



中国科学技术大学

UNIVERSITY OF SCIENCE AND TECHNOLOGY OF CHINA

Atomic-orbital Based Ab-initio simulations

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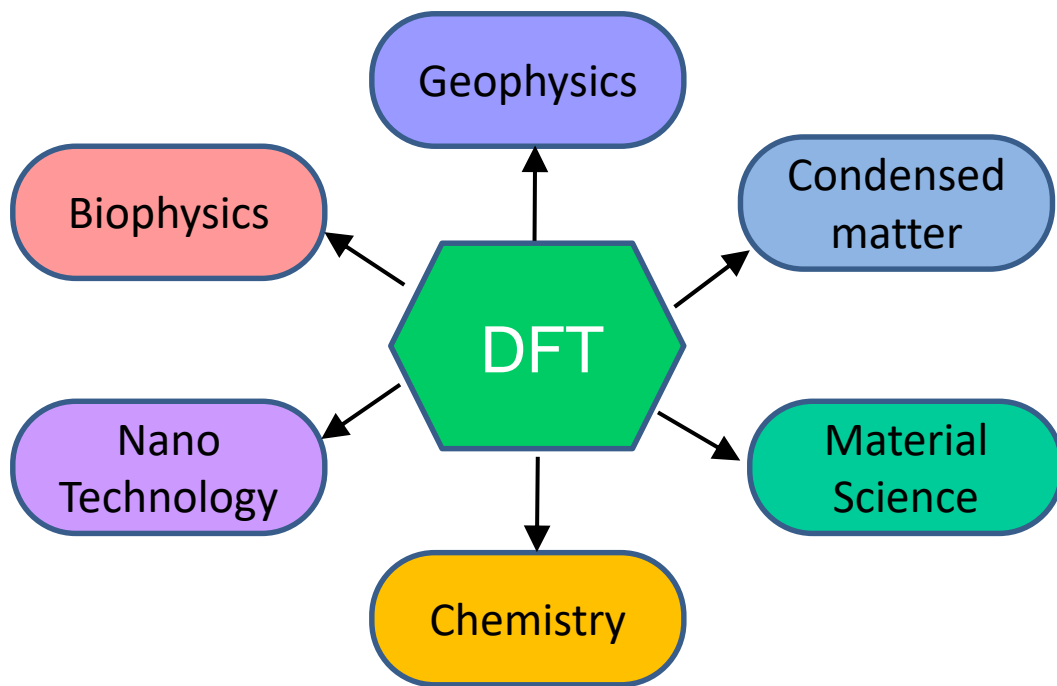
University of Science and Technology of China

The **Nobel Prize in Chemistry 1998** was awarded to W. Kohn "for his development of the density-functional theory".



1923/3/9-2016/4/19

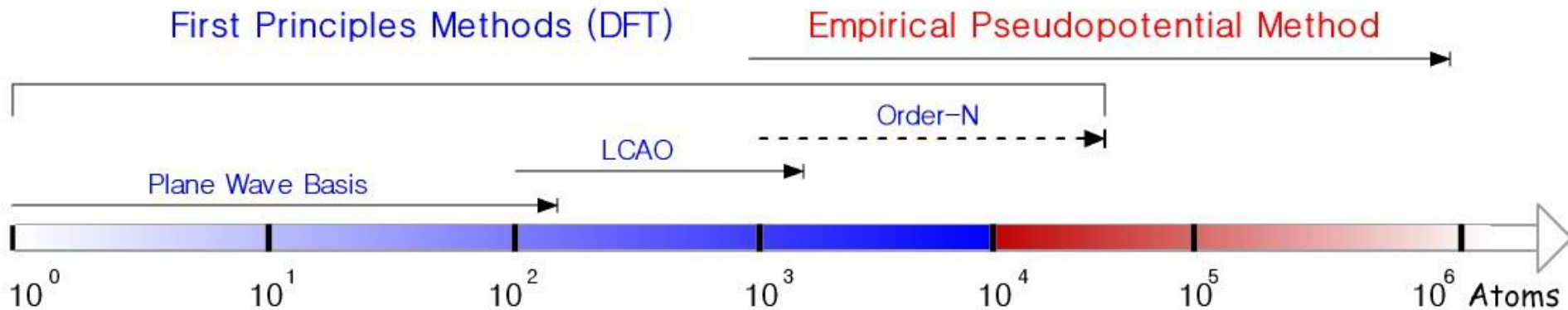
◆ Applications:



◆ New developments:

New functionals: more accurate
Numerical methods: larger systems

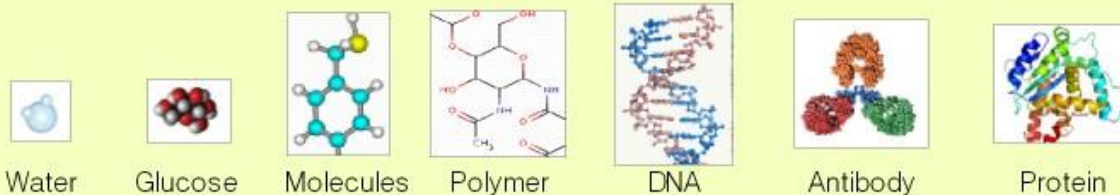
The development of first-principles methods



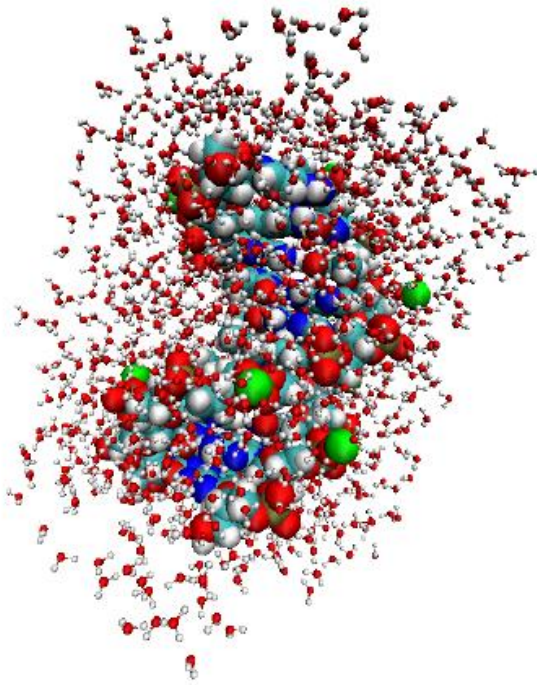
Physics



Chemistry/Biology

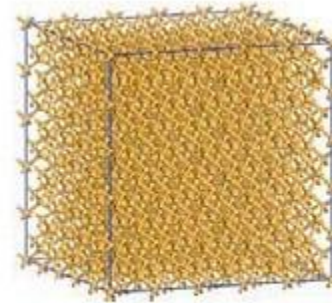


Application of DFT for thousands of atoms

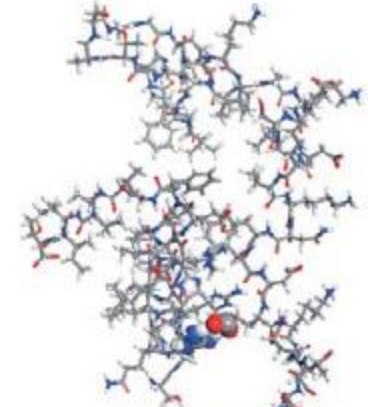


DNA containing 3439 atoms

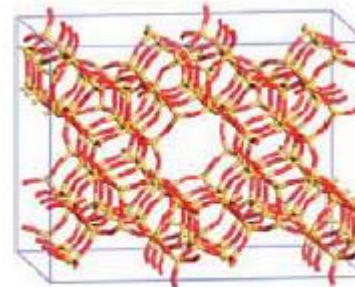
J. Phys. : Condens. Matter
20 (2008) 294201
(CONQUEST)



Crystalline silicon
(1000 atoms)

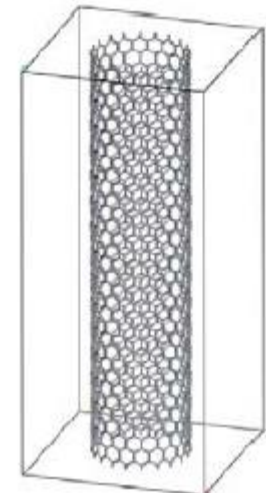


Protein
(988 atoms)



ZSM5 zeolite
(576 atoms)

Phys. stat. sol. (b)
243, No.5 (2006) 977
(ONETEP)



(20, 0) Nanotube
(1280 atoms)

Kohn-Sham equation

$$\left[-\frac{\nabla^2}{2m} + V_{Hxc}[n](r) + V_{ext}(r) \right] \varphi_i(r) = \varepsilon_i \varphi_i(r)$$

$$V_{Hxc}[n] = \frac{\delta E_{Hxc}[n]}{\delta n} \quad n(r) = \sum_i^{\text{occ}} |\varphi_i(r)|^2$$

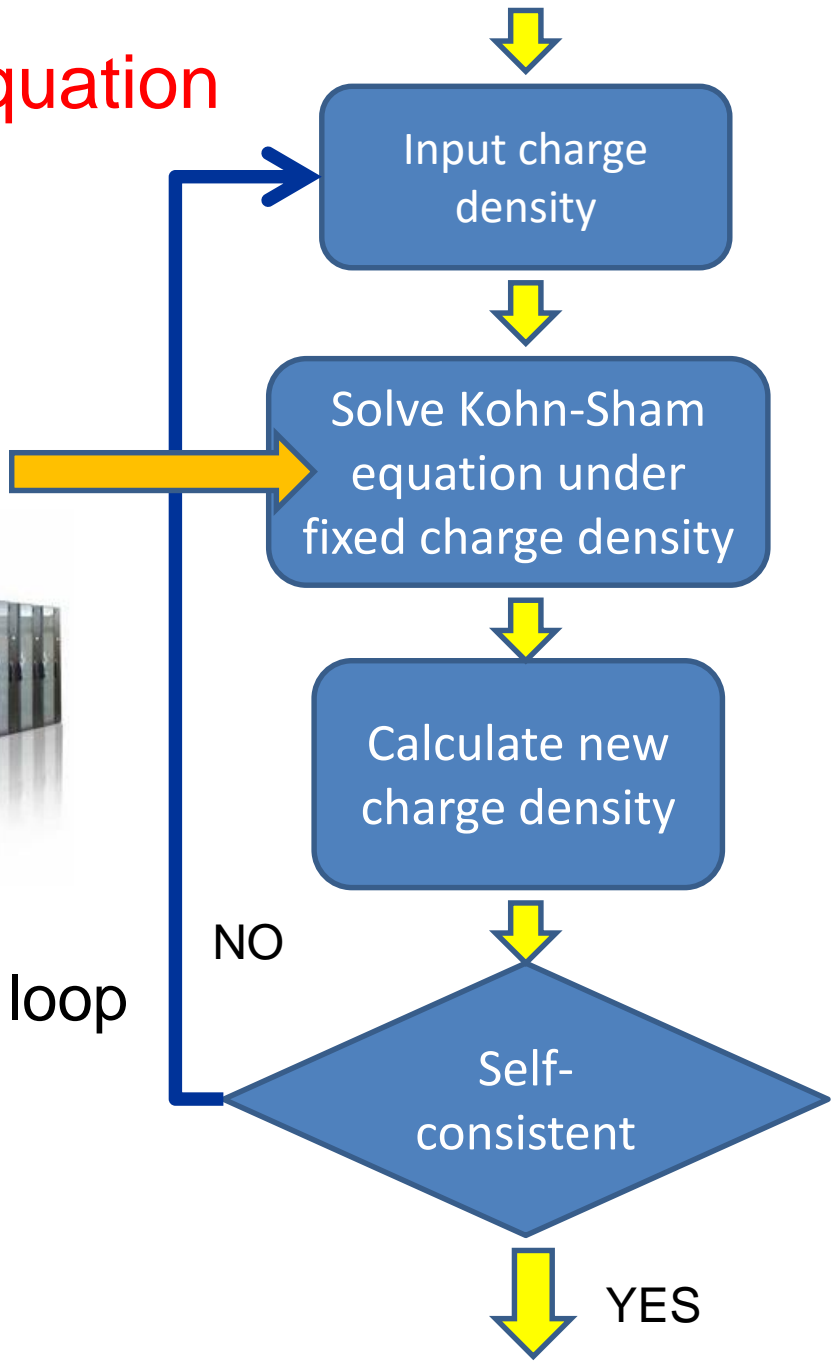
KS equation maps the many-particle problem into a single-particle problem. The mapping is exact, if the energy functional is exact!

Solving Kohn-Sham equation

Most time consuming step



Self-consistent loop



Solve Kohn-Sham equation

By choosing proper **basis set**, one can turn the Kohn-Sham equation to a **matrix eigenvalue problem**

$$|\psi_i\rangle = \sum_{q'} c_{i,q'} |q'\rangle$$

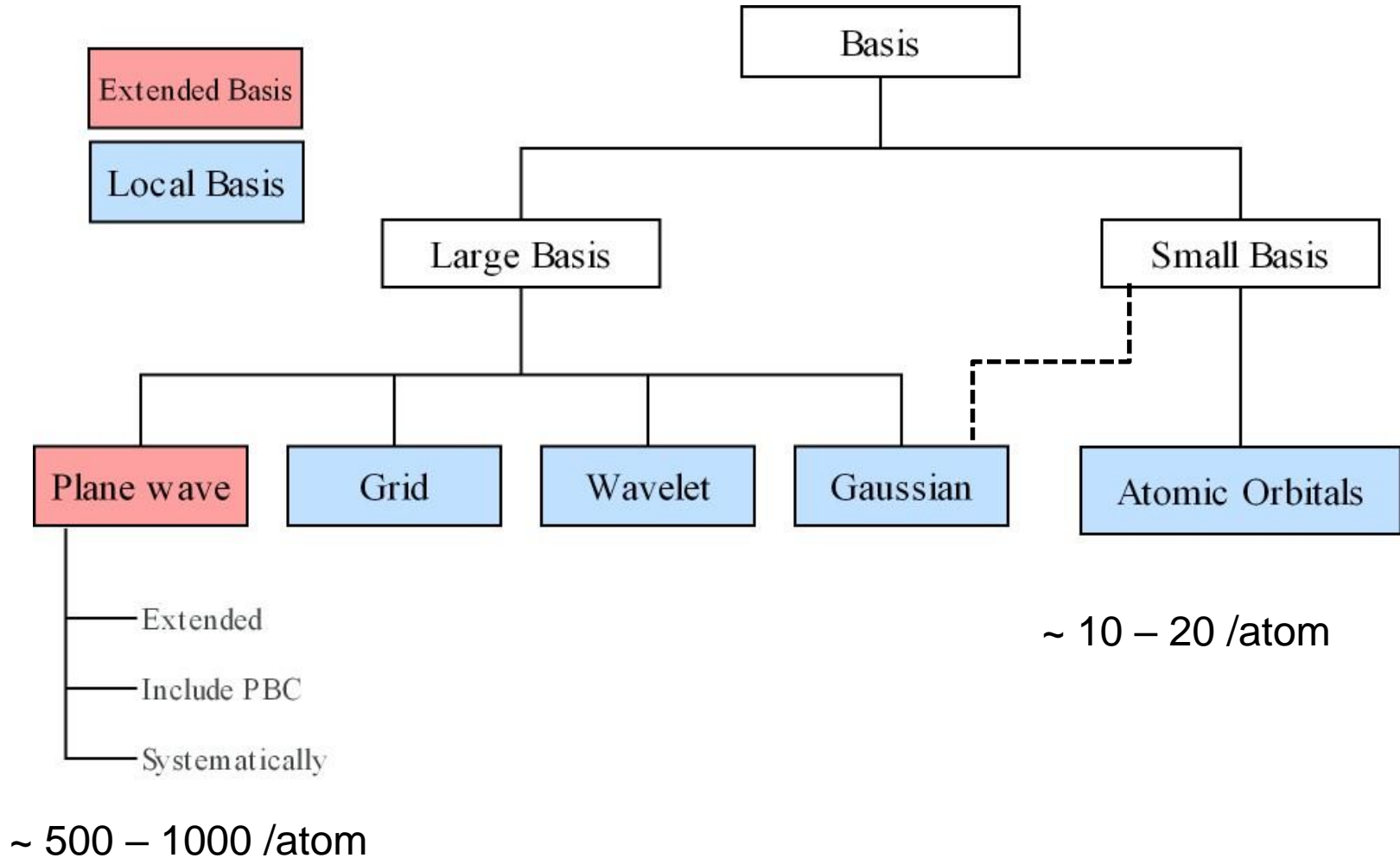
$$\sum_{q'} \langle q|H|q'\rangle c_{i,q'} = \varepsilon_i \sum_{q'} c_{i,q'} \langle q|q'\rangle$$

$$\rightarrow \sum_{q'} H(q, q') c_{i,q'} = \varepsilon_i \sum_{q'} S(q, q') c_{i,q'}$$

One of the central problems for designing efficient electronic structure algorithm is to find proper basis sets

- All electron: APW, MTO
- Pseudopotential: Plane wave, atomic orbital, etc.

Popular basis sets





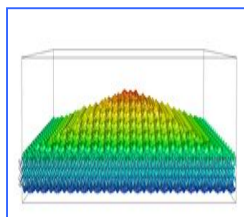
➤ Plane wave basis (PWB)

- 😊 Independent of atom positions, easy to implement.
- 😊 Forces can be easily calculated.
- 😊 Can be systematically improved by a single parameter (Ecut)
- 😊 Efficient for <100 atoms.
- 😞 Inefficient for large systems: $O(N^3)$.



