



## 6. Fluid mechanics: fluid statics; fluid dynamics (internal flows, external flows)

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### 6.1 Fluid statics



# Fluid statics, static pressure /1

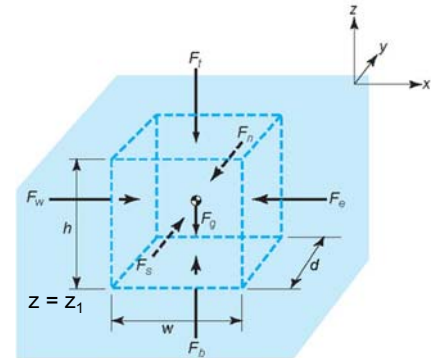
- In engineering applications, a **fluid** (sv: *fluid*) is a **liquid or a gas**
- The behaviour of **stationary fluids** is described by **fluid statics**
- A **liquid** in a container forms a layer with a distinct surface, and exerts forces on the walls supporting it, while a **gas** will fill the whole container.

- **Two types of forces** act on a fluid volume element:

**surface (pressure) forces** and **body (gravitational) forces**: see Figure →

- **Pressure** (a scalar!) is defined as **surface force / area**, for example

$$p_b = F_b / (d \cdot w) = p @ z = z_1$$



Fluid volume  $h \cdot d \cdot w$   
with density  $\rho$  and  
mass  $m = h \cdot d \cdot w \cdot \rho$

Picture: KJ05

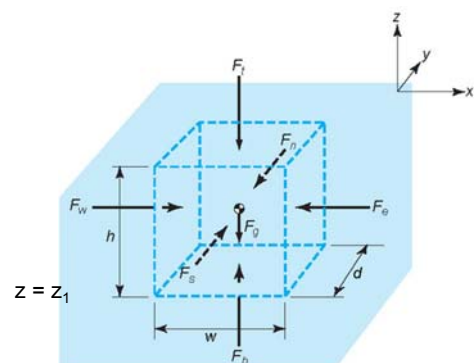
# Fluid statics, static pressure /2

- For the horizontal forces  $F_n + F_s = 0$  or  
 $-p_y \cdot h \cdot w + p_y \cdot h \cdot w = 0 \rightarrow p_y = 0$
- Similarly  $F_w + F_e = 0$  gives  $p_x = 0$ ,
- There are three vertical forces:  
 $-F_t \cdot h \cdot d - m \cdot g + F_b \cdot h \cdot d = 0$  (gravity  $g$ )
- The **pressure difference** between  
 $z = z_1$  and  $z = z_1 + h$  follows from  
 $-F_t - \rho \cdot h \cdot d \cdot w \cdot g = -F_b$ , with  
 $-F_b / (d \cdot w) = -p_z @ z = z_1$ ; and  
 $F_t / (d \cdot w) = -p_z @ z = z_1 + h$ ; gives

$$p_z(z_1) = p_z(z_1 + h) + \rho \cdot h \cdot g$$

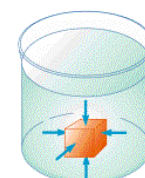
- If  $z = z_1 + h$  is at the fluid surface exposed to atmospheric pressure  $p^0$  then

$$p_z(z_1) = p^0 + \rho \cdot h \cdot g$$



Fluid volume  $h \cdot d \cdot w$   
with density  $\rho$  and  
mass  $m = h \cdot d \cdot w \cdot \rho$

Picture:  
KJ05



# U-tube manometer

- The **U-tube manometer** is based on the relation between depth and pressure in static fluids, with one end open to the atmosphere at  $p_{\text{atm}}$
- For the Figure, with gravity  $g$  and densities  $\rho_g$  and  $\rho_l$  for gas and liquid:

$$p_C = \rho_g \cdot h_1 \cdot g + p_B$$

$$p_D = \rho_l \cdot h_2 \cdot g + p_C = \rho_l \cdot h_2 \cdot g + \rho_g \cdot h_1 \cdot g + p_B$$

and also, from the other side

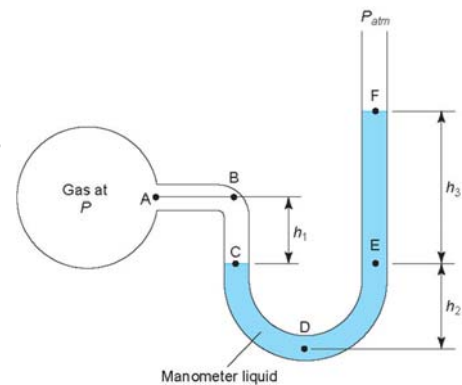
$$p_D = \rho_l \cdot (h_3 + h_2) \cdot g + p_F = \rho_l \cdot (h_3 + h_2) \cdot g + p_{\text{atm}}$$

which gives, with  $p_B = p_A$

$$\rho_l \cdot h_2 \cdot g + \rho_g \cdot h_1 \cdot g + p_A = \rho_l \cdot (h_3 + h_2) \cdot g + p_{\text{atm}}$$

$$p_A - p_{\text{atm}} = \rho_l \cdot h_3 \cdot g - \rho_g \cdot h_1 \cdot g$$

and noting that  $\rho_l \gg \rho_g$  :  $p_A - p_{\text{atm}} = \rho_l \cdot h_3 \cdot g$



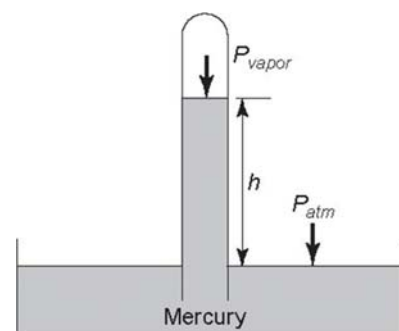
Note that the U-tube manometer measures **pressure differences**

Picture: KJ05

# Barometer

- A device for measuring atmospheric pressure (which cannot be done using an U-tube manometer) is referred to as **barometer**
- A closed tube filled with mercury (Hg) is quickly put upside-down in an open container filled with Hg
- Gravity causes the Hg level in the tube to fall, but no air can enter the tube. The small gas volume trapped is Hg vapour at equilibrium with liquid Hg.
- For the tube  $p_{\text{vapor,Hg}} + \rho_{\text{Hg}} \cdot h_{\text{Hg}} \cdot g = p_{\text{atm}}$
- At 20°C,  $p_{\text{vapor,Hg}} = 0.158 \text{ Pa} \ll p_{\text{atm}}$ , thus

$$p_{\text{atm}} \approx \rho_{\text{Hg}} \cdot h_{\text{Hg}} \cdot g$$



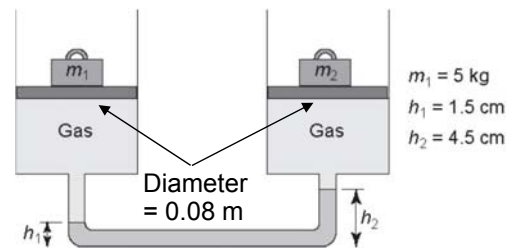
the density of liquid Hg is  
13546.2 kg/m³ at 20°C

after Torricelli:  
1 torr = 1 mm Hg pressure  
1 atm = 760 torr at 0°C

Picture: KJ05

## Example: a manometer

- Two piston-cylinder assemblies are connected by a tube filled with mercury (Hg) at 20°C (density 13546 kg/m<sup>3</sup>)
- The diameter of each piston is 0.08 m, the mass of each piston is 0.40 kg. Mass  $m_1 = 5.00$  kg
- Use the data to calculate mass  $m_2$ .

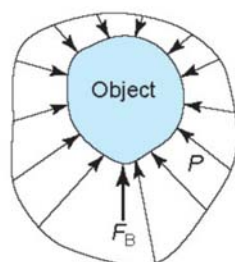


Picture: KJ05

## Buoyancy //

- Buoyancy** (sv: flytkraft, fi: nostovoima) or **buoyant force** acts on all objects immersed or submerged (sv: sänkad) in a fluid
- It is an overall **upwards force** as the result of the fact that pressure  $p$  in a static fluid increases with depth

surface \_\_\_\_\_



Picture: KJ05

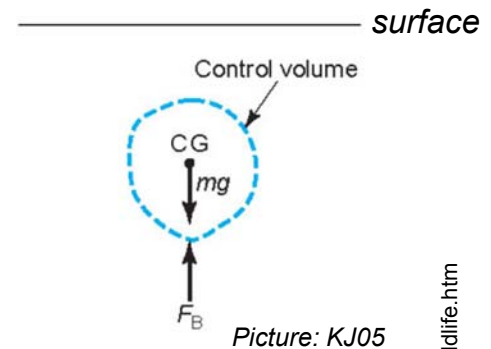


Picture: <http://www.math.nyu.edu/~corres/Archimedes/Crown/Vitruvius.html>  
Picture: [http://www.physics.lsa.umich.edu/demolab/graphics/2b40\\_u2.jpg](http://www.physics.lsa.umich.edu/demolab/graphics/2b40_u2.jpg)

## Buoyancy /2

- For an immersed object, horizontal forces cancel each other, and the two vertical forces are gravity and buoyancy.
- The forces on the surface of the object are the same as when that surface would be filled with the fluid
- Thus, the **buoyant force** on a mass with volume  $V$  is equal (but opposite in sign) to the weight of the fluid in the volume  $V$ , and acts on the same **centre of gravity (CG)**:

$$F_B = - m_{\text{fluid}} \cdot g = - \rho_{\text{fluid}} \cdot V \cdot g$$



## Buoyancy /3

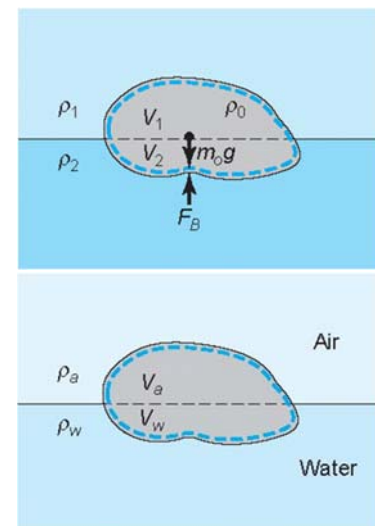
- For any object the buoyancy force it experiences may be less than, equal to or larger than its weight
- If  $F_B > \text{weight}$ , the object will rise / float  
If  $F_B < \text{weight}$ , the object will sink  
If  $F_B = \text{weight}$ , the will float in suspension
- For example, for the two fluids geometry  $\rightarrow$   

$$F_B = (\rho_1 \cdot V_1 + \rho_2 \cdot V_2) \cdot g$$
 in equilibrium with  $F_{\text{gravity}} = m_0 \cdot g = \rho_0 \cdot V_{\text{tot}} \cdot g$   
 for object mass  $m_0$  (kg).  

$$\rightarrow \rho_0 \cdot V_{\text{tot}} = \rho_1 \cdot V_1 + \rho_2 \cdot V_2 \quad \text{and} \quad V_{\text{tot}} = V_1 + V_2$$
- For example, for cases with water + air  $\rightarrow$   

$$F_B = (\rho_a \cdot V_a + \rho_w \cdot V_w) \cdot g \approx \rho_w \cdot V_w \cdot g \quad (\rho_a \gg \rho_w)$$

$$\rightarrow \rho_0 \cdot V_{\text{tot}} = \rho_w \cdot V_w, \quad \text{or} : \rho_0 / \rho_w = V_w / V_{\text{tot}}$$



Pictures: KJ05

## Example: buoyancy

- The tip of a certain iceberg (which is the volume of the iceberg above the water surface) is  $V_{\text{tip}} = 79 \text{ m}^3$ , in seawater of with density  $\rho_{\text{sea}} = 1027 \text{ kg/m}^3$ . Calculate the submerged (i.e. under water) volume of the iceberg. For ice the density is  $\rho_{\text{ice}} = 920 \text{ kg/m}^3$ .



picture: <http://www.sgisland.org/pages/zone/download.htm>

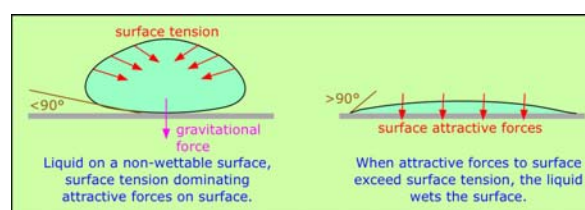
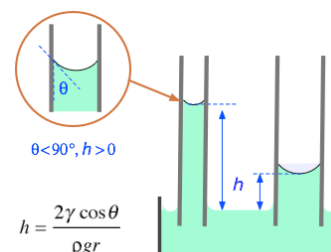
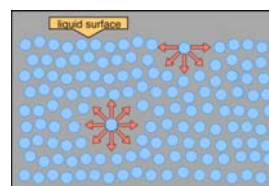
Source: KJ05

## Surface tension

- A liquid at a material interface, usually liquid-gas, exerts **a force  $F_{\text{int}}$  per unit length  $L$  along the surface.**
- It is the result of molecular attraction at a liquid surface being different from that "in" the liquid  $\rightarrow$  the surface acts like a stretched membrane
- Surface tension** ( $\sigma$  or  $\gamma$ , unit: N/m) **quantifies this force:**

$$F_{\text{int}} = \gamma \cdot L$$

- Result phenomena:
  - Contact angle
  - Capillary action (rise or drop)
  - Bubbles, droplets



For ambient water-air:  $\gamma = 0.073 \text{ N/m}$

<http://www.chem1.com/acad/sci/aboutwater.html>  
<http://www.chem1.com/acad/webtext/states/state-images/caprise1.png>  
[http://3.bp.blogspot.com/\\_ADtjVgTw6-Q/TJE\\_2usgl/AAAAAAAAAAC/RFSv-ttCeuc/s1600/surface\\_tension\\_c\\_ph\\_784.jpg](http://3.bp.blogspot.com/_ADtjVgTw6-Q/TJE_2usgl/AAAAAAAAAAC/RFSv-ttCeuc/s1600/surface_tension_c_ph_784.jpg)

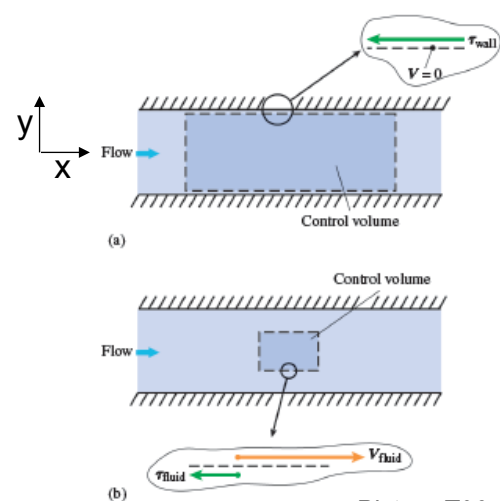


## 6.2 Fluid dynamics: viscosity, laminar, turbulent flow, boundary layer



### Internal friction in fluid flow //

- Fluids will (try to) resist a change in shape, as will occur in fluid flow situations where different fluid elements have different velocities
- Note the **definition** of a fluid:  
**a fluid is a substance that deforms continuously under the application of a shear stress** (sv: skjuvspänning)
- Consider fluid flow between plates:
  - The **no-slip condition** says that at the wall the velocity of the fluid is the same as the wall velocity \*), for a fixed wall  $v_{\text{fluid}} = 0$  at the wall
  - Between the plates a **velocity profile** exists: it can be described as  $v_x = v_x(y)$
  - Shear stresses,  $\tau_{\text{fluid}}$ , arise due to velocity differences between different fluid elements



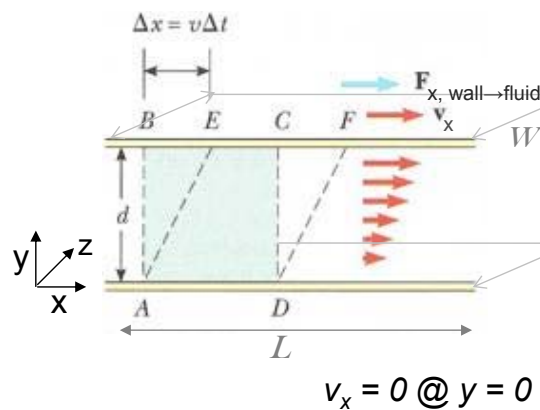
Picture T06

\*) this applies always except for very low pressure gases, for example in the upper atmosphere



## Internal friction in fluid flow /2

- For a fluid between plates with width  $W$  (m), distance  $d$  (m) the shear force  $\underline{F} = (F_x, F_y, F_z) = (F_x, 0, 0)$  (unit: N) to pull the fluid at velocity  $\underline{v} = (v_x, v_y, v_z) = (v_x, 0, 0)$  gives a **shear stress**  $\tau_{yx}$  (unit: N/m<sup>2</sup>) in the fluid at  $y = d$  that is equal to:



$$\frac{F_{x, \text{fluid} \rightarrow \text{wall}}}{\text{surface}} = \frac{-F_{x, \text{wall} \rightarrow \text{fluid}}}{W \cdot L} = \tau_{yx} \Big|_{y=d} = -\eta \frac{dv_x}{dy} \approx -\eta \frac{\Delta v_x}{\Delta y}$$

with  $\tau_{yx}$  as stress in direction "x" in a plane for constant "y"

- This defines the **dynamic viscosity**  $\eta$  (unit: Pa.s = kg.m<sup>-1</sup>.s<sup>-1</sup>)
- ! Note:**  $\tau_{yx}$  at  $y = y_0$  is the shear stress of fluid elements with  $y < y_0$  on the fluid elements with  $y > y_0$ . As a result  $F_x > 0$  if  $dv_x/dy < 0$ !