Some DFT packages for thousands of atoms

CONQUEST
(England and Japan)



Basis : Numerical orbitals
/ finite element basis

ONETEP
(England)



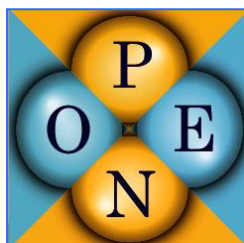
Basis : real space grid/
Wannier functions

SIESTA
(Spain)



Basis : Numerical orbitals

OpenMX
(Japan)



Basis : Numerical orbitals

FHI-aims
(Germany, FHI)



Basis : Numerical orbitals
All electron code

Bloch wave function in plane wave bases

$$\psi_{n\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{G}} e^{i(\mathbf{k}+\mathbf{G})\cdot\mathbf{r}} u_{n\mathbf{k}}(\mathbf{G})$$

Bloch wave function in numerical atomic orbital bases (LCAO)

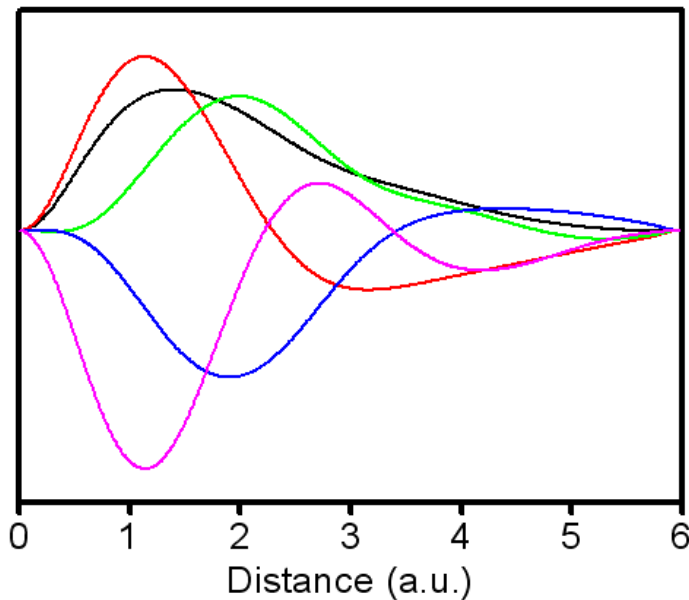
$$\psi_{n\mathbf{k}}(\mathbf{r}) = \frac{1}{\sqrt{N}} \sum_{\mathbf{R}} \sum_{\alpha} c_{n\alpha,\mathbf{k}} e^{i\mathbf{k}\cdot\mathbf{R}} \phi_{\alpha}(\mathbf{r} - \boldsymbol{\tau}_{\alpha} - \mathbf{R})$$

↑
NAO

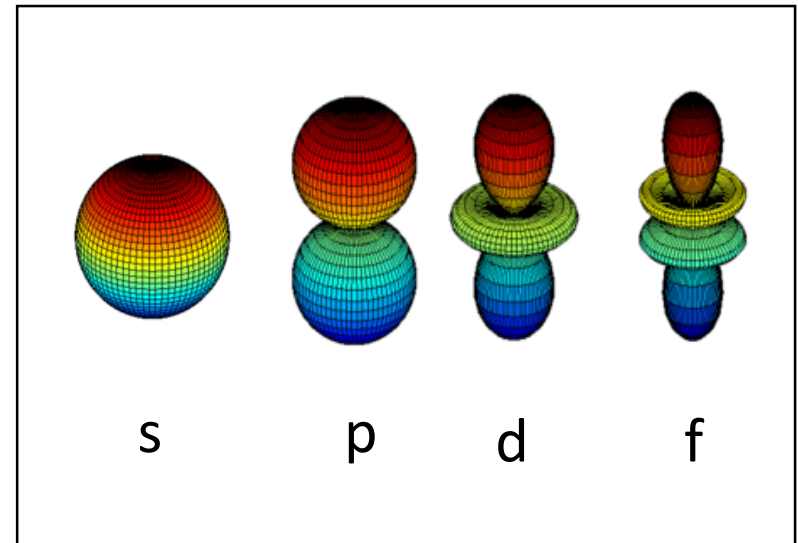
Atomic orbitals

$$\phi_{I,l,m,n}(\vec{r}) = R_{I,l,n}(r)Y_{l,m}(\hat{r}_I)$$

Radial wave functions
(size, range, shape)

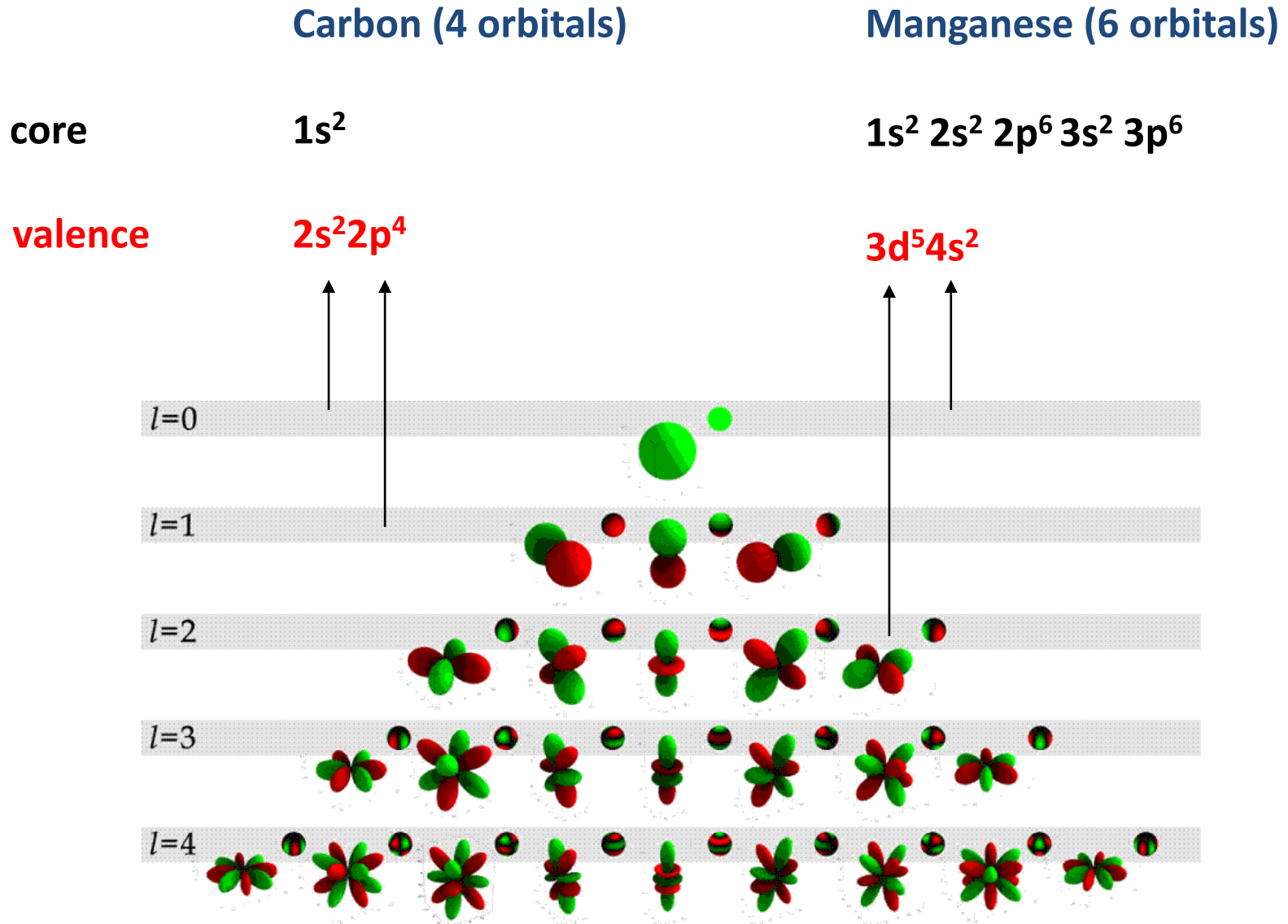


Spherical harmonics



Adapt to any symmetry!

Examples of atomic orbitals



How to improve the atomic basis sets

Low Accuracy

Minimal basis
set



High Accuracy

(need more computational resource)

Multiple-zeta orbitals
+
Polar orbitals

Advantages

- Small basis sets.
- Local basis, suitable for $O(N)$ algorithms.

Two problems

- Accuracy
- Transferability

Systematically improvable?

Atomic basis sets

- ◆ **Multi-zeta orbitals:** Different radial wave functions for the same angular momentum
- ◆ **Polar orbitals:** Orbitals with higher angular momentum than valence orbitals.
- **Systematically improvement:** 'SZ', 'DZ', 'DZP', 'TZDP'...

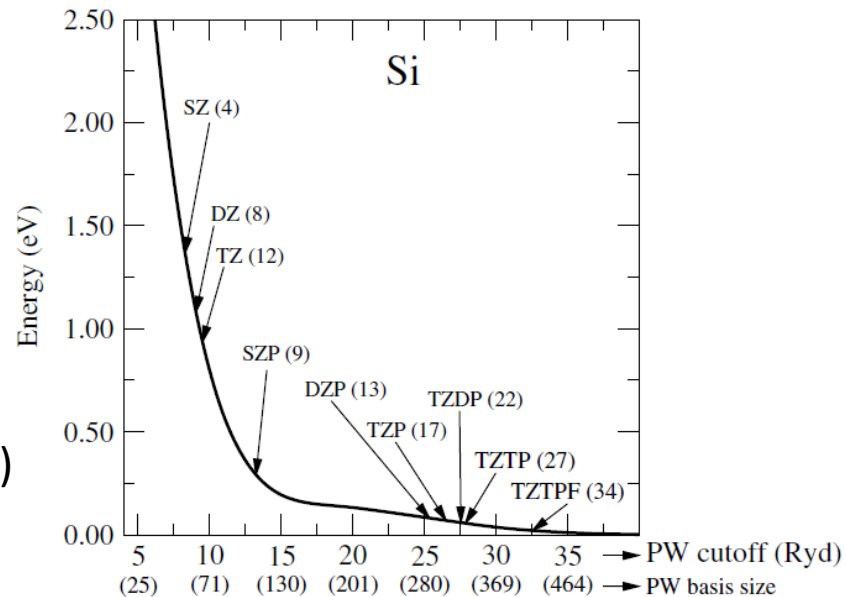
SZ: Single-zeta (minimal basis set sp)

DZ: Double-zeta (2s2p)

DZP: Double-zeta + Polar orbital (2s2p1d)

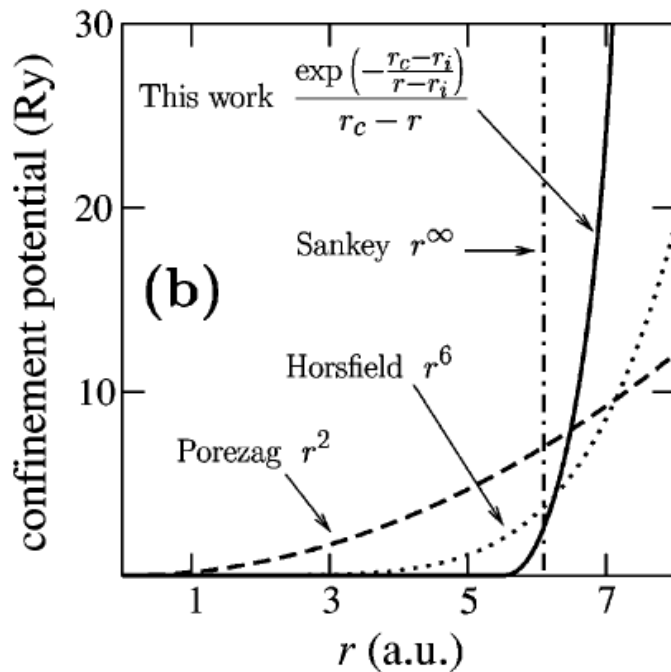
TZDP: Triple-zeta + double Polar orbital (3s3p2d)

Taken Si as an example

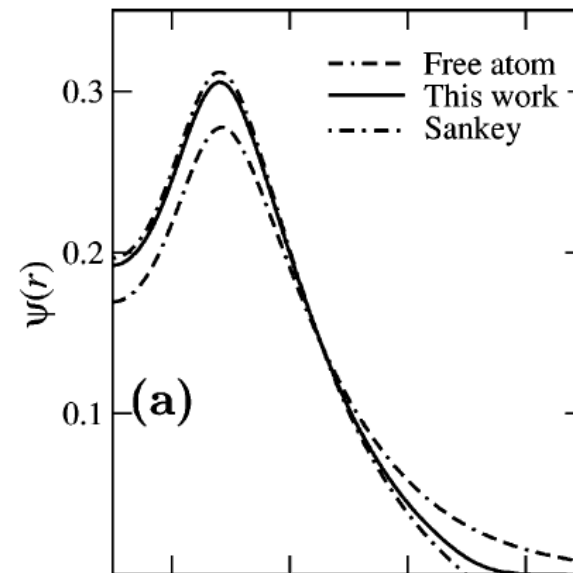


How to construct atomic orbitals

$$\left[-\frac{\nabla^2}{2} + V_{atom}(r) + V_{confine}(r) \right] \phi_{lm}(r) = \varepsilon_n \phi_{lm}(r)$$



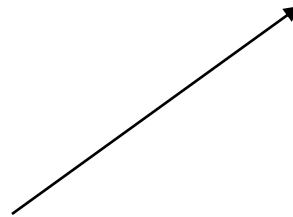
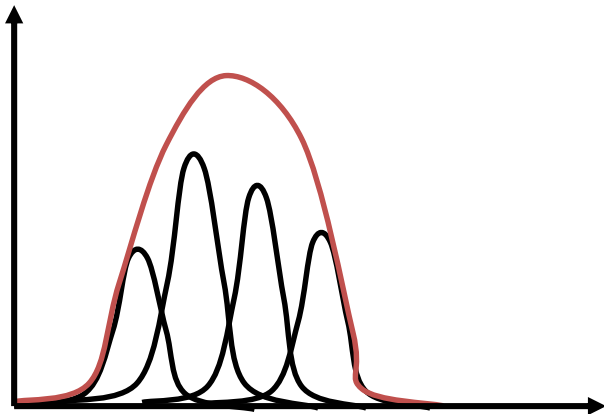
Confinement potential



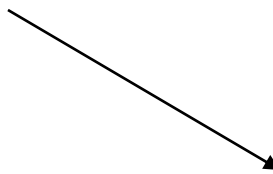
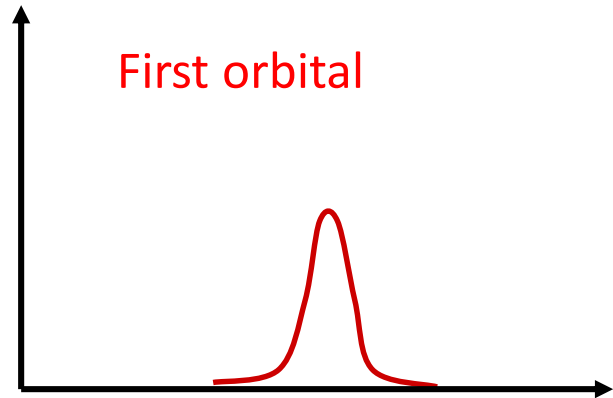
The shape of orbitals

How to construct multi-zeta orbitals

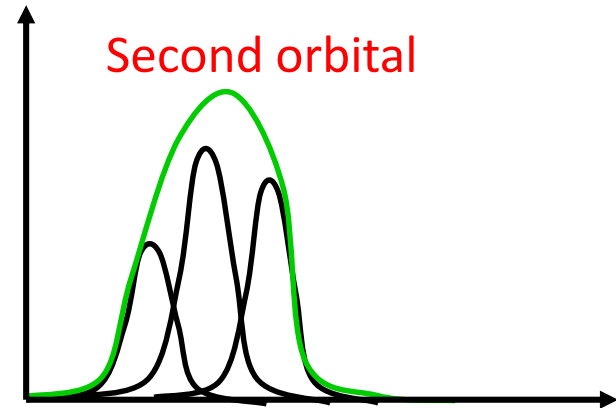
Split Valence (Gaussians)



First orbital



Second orbital



How to construct polarized orbitals

Method 2: Apply electric field

$$\left[-\frac{\nabla^2}{2} + V_{atom}(r) + V_{confine}(r) + E \cdot r \right] \psi_n = \varepsilon_n \psi_n$$



Selection rule

$$\left[\frac{1}{2r} \frac{d^2}{dr^2} r + \frac{(l+1)(l+2)}{2r^2} + V_r(l) - E_l \right] \phi_{l+1} = -\epsilon r \phi_l(r) c_{l+1}$$

How to construct polar orbitals

Compare method 1 and method 2

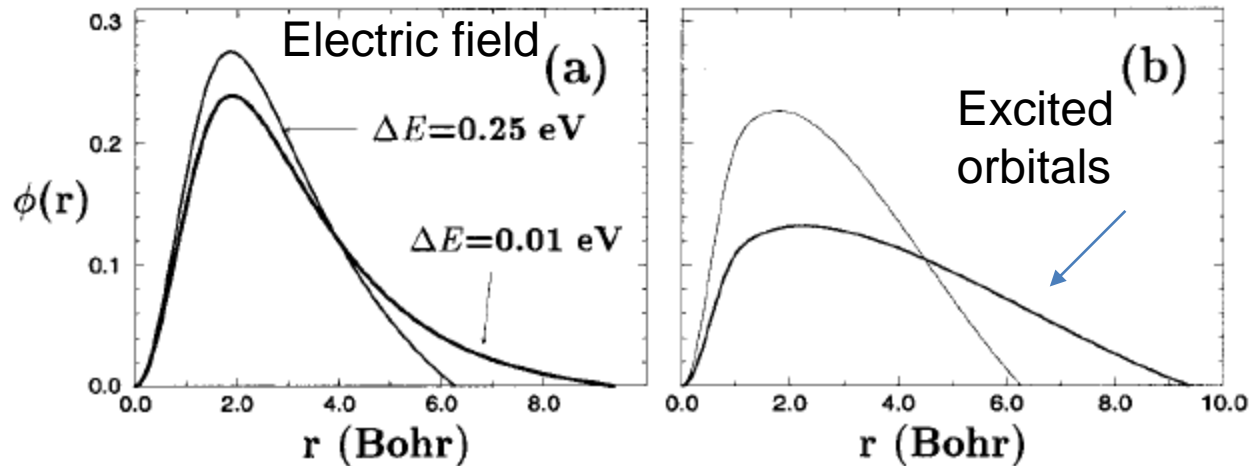
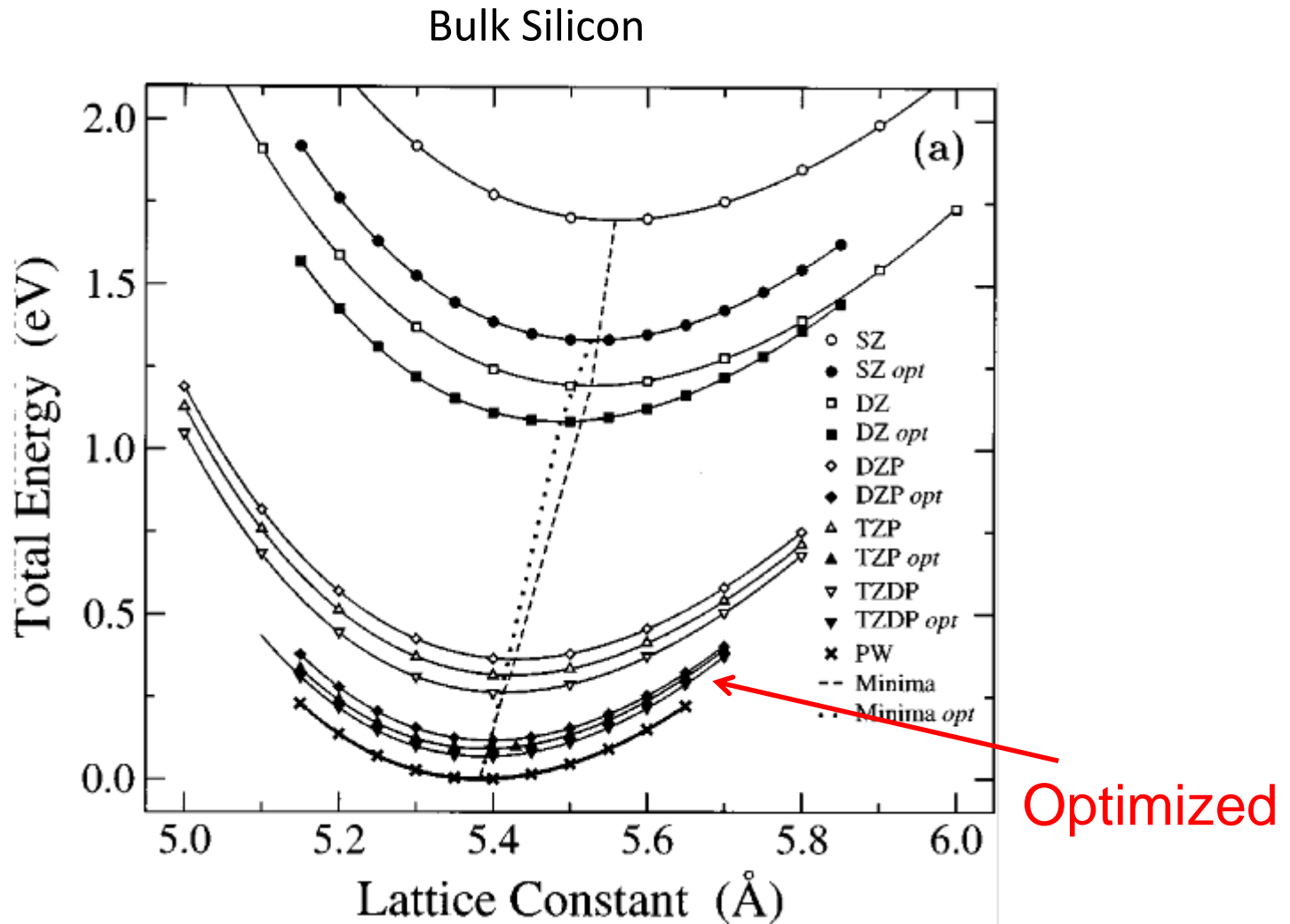


Fig. 3. d polarization orbitals for silicon for two different confinement conditions. a) Obtained with the electric-field polarization method and b) the confined d PAOs

Scheme (2)

Scheme (1)

Test atomic basis on bulk Si



Problems of previous orbital generation methods

1. Complicated especially for multi-zeta and polar orbitals.

2. The orbitals are not fully optimized.

Our Method: spillage minimization

What is "Spillage": $\mathcal{S} = \frac{1}{N_n} \sum_{n=1}^{N_n} \langle \Psi_n | 1 - \hat{P} | \Psi_n \rangle$ → Reference states

$$\hat{P} = \sum_{\mu\nu} | \phi_\mu \rangle \mathcal{S}_{\mu\nu}^{-1} \langle \phi_\nu | \quad \text{Local orbital space}$$



Cup: Hilbert Space spanned by reference states (plane wave).