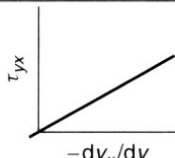
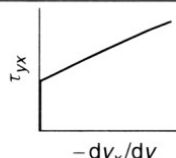
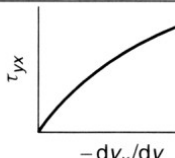
**SIGN:**

Picture: <http://www.physics.uc.edu/~sitko/CollegePhysicsIII/9-Solids&Fluids&Fluids.htm>

## Internal friction in fluid flow /3

- The linear relation between  $\tau_{yx}$  and  $dv_x/dy$  is referred to as **Newton's Law** which holds for so-called **Newtonian fluids**
- For **non-Newtonian fluids**, other relations between shear force and velocity gradient hold, for example Bingham fluids (toothpaste, clay) or pseudo-plastic (Ostwald) fluids (blood, yoghurt). For those, *viscosity is a function of the velocity gradient*:  $\tau_{yx} = \eta(dv_x/dy) \cdot dv_x/dy$

**Newtonian vs non-Newtonian fluids**

		
Newton	Bingham	Ostwald
$\tau_{yx} = -\eta \frac{dv_x}{dy}$	$\tau_{yx} = \pm \tau_0 - \eta \frac{dv_x}{dy}$	$\tau_{yx} = -K \left( \frac{dv_x}{dy} \right)^{n-1} \frac{dv_x}{dy}$

Picture: BMH99

- Note:** The flow of a fluid between plates, or in a tube or on a surface doesn't necessarily require moving walls:

**usually the driving force is gravity, or a static pressure difference**

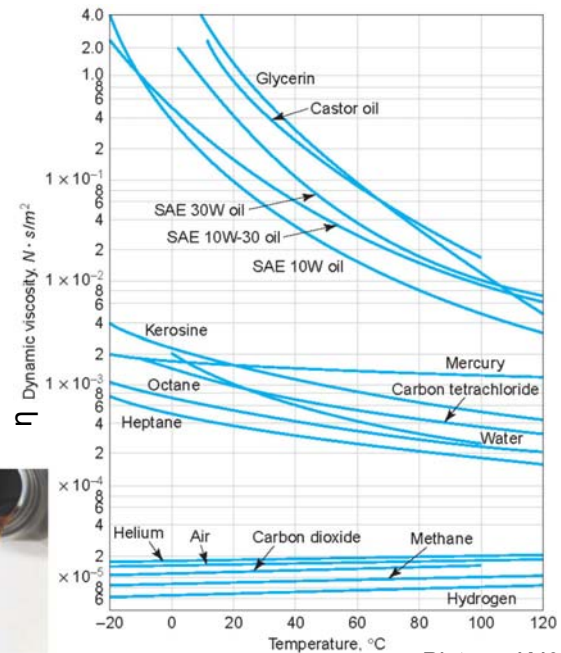


# Viscosity

- **Viscosity** (sv: viskositet) is a measure of a fluid's resistance to flow; it describes the internal friction of a moving fluid.
- More specifically, it defines the **rate of momentum transfer in a fluid as a result of a velocity gradient**.
- **Dynamic viscosity**  $\eta$  (unit: Pa.s) is related to a **kinematic viscosity**,  $\nu$  (unit: m<sup>2</sup>/s) via fluid density  $\rho$  (kg/m<sup>3</sup>) as:  $\nu = \eta/\rho$



Picture T06



Picture: KJ05

## Internal friction in fluid flow /5

- Concentration,  $c$ , temperature,  $T$ , and energy,  $E$ , are scalars, and their gradient is a **vector** such as  $dT/dx$  or  $\nabla T = (\partial T/\partial x, \partial T/\partial y, \partial T/\partial z)$ , etc.
- Velocity is a vector  $\underline{v}$ , for example  $\underline{v} = (v_x, v_y, v_z)$  and it's gradient is a (second order) **tensor** with elements such as  $dv_x/dy$  (gradient of  $v_x$  in  $y$ -direction)

$$\nabla \underline{v} = \begin{pmatrix} \frac{\partial v_x}{\partial x} & \frac{\partial v_y}{\partial x} & \frac{\partial v_z}{\partial x} \\ \frac{\partial v_x}{\partial y} & \frac{\partial v_y}{\partial y} & \frac{\partial v_z}{\partial y} \\ \frac{\partial v_x}{\partial z} & \frac{\partial v_y}{\partial z} & \frac{\partial v_z}{\partial z} \end{pmatrix}$$

note :

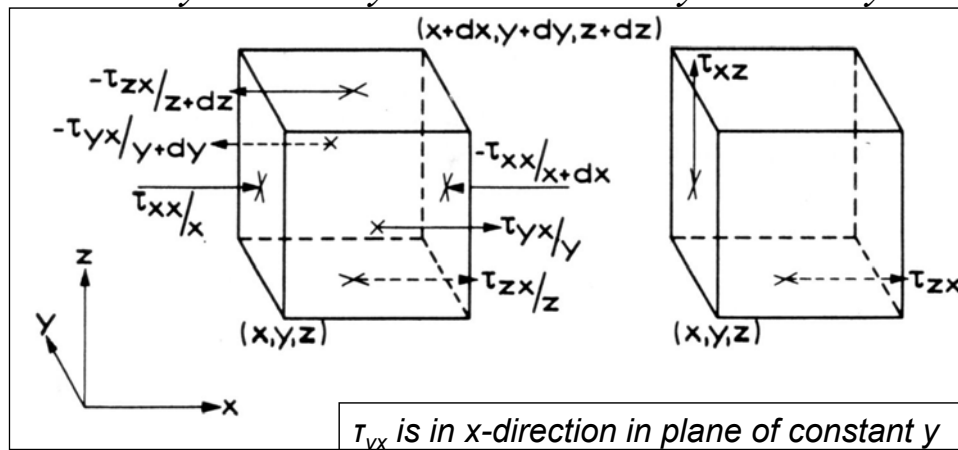
$$\nabla \cdot \underline{v} = \left( \frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} + \frac{\partial v_z}{\partial z} \right)$$

Gradients of a scalar property give a vector (or 1<sup>st</sup> order tensor);  
gradients of a vector property give a 2<sup>nd</sup> order tensor, etc.

# Internal friction in fluid flow 16

- $\nabla \underline{v}$  results in **3 compressive stresses** (sv: tryckspänningar)  $\tau_{xx}$ ,  $\tau_{yy}$  and  $\tau_{zz}$  and **6 shear stresses** (sv: skjuvspänningar)  $\tau_{xy}$ ,  $\tau_{xz}$ ,  $\tau_{yz}$ ,  $\tau_{zx}$ ,  $\tau_{yx}$  and  $\tau_{zy}$ :

$$\tau_{yx} = -\eta \frac{dv_x}{dy} = -\nu \frac{d\rho v_x}{dy}; \quad \tau_{yz} = -\nu \frac{dv_z}{dy} = -\nu \frac{d\rho v_z}{dy}; \quad \text{etc.}$$



Picture: SSJ84

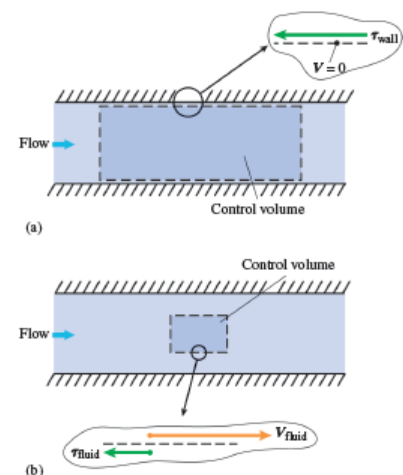
## Viscous work

Vector/tensor calculations like this are beyond this course

- The shear stresses can be expressed as **tensor**  $\underline{\underline{\tau}}$ , resulting in a viscous shear force on a certain area  $A$  that is equal to  $\underline{F}_{\text{visc}} = \underline{\underline{\tau}} \cdot \underline{A}$ , with  $\underline{A} = A \underline{n}$  with normal vector  $\underline{n}$
- If the velocity  $\underline{v}$  at surface  $\underline{A}$  the rate of **viscous work** done by the fluid at surface  $A$  equals  $\dot{W}_{\text{visc}} = \underline{F}_{\text{visc}} \cdot \underline{v} = \underline{\underline{\tau}} \cdot \underline{A} \cdot \underline{v}$ , which for a certain volume element of control volume (inside which  $\underline{v}$  and  $\underline{\underline{\tau}}$  can vary) with total outside surface  $A$  gives the rate of work done:

$$\dot{W}_{\text{visc}} = \int_A (\underline{\underline{\tau}} \cdot \underline{v}) \cdot d\underline{A}$$

- Note: at the wall  $\underline{v} = 0$  so no work is done; also at points where velocity and shear are perpendicular  $\underline{\underline{\tau}} \cdot \underline{v} = 0$  and no work is done.

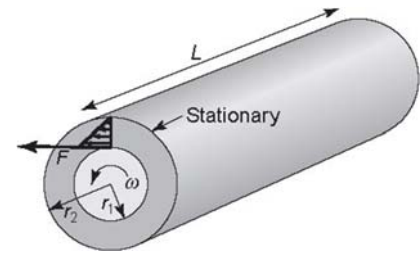


Picture T06

**The friction work is dissipated as HEAT**

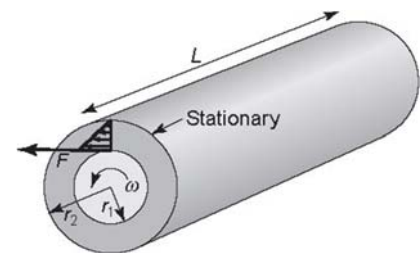
## Example: shear stress concentric cylinders / I

- Oil with viscosity  $\eta = 0.05 \text{ Pa}\cdot\text{s}$  fills a 0.4 mm gap between two cylinders of which the inner one rotates whilst the outer one is fixed.
- The diameter of the inner cylinder is 8 cm, the length is 20 cm.
- Question: How much power is required to rotate the inner cylinder at 300 rpm?



Picture: KJ05  
Question ÖS96-4.1

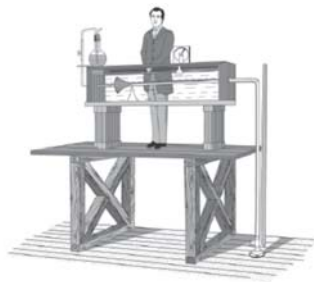
## Example: shear stress concentric cylinders / 2



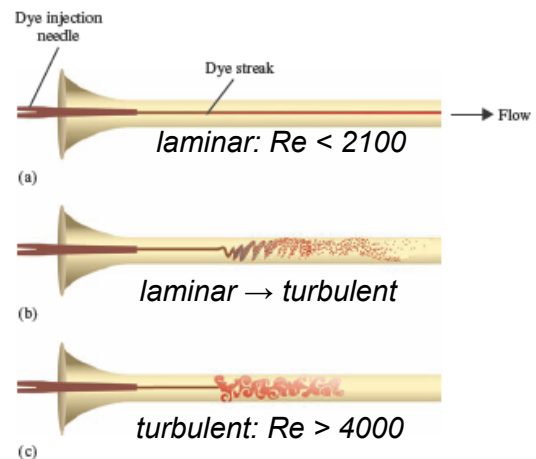
\*) The space between the two cylinders is very small and may be treated as a flat plate

Picture: KJ05  
Question ÖS96-4.1

# Laminar ↔ turbulent fluid flow



Osborne Reynolds's dye-streak experiment (1883) for measuring laminar → turbulent flow transition



Pictures: T06

- For circular tube flow, the laminar → turbulent flow transition occurs at **Reynolds number Re 2100 - 2300**, with the **dimensionless number** defined as  $Re = \rho \langle v \rangle \cdot d / \eta$  for  $\rho$  = fluid's density ( $\text{kg/m}^3$ ),  $\langle v \rangle$  = fluid's **average** velocity ( $\text{m/s}$ ),  $d$  = tube diameter ( $\text{m}$ ) and  $\eta$  = fluid's dynamic viscosity ( $\text{Pa} \cdot \text{s}$ )



## Example: a liquid film on a vertical wall / I

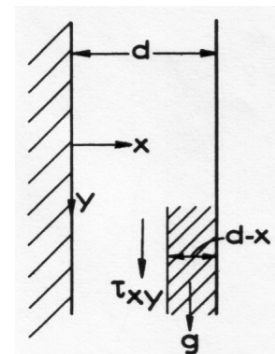
- A stationary laminar flow of water (at 1200 kg/h) runs down a vertical surface (with width  $W = 1 \text{ m}$ ).

Give

- the expression for the shear stress distribution,
- the expression for the velocity profile, and
- the expression for volumetric flow rate  $V$  ( $\text{m}^3/\text{s}$ )

and calculate

- film thickness  $d$
- velocity  $\langle v_y \rangle$  averaged over the film thickness
- maximum velocity  $v_{y,\text{max}}$



Data: dynamic viscosity for water  $\eta = 10^{-3} \text{ Pa} \cdot \text{s}$   
density for water  $\rho = 1000 \text{ kg/m}^3$   
gravity  $g = 9.8 \text{ m/s}^2$

Source: SSJ84



## Example: a liquid film on a vertical wall /2

Answer: For this steady-state process:

- The vertical force balance for a volume element with length  $dy$  as shown gives  $F_{\text{gravity}} = F_{\text{shear}}$

$$\rho \cdot (d-x) \cdot W \cdot dy \cdot g + \tau_{xy} \cdot W \cdot dy = 0 \Rightarrow \rho \cdot (d-x) \cdot g + \tau_{xy} = 0$$

with  $\tau_{xy} = -\eta \frac{dv_y}{dx} = -\rho \cdot (d-x) \cdot g \Rightarrow \frac{dv_y}{dx} = \frac{\rho \cdot (d-x)g}{\eta}$ , integrating :

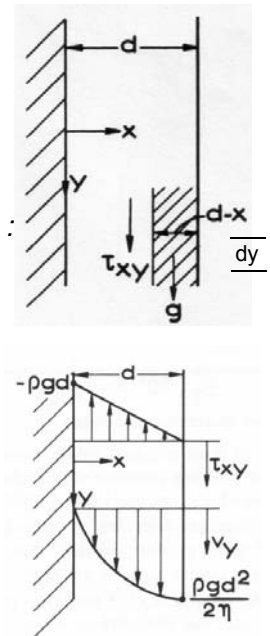
$$v_y(x) = \int_0^x \frac{dv_y}{dx} dx = \int_0^x \frac{\rho \cdot (d-x)g}{\eta} dx = \frac{\rho \cdot g}{\eta} \cdot (xd - \frac{1}{2}x^2)$$

with  $v_y = v_{y,\text{max}}$  @  $x = d$ :  $v_{y,\text{max}} = \frac{1}{2} \rho g d^2 / \eta$

For the average velocity  $\langle v \rangle$  with  $\dot{V} = \langle v \rangle \cdot d \cdot W$ :

$$\langle v_y \rangle = \frac{1}{d} \int_0^d v_y(x) \cdot dx = \frac{1}{d} \int_0^d \frac{\rho g}{\eta} \cdot (xd - \frac{1}{2}x^2) \cdot dx = \frac{\rho g d^2}{3\eta}$$

and  $\langle v_y \rangle = \frac{\dot{V}}{W \cdot d}$  gives  $d = \sqrt[3]{\frac{3\eta \dot{V}}{\rho g}}$



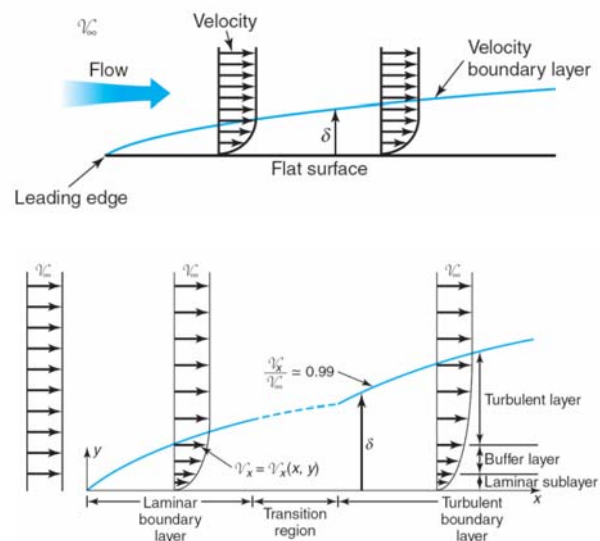
Source: SSJ84

The data gives:  $d = 0.47 \text{ mm}$ ,  $\langle v_y \rangle = 0.71 \text{ m/s}$ ;  $v_{y,\text{max}} = 1.07 \text{ m/s}$

## Boundary layers

- At the interface of a surface\* and a flowing medium, a **thin** ( $\sim 0.01 - 1 \text{ mm}$ ) layer of fluid is created in which the velocity increases from  $v = 0$  at the interface to the free-flow velocity  $v = v_\infty$  (or  $0.99 \cdot v_\infty$ )
- In this **boundary layer** (sv: gränsskikt) all the **thermal and/or viscous effects** of the surface are concentrated
- The boundary layer can develop from **laminar to turbulent** flow

\* This can be a solid surface or another flowing medium



Growth of the velocity boundary layer on a flat surface.

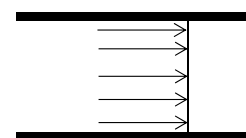
Pictures: KJ05

## 6.3 Fluid dynamics: internal flows / tube flow

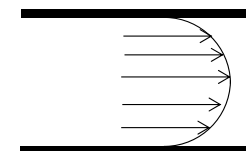


### Internal flows; velocity profiles

- Fluid flow in a tube or other confinement (*sv: inspärning*) will show:
  - zero velocity (the no-slip condition) at the walls; and
  - maximum velocity furthest from the walls (i.e. at a tube flow centre line or at a free surface)
- The velocity profile is the result of viscous friction, and for turbulent flow, "**eddy**" currents (→ so-called "eddy viscosity":  $\eta = \eta_{\text{viscous}} + \eta_{\text{eddy}}$ )
- In many applications a plug flow idealisation may be used described by an **average velocity**  $\langle v \rangle$



*Plug flow idealisation*



*Velocity profile due to viscous friction*



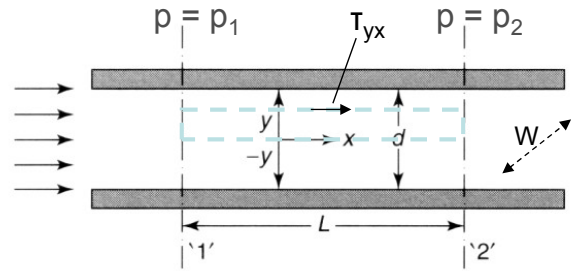
*Velocity profile due to turbulent "eddies"*





## Laminar flow between two plates /1

- For a steady-state fluid flow between two stagnant parallel plates, the forces on a volume element between point "1" and "2" and between  $y =$  centre line and  $y = y$  are (for plate width  $W$ ) :



Picture: BMH99

@ "1" pressure force  $= p_1 \cdot y \cdot W$  ; @ "2" pressure force  $= -p_2 \cdot y \cdot W$   
 shear force on volume element  $= -\tau_{yx} \cdot L \cdot W$

The force balance gives  $p_1 \cdot y - p_2 \cdot y - \tau_{yx} \cdot L = 0 \Rightarrow \tau_{yx} = \frac{p_1 - p_2}{L} \cdot y$

With  $\tau_{yx} = -\eta \cdot \frac{dv_x}{dy} \Rightarrow \frac{dv_x}{dy} = -\frac{p_1 - p_2}{\eta \cdot L} \cdot y$  with  $v_x = 0$  @  $y = \pm \frac{1}{2}d$

$\tau_{yx}$  acts on fluid  $y > y$ , so  $-\tau_{yx}$  acts on fluid  $y < y$  which is the fluid element

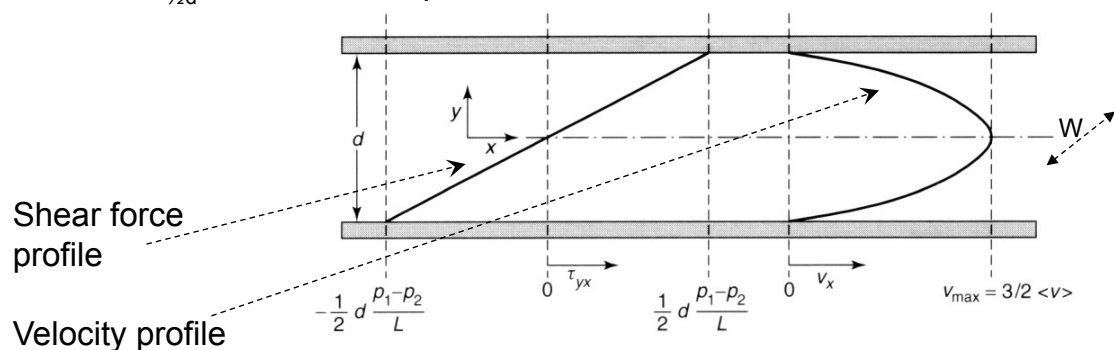
## Laminar flow between two plates /2

Calculation of the velocity profile and maximum velocity :

Integration :  $v_x = \frac{p_1 - p_2}{2 \cdot \eta \cdot L} \left( \frac{d^2}{4} - y^2 \right)$  and  $v_{x, \max} = \frac{p_1 - p_2}{8 \cdot \eta \cdot L} d^2$  @  $y = 0$

Calculation of the flow rate  $\dot{V} (m^3/s)$  :

$\dot{V} = W \cdot \int_{-\frac{1}{2}d}^{\frac{1}{2}d} v_x dy = W \cdot \frac{p_1 - p_2}{12 \cdot \eta \cdot L} d^3 = W \cdot d \cdot \langle v_x \rangle$  ,  $\Rightarrow \langle v_x \rangle = \frac{2}{3} v_{x, \max}$



Picture: BMH99

# Stationary laminar tube flow

@ "1" pressure force =  $p_1 \cdot \pi \cdot r^2$  ;

@ "2" pressure force =  $-p_2 \cdot \pi \cdot r^2$

shear force on volume element =  $-\tau_{rx} \cdot (x_2 - x_1) \cdot 2\pi \cdot r$

Force balance :  $p_1 \cdot \pi \cdot r^2 - p_2 \cdot \pi \cdot r^2 - \tau_{rx} \cdot (x_2 - x_1) \cdot 2\pi \cdot r = 0$

$$\Rightarrow \tau_{rx} = \frac{p_1 - p_2}{2 \cdot (x_2 - x_1)} \cdot r = \frac{1}{2} r \cdot \left( -\frac{dp}{dx} \right)$$

With  $\tau_{rx} = -\eta \cdot \frac{dv_x}{dr} \Rightarrow \frac{dv_x}{dr} = \frac{-1/2 r}{\eta} \cdot \left( -\frac{dp}{dx} \right)$  with  $v_x = 0$  @  $r = R$

Velocity profile and maximum velocity :

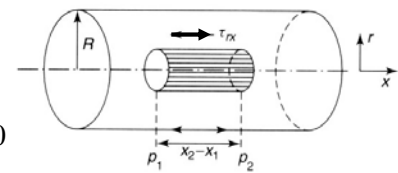
Integration :  $v_x(r) = \frac{1}{4 \cdot \eta} \cdot \left( -\frac{dp}{dx} \right) \cdot (R^2 - r^2)$  and

$$v_{x,max} = \frac{1}{4 \cdot \eta} \cdot \left( -\frac{dp}{dx} \right) \cdot R^2 \text{ @ } r = 0$$

Calculation of the flow rate  $\dot{V}(\text{m}^3/\text{s})$  :

$$\dot{V} = \int_0^R 2\pi r \cdot v_x dr = \frac{\pi R^4}{8 \cdot \eta} \cdot \left( -\frac{dp}{dx} \right) = \pi R^2 \cdot \langle v_x \rangle, \Rightarrow \langle v_x \rangle = \frac{1}{2} v_{x,max}$$

"Hagen-Poiseuille relationship"



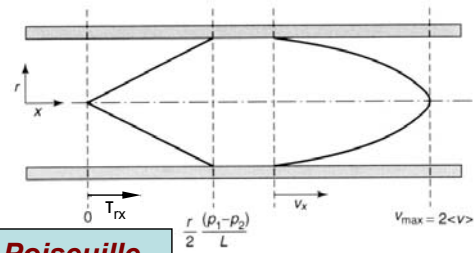
at the centre line :

$$\tau_{rx} = 0$$

at the wall :

$$\tau_{rx} = \frac{1}{2} R \cdot \left( -\frac{dp}{dx} \right) = \frac{4 \cdot \langle v_x \rangle \cdot \eta}{R}$$

$$= \tau_{\text{fluid} \rightarrow \text{wall}} = -\tau_{\text{wall} \rightarrow \text{fluid}}$$



Pictures: BMH99

## Tube flow velocity profiles

- **Laminar and turbulent tube flows** show different velocity profiles

Laminar:

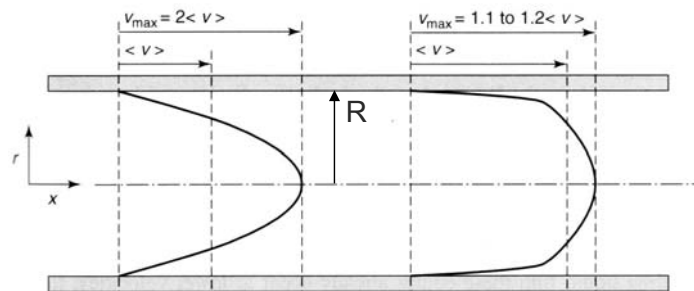
$$v_x(r) = (1 - r^2/R^2) \cdot v_{max}$$

cross-sectional average velocity  $\langle v \rangle = \frac{1}{2} \cdot v_{max}$

Turbulent:

$$v_x(r) \approx (1 - r/R)^{1/7} \cdot v_{max}$$

cross-sectional average velocity  $\langle v \rangle = 0.875 \cdot v_{max}$



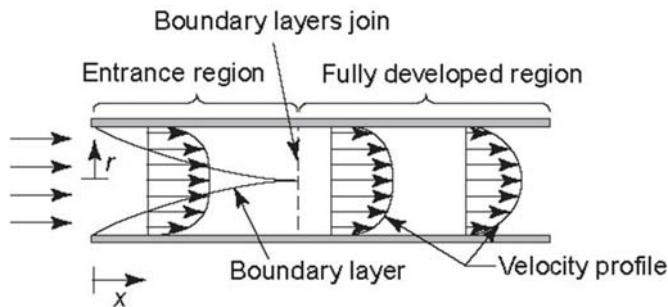
- The **cross-sectional average velocity  $\langle v \rangle$**  is used in dimensional analysis or the resulting dimensionless groups (Re, and others)

$$\langle v \rangle = \frac{\int_0^R v_x(r) \cdot 2 \cdot \pi \cdot r \cdot dr}{\pi R^2} = \frac{\dot{V}(\text{m}^3/\text{s})}{\pi R^2 (\text{m}^2)}$$

Picture: BMH99

# Tube flow entrance region

- Flow entering a tube requires a certain distance to produce a "developed flow" with a constant boundary layer: **the entrance region**
- For the entrance region in **laminar tube flow**, the **Graetz number** quantifies for the boundary layer build-up (see also section 5.2 - Convective heat transfer)



Picture: KJ05

- The entrance length  $L_{\text{ent}}$  for a hydrodynamically developed tube flow (tube diameter  $D$ ) is

$$L_{\text{ent}} \approx 0.065 \cdot \text{Re} \cdot D$$

for laminar flow  $\text{Re} < 2100$

$$L_{\text{ent}} \approx 4.4 \cdot \text{Re}^{1/6} \cdot D$$

for turbulent flow  $\text{Re} > 4000$

## 6.4 Fluid dynamics: pressure drop & energy dissipation in tube systems

## Tube systems //

- In a tube system, **pressure drop losses** resulting from fluid internal friction and wall friction in straight and curved tube sections, valves, inlet/outlet sections, diameter changes etc. etc. **must be compensated for by adding mechanical energy** via pumps, compressors, turbines, ventilators (sv: *pumpar, kompressorer, turbiner, fläktar*) etc.
- Additional effects that must be compensated for are **kinetic energy** (if flow velocities change) and **potential energy** (for non-horizontal tube sections)



Picture: <http://www.raneng.com/Arco%20EH.htm>



## Tube systems /2

- For a flow tube system from point "1" at height  $z_1$ , average velocity  $\langle v \rangle_1$ , pressure  $p_1$ , volume flow  $\dot{V}_1$ , to point "2" at height  $z_2$ , velocity  $\langle v \rangle_2$ , pressure  $p_2$ , volume flow  $\dot{V}_2$ , **pumping power** (sv: *pumpeffekt*)  
 $P_{\text{pump}}$  compensates for flow friction losses  $P_{\text{losses}}$  :