Hilbert Space spanned by Local orbitals

Hilbert space left : spillage

Ice Block: atomic orbitals

D. S. Portal, et al. Solid State Communications **95**, 685 (1995)

Radial functions

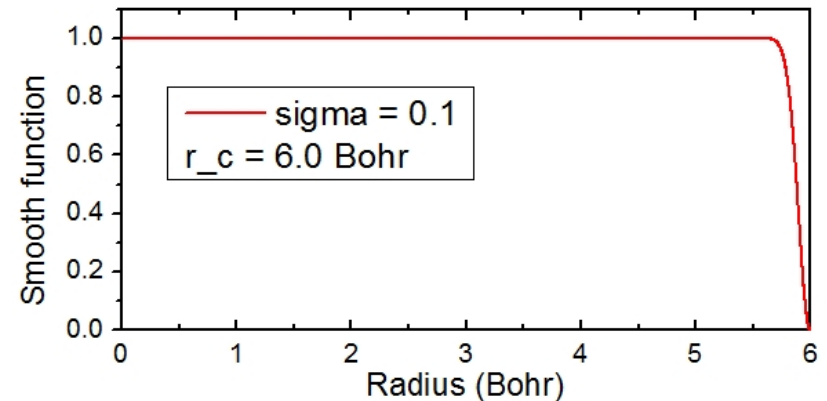
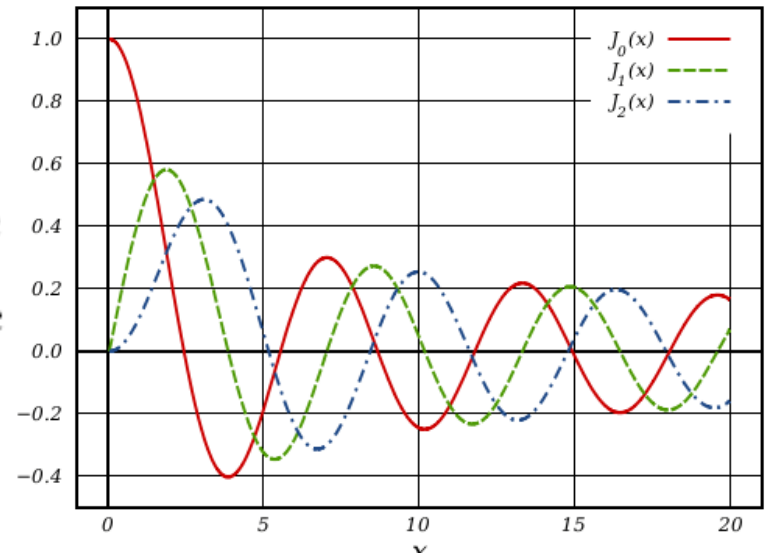
Spherical Bessel functions:

$$f_{\mu,l}(\mathbf{r}) = \begin{cases} \sum_q c_{\mu q} j_l(qr), & r < r_c \\ 0 & r \geq r_c \end{cases}$$

Very flexible!

To ensure continuity, multiply f by a smooth function:

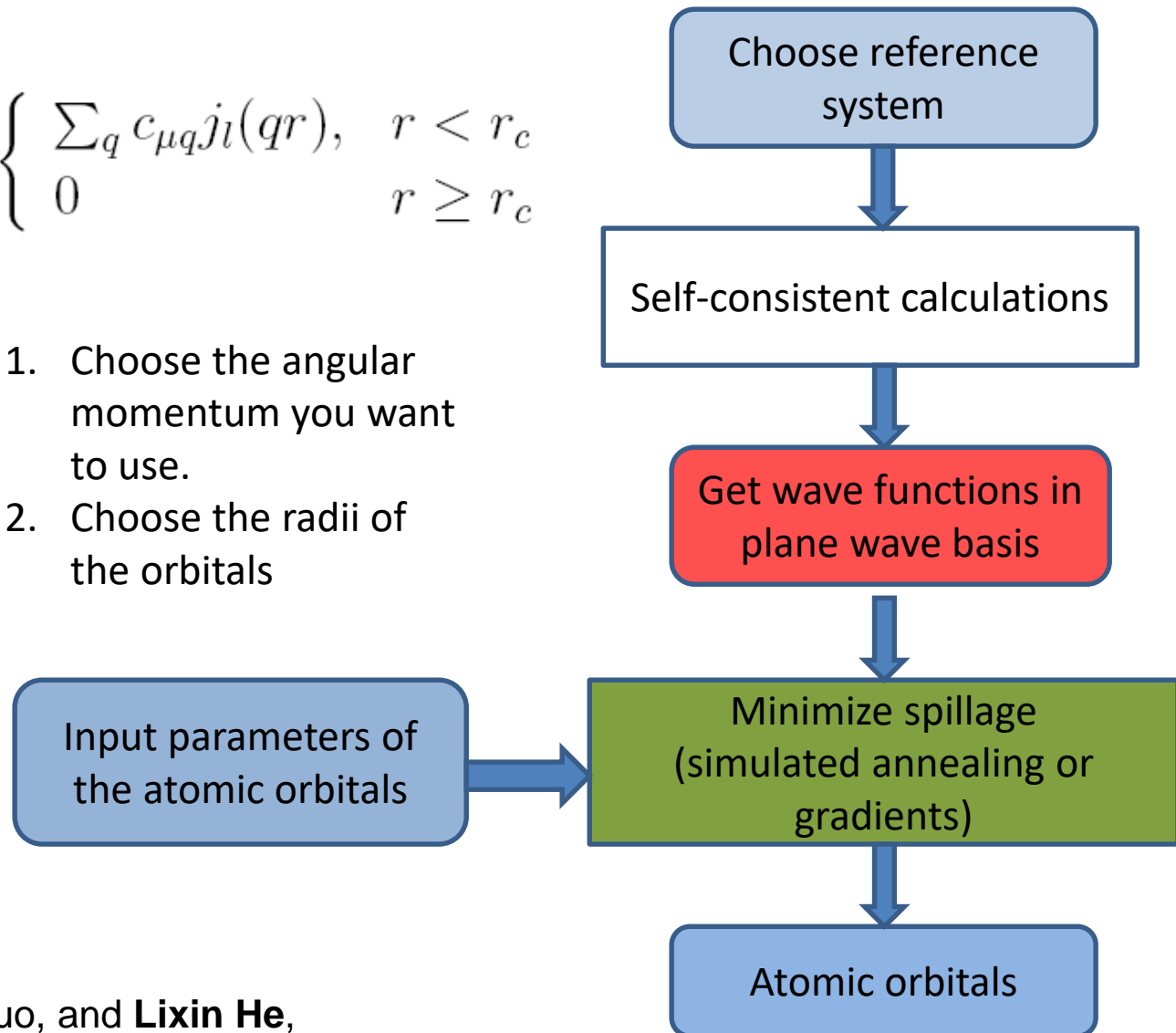
$$g(r) = 1 - \exp\left[-\frac{(r - r_{\text{cut}})^2}{2\sigma^2}\right]$$



Steps to get the orbitals

$$f_{\mu,l}(\mathbf{r}) = \begin{cases} \sum_q c_{\mu q} j_l(qr), & r < r_c \\ 0 & r \geq r_c \end{cases}$$

1. Choose the angular momentum you want to use.
2. Choose the radii of the orbitals

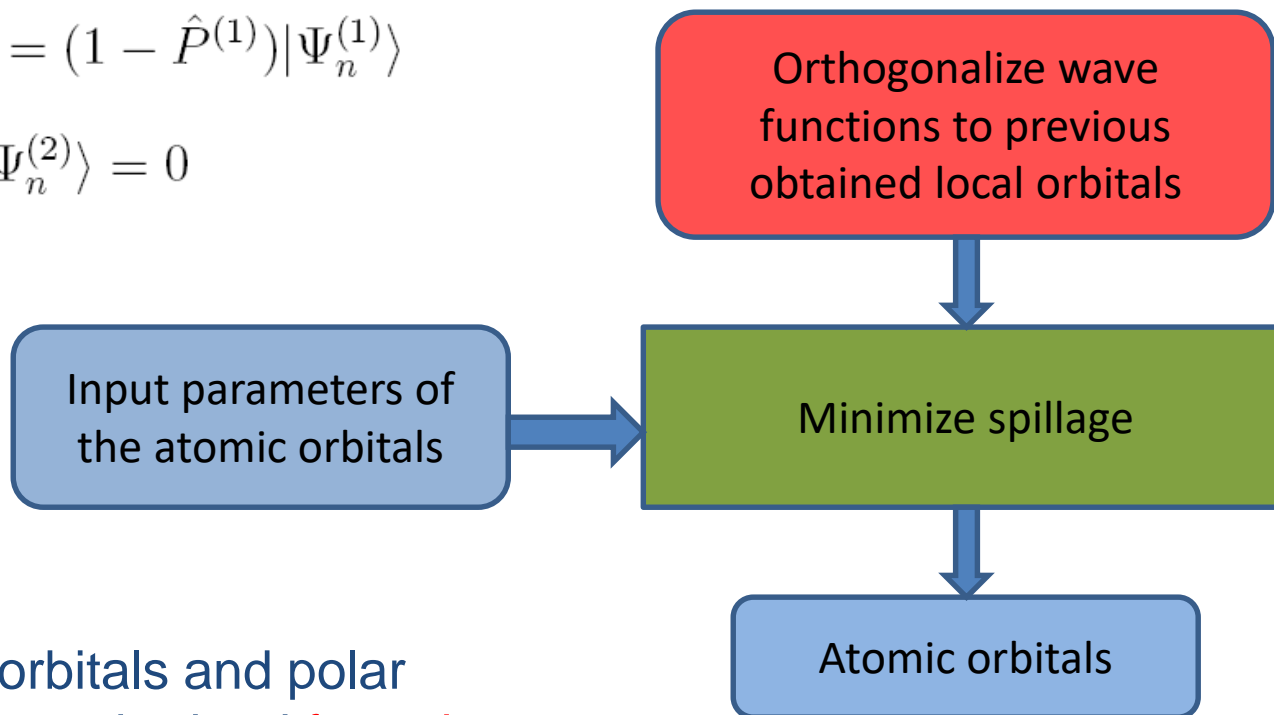


Systematically generate atomic basis

$$\hat{P}^{(1)} = \sum_{\mu\nu} |\phi_{\mu}^{(1)}\rangle S_{\mu\nu}^{-1} \langle\phi_{\nu}^{(1)}|$$

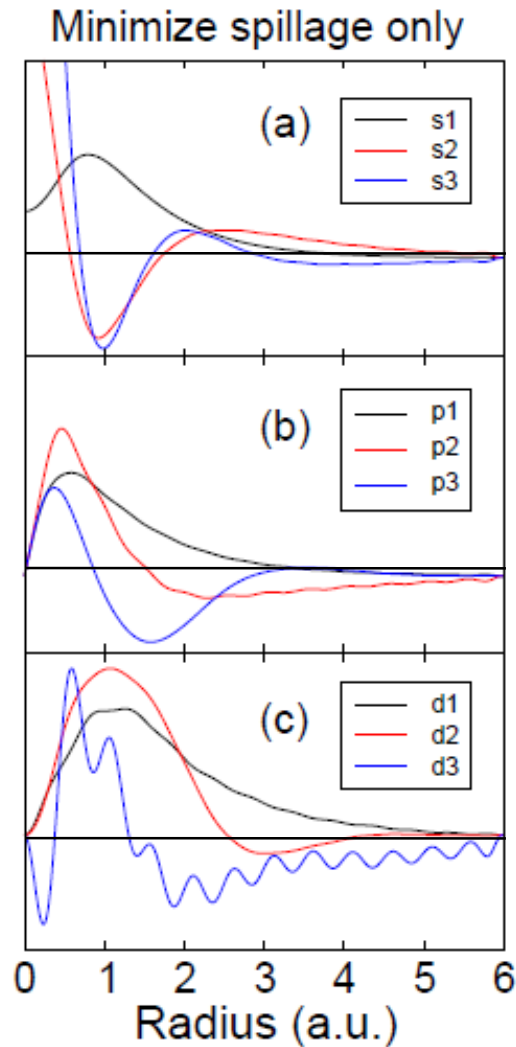
$$|\Psi_n^{(2)}\rangle = (1 - \hat{P}^{(1)})|\Psi_n^{(1)}\rangle$$

$$\hat{P}^{(1)}|\Psi_n^{(2)}\rangle = 0$$



All zeta orbitals and polar orbitals are obtained **from the same procedure!**

The radial functions of Carbon (6.0 Bohr)



SCF with:

- ◆ Carbon dimers of 5 bond lengths.
- ◆ LDA
- ◆ Energy cutoff = 100 Ry

Radial functions have oscillation!

Unphysical + worsen the transferability of the basis

Optimize the orbital shape

- ◆ Minimize “kinetic energy” of the orbitals

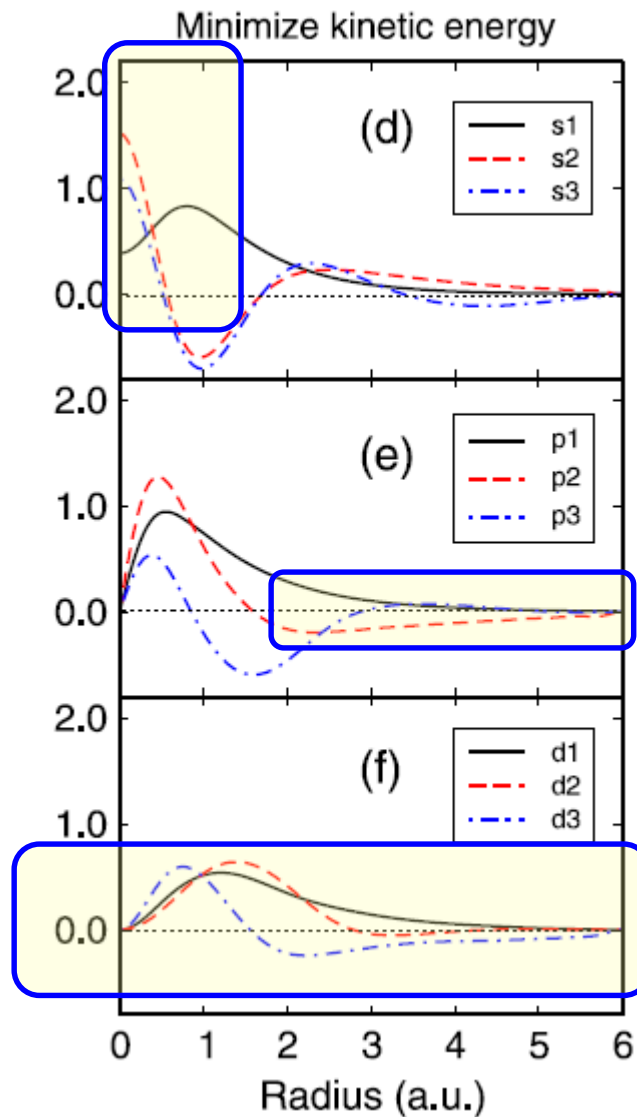
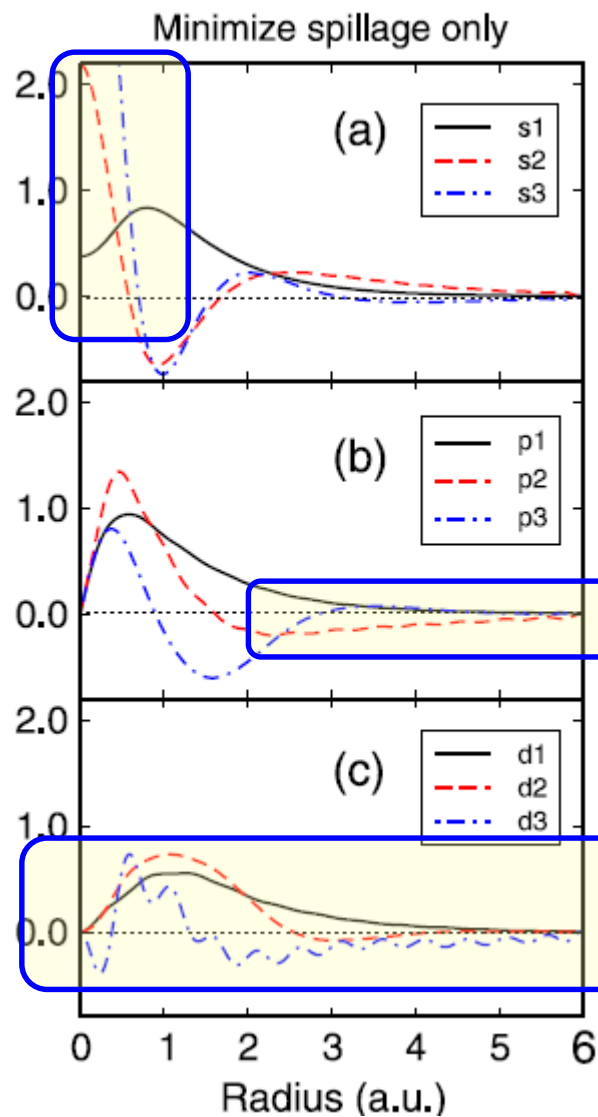
$$\min[\mathbf{T}_\mu(c_{\mu q})] = \min\left[\sum_q c_{\mu q}^2 q^2 / 2 + \kappa\right]$$

$$\kappa = \begin{cases} 0, & \mathcal{S}/\mathcal{S}_0 - 1 < \Delta \\ \infty, & \mathcal{S}/\mathcal{S}_0 - 1 > \Delta \end{cases} \quad (0.002 \sim 0.005 \text{ is enough})$$