General energy balance with heat input  $\dot{Q}$ ,

work input  $\dot{W}$ , potential and kinetic energy and "flow work" :

$$\dot{m}_1 \cdot (u_1 + gz_1 + \frac{1}{2} \langle v \rangle_1^2) + p_1 \dot{V}_1 + \dot{Q} + \dot{W} = \dot{m}_2 \cdot (u_2 + gz_2 + \frac{1}{2} \langle v \rangle_2^2) + p_2 \dot{V}_2$$

For isothermal flows, no heat effect ( $\dot{Q} = 0$ ), no work ( $\dot{W} = 0$ ) :

$$\dot{m}_1 \cdot (gz_1 + \frac{1}{2} \langle v \rangle_1^2) + p_1 \dot{V}_1 = \dot{m}_2 \cdot (gz_2 + \frac{1}{2} \langle v \rangle_2^2) + p_2 \dot{V}_2$$

With work input to compensate for flow friction losses  $P_{\text{losses}}$

for example  $\dot{W} = P_{\text{pump}} = P_{\text{losses}}$  ( $= -\dot{Q}$ , but assuming  $\dot{Q} \approx 0$ ) :

$$\dot{m}_1 \cdot (gz_1 + \frac{1}{2} \langle v \rangle_1^2) + p_1 \dot{V}_1 + P_{\text{pump}} = \dot{m}_2 \cdot (gz_2 + \frac{1}{2} \langle v \rangle_2^2) + p_2 \dot{V}_2 + P_{\text{losses}}$$

Picture: [http://www.pharmaceutical-technology.com/contractors/water\\_treatment/fischer/fischer2.html](http://www.pharmaceutical-technology.com/contractors/water_treatment/fischer/fischer2.html)



## Tube systems /3

- Flow through pipes and conduits (sv: rör, ledning, kanal) with height  $z_1$ , velocity  $v_1$ , pressure  $p_1$ , volume flow  $\dot{V}_1 \rightarrow$  height  $z_2$ , velocity  $v_2$ , pressure  $p_2$ , volume flow  $\dot{V}_2$



$$\dot{m}_1 \cdot (u_1 + gz_1 + \frac{1}{2} \langle v \rangle_1^2) + p_1 \dot{V}_1 + P_{pump} = \dot{m}_2 \cdot (u_2 + gz_2 + \frac{1}{2} \langle v \rangle_2^2) + p_2 \dot{V}_2 + P_{losses}$$

Special case 1: for an isothermal inviscid fluid (negligible viscosity),

$\rightarrow P_{pump} = P_{losses} \approx 0$ ; this gives Bernoulli's equation:

$$\dot{m}_1 \cdot (gz_1 + \frac{1}{2} \langle v \rangle_1^2) + p_1 \dot{V}_1 = \dot{m}_2 \cdot (gz_2 + \frac{1}{2} \langle v \rangle_2^2) + p_2 \dot{V}_2$$

Picture: <http://www.chemicals-technology.com/contractors/pipes/fischer/fischer2.html>

## Tube systems /4

- Flow through pipes and conduits (sv: rör, ledning, kanal) with height  $z_1$ , velocity  $v_1$ , pressure  $p_1$ , volume flow  $\dot{V}_1 \rightarrow$  height  $z_2$ , velocity  $v_2$ , pressure  $p_2$ , volume flow  $\dot{V}_2$



$$\dot{m}_1 \cdot (u_1 + gz_1 + \frac{1}{2} \langle v \rangle_1^2) + p_1 \dot{V}_1 + P_{pump} = \dot{m}_2 \cdot (u_2 + gz_2 + \frac{1}{2} \langle v \rangle_2^2) + p_2 \dot{V}_2 + P_{losses}$$

Special case 2: correcting for velocity profiles in stream cross - section:

$$\dot{m}_1 \cdot (gz_1 + \frac{1}{2} \xi_1 \langle v \rangle_1^2) + p_1 \dot{V}_1 + P_{pump} = \dot{m}_2 \cdot (gz_2 + \frac{1}{2} \xi_2 \langle v \rangle_2^2) + p_2 \dot{V}_2 + P_{losses}$$

with kinetic energy correction factor  $\xi$ , for stream cross – sectional area A:

$$\xi = \frac{\dot{E}_{kinetic}}{\frac{1}{2} \cdot \dot{m} \cdot \langle v \rangle^2} = \frac{\int \frac{1}{2} \cdot \dot{m} \cdot v^2 dA}{\frac{1}{2} \cdot \rho \cdot A \cdot \langle v \rangle^3} = \frac{\frac{1}{2} \cdot \rho \int v^3 dA}{\frac{1}{2} \cdot \rho \cdot A \cdot \langle v \rangle^3} = \frac{\frac{1}{A} \int v^3 dA}{\langle v \rangle^3}$$

$\xi \approx 2$  for laminar flows, and  $\xi \approx 1.05 - 1.10$  for turbulent flows

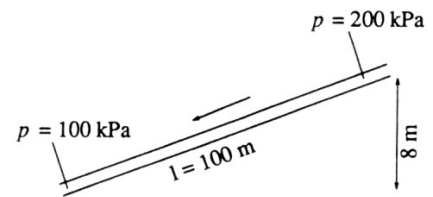
See  
also  
CEWR10  
p. 222

Picture: <http://www.chemicals-technology.com/contractors/pipes/fischer/fischer2.html>



## Example: friction losses (ÖS96-4.6)

- 1 liter/s ethanol (density  $\rho = 791 \text{ kg/m}^3$ ) is pumped through a tube (diameter  $d = 25 \text{ mm}$ ) with a downwards slope. Pressure is measured at 2 points 100 m apart, as shown. Calculate the friction losses per meter tube,  $P_{\text{losses}} / l$  (W/m)



Picture: ÖS96



## Tube systems /5

- For a tubing network (sv: rörsystem), a design calculation can involve
  - **Calculation of power losses**, primarily pressure drop losses that must be compensated for with pumps etc. in a given process tubing situation
  - **Calculation of flow velocities** or volume streams that will result when applying a certain pumping power to a certain tube system flow situation
  - **Calculation of tube diameters**, lengths and tubing lay-out for a certain process situation, often based on given pumps or pressure drop data etc.



Picture <http://www.pipetuff.com/>

*Sometimes iterative calculations are needed:*

$P_{\text{pump}} \rightarrow p_2 \text{ and } v_2 ; \rightarrow$   
*adjust  $p_2 \rightarrow$  new value for  $P_{\text{pump}}$  etc.*

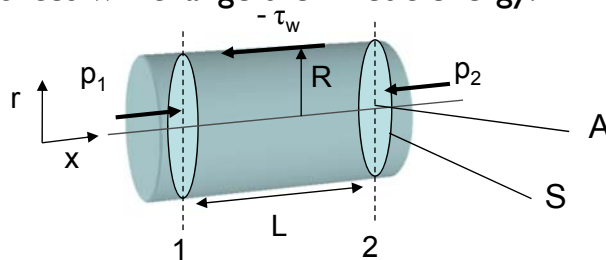
(see also ÖS96 p. 41)





## Pressure drop //

- The pressure drop in a tube flow system can be predicted **if the shear force at the wall  $\tau_w$  is known**
- For example for **laminar** tube flow (tube diameter  $d = 2R$ , flow direction "x"),  $(-dp/dx) = -2 \cdot \tau_w / R$  where  $\tau_w = \tau_{\text{fluid} \rightarrow \text{wall}}$  can be related to  $dv_x/dr$ , but for **turbulent** flow such **information is not available**
- Force analysis shows **3 forces acting on a flow volume element**: surface forces (pressure and surface shear), and body force (gravity). These can **change the kinetic energy  $E_k = \frac{1}{2}mv^2$  and potential energy  $E_p = mgz$** . For a horizontal tube the body forces cannot change, but surface forces will change the kinetic energy.



Volume element with length  $L$  (m), cross-section  $A$  (m<sup>2</sup>), circumference  $S$  (m), density  $\rho$  (kg/m<sup>3</sup>)

## Pressure drop //2 friction factor

- The surface shear force acting on the surface of a moving fluid element can be expressed as

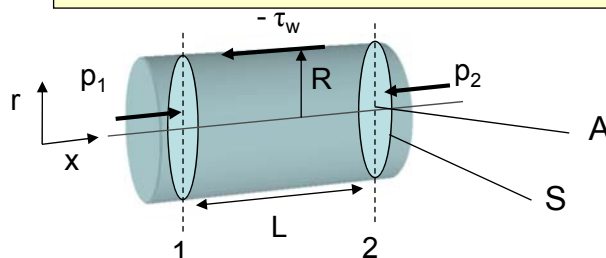
$$\tau_w = \text{friction factor} \cdot (E_{\text{kinetic}}/\text{volume}) = f \cdot \frac{1}{2} \rho \langle v \rangle^2$$

= dynamic pressure or "thrust" (sv: stöt)

- For flow in a horizontal tube with radius  $R$  the force balance at the wall for length section  $L$  gives

$$p_1 \cdot A - p_2 \cdot A - \tau_w \cdot S \cdot L = 0, \text{ with } \tau_w = \tau_{\text{fluid} \rightarrow \text{wall}} = -\tau_{\text{wall} \rightarrow \text{fluid}}$$

$$\rightarrow (p_1 - p_2) = \tau_w \cdot L \cdot S/A = f \cdot \frac{1}{2} \rho \langle v \rangle^2 \cdot L \cdot S/A = -\Delta p$$



with for a round tube cross-section  $A = \pi R^2$ , circumference  $S = 2\pi R$

# Pressure drop /3 friction factor

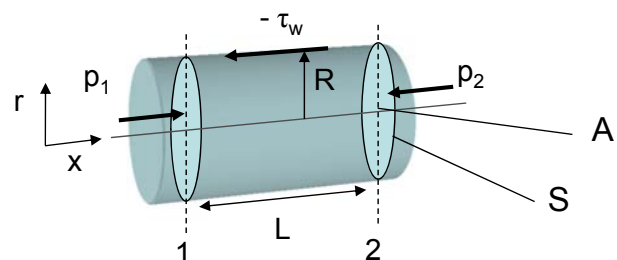
Be careful with literature data !

- This defines the **Fanning friction factor  $f$** ;  
also used is **Darcy or Blasius friction factor  $\zeta = 4f$**
- The group  $\frac{1}{2} \cdot \rho \cdot \langle v \rangle^2$  (unit: N/m<sup>2</sup>) follows also from dimensional analysis, reasoning that  $\tau_w = \tau_w(\rho, \eta, \langle v_x \rangle, \text{geometry})$ , which for a tube with diameter  $D$  gives  $\tau_w = \tau_w(\rho, \eta, \langle v_x \rangle, D)$ .
- It is found that

$$\tau_w / (\rho \cdot \langle v \rangle^2) = f(\text{Re}),$$

which is usually written as








$$\tau_w = \frac{1}{2} \cdot f \cdot \rho \cdot \langle v \rangle^2$$



# Hydraulic diameter

- The ratio  **$A/S$**  (unit: m) is a characteristic dimension of the tube, pipe, duct or channel known as hydraulic radius, while  **$4 \cdot A/S$**  is known as **hydraulic diameter  $D_h$**  (see Figure below) with  **$A$  = cross-sectional area** (sv: tvärsnitt);  **$S$  = perimeter** (sv: omkrets) **touched by fluid**

- For example for a round tube with diameter  $D$ , completely filled with fluid:  $D_h = D$ ;  
for a square channel with width  $W$ , fluid height  $H$ :  
 $D_h = 4 \cdot A/S = 4 \cdot (H \cdot W)/(2H + W)$

Flow situation	Hydraulic diameter $D_h = 4A/S$	$A$
 Circular pipe	$D$	$\frac{\pi}{4} D^2$
 Concentric pipe or slit	$D_2 - D_1 = 2\delta$	$\frac{\pi}{4} (D_2^2 - D_1^2)$
 Rectangular pipe	$\frac{2WB}{W+B}$	$WB$
 Open channel	$\frac{4WH}{W+2H}$	$WH$
 Open channel	$\frac{2H}{\sqrt{2}}$	$H^2$
 Half-filled	$D$	$\frac{\pi}{8} D^2$
 Liquid film in a tube	$4\delta$	$\delta \pi D$

Picture: BMH99

## Pressure drop /4 laminar tube flow

- Thus for the pressure drop for flow in a tube or duct with hydraulic diameter  $D_h = 4 \cdot A/S$ :

$$(p_1 - p_2) = -\Delta p = \tau_w \cdot L \cdot (4/D_h) = 4f \cdot \frac{1}{2} \cdot \rho \cdot \langle v \rangle^2 \cdot L/D_h$$

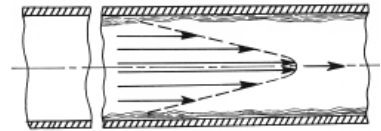
- For a **laminar flow in a round tube** (Hagen - Poiseuille flow, with  $D_h = \text{diameter } D = 2R$ ):

$$-\tau_w = \tau_{\text{wall} \rightarrow \text{fluid}} = \frac{1}{2}R \cdot (-\Delta p/L)$$

$$\rightarrow -\tau_w = 4\eta \langle v \rangle / R = 8\eta \langle v \rangle / D = f \cdot \frac{1}{2} \cdot \rho \cdot \langle v \rangle^2$$

$$\rightarrow \boxed{f = 16\eta / (\rho \langle v \rangle D) = 16 / \text{Re} ; \quad 4f = \zeta = 64 / \text{Re}}$$

with  $\text{Re} < 2100$



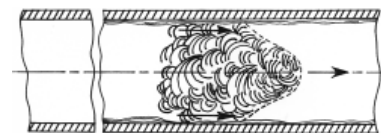
- For non-circular ducts another proportionality constant is needed !

Picture: <http://www.mbamanufacturing.com/Publication/LamTurbFlow.htm>

## Pressure drop /5 turbulent tube flow

- Pressure drop for flow in a tube or duct with hydraulic diameter  $D_h = 4 \cdot A/S$  :  $(p_1 - p_2) = -\Delta p = \tau_w \cdot L \cdot (4/D_h) = 4f \cdot \frac{1}{2} \cdot \rho \cdot \langle v \rangle^2 \cdot L/D_h$

- For a **turbulent flow in a tube of duct** it is found that  $f \sim \text{Re}^{-0.25} \dots^0$  (less direct influence of viscosity than in laminar flow) and  $\Delta p \sim v^{1.75..2}$



- For smooth pipes

$$\boxed{f = 0.0791 \cdot \text{Re}^{-0.25} ; \quad 4f = \zeta = 0.316 \cdot \text{Re}^{-0.25}}$$

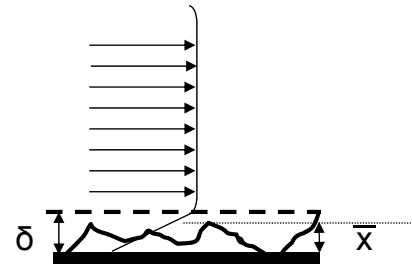
(Blasius' equation) with  $4000 < \text{Re} < 10^5$

can be used for any cross-sectional shape using characteristic diameter = hydraulic diameter  $D_h$

Picture: <http://www.mbamanufacturing.com/Publication/LamTurbFlow.htm>

# Pressure drop /6 wall roughness

- For rough pipes, **wall surface roughness** (sv: väggsbäckslighet)  $\bar{x}$  is important if it is of the same order as the thickness of the laminar boundary layer,  $\delta$ ;
- Important at great wall roughness or high Re numbers.
- Roughness data is found in tables
- Important is the **relative roughness**  $\bar{x}/D$ , with tube diameter D
- Not important for laminar flows
- The friction factor  $f$  or  $\zeta$  can be read from a **friction factor chart** or **Moody chart** as function of Re and relative wall roughness



APPROXIMATION for MOODY CHART

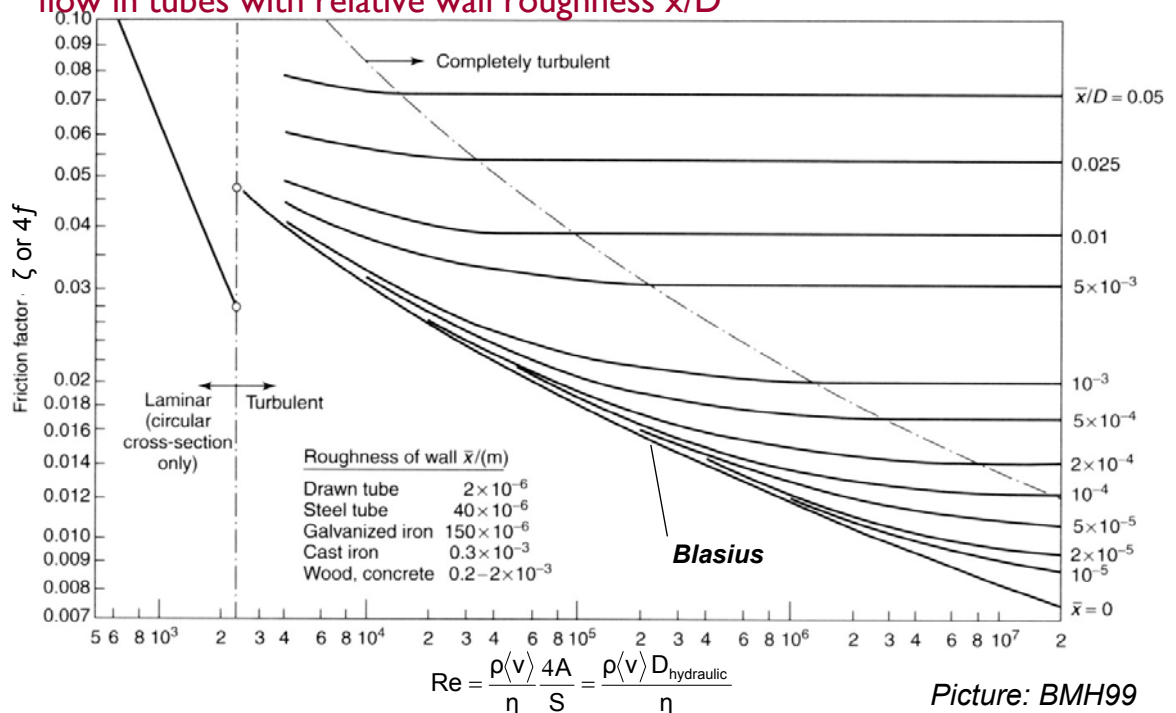
$$4f = \zeta = \frac{0.25}{\left(10 \log \left( \frac{\bar{x}}{3.7D} + \frac{5.74}{Re_D^{0.9}} \right)\right)^2}$$

$$5000 \leq Re_D \leq 10^8 \quad \text{and} \quad 10^{-6} \leq \frac{\bar{x}}{D} \leq 10^{-2}$$

## MOODY CHART

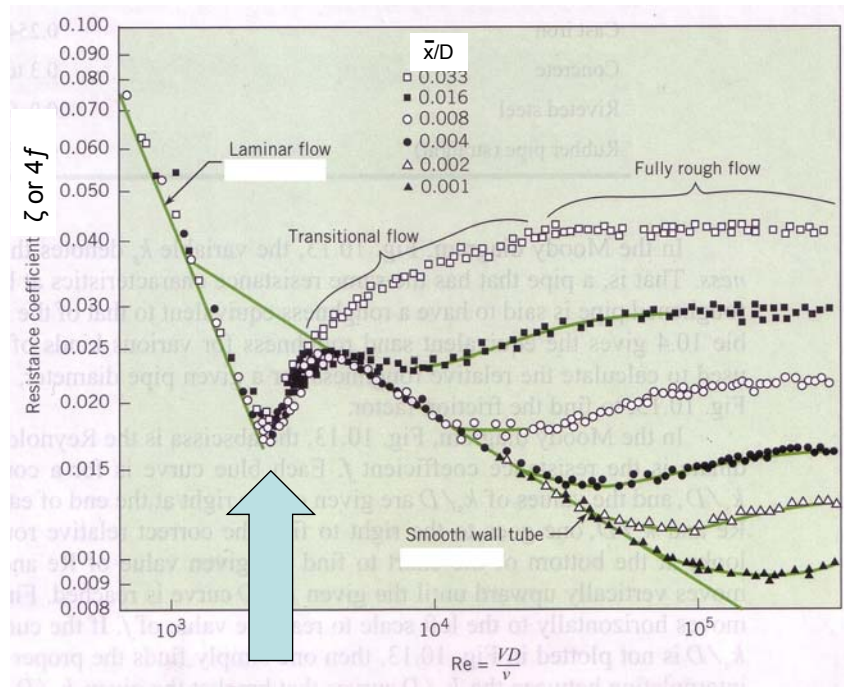
# Tube flow friction factor

flow in tubes with relative wall roughness  $\bar{x}/D$



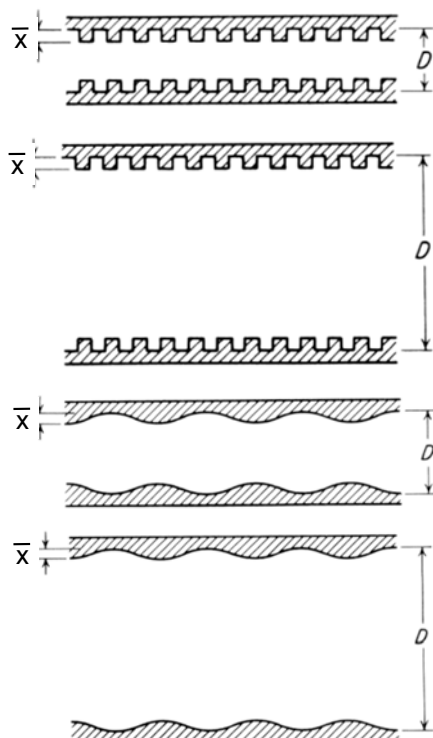
# Tube flow friction factor

flow in tubes with relative wall roughness  $\bar{x}/D$  - **the transition region**



Picture:  
CEWR10

## Wall roughness data

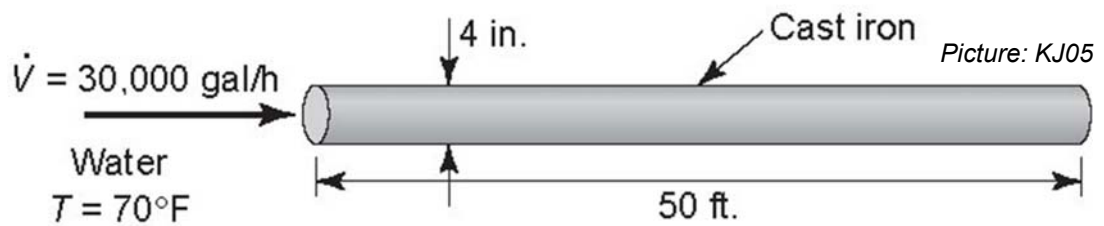


Material	Condition	Roughness Height, $\bar{x}$ (mm)	Uncertainty (%)
Steel	Sheet metal, new	0.05	$\pm 60$
	Stainless, new	0.002	$\pm 50$
	Commercial, new	0.046	$\pm 30$
	Riveted	3.0	$\pm 70$
	Rusted	2.0	$\pm 50$
Iron	Cast, new	0.26	$\pm 50$
	Wrought, new	0.046	$\pm 20$
	Galvanized, new	0.15	$\pm 40$
	Asphalted cast	0.12	$\pm 50$
Brass	Drawn, new	0.002	$\pm 50$
Plastic	Drawn tubing	0.0015	$\pm 60$
Glass	—	Smooth	
Concrete	Smoothed	0.04	$\pm 60$
	Rough	2.0	$\pm 50$
Rubber	Smoothed	0.01	$\pm 60$
Wood	Stave	0.5	$\pm 40$

← Relative wall roughness, small or large diameter tubes

Table: T06  
Pictures: MSH93

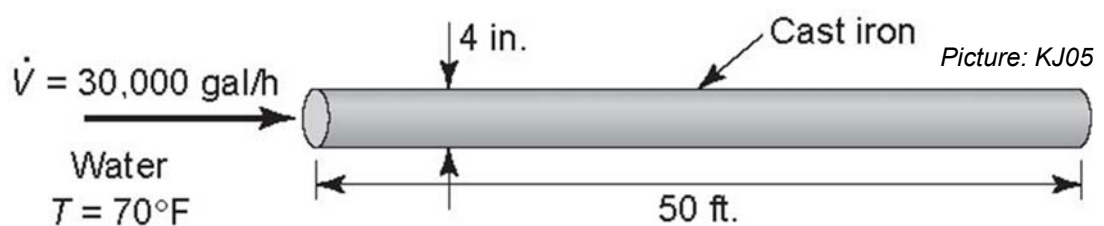
## Example: pipe flow friction /1



- A horizontal cast-iron pipe with diameter 4" carries 30000 (US) gal/h water at  $70^\circ\text{F}$ . Pipe length is 50 ft. Calculate the pressure drop. The water's density is  $62.2 \text{ lbm/ft}^3$ ; dynamic viscosity is  $65.8 \cdot 10^{-5} \text{ lbm/(ft} \cdot \text{s)}$



## Example: pipe flow friction /2





# Pressure drop /7 Fittings and valves

- Pressure drop across a tube section can be expressed as  

$$-\Delta p = 4f \cdot \frac{1}{2} \cdot \rho \cdot \langle v \rangle^2 \cdot L/D_h = \zeta \cdot \frac{1}{2} \cdot \rho \cdot \langle v \rangle^2 \cdot L/D_h$$
- Similarly, for the sudden local pressure drop caused over a very short distance by, for example
  - A change in tube diameter, or a bend or curve, or a T-junction
  - A valve (sv: ventil, klaff) or other fitting (sv: rörellement)
  - An inlet or outlet (sharp or smooth)

For these, pressure drop can be expressed as

$$-\Delta p = K_w \cdot \frac{1}{2} \cdot \rho \cdot \langle v \rangle^2 \quad \text{or}$$

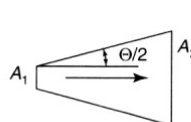
$$-\Delta p = \zeta' \cdot \frac{1}{2} \cdot \rho \cdot \langle v \rangle^2$$

with coefficients  $K_w$  or  $\zeta'$  independent of flow Reynolds number for  $Re > 10^5$



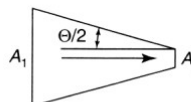
Picture: <http://www.chicagobrassworks.com/gs.htm>

## Friction loss factors $K_w$ (or $\zeta'$ ) for flow tube elements / 1 of 4

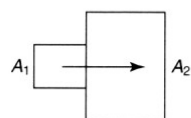


$$K_w = k \left( 1 - \frac{A_1}{A_2} \right)^2$$

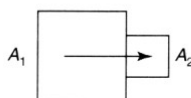
$\Theta$	$< 10^\circ$	$10^\circ$	$20^\circ$	$30^\circ$	$40^\circ$	$50^\circ$	$60^\circ$	$70^\circ$	$80^\circ$	$90^\circ$
$k$	0	0.17	0.41	0.71	0.90	1.03	1.12	1.13	1.10	1.05



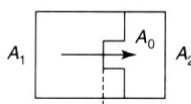
$\Theta$	$10^\circ$	$20^\circ$	$30^\circ$	$40^\circ$	$50^\circ$	$60^\circ$	$70^\circ$	$80^\circ$
$K_w$	0.16	0.20	0.24	0.28	0.31	0.32	0.34	0.35



$$K_w = \left( 1 - \frac{A_1}{A_2} \right)^2 \quad (\text{corrected!})$$



$$K_w = 0.45 \left( 1 - \frac{A_2}{A_1} \right)$$

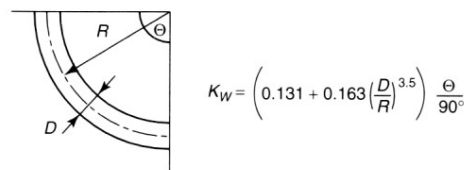
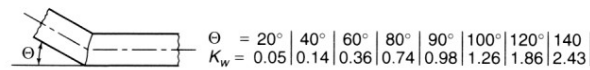
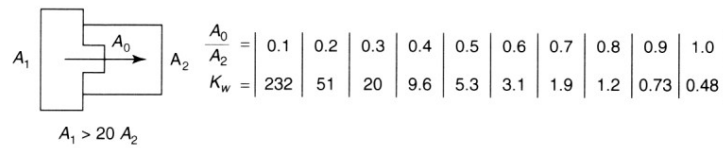


$\frac{A_0}{A_1}$	0.1	0.2	0.3	0.4	0.5	0.6	0.7	0.8	0.9
$K_w$	226	47.8	17.5	7.8	3.75	1.80	0.80	0.30	0.06

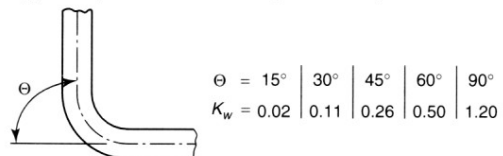
Downstream velocity,  
 $Re > 10^5$

Picture: BMH99

# Friction loss factors $K_w$ (or $\zeta'$ ) for flow tube elements / 2 of 4



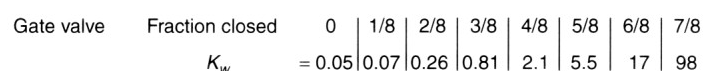
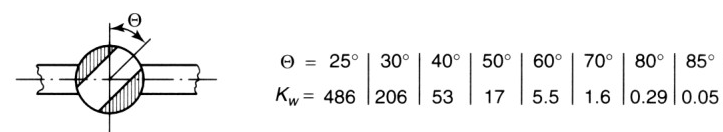
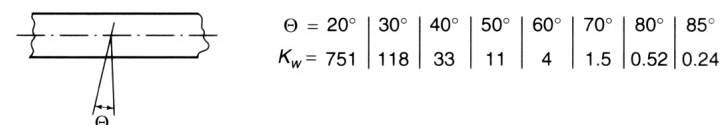
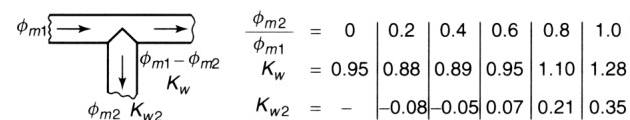
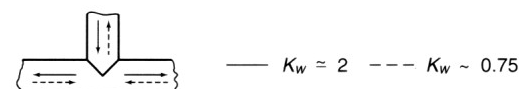
$K_w$  (referring to downstream velocity for  $Re > 10^5$ )