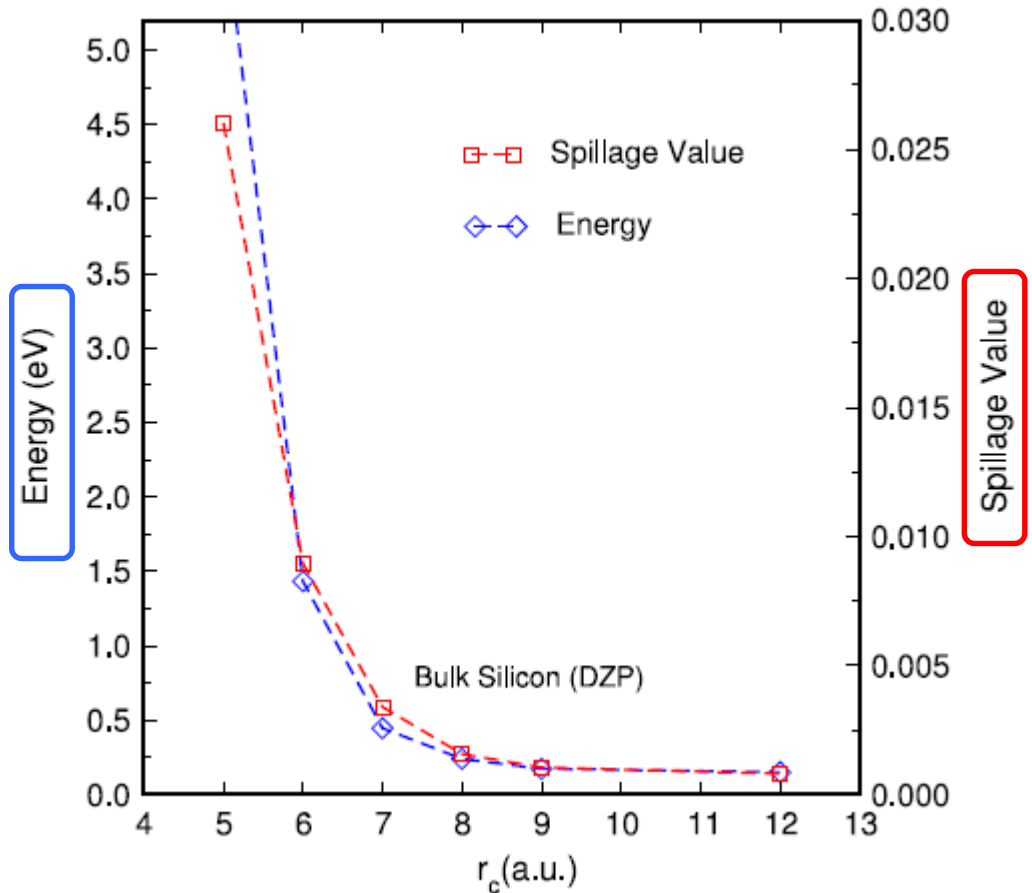
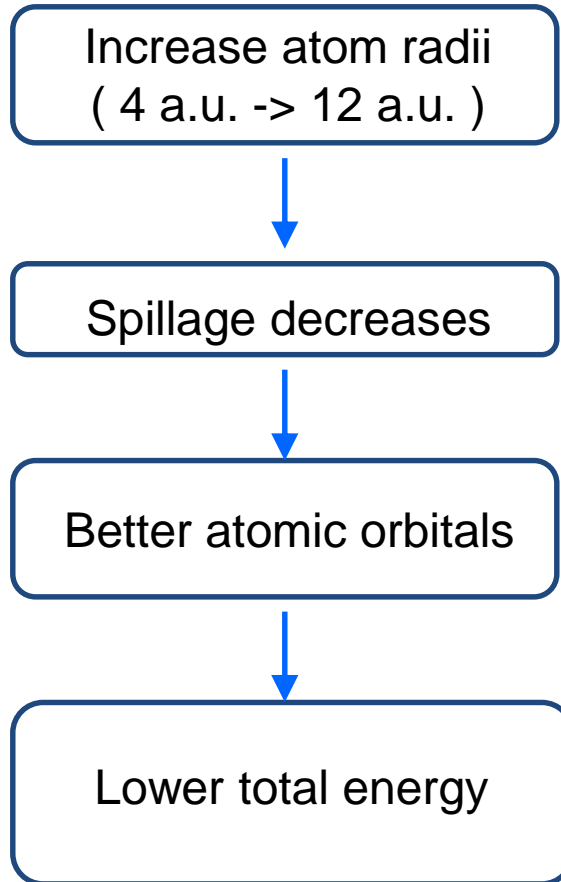
S: Current Spillage

S0: Spillage before minimize kinetic energy.

Optimized radial functions for Carbon

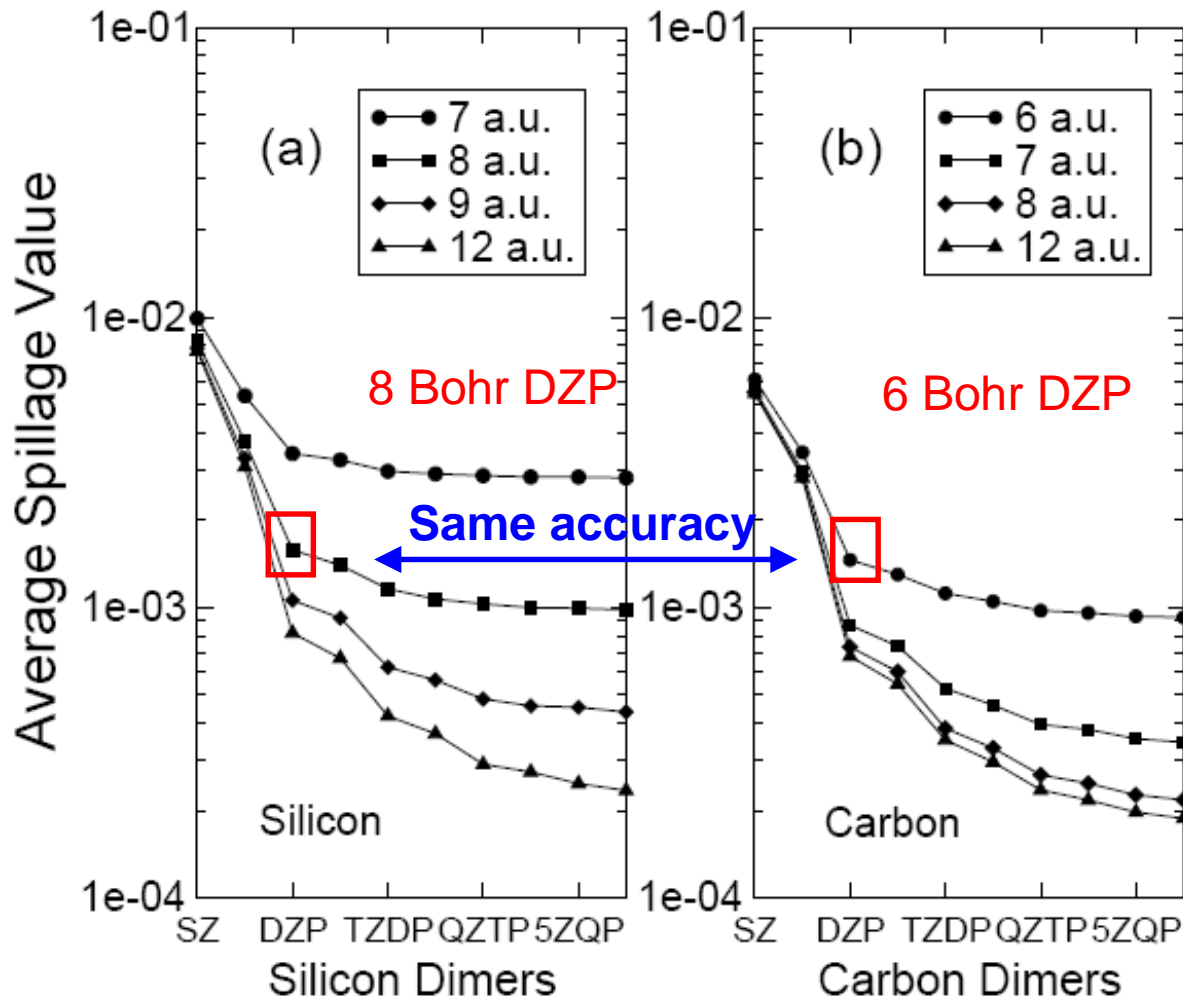


Spillage vs. Total energy



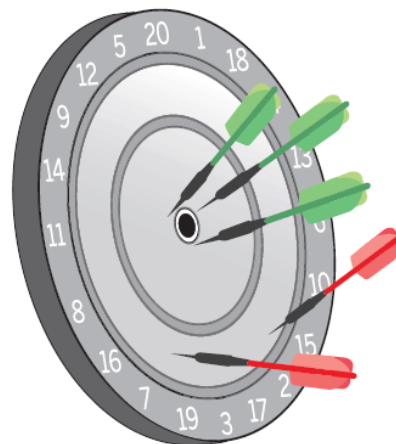
Spillage can describe very well the convergence of the total energy!

Balance the accuracy of different elements

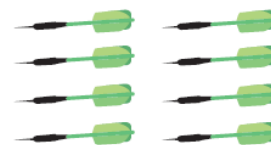


Reproducibility in density functional theory calculations of solids

Kurt Lejaeghere,^{1*} Gustav Bihlmayer,² Torbjörn Björkman,^{3,4} Peter Blaha,⁵ Stefan Blügel,² Volker Blum,⁶ Damien Caliste,^{7,8} Ivano E. Castelli,⁹ Stewart J. Clark,¹⁰ Andrea Dal Corso,¹¹ Stefano de Gironcoli,¹¹ Thierry Deutsch,^{7,8} John Kay Dewhurst,¹² Igor Di Marco,¹³ Claudia Draxl,^{14,15} Marcin Dułak,¹⁶ Olle Eriksson,¹³ José A. Flores-Livas,¹² Kevin F. Garrity,¹⁷ Luigi Genovese,^{7,8} Paolo Giannozzi,¹⁸ Matteo Giantomassi,¹⁹ Stefan Goedecker,²⁰ Xavier Gonze,¹⁹ Oscar Grånäs,^{13,21} E. K. U. Gross,¹² Andris Gulans,^{14,15} François Gygi,²² D. R. Hamann,^{23,24} Phil J. Hasnip,²⁵ N. A. W. Holzwarth,²⁶ Diana Iușan,¹³ Dominik B. Jochym,²⁷ François Jollet,²⁸ Daniel Jones,²⁹ Georg Kresse,³⁰ Klaus Koepfner,^{31,32} Emine Küçükbenli,^{9,11} Yaroslav O. Kvashnin,¹³ Inka L. M. Locht,^{13,33} Sven Lube,³⁴ Martijn Marsman,³⁰ Nicola Marzari,⁹ Ulrike Nitzsche,³¹ Lars Nordström,¹³ Taisuke Ozaki,³⁴ Lorenzo Paulatto,³⁵ Chris J. Pickard,³⁶ Ward Poelmans,^{1,37} Matt I. J. Probert,²⁵ Keith Refson,^{38,39} Manuel Richter,^{31,32} Gian-Marco Rignanese,⁴⁰ Santanu Saha,²⁰ Matthias Scheffler,^{15,40} Martin Schlipf,²² Karlheinz Schwarz,⁵ Sangeeta Sharma,¹² Francesca Tavazza,¹⁷ Patrik Thunström,⁴¹ Alexandre Tkatchenko,^{15,42} Marc Torrent,²⁸ David Vanderbilt,²³ Michiel J. van Setten,¹⁹ Veronique Van Speybroeck,¹ John M. Wills,⁴³ Jonathan R. Yates,²⁹ Guo-Xu Zhang,⁴⁴ Stefaan Cottenier^{1,45*}



New methods
Mutual agreement

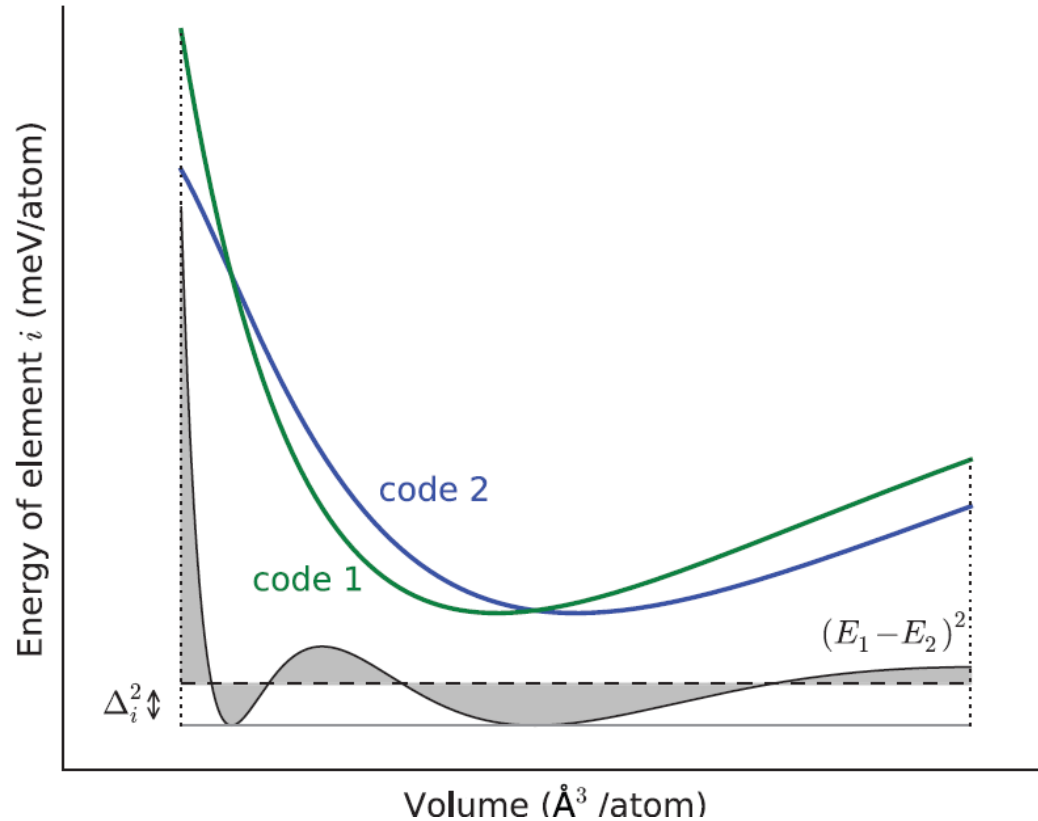


Old methods
Different values



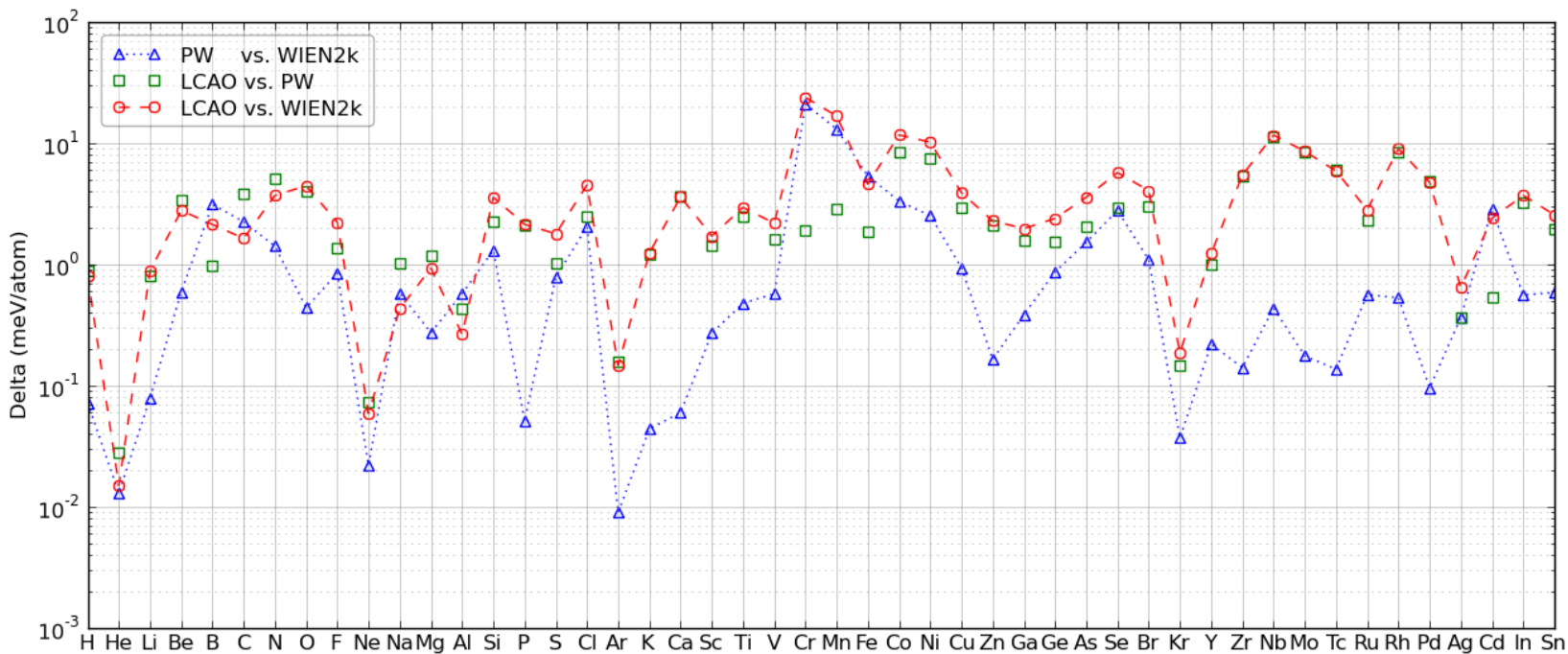
Stefaan Cottenier et al., *Science* **351**, 1415 (2016)

Fig. 3. Graphical representation of the Δ gauge. The black curve depicts the quadratic energy difference between two EOS $[(E_1 - E_2)^2]$, where the subscripts correspond to the two codes shown], and Δ_i corresponds to the root-mean-square average. This is demonstrated by the shaded area, which is equally large above and below the Δ_i^2 line.



$$\Delta_i(a, b) = \sqrt{\frac{\int_{0.94V_{0,i}}^{1.06V_{0,i}} (E_{b,i}(V) - E_{a,i}(V))^2 dV}{0.12V_{0,i}}}$$

DZP bases



LCAO-DZP vs. PW $\langle \Delta \rangle \approx 2.9 \text{ meV}$

NC(SG15) $\langle \Delta \rangle \approx 1.4 \text{ meV}$

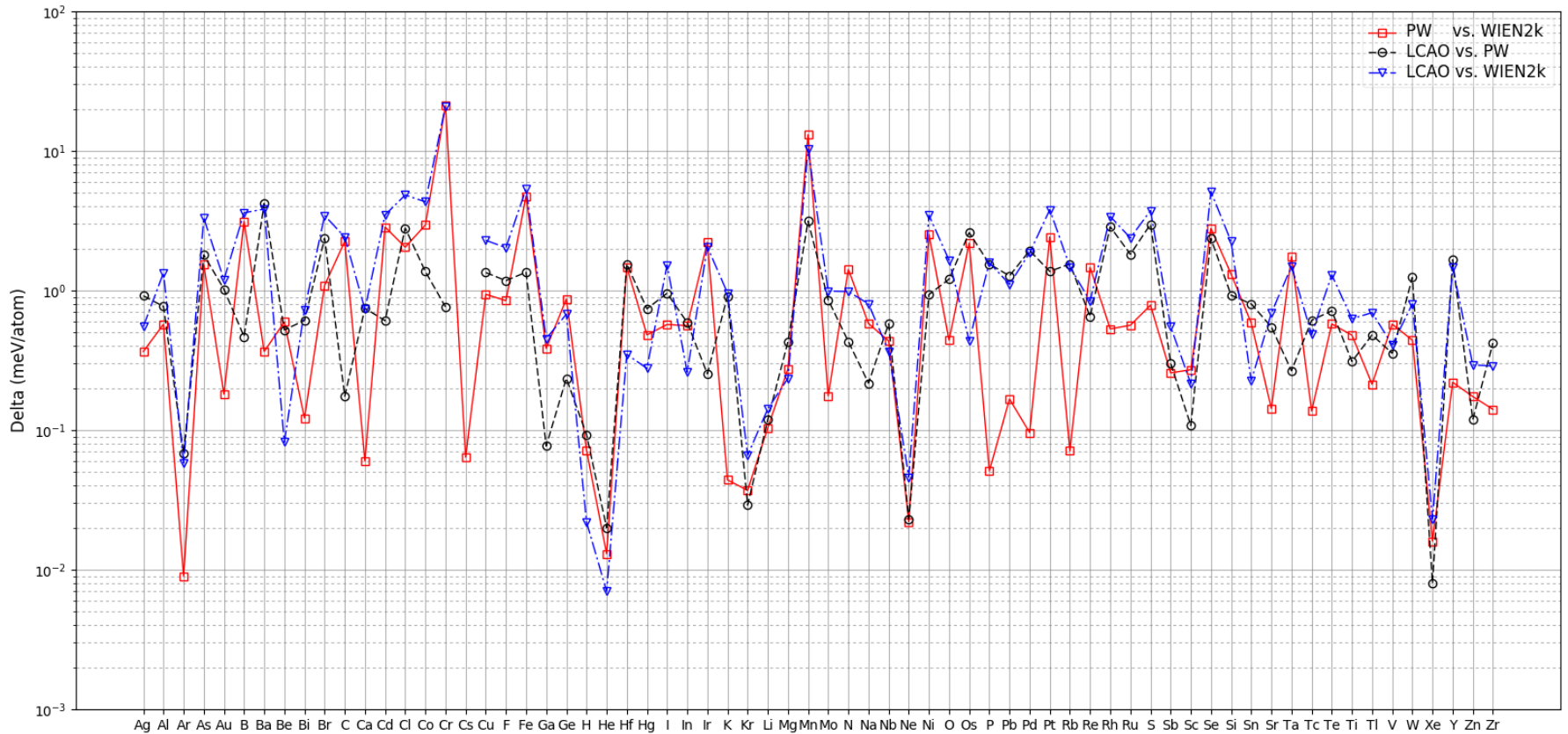
LCAO-DZP vs. WIEN2K $\langle \Delta \rangle \approx 3.0 \text{ meV}$