Downstream velocity,  
 $Re > 10^5$

Picture: BMH99

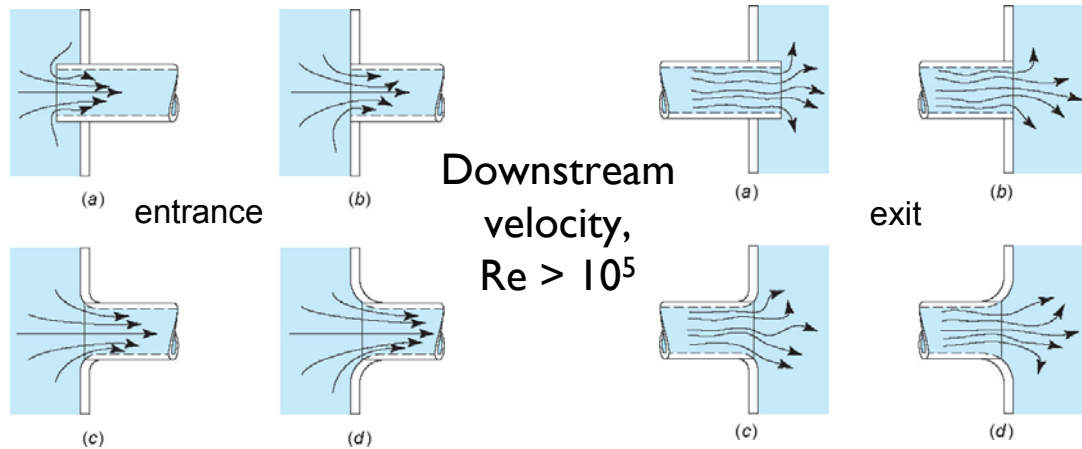
# Friction loss factors $K_w$ (or $\zeta'$ ) for flow tube elements / 3 of 4



Downstream velocity,  
 $Re > 10^5$

Picture: BMH99

# Friction loss factors $K_w$ (or $\zeta'$ ) for flow tube elements / 4 of 4



Entrance / exit flow conditions & loss coefficient:

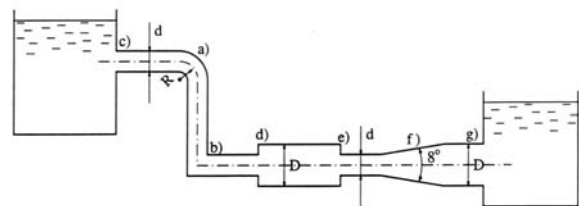
(a) reentrant,	entrance $K_w = 0.8$ ,	exit $K_w = 1.0$
(b) sharp-edged,	entrance $K_w = 0.5$ ,	exit $K_w = 1.0$
(c) slightly rounded,	entrance $K_w = 0.2$ ,	exit $K_w = 1.0$
(d) well-rounded,	entrance $K_w = 0.04$ ,	exit $K_w = 1.0$

Pictures: KJ05

## Tube elements: example

- Friction coefficients  $K_w$  or  $\zeta'$  for several tube sections and fitting elements:

- Bend  $90^\circ$ ,  $R/d = 1$   $\zeta' = 0.5$
- Sharp bend  $90^\circ$   $\zeta' = 0.98$  or elbow  $\zeta' = 1.2$
- Tube inlet, sharp  $\zeta' = 0.5$  or smooth  $\zeta' = 0.20$
- Diameter increase, sharp  $\zeta' = (1 - d^2/D^2)^2$
- Diameter decrease, sharp  $\zeta' = 0.45 \cdot (1 - d^2/D^2)$
- Diameter increase, diffusor with  $\theta/2 < 10^\circ$   $\zeta' \approx 0$
- Tube outlet, turbulent  $\zeta' = 1$  or laminar  $\zeta' = 2$



For this set-up if for example  $D = 80$  mm,  $d = 50$  mm, for turbulent flow:  
 $\sum \zeta' = 0.50 + 0.50 + 0.98 + 0.37 + 0.27 + 0 + 1.1 = 3.72$  for the fittings, bends and diameter changes only.

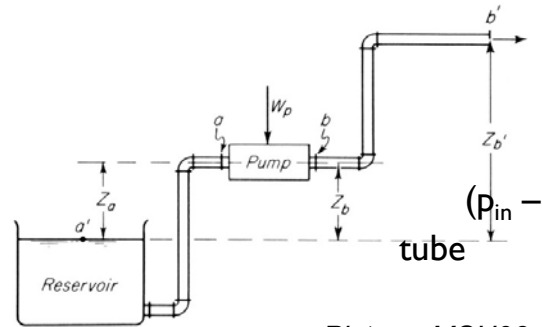
Picture: ÖS96

## Pressure drop, pressure loss, power loss, energy dissipation //

- For fluid flow with viscous friction through a channel the power loss (energy dissipation)  $P_{\text{loss}}$  (sv: effektförlust) can be related to pressure loss  $-\Delta p_{\text{loss}}$  for a given volume stream  $\dot{V}$ :

$$-\Delta p_{\text{loss}} = \frac{P_{\text{losses}}}{\dot{V}}$$

(unit: Pa) which is equal to  $p_{\text{out}}$  **only** for a horizontal without diameter changes.



Picture: MSH93

- For the energy equation for a tube system (with  $\dot{Q} = 0$ ), dividing by  $\dot{V}$  (noting that  $\dot{m} = \rho \cdot \dot{V}$  requires  $\rho = \text{constant}$ ) this gives

$$\rho g(z_1 - z_2) + \frac{1}{2}(\xi_1 \rho \langle v \rangle_1^2 - \xi_2 \rho \langle v \rangle_2^2) + (p_1 - p_2) + (-\Delta p)_{\text{pump}} = (-\Delta p)_{\text{losses}}$$

## Pressure drop, pressure loss, power loss, energy dissipation /2

- If density changes are significant (typical for gases) then  $\dot{V}_1 \neq \dot{V}_2$  and that must be accounted for:

$$-\Delta p_{\text{loss}} = \frac{P_{\text{losses}}}{\dot{V}} = \frac{P_{\text{losses}}}{\dot{m}} \int_1^2 dp = \int_1^2 -dp_{\text{loss}}$$

and

$$g(z_1 - z_2) + \frac{1}{2}(\xi_1 \langle v \rangle_1^2 - \xi_2 \langle v \rangle_2^2) + \int_1^2 \frac{-dp}{\rho} + (-\Delta p)_{\text{pump}} = (-\Delta p)_{\text{losses}}$$

- With pressure drop  $\Delta p \sim \text{shear force}$  it follows that  **$\Delta p \sim \text{velocity}$  for laminar flow**, and  **$\Delta p \sim \text{velocity}^{1.75 \dots 2}$  for turbulent flow**. Note: for laminar:  $\Delta p \sim v$  with  $4f \sim 1/\text{Re} \sim 1/v$
- With viscous work  $\sim \text{shear force} \times \text{velocity}$ ,  $P_{\text{loss}} \sim \Delta p \cdot \dot{V} \sim \text{velocity} \cdot \Delta p$  this gives  **$P_{\text{loss}} \sim \text{velocity}^2$  for laminar flow**, and  **$P_{\text{loss}} \sim \text{velocity}^{2.75 \dots 3}$  for turbulent flow**.

## Pressure drop, pressure loss, power loss, energy dissipation /3

- For the power loss (energy dissipation) for a flow channel with **total** pressure losses  $\Delta p_{\text{loss}}$ , composed of
  - $\Delta p_{\text{loss}} (\zeta, L, D)$  for the **straight sections** and
  - $\Delta p_{\text{loss}} (\zeta')$  for the **fittings, valves, diameter changes, in-/outlet, ...** :

$$4f = \zeta = \frac{-\Delta p_{\text{losses}}}{\frac{1}{2}\rho \langle v \rangle^2} \cdot \frac{D_h}{L} = \frac{P_{\text{losses}}}{\frac{1}{2}\rho \langle v \rangle^2 \cdot \dot{V}} \cdot \frac{D_h}{L} = \frac{P_{\text{losses}}}{\frac{1}{2}\dot{m} \langle v \rangle^2} \cdot \frac{D_h}{L} \quad \text{for tube sections}$$

$$K_w = \zeta' = \frac{-\Delta p_{\text{losses}}}{\frac{1}{2}\rho \langle v \rangle^2} = \frac{P_{\text{losses}}}{\frac{1}{2}\rho \langle v \rangle^2 \cdot \dot{V}} \quad \text{for valves, fittings, diameter changes, .....}$$

which gives for the total tubing system including fittings etc:

$$-\Delta p_{\text{loss}} = \frac{1}{2}\rho \langle v \rangle^2 \cdot \left( \zeta \cdot \frac{L}{D_h} + \sum \zeta' \right) \quad \text{and} \quad P_{\text{losses}} = \frac{1}{2}\rho \dot{V} \langle v \rangle^2 \cdot \left( \zeta \cdot \frac{L}{D_h} + \sum \zeta' \right)$$

*Note: kinetic energy correction factor  $\xi$  is now included in  $\zeta$  or  $4f$  !!!!*

## Calculation of volume flow or tube diameter

- Calculation of pressure drop  $-\Delta p$  or power loss  $P_{\text{loss}}$  from flow channel diameters and friction factors is relatively straight-forward; more complicated, however, is to determine volume stream  $\dot{V}$  or channel diameter  $D_h$  based on  $-\Delta p$  or  $P_{\text{loss}}$
- An **iterative procedure** can be used, using  $\dot{V} = A \cdot \langle v \rangle$  for flow cross-section  $A$  and the expressions given above; for tube system based on a round tube with  $A = \frac{1}{4}\pi D^2$  this gives

$$\dot{V} = \sqrt{\frac{\pi^2 (-\Delta p)_{\text{loss}} D^4}{8\rho \left( \zeta \frac{L}{D} + \sum \zeta' \right)}} \quad \text{and} \quad D = \sqrt[4]{\frac{8\rho \dot{V}^2 \left( \zeta \frac{L}{D} + \sum \zeta' \right)}{\pi^2 (-\Delta p)_{\text{loss}}}}$$

where  $\zeta$  (or  $4f$ ) and  $\zeta'$  (or  $K_w$ ) are functions of  $\langle v \rangle$ ,  $D$  and/or  $Re$  !

(see also ÖS96 p. 48)

## Example: old exam question /question

- Calculate what the inner diameter  $d$  (in m) of a well heat-insulated steel tube should be for transporting  $\dot{m} = 3,2$  kg/s steam with temperature  $180^\circ\text{C}$  and pressure 300 kPa (density  $\rho = 1,464$  kg/m<sup>3</sup>, dynamic viscosity  $\eta = 15,1 \times 10^{-6}$  Pa · s), if the pressure drop in straight tube sections may not be more than 250 Pa per meter. Wall roughness is  $k = \bar{x} = 0,4$  mm.
- Note that for round tubes: 
$$\text{Re} = \frac{4 \cdot \dot{m}}{\pi \cdot \eta \cdot d}$$
- Advice: develop an expression  $d = f(\langle v \rangle, \zeta, \dots)$  and iterate a few times to find a result for  $d$  (m).



## Example: old exam question /answer





## Calculation of volume flow or tube diameter

- Two expressions for this are given in CEWR10, p. 332

$$\dot{V} = -2.22 \cdot D^{5/2} \cdot \frac{(-\Delta p)_{loss}}{\rho \cdot L} \cdot {}^{.10} \log \left( \frac{\bar{x}}{3.7 \cdot D} + \frac{1.78 \cdot \eta}{D^{3/2} \cdot \rho \cdot \sqrt{\frac{(-\Delta p)_{loss}}{\rho \cdot L}}} \right)$$

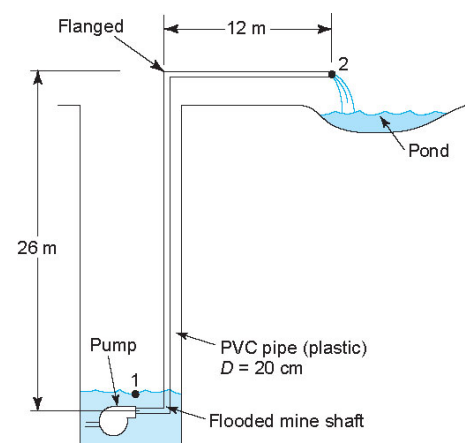
$$D = 0.66 \cdot \left( \left( \frac{\bar{x}}{D} \right)^{1.25} \cdot \left( \frac{L \cdot \dot{V}^2 \cdot \rho}{(-\Delta p)_{loss}} \right)^{4.75} + \frac{\eta}{\rho} \cdot \dot{V}^{9.4} \cdot \left( \frac{L \cdot \rho}{(-\Delta p)_{loss}} \right)^{5.2} \right)^{0.04}$$

for  $Re > 3000$ ,  $\frac{\bar{x}}{D} < 0.02$

which should be used with caution.

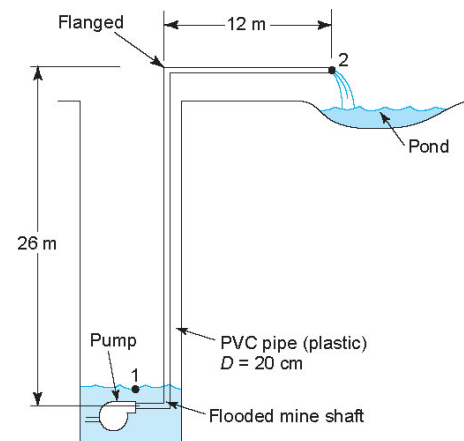
## Example: water pumping system //

- A pump is used to remove water from a mine shaft – see Figure. How much pump power  $P_{\text{pump}}$  (in kW) is needed to remove water at a rate of 65.0 kg/s? Assume an ideal pump (efficiency 100%). Assume density  $\rho = 997 \text{ kg/m}^3$ , viscosity  $\eta = 1.12 \cdot 10^{-3} \text{ Pa}\cdot\text{s}$



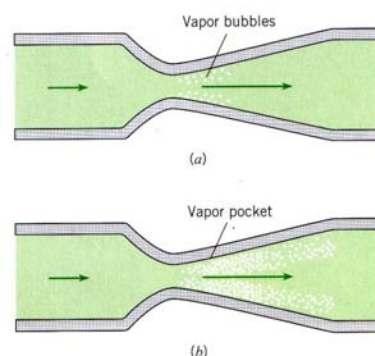
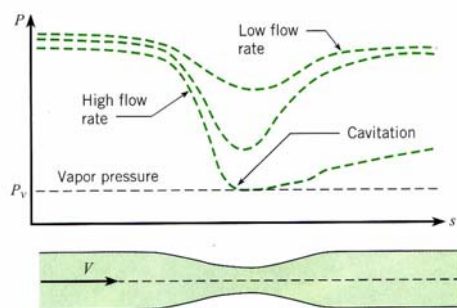
Picture: KJ05

## Example: water pumping system /2



Picture: KJ05

## Cavitation



Pictures:  
CEWR10

- Cavitation occurs if fluid pressure is reduced to the vapour pressure (at the given temperature) so that **boiling** occurs.
- The formation and collapse of bubbles gives shock waves, noise, and other problematic dynamic effects that can result in **reduced performance, failure and damage**.
- Typically occurs at high velocity locations in, for example, pumps or valves, but can damage also tube walls.

## 6.5 Flow systems with negligible losses, flow measurement



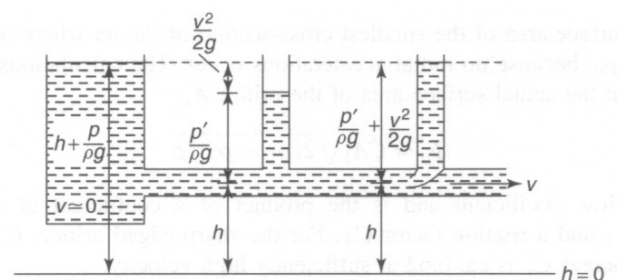
### Flow systems with negligible losses //

- Often the energy dissipation  $P_{\text{loss}}$  can be neglected in comparison with the (mechanical) energy changes in a flow system.
- If the fluid density can be considered constant this gives the **Bernoulli's equation**, which can be written as

$$\frac{p}{\rho g} + z + \frac{1/2 v^2}{g} = \text{constant}$$

where the three terms  
(unit: m) are referred to as

- **pressure head**,
- **static head** and
- **velocity head**



Picture: BMH99



## Flow systems with negligible losses /2

- This is used when measuring fluid velocities with a so-called **Pitot tube**: in the Figure →

$$P_{@b} - P_{@a} = \frac{1}{2}\rho \langle v \rangle^2 = \rho gh$$

- In a **venturi flowmeter**, the pressure difference between main flow and the throat as shown in Figure → equals

$$P_{@A1} - P_{@A2} = \frac{1}{2}\rho \langle v \rangle^2_{@A2} - \frac{1}{2}\rho \langle v \rangle^2_{@A1}$$

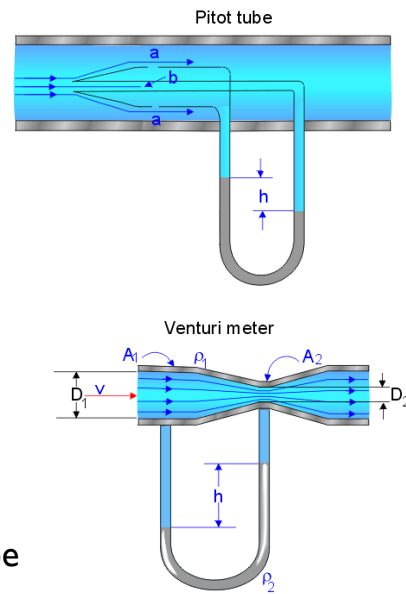
(which gives  $p_{@1} > p_{@2}$  !)

$$\text{with } \langle v \rangle_1 \cdot A_1 = \langle v \rangle_2 \cdot A_2 \text{ and}$$

$P_{@A1} - P_{@A2} = \rho gh$  the flow  $\dot{V}$  at  $A_2$  can be calculated for a liquid:

$$\dot{V} = A_2 \sqrt{\frac{2(p_1 - p_2)}{\rho}} \bigg/ \sqrt{1 - \frac{A_2^2}{A_1^2}}$$

For a gas: (ideal, adiabatic process):  
use  $p \cdot \rho^{-\gamma} = \text{constant}$ ,  $\gamma = c_p/c_v$



Pictures: <http://www1.uts.com/Physics/flowmetering/flowmeter.htm>

## Flow systems with negligible losses /3

- For flow of **liquid** from an **orifice**  
(sv: mynning, öppning) friction losses can be neglected
- At some distance from the opening, (at cross-sectional area  $A_1$ ), the velocity is much smaller than the velocity  $\langle v \rangle$  in the opening (area  $A_1'$ ):

$$p_0 + \frac{1}{2}\rho \langle v \rangle^2 \approx p_1 \text{ this gives}$$

$$\langle v \rangle \approx \sqrt{2(p_1 - p_0)/\rho}$$

$$\dot{V} = A_1' \langle v \rangle = C_f A_1 \sqrt{2(p_1 - p_0)/\rho}$$

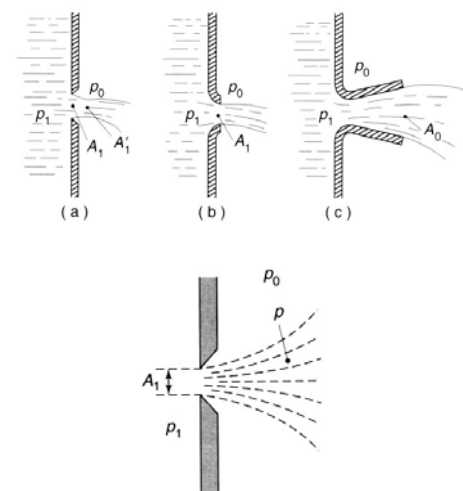
with friction factor  $C_f$

$C_f \approx 1$  for a sharp edge (a),

$C_f \approx 0.95-0.99$  for a rounded outlet (b).

For a diffuser (c) with angle  $< 8^\circ$ ,

$$\dot{V} = C_f A_0 \sqrt{2(p_1 - p_0)/\rho} \text{ with } C_f \approx 1$$



For a gas :  
(ideal, adiabatic process):  
 $p_0 < p \text{ in jet} < p_1$   
use  $p \cdot \rho^{-\gamma} = \text{constant}$ ,  $\gamma = c_p/c_v$

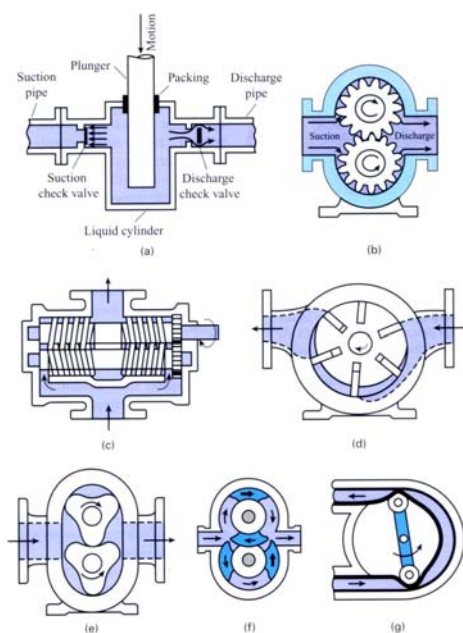
Pictures: BMH99

## 6.6 Pumps, compressors, fans



## Pumps, compressors, fans //

- **Creating a flow and/or increasing the pressure of a fluid, or compensating for pressure losses is accomplished with pumps** (sv: *pumpar*) for liquids, or with **compressors** or **fans** (sv: *kompressorer, fläktar*) for gases
- Usually a fan creates flow with minimal pressure change; if a fan creates a higher outlet pressure then it is generally referred to as a **blower** (sv: *bläster*)



*Positive-displacement pumps*

Picture: T06

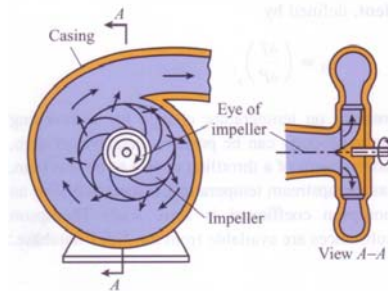


# Pumps, compressors, fans /2

- Pumps, compressors and fans can be divided into **two major categories**:
  - Positive displacement devices** based on "pushing" the fluid through the device (see *previous slide*)
  - Dynamic devices** based on transfer of energy as momentum (sv: *rörelsemängd*) from rotary blades or vanes, or from a high-speed fluid stream (*for example, centrifugal pumps and rotodynamic compressors and fans*)

Centrifugal pump

pictures: TO6



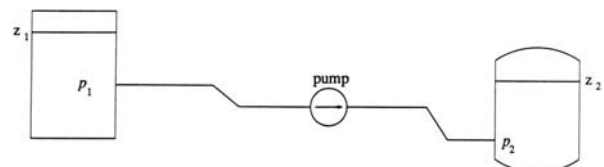
# Pumps /1

- The general relation between pump (or compressor) power and the pressure difference  $\Delta p_{\text{pump}}$  (sv: *uppföringstryck*) **for a given flow tubing system situation** follows from the mechanical energy balance ( $\dot{Q} = 0$ , no heat transfer or significant temperature changes), assuming also that  $\Delta \dot{E}_{\text{kinetic}} = 0$ :

$$-\Delta p_{\text{pump}} = (p_2 - p_1) + \rho g(z_2 - z_1) + \frac{1}{2}\rho \langle v \rangle^2 \cdot \left( \zeta \cdot \frac{L}{D_h} + \sum \zeta' \right) \quad \text{with} \quad \langle v \rangle^2 = \frac{\dot{V}^2}{A^2}$$

$$-\Delta p_{\text{pump}} = \frac{P_{\text{pump}}}{\dot{V}} = \frac{\dot{H}_2 - \dot{H}_1}{\dot{V}} = \rho(h_2 - h_1) = \rho g \cdot \Delta z_{\text{pump}}, \quad \text{with "pump head"} \Delta z_{\text{pump}}$$

- The **pump head** (unit: m) is the pressure rise across the pump equivalent height fluid



Picture: ÖS96

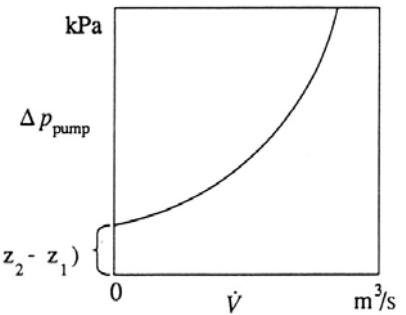


## Pumps /2

- The relation between  $-\Delta p_{\text{pump}}$  and  $\dot{V}$  is a **characteristic for the flow tubing system**

(sv: rörledningskaraktersistika)

for example →

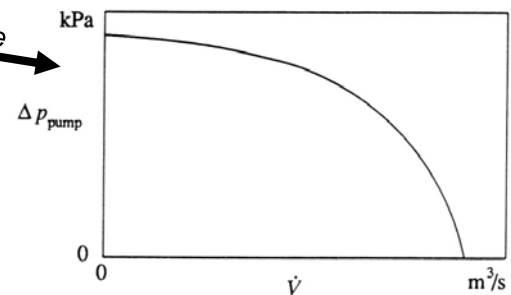


- For the pump itself, the **pump characteristic**

(sv: pumpkaraktersistika) gives the performance

$-\Delta p_{\text{pump}}$  versus  $\dot{V}$

for example →

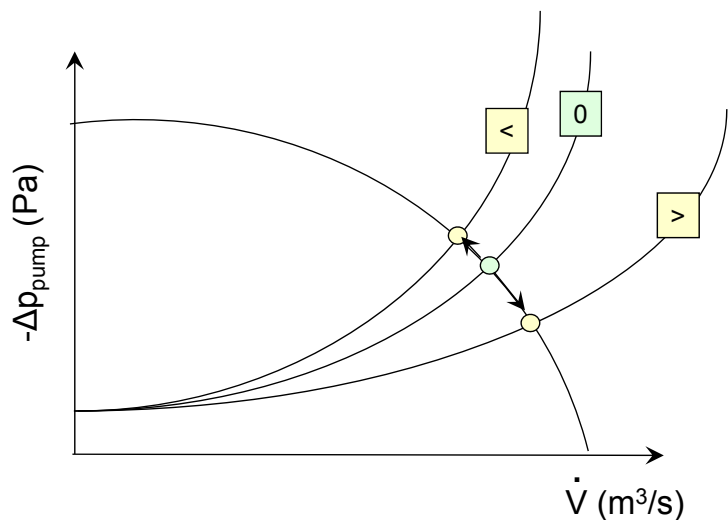


- Combining the two lines in one diagram gives the **working point** (sv: arbetspunkt) at the point where the lines cross

Pictures: ÖS96

## Pumps /3

- Changing the flow resistance in the tube network will give another system characteristic line
- The new working point will give another fluid flow throughput



For example, closing or opening a valve

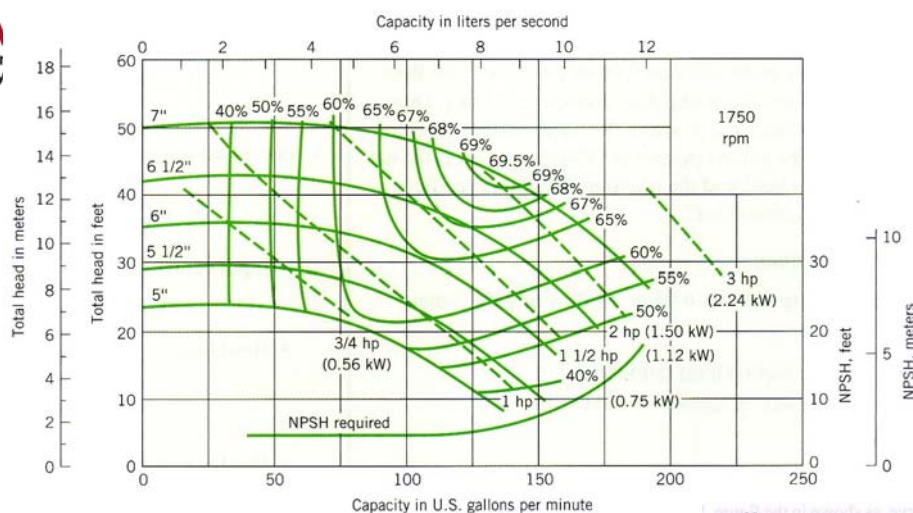
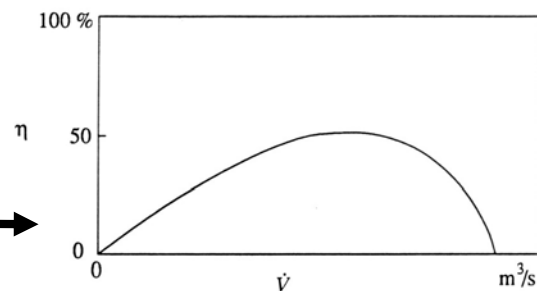
- The pump itself generates a viscous friction effect in the fluid, and as a result not all pump power  $P_{\text{pump}}$  will be available to give a pressure increase  $-\Delta p_{\text{pump}}$  in flow  $\dot{V}$ . The **pump efficiency** (sv: *pumpverkningsgrad*)  $\eta_{\text{pump}}$  quantifies for this:

$$\eta_{\text{pump}} = \frac{P_{\text{pump}}}{\text{Power input}} = \frac{-\Delta p_{\text{pump}} \cdot \dot{V}}{\dot{W}_{\text{in}}}$$

Picture: ÖS96

- For a given pump the efficiency depends on the fluid that is pumped and the volume stream  $\dot{V}$

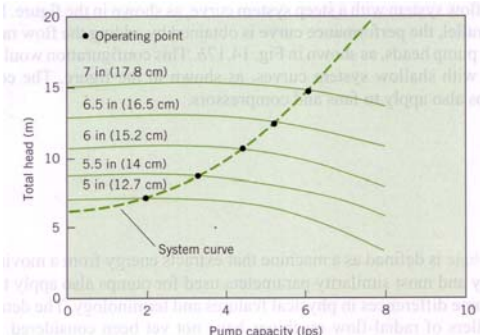
for example →



Picture: CEWR10

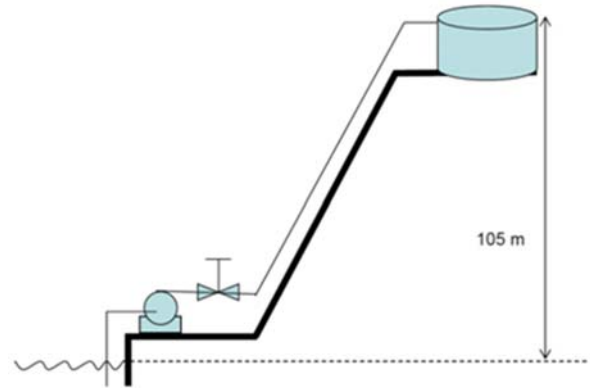
## Pump characteristics & performance

- examples



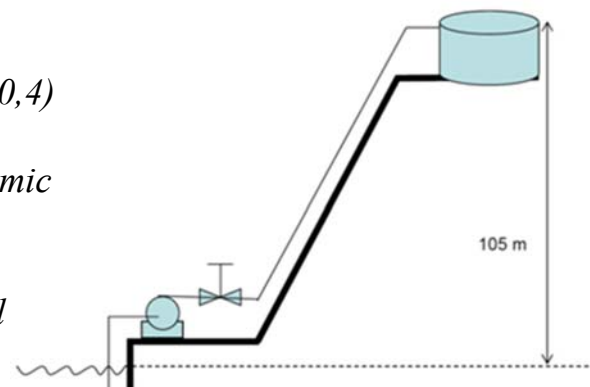
## Old exam question /q 1 of 2

- Cooling water must be pumped from a reservoir up to a process through a tube system with a few bends and a valve as shown in the figure. At both liquid surfaces the pressure equals ambient atmospheric pressure. The cooling water (20°C) flow is 135 m<sup>3</sup>/h. The height difference between reservoir and process is 105 m and the total tube length is 166 m, of which 16 m is upstream (“before”) of the pump.
- a) What tube diameter must be chosen so that the flow velocity does not exceed 2 m/s, and what is the Reynolds number of the flow then?
- ..... continues



## Old exam question /q 2 of 2

- b) What pressure head should the pump be able to produce so that the flow objective is achieved?
- c) Calculate the pump power that is needed for a pump with an efficiency of 80%.
- d) Is there a risk of so-called “cavitation” somewhere in this tube system?
- Assume that the friction coefficient for the valve is  $\zeta' = 2,0$ , assume two 45° elbow bends ( $\zeta' = 0,4$ ) and an 90° elbow bend ( $\zeta' = 0,9$ ). Density water = 1000 kg/m<sup>3</sup>; dynamic viscosity water = 0,001 Pa·s. Water vapour pressure at 20 °C is 2336,8 Pa. Assume the tube wall roughness to be  $= 4,7 \cdot 10^{-4}$  m.