PAW $\langle \Delta \rangle \approx 0.6 - 1.6 \text{ meV}$

(no f orbitals for 3d, 4d elements)

USPP $\langle \Delta \rangle \approx 0.5 - 6.3 \text{ meV}$

TZDP bases



LCAO-TDZP vs. PW $\langle \Delta \rangle \approx 0.9 \text{ meV}$

SG15 $\langle \Delta \rangle \approx 1.4 \text{ meV}$

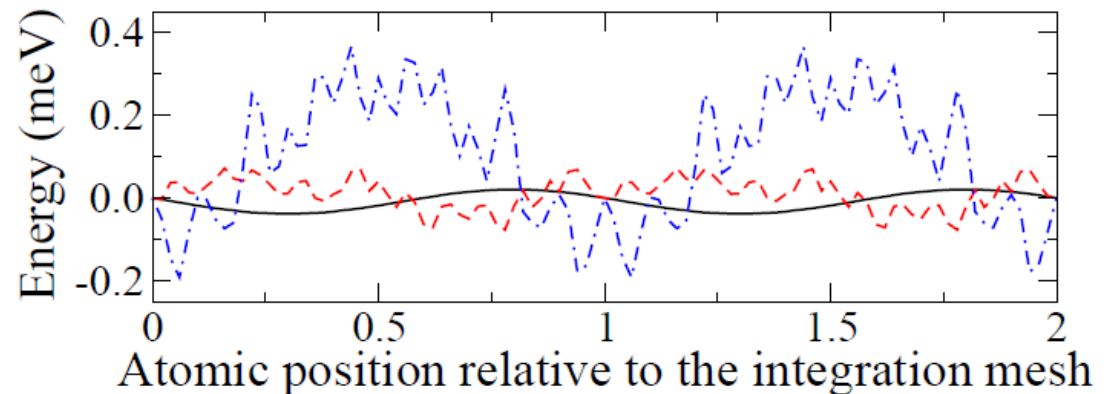
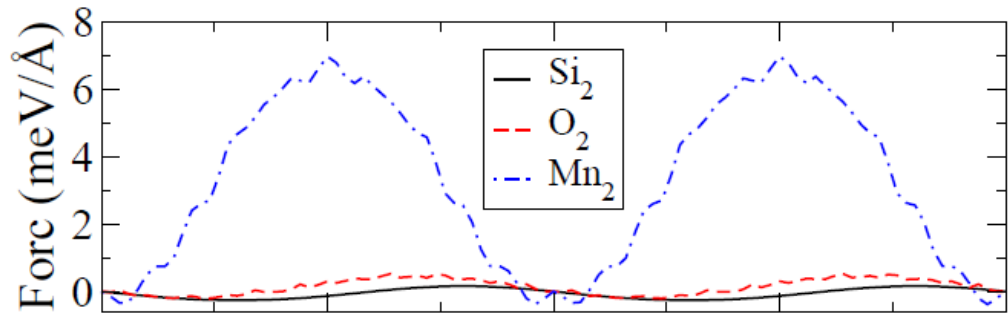
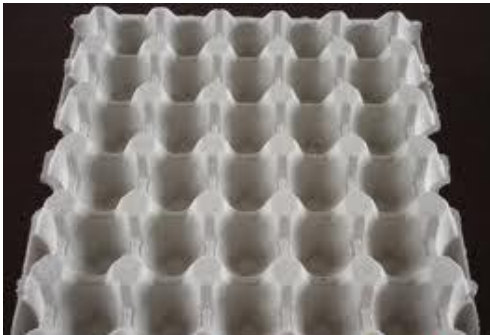
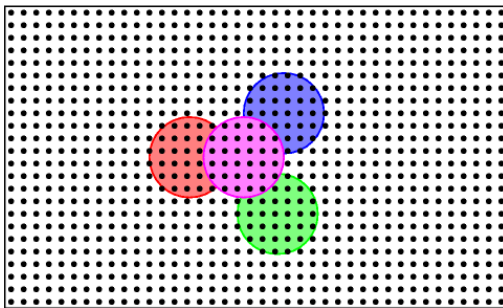
LCAO-TDZP vs. WIEN2K $\langle \Delta \rangle \approx 1.9 \text{ meV}$

PAW $\langle \Delta \rangle \approx 0.6 - 1.6 \text{ meV}$

USPP $\langle \Delta \rangle \approx 0.5 - 6.3 \text{ meV}$

Egg-Box effects

Translation symmetry broken by grid points!



SIESTA used a special filtering procedure to reduce the egg-box effects

P-F Li et al. Computational Materials Science 112 (2016) 503–517

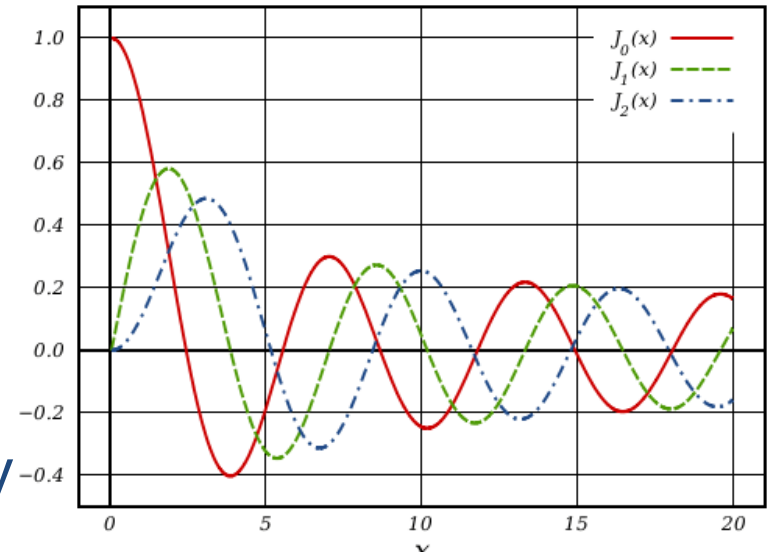
Egg-Box effects – why?

$$f_{\mu,l}(\mathbf{r}) = \begin{cases} \sum_q c_{\mu q} j_l(qr), & r < r_c \\ 0 & r \geq r_c \end{cases}$$

Strictly Localized orbitals usually have high \mathbf{G} components, resulting in egg box effects.

Our scheme allow the atomic orbitals localized both in real space and in \mathbf{G} space!

Spherical Bessel functions:



Tested Systems

- ✓ Molecules

 - Bond length, vibration frequency, atomization energy

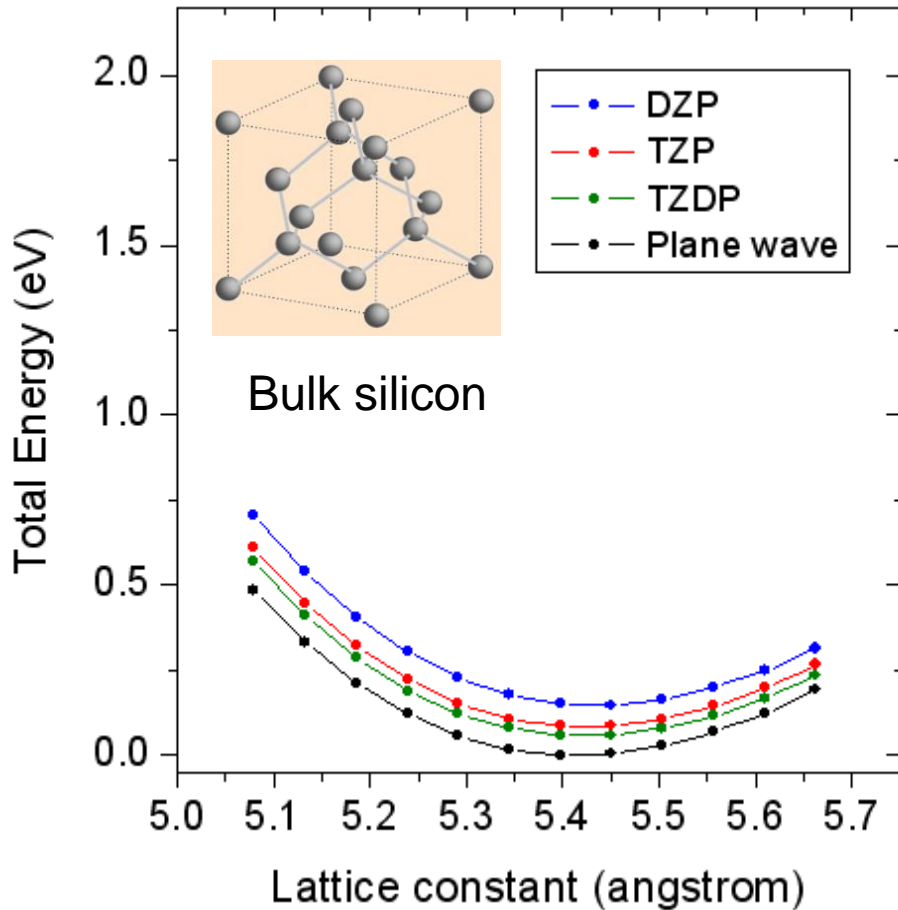
- ✓ Solid (bulk)

 - Lattice constants, cohesive energy, bulk modulus, structure

- ✓ Surface reconstruction

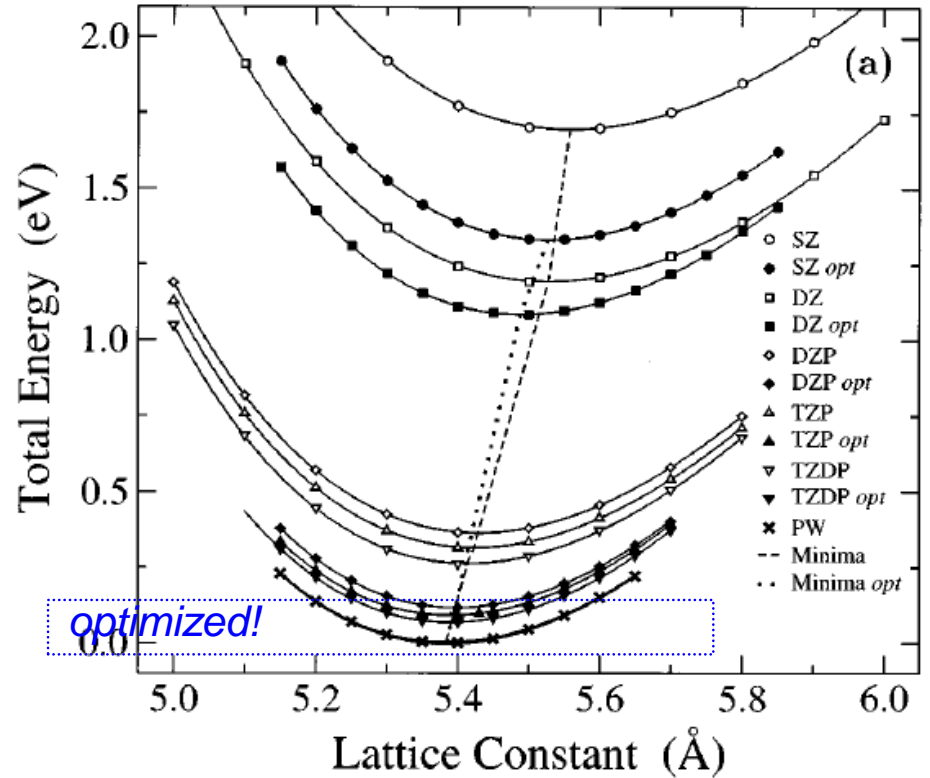
Test on bulk Si

Our orbitals



◆ References: dimers

SIESTA



◆ references: bulk Si.

Semiconductors

Compound	SZ	<i>a</i>			Expr.	SZ	<i>B</i>			Expr.
		DZP	TZDP	PW			DZP	TZDP	PW	
GaAs	10.67	10.50	10.49	10.48	10.68 ^a	69	78	77	77	75.57 ^c
GaP	10.28	10.11	10.11	10.10	10.30 ^a	82	92	93	93	89 ^d
GaSb	11.54	11.38	11.37	11.36	11.52 ^a	49	59	58	57	57 ^d
InAs	11.40	11.27	11.28	11.28	11.45 ^a	63	66	65	65	60 ^d
InP	11.07	10.94	10.94	10.93	11.09 ^a	78	79	79	80	71 ^d
InSb	12.33	12.05	12.05	12.07	12.24 ^a	41	50	49	50	47 ^d
AlAs	10.88	10.63	10.62	10.59	10.70 ^a	67	76	76	76	77 ^d
AlP	10.50	10.26	10.25	10.21	10.33 ^a	64	87	88	89	86 ^d
AlSb	11.83	11.58	11.57	11.54	11.59 ^a	48	55	56	56	58 ^d
Ge	10.82	10.68	10.61	10.61	10.70 ^c	57	67	71	71	77.20 ^c
Si	10.59	10.28	10.25	10.23	10.26 ^b	74	94	94	94	99 ^b
C ^e	6.78	6.67	6.67	6.67	6.75 ^b	436	470	467	466	442 ^b

B	C	N
Al	Si	P
Ga	Ge	As
In	Sn	Sb

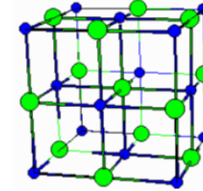
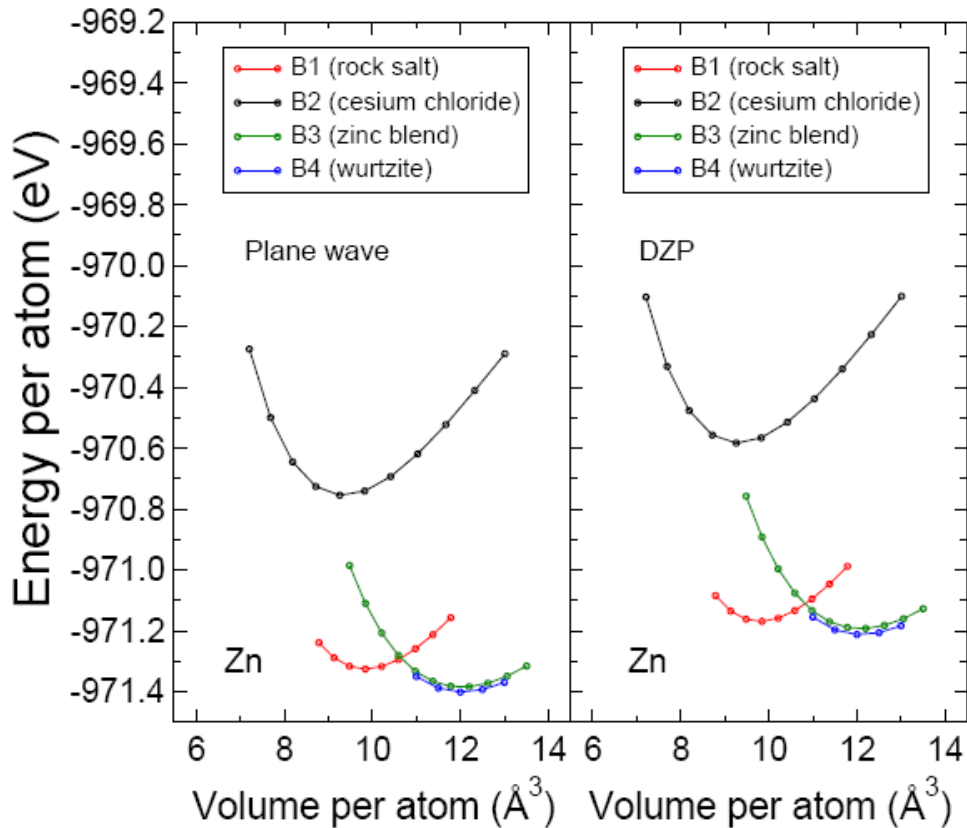
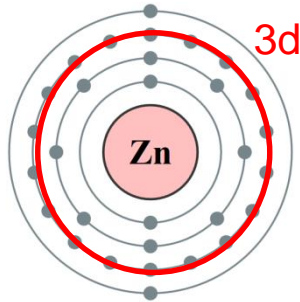
晶格常数平均误差(Bohr) : 0.218 (SZ), 0.026 (DZP), 0.015 (TZDP)

体弹性模量平均误差(GPa) : 12.3 (SZ), 1.4 (DZP), 0.4 (TZDP)

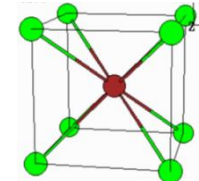
DZP 可以准确预测半导体的结构性质

Tests on ZnO

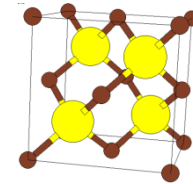
- ◆ Include 3d electrons
- ◆ Four phases



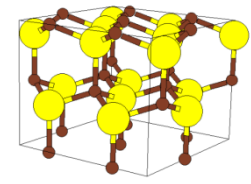
B1 (NaCl)



B2 (CsCl)



B3 (Zinc Blend)



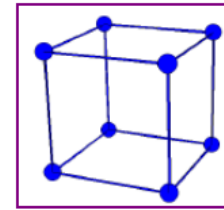
B4 (Wurtzite)

Correct energy orders and energy differences!

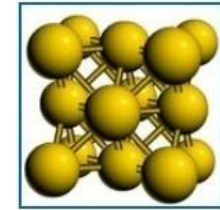
Good transferability!

Tests on Al (metal)

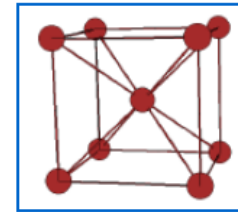
- ◆ Metal
- ◆ Four different structures



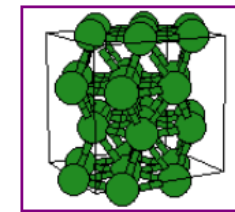
sc



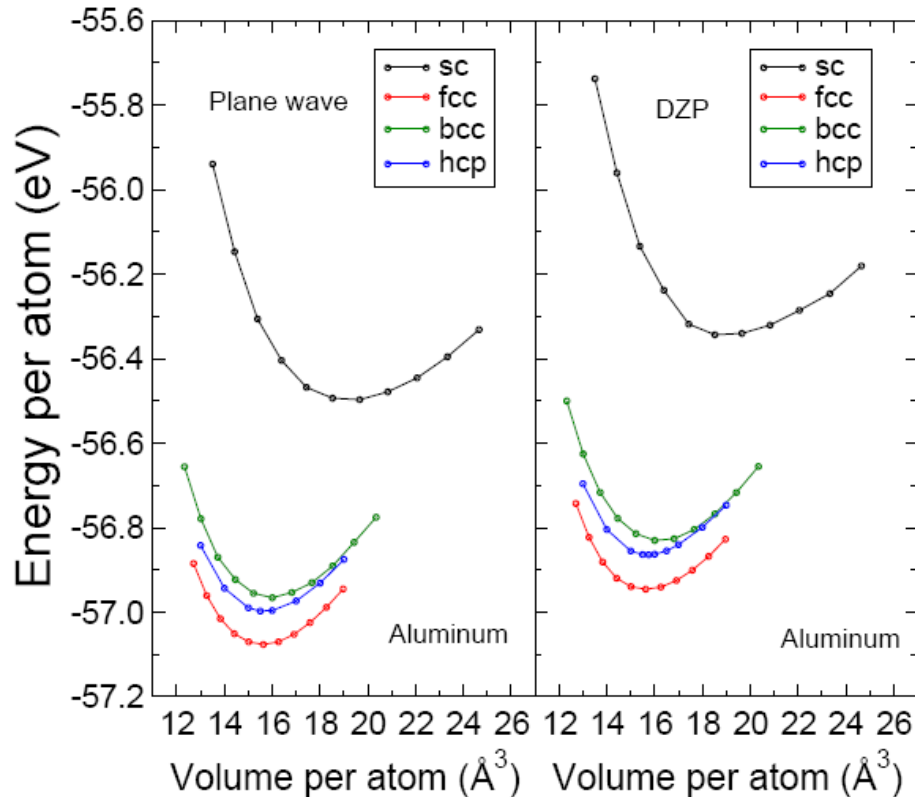
fcc



bcc



hcp



Correct energy orders and energy differences!

LCAO works well for metals and insulators.

Electronic structure interpolation via atomic orbitals

Mohan Chen, G-C Guo, and **Lixin He**, *J. Phys.: Cond. Mat.* **23** 325501 (2011)

Motivation : Large k sampling problem

The number of k points used to calculate anomalous Hall conductivity in bcc Fe:

Self consistent – 20, 000 k points

Berry curvature , larger than 400,000 – 2,000,000 k points

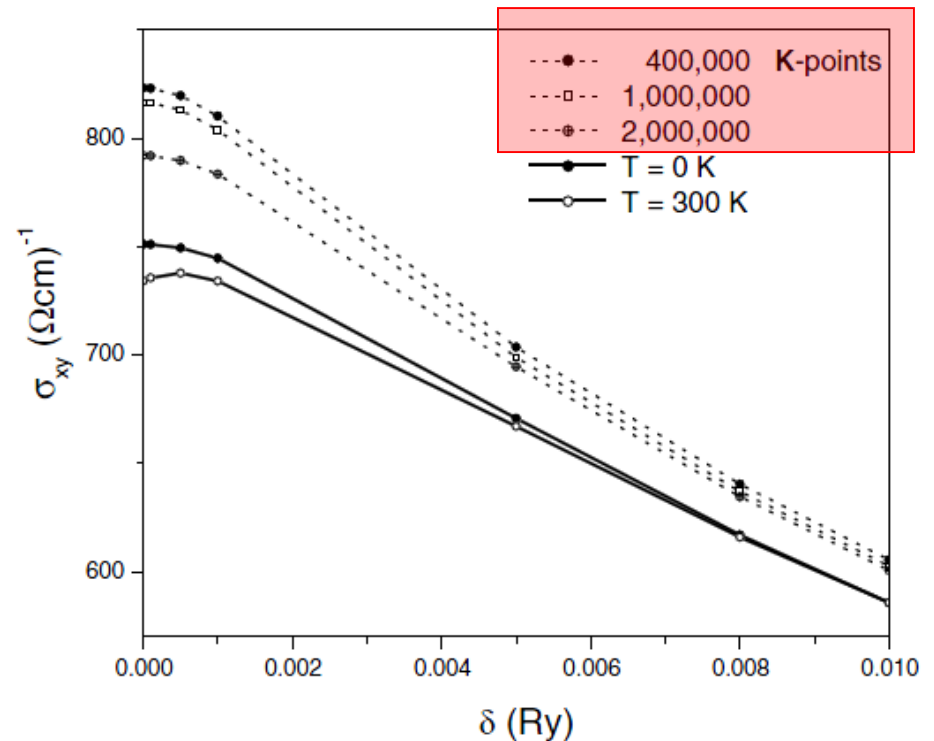
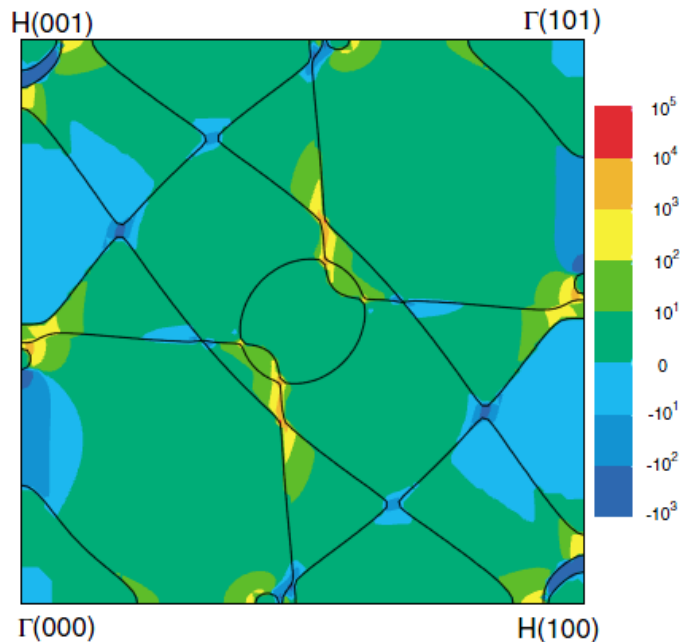


FIG. 3 (color). Fermi surface in (010) plane (solid lines) and Berry curvature $-\Omega^z(\mathbf{k})$ in atomic units (color map).

Yugui Yao et al, Phys. Rev. Lett **92**, 037204 (2004)

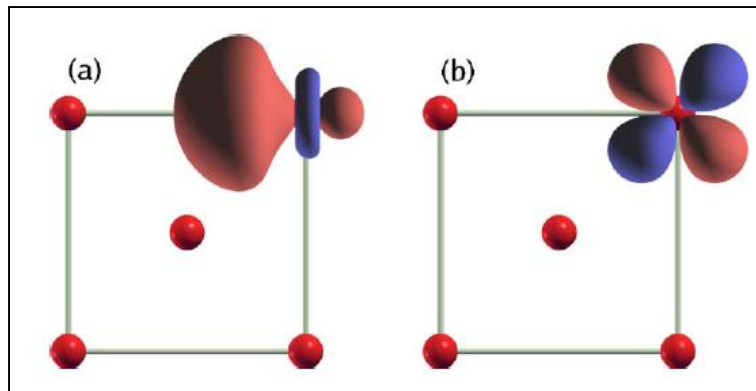
First Principle Calculation of Anomalous Hall Conductivity in Ferromagnetic bcc Fe

Band interpolation via Wannier functions

◆ Maximally localized wannier functions (MLWF)

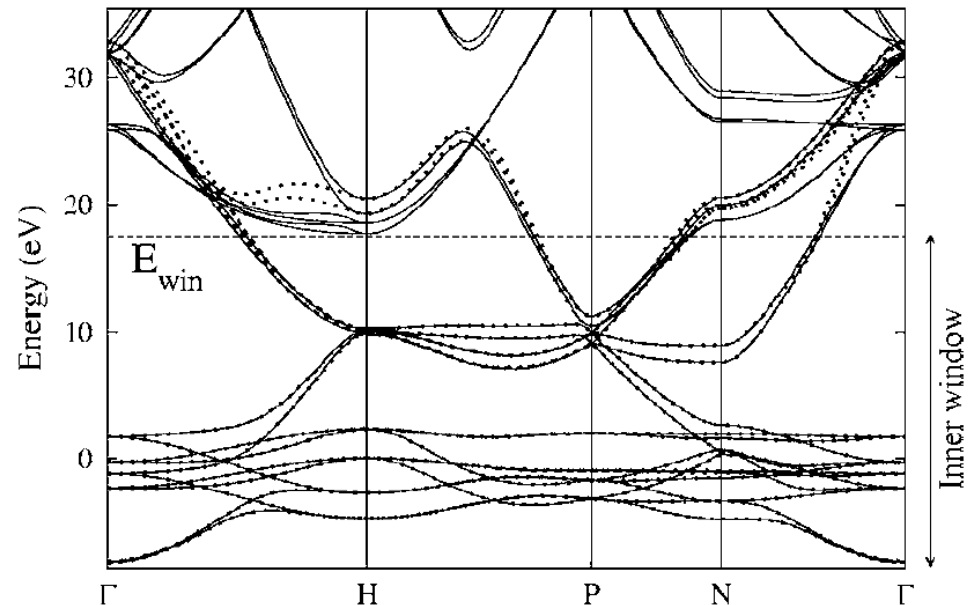
Minimize the spread functional:

$$\Omega = \sum_n [\langle r^2 \rangle_n - \bar{\mathbf{r}}_n^2]$$



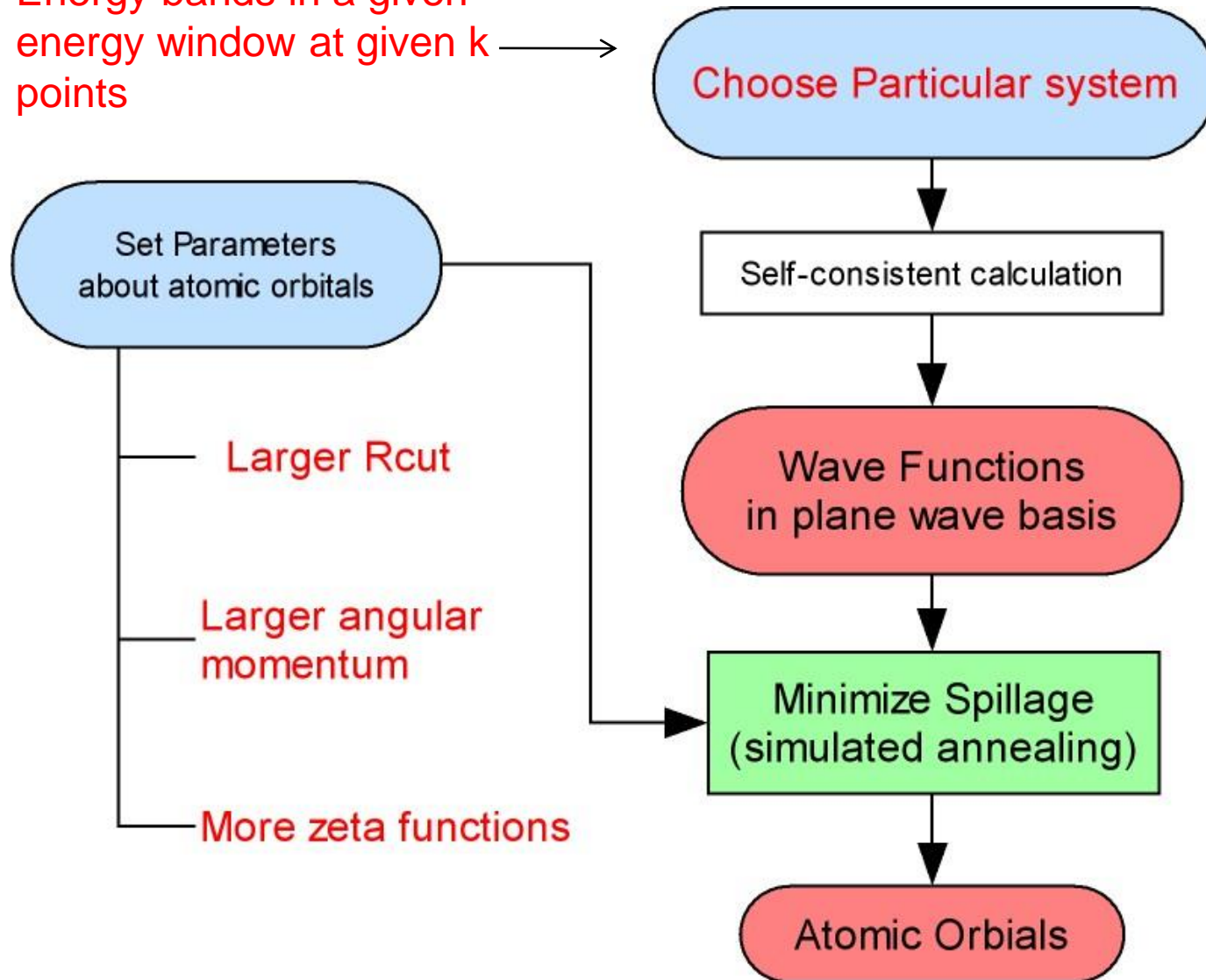
MLWF

Band structure of bcc Fe



Our Method: atomic basis

Energy bands in a given energy window at given k points



Interpolation of operators

$$H_{\mu\nu}(\mathbf{R}_n) = \langle \phi_{\mu 0} | \hat{H} | \phi_{\nu n} \rangle, \quad \text{Calculate once for all k-points}$$



$$H_{\mu\nu}(\mathbf{k}) = \sum_{\mathbf{R}_n} e^{i\mathbf{k}\cdot\mathbf{R}_n} H_{\mu\nu}(\mathbf{R}_n),$$



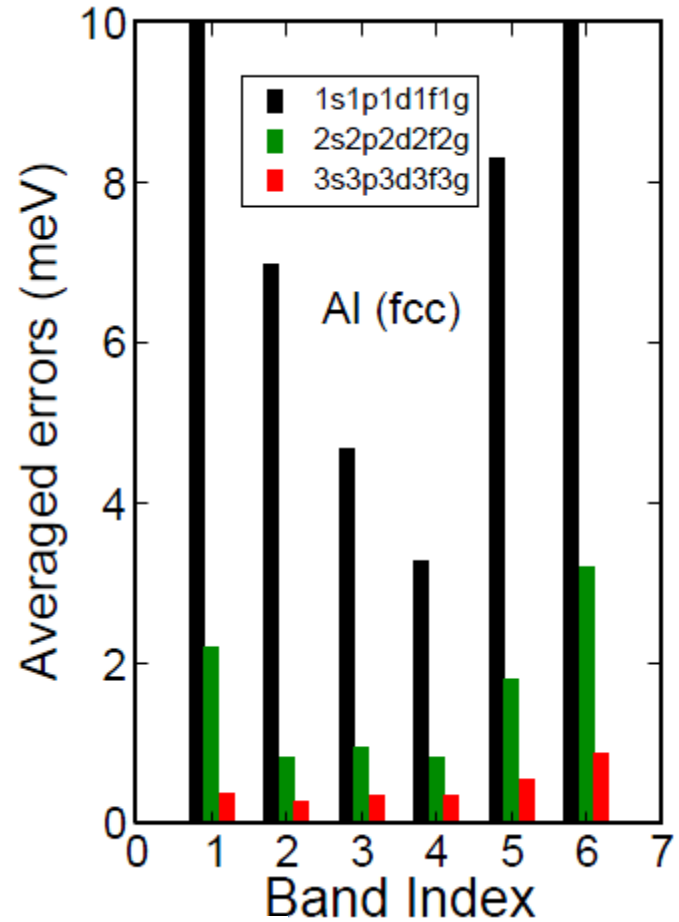
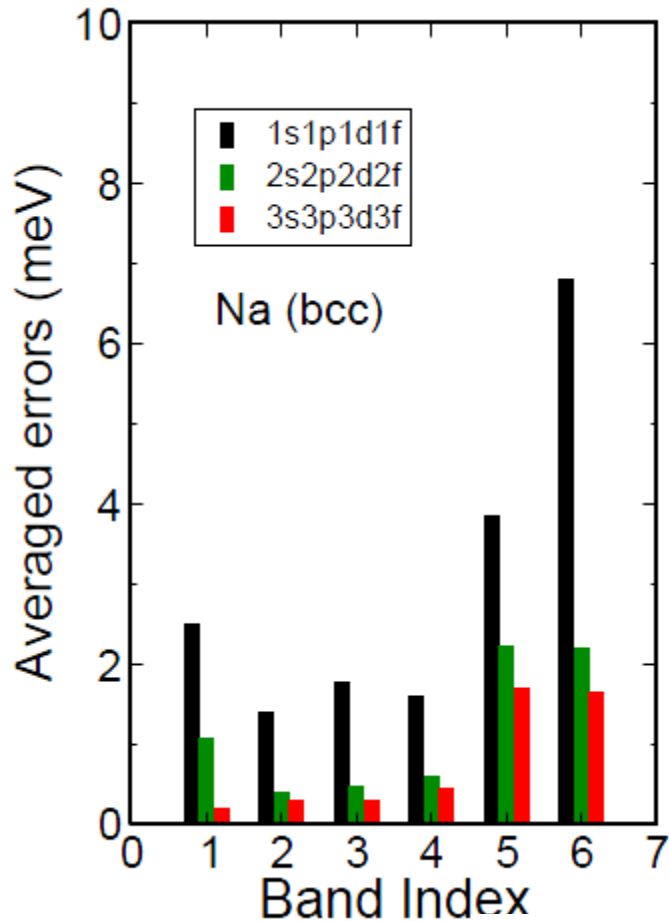
$$\Psi_{i\mathbf{k}}(\mathbf{r}) = \sum_{\mathbf{R}_n} c_{i\mu}(\mathbf{k}) \phi_{\mu}(\mathbf{r} - \mathbf{r}_{\mu} - \mathbf{R}_n) e^{i\mathbf{k}\cdot\mathbf{R}_n}.$$



$$O_{ij}(\mathbf{k}, \mathbf{q}) = \langle \Psi_{i,\mathbf{k}} | \hat{O} | \Psi_{j,\mathbf{q}} \rangle$$

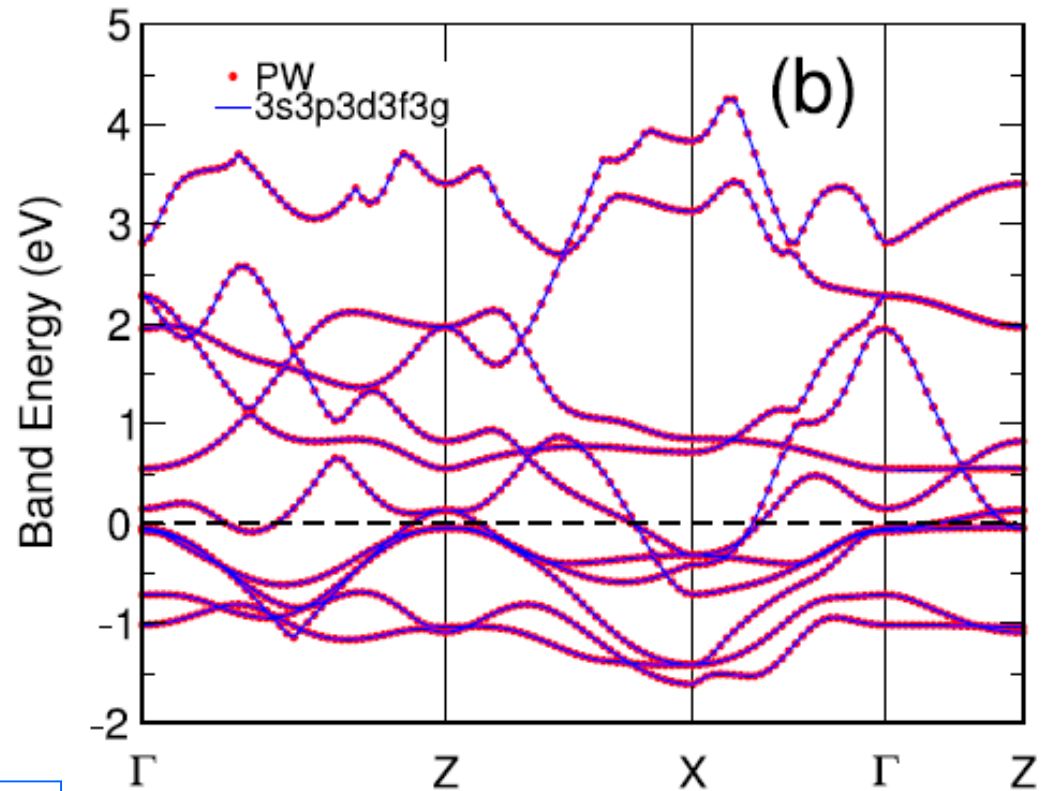
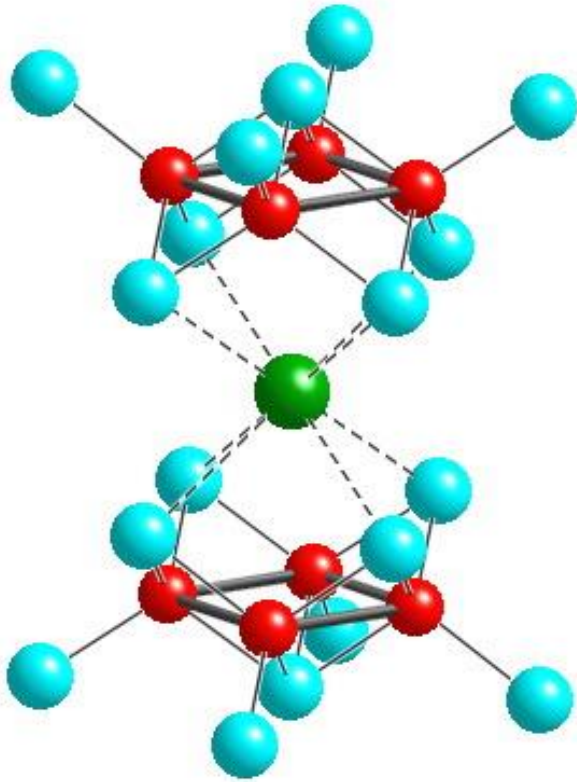
$$O_{ij}(\mathbf{k}, \mathbf{q}) = \sum_{\mathbf{R}_n \mathbf{R}_m} e^{-i\mathbf{k}\cdot\mathbf{R}_n + i\mathbf{q}\cdot\mathbf{R}_m} c_{i\mu}^*(\mathbf{k}) c_{j\nu}(\mathbf{q}) \langle \phi_{\mu n} | \hat{O} | \phi_{\nu m} \rangle.$$

Test: Na (bcc) , Al (fcc)



$$\text{RMSE} = \sqrt{\frac{\sum_{i=1}^N (\epsilon_i^{\text{PW}} - \epsilon_i^{\text{LCAO}})^2}{N}}$$

Test: BaFe_2As_2



平面波: 5319

1s1p1d1f1g: 125(快884倍)

2s2p2d2f2g: 250 (快212倍)

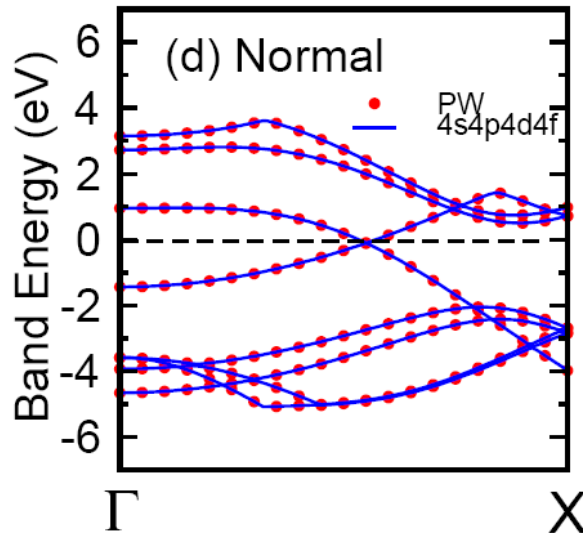
3s3p3d3f3g: 375 (快81倍)

2s2p2d2f2g is accurate
enough! ($< 3\text{meV}$)

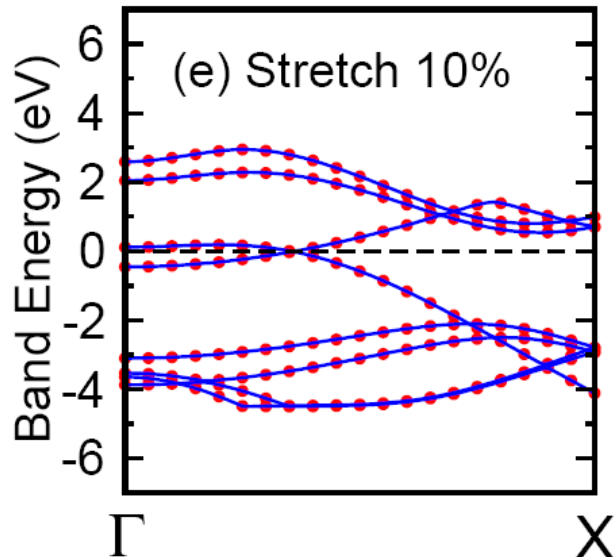
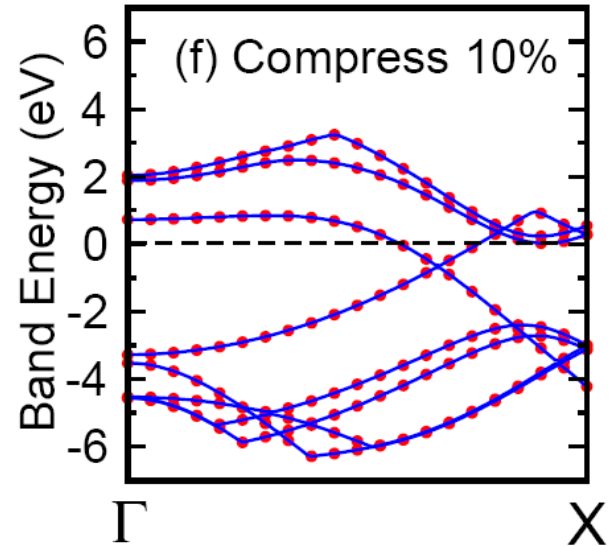
Transferability of the bases

- ◆ Atomic orbitals generated under normal conditions (10 Bohr, 4s4p4d4f)

planewaves: 17221
Atomic bases(4s4p4d4f): 768



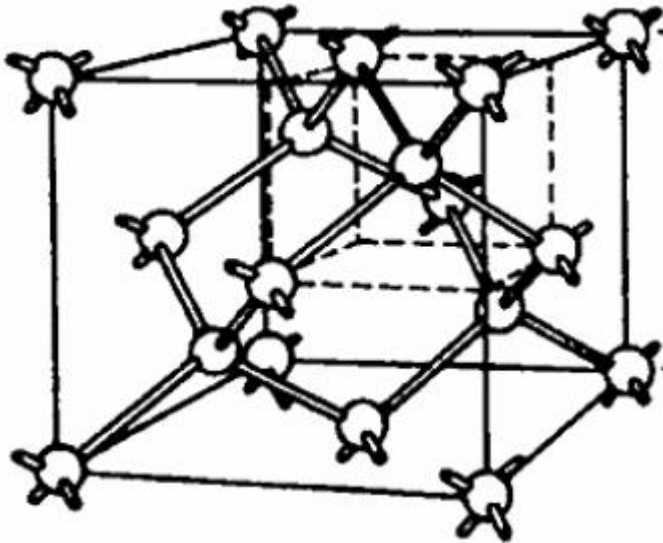
Compressed by 10%



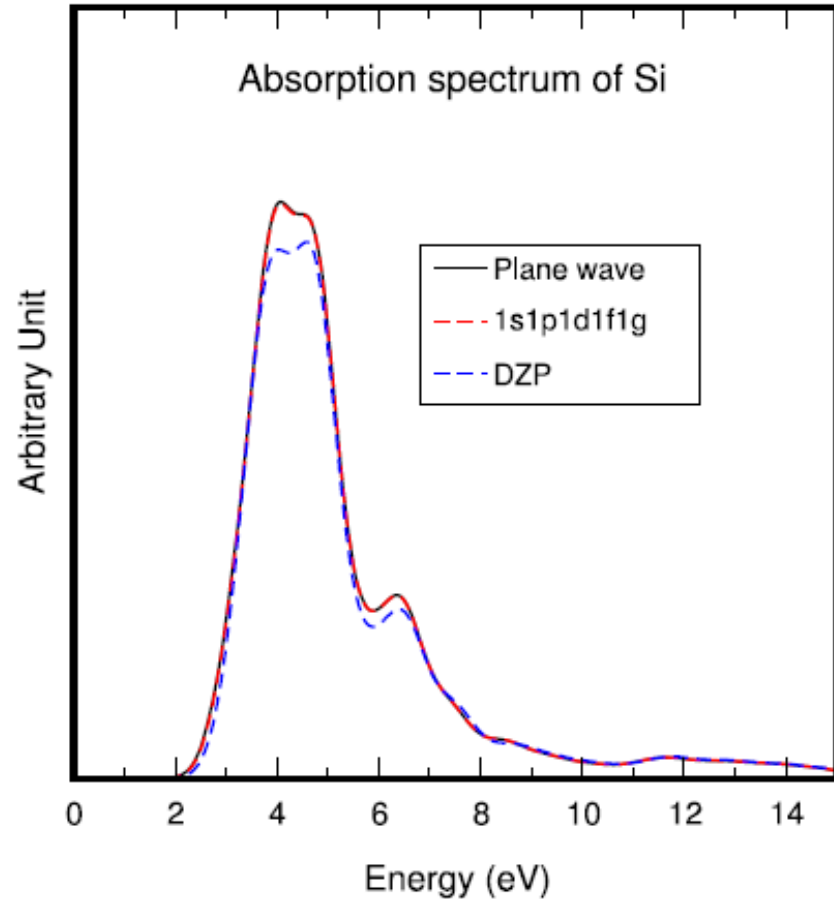
Good transferability!

Stretched by 10%

Absorption spectrum of Si



xc: LDA
K points: 13*13*13

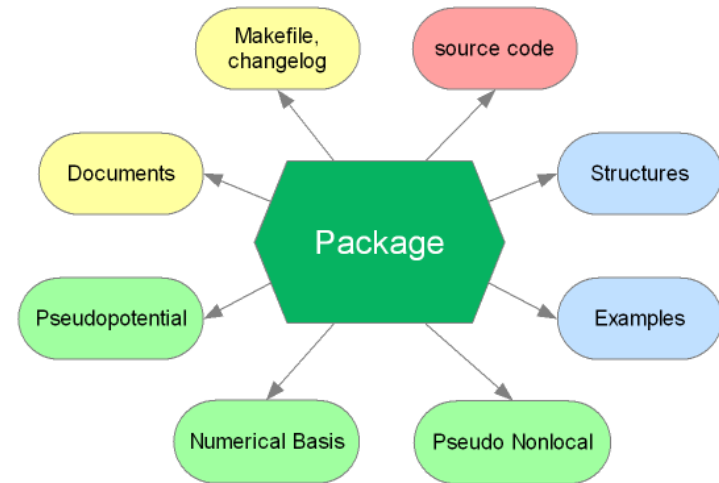


Dielectric function:

$$\epsilon_2(\omega) = \frac{2e^2\pi}{\Omega\epsilon_0} \sum_{\mathbf{k},c,v} |\langle \Psi_{\mathbf{k},c} | \hat{\mathbf{e}} \cdot \mathbf{r} | \Psi_{\mathbf{k},v} \rangle|^2 \delta[E_c(\mathbf{k}) - E_v(\mathbf{k}) - \hbar\omega]$$

ABACUS project

- ◆ Atomic-orbital Based Ab-initio Computation at UStc
- ◆ Written in C++
- ◆ Start at 2007
- ◆ first version v1.0.0 released at 2015.10
- ◆ v1.0.1 released at 2016.12
- ◆ v1.1 will be released in 2018



<http://abacus.ustc.edu.cn/>

http://abacus.ustc.edu.cn

ABACUS

ATOMIC-ORBITAL BASED AB-INITIO COMPUTATION AT USTC

Today is: Tuesday the 10th of October, 2017

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16.12.19

ABACUS V1.0.1

ABACUS version 1.0.1 is updated for download!

ABACUS (Atomic-orbital Based Ab-initio Computation at USTC) is an open-source computer code package aiming for large-scale electronic-structure simulations from first principles, developed at the Key Laboratory of Quantum Information, University of Science and Technology of China (USTC) - Computer Network and Information Center, Chinese of Academy (CNIC of CAS).

ABACUS currently provides the following features and functionalities:

1. Ground-state total energy calculations using Kohn-Sham (KS) density functional theory (DFT) with local-density or generalized gradient approximations (LDA/GGAs).
2. Brillouin zone sampling using the Monkhorst-Pack special k-points.
3. Geometry relaxation with both Conjugated Gradient (CG) and BFGS methods.
4. Semi-empirical van der Waals energy correction using the Grimme DFT-D2 scheme.
5. NVT molecular dynamics simulation using the Nose-Hoover thermostat.

[More about >>](#)



user's guide



pseudopotential
library



download



已
完
成

- ◆ Plane waves/Atomic orbital dual bases sets
- ◆ Normal conserving pseudopotential (UPF format)
- ◆ Functionals: LDA, GGA, spins, van der Waals (DFT-D2)
- ◆ Electronic structures, structural relaxation
- ◆ Molecular dynamics (NVE, NVT)

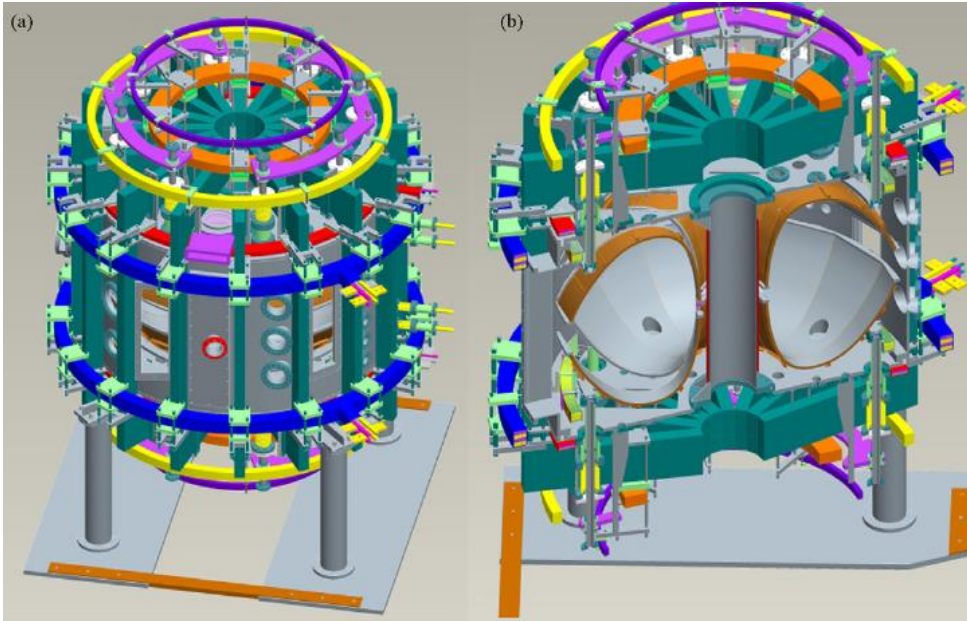
测
试
中

- ◆ stress and cell optimization
 - ◆ Real-time TDDFT (cooperate with 孟胜)
 - ◆ Berry phase
 - ◆ Interface to Wannier90
 - ◆ spin-orbital coupling
 - ◆ Hybrid functional (HSE and PBE0)
- } Will be included in the next release

进
行
中

- ◆ RPA, GW, optical properties
- ◆ Transports (cooperate with 夏柯, 袁喆)

Example: LiSn metal liquid



Lithium Tokamak Experiment (LTX),
Princeton Plasma Physics Laboratory (PPPL)

Liquid metals: promising materials are lithium, tin, gallium, etc.

Advantages:

- no irreversible erosion
- no neutron damage
- no heat overload

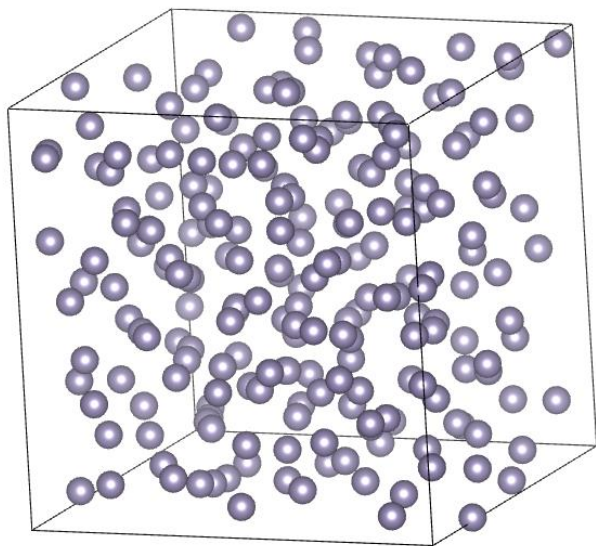
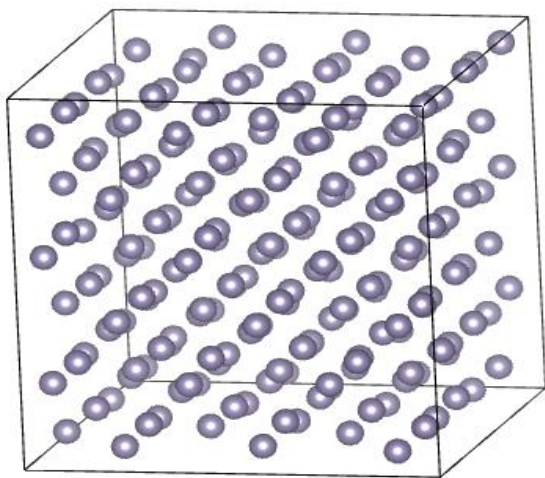
Nuclear fusion reactor wall protection materials

MD simulation system size: 200 – 500 atoms



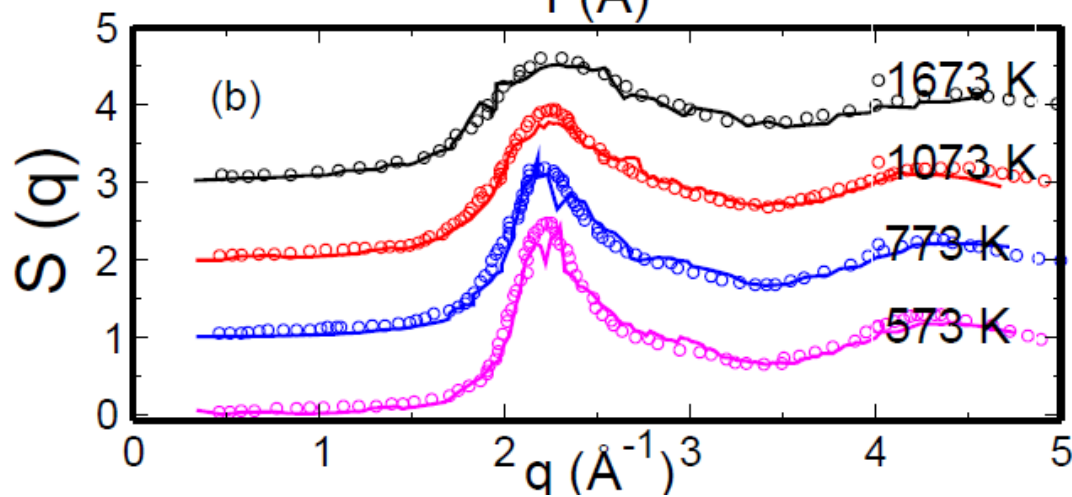
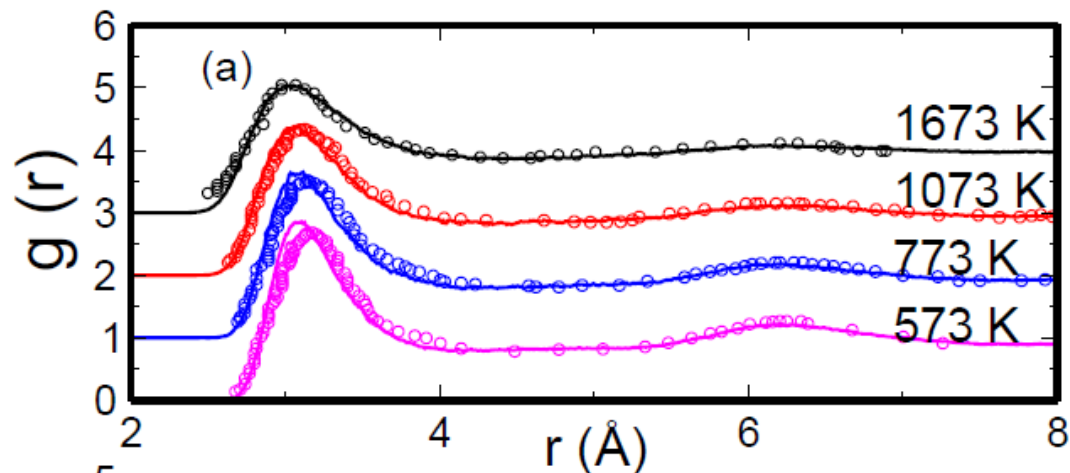
Structure properties of solid tin

	a_0	c_0/a_0	V_0	ΔE	B_0	Method
β -tin	5.780	0.537	26.15	0.000	58	FP (PW)
	5.786	0.538	26.23	0.000	57	FP (LCAO)
	5.831	0.546	27.07	-	-	EXP ⁴⁵
	5.8119	0.543	26.65	-	-	EXP ⁴⁶
	-	-	-	-	57.037	EXP ⁴⁷
	-	-	-	-	57.9	EXP ⁴⁸
α -tin	6.442	-	33.41	-0.019	43	FP (PW)
	6.445	-	33.47	-0.055	43	FP (LCAO)
	6.483	-	34.05	-	-	EXP ⁴⁵
	-	-	-	-	42.617	EXP ⁴⁷
	-	-	-	-	54	EXP ⁴⁹
bct	3.933	0.846	25.73	0.045	53	FP (PW)
	3.920	0.844	25.42	0.047	54	FP (LCAO)
bcc	3.664		24.60	0.095	70	FP (PW)
	3.658		24.47	0.130	71	FP (LCAO)



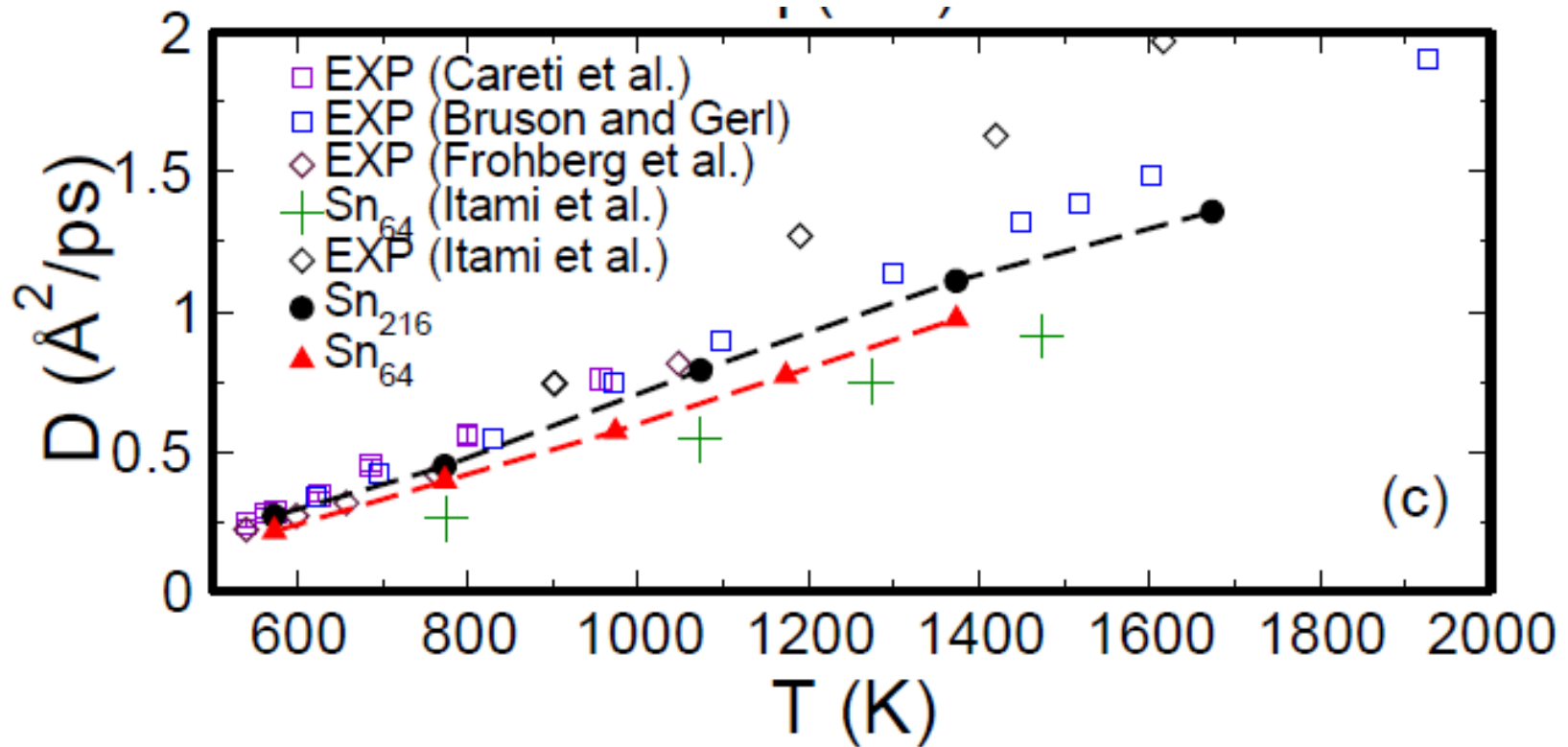
216 atoms, 32 cores, ~30 s/MD step

$$g(r) = \frac{1}{\rho N} \left\langle \sum_{i=1}^N \sum_{j=1, j \neq i}^N \delta(\mathbf{r} - \mathbf{R}_i + \mathbf{R}_j) \right\rangle$$



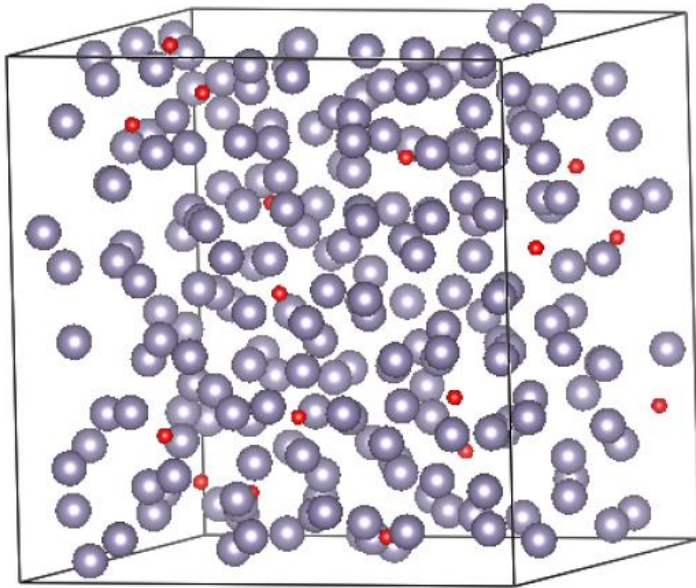
$$S(q) = \frac{1}{N} \left\langle \sum_{i=1}^N \sum_{j=1}^N e^{i\mathbf{q} \cdot (\mathbf{R}_i - \mathbf{R}_j)} \right\rangle$$

Diffusion coefficient of pure Sn metal liquid

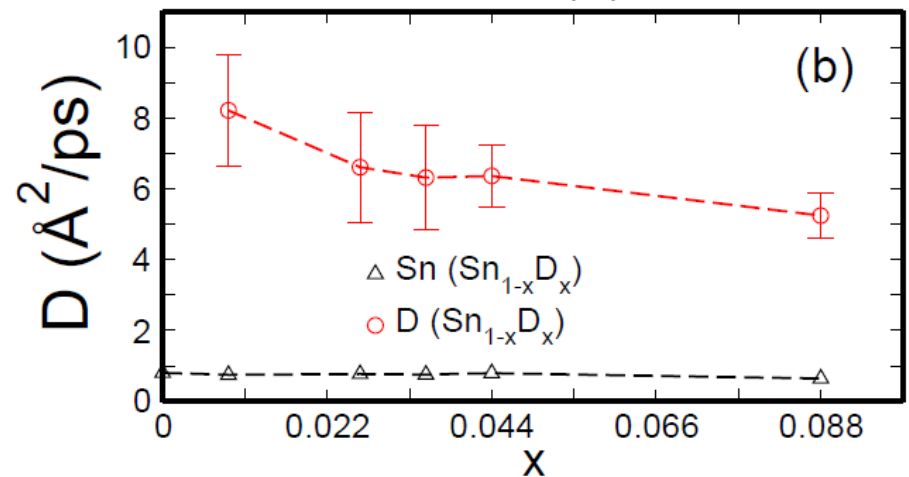
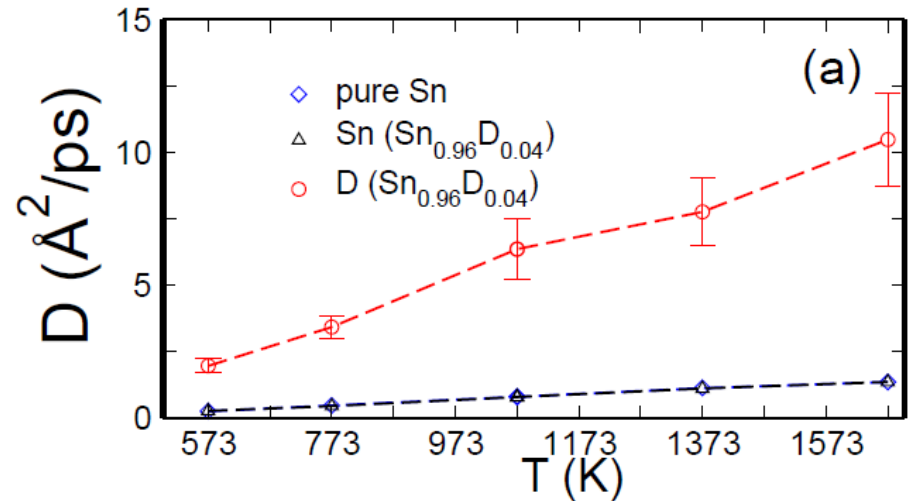


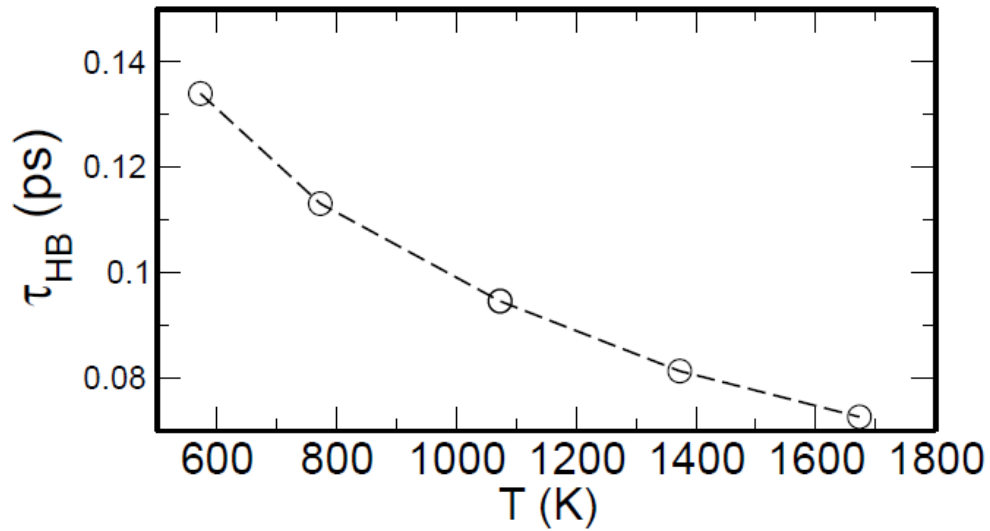
$$D = \frac{1}{6} \frac{d}{dt} \langle \Delta r(t)^2 \rangle$$

Deuterium diffusion in Sn metal liquid

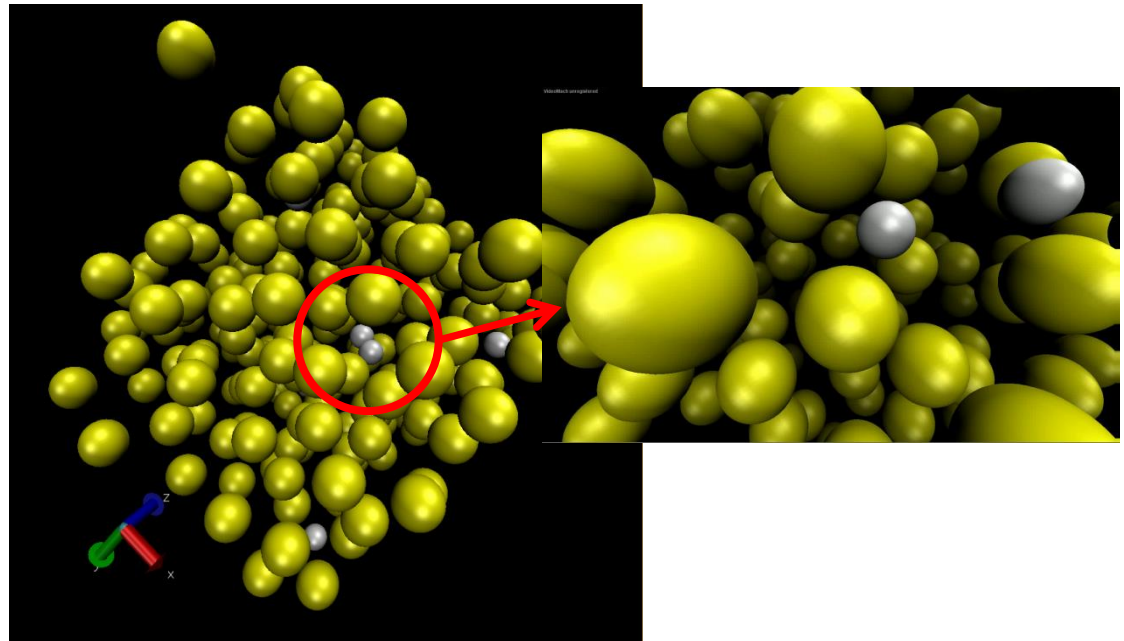


~2ps, 10^5 MD steps





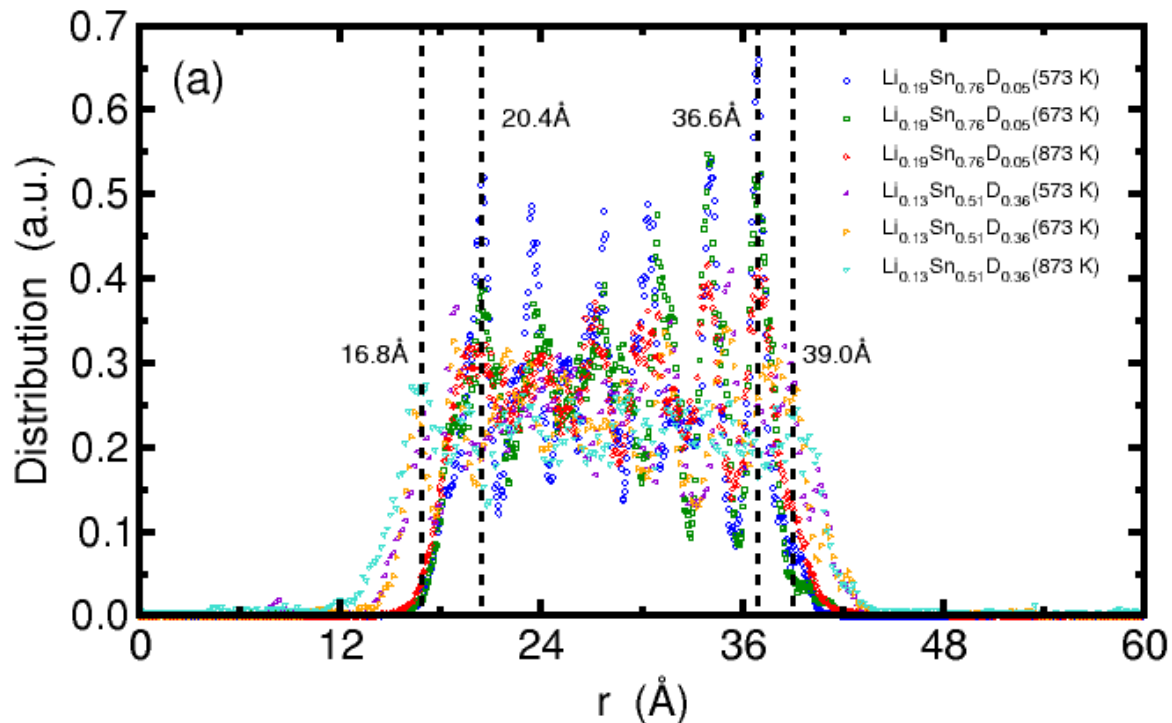
1. D does not form bonds with Sn;
2. D does not form bonds with other D;
3. Good for D distraction.





LiSnD liquid slab

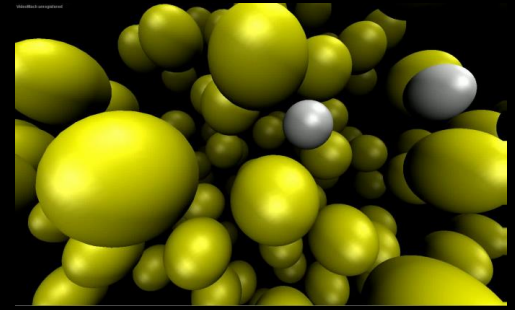