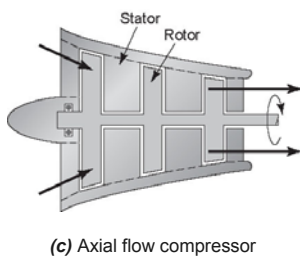
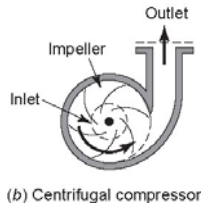
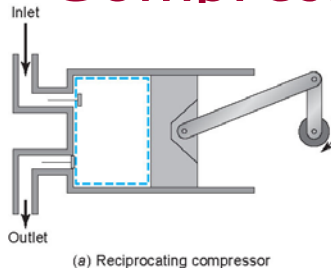
# Old exam question /a 1 of 2



# Old exam question /a 2 of 2



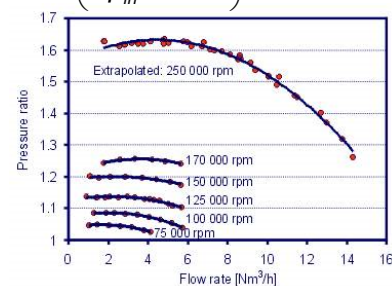
# Compressors



- A compressor increases pressure of a gas
- Most important are (a) reciprocating, (b) centrifugal and (c) axial flow types
- If  $p_{out} \leq 1.1 \times p_{in}$  the calculations are similar to those for a pump, **otherwise** gas compressibility must be considered → calculate as polytropic process with

$$P_{compr} = \dot{W}_{in} \cdot \eta_{compr} = \dot{H}_{out} - \dot{H}_{in}$$

$$\dot{W}_{in} = P_{compressor,theor} = \frac{\kappa}{\kappa - 1} \cdot \dot{V}_{in} \cdot p_{in} \cdot \left( \left( \frac{p_{out}}{p_{in}} \right)^{\frac{\kappa-1}{\kappa}} - 1 \right), \text{ and } \kappa = \frac{c_p}{c_v}$$



Pictures: JK05

A compressor characteristic

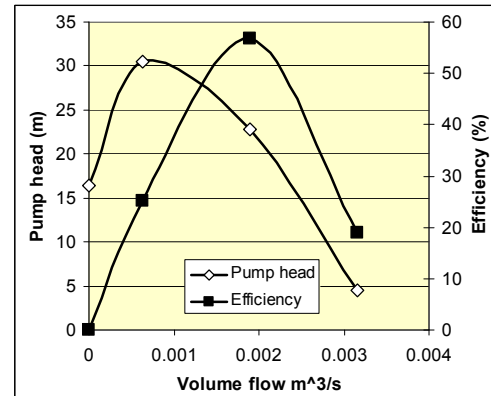
Picture: <http://www.mech.kuleuven.be/micro/topics/turbine/>

## Example: pump //

- A 1 hp (746 W) electrical motor drives a centrifugal pump for which the catalogue gives some *tabelised data*.
- Calculate the pumping power and efficiency for pumping water ( $\rho = 996 \text{ kg/m}^3$ ) with this pump, and plot these as a function of the flow rate  $\dot{V}$ .

Flow (m³/s)	Pump head (m)
$3.16 \times 10^{-3}$	4.6
$1.89 \times 10^{-3}$	22.9
$0.63 \times 10^{-3}$	30.5
0	16.5

## Example: pump /2



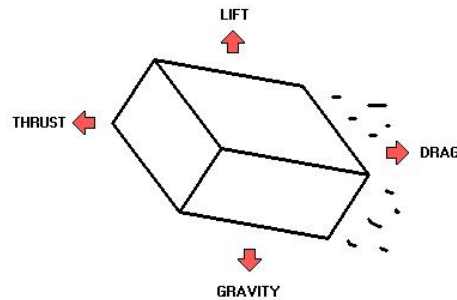
Source: T06

## 6.7 Fluid dynamics: external flows



# Fluid flow around objects

- In the cases of
  - an object moving through a fluid
  - a fluid flow around an object
 the **velocity difference** generates forces
- Forces acting parallel to the flow direction are **drag forces**; forces acting perpendicular to the flow direction are **lift forces**
- The flow field around an object can be divided in two parts: the **boundary layer** where the viscous forces are active, and the **free-stream velocity** (or the stagnant surrounding fluid)



\* See section 6.2

Picture: <http://www.weirdrichard.com/images/forces.jpg>



## Flow around a flat plate / I

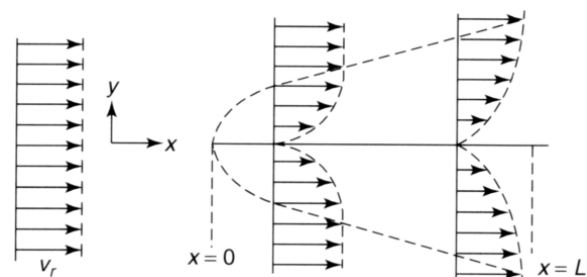
- For flow along a flat plate, the forces on the plate are friction forces. The **shear stress on each side of the surface** is

$$\tau_{yx}|_{y=0} = \tau_w = \eta_{fluid} \frac{v_r}{\delta} = 0.664 \cdot \left( \frac{v_r x \rho_{fluid}}{\eta_{fluid}} \right)^{-1/2} \cdot \frac{1}{2} \rho_{fluid} v_r^2, \quad \text{for } Re_x = \frac{v_r \cdot x \cdot \rho_{fluid}}{\eta_{fluid}} < 3 \times 10^5$$

with (laminar) boundary layer thickness  $\delta$  and **relative velocity**  $v_r$

- The **drag force** on each side of a plate with length  $L$  and width  $b$  is then given by

$$\begin{aligned} F_{drag} = F_D &= b \int_0^L \tau_w dx = \\ &= 1.33 \left( \frac{v_r L \rho_{fluid}}{\eta_{fluid}} \right)^{-1/2} \cdot bL \cdot \frac{1}{2} \rho_{fluid} v_r^2 \\ \text{for } Re_L &= \frac{v_r \cdot L \cdot \rho_{fluid}}{\eta_{fluid}} < 3 \times 10^5 \end{aligned}$$



- The pressure  $\frac{1}{2} \rho v^2$  is known as **THRUST** (sv: stöt)

Picture: BMH99



## Flow around a flat plate /2

- This defines the (length-averaged) **drag coefficient**  $C_D$  as

$$F_D = C_D \cdot A \cdot \frac{1}{2} \rho V_r^2 \quad \text{with } C_D = \frac{1.33}{\sqrt{Re_L}} \quad \text{for } Re_L < 3 \times 10^5$$

where  $A$  ( $m^2$ ) is the area (one side) of the plate

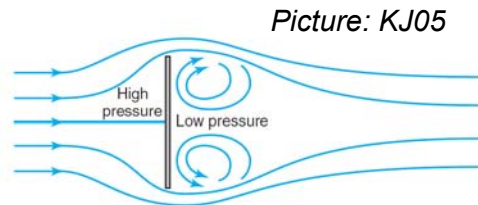
- For turbulent cases, experimental results give

$$C_D = \frac{0.074}{Re_L^{1/5}} \quad \text{for } 10^5 < Re_L < 10^7; \quad C_D = \frac{0.445}{10 \log(Re_L)^{2.58}} \quad \text{for } 10^7 < Re_L < 10^9$$

- For a flat surface with a laminar region followed by a turbulent region, a "composite" drag composition can be calculated with

$$C_D = \frac{0.074}{Re_L^{1/5}} - \frac{1740}{Re_L}$$

- For a flat plate **perpendicular** to fluid the drag coefficient equals  $\sim 2$ , largely independent of Re-number



Picture: KJ05

## Flow around cylinders, spheres /1

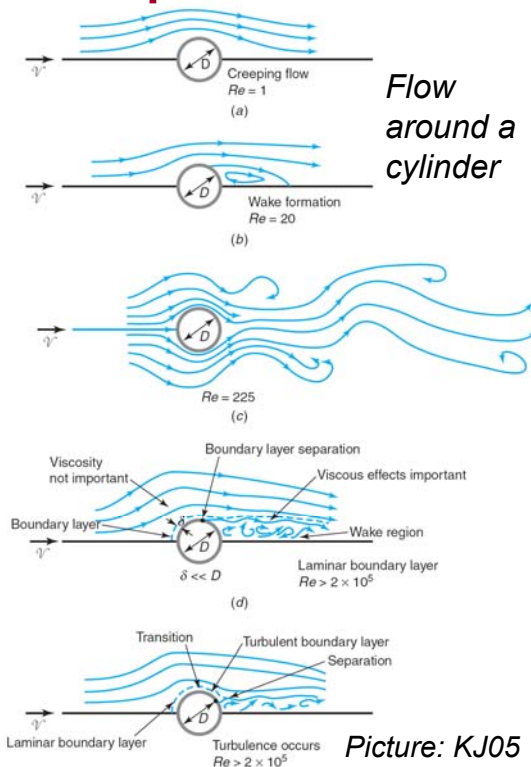
- For a general surface area  $A_\perp$  ( $m^2$ ) perpendicular to the flow, the **drag force** is

$$F_D = C_D \cdot A_\perp \cdot \frac{1}{2} \rho V_r^2$$

(where  $\frac{1}{2} \rho V_r^2$  is actually the pressure difference between the front and the back of the object)

- With increasing Re-numbers, **boundary layer separation** occurs, and

a **wake region** (sv: köl(vatten)) arises where kinetic energy is only partly converted into pressure



Picture: KJ05

# Flow around cylinders, spheres /2

- For spherical particles the drag coefficient equals
- For flow at  $Re < 0.1$  around a sphere, the relation  $C_D = 24/Re$

$$C_D = \frac{24}{Re}$$

for  $Re \ll 1$  or  $< 0.2$

$$C_D = \frac{24}{Re} \left( 1 + \frac{3}{16} Re \right)$$

for  $0.2 < Re < 2$

$$C_D = \frac{24}{Re} \left( 1 + \frac{1}{6} Re^{2/3} \right)$$

for  $2 < Re < 800$

$$C_D = 0.44$$

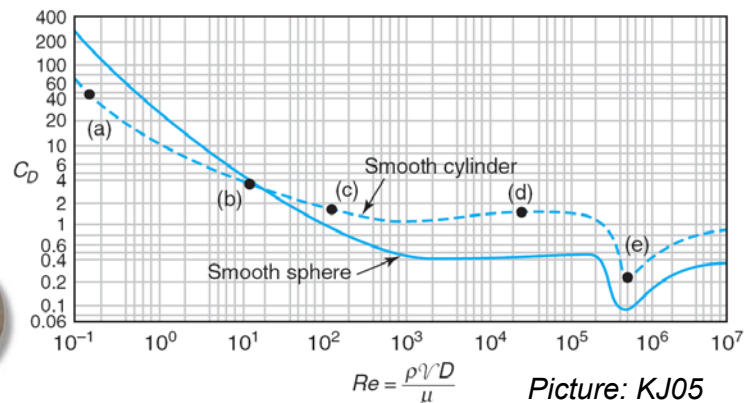
for  $800 < Re < 10^5$



follows also from **Stokes' law**

$$F_{\text{drag}} = 3\pi\eta v_r D$$

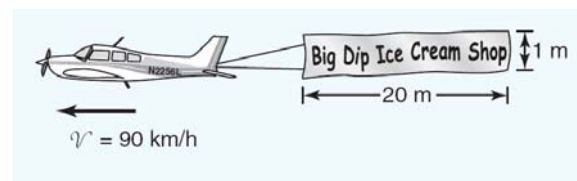
for a sphere with diameter  $D$  and relative velocity  $v_r$



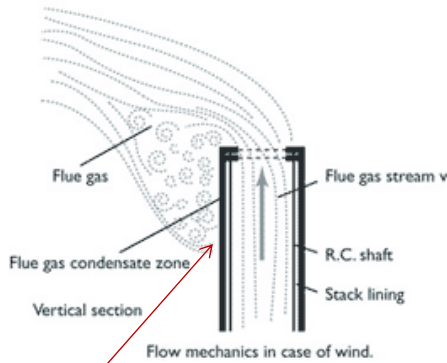
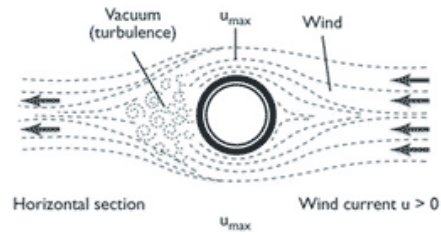
Picture: [http://www.school-for-champions.com/science/friction\\_changing\\_fluid.htm](http://www.school-for-champions.com/science/friction_changing_fluid.htm)

## Example: drag on a flat plate

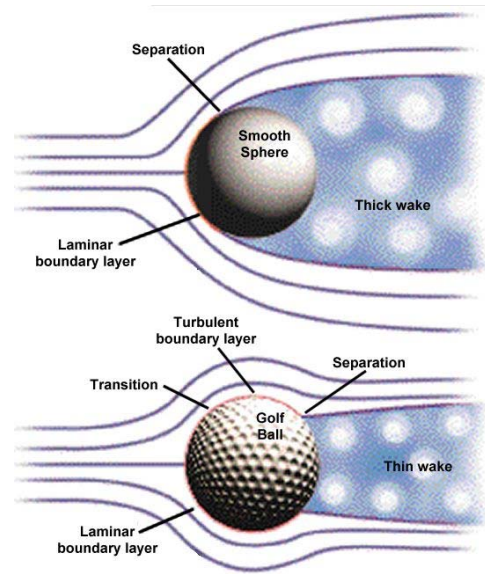
- An advertising banner (1 m x 20 m) is towed behind an aeroplane at 90 km/h, in air at 32°C.
- Calculate the power (in kW) needed to pull the banner.



# Boundary layer separation examples



*erosion, corrosion*



Picture: <http://www.aerospaceweb.org/question/aerodynamics/q0215.shtml>

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Picture: [http://www.ecotrust.org/copperriver/crks\\_cd/content/pages/photographs/images/](http://www.ecotrust.org/copperriver/crks_cd/content/pages/photographs/images/)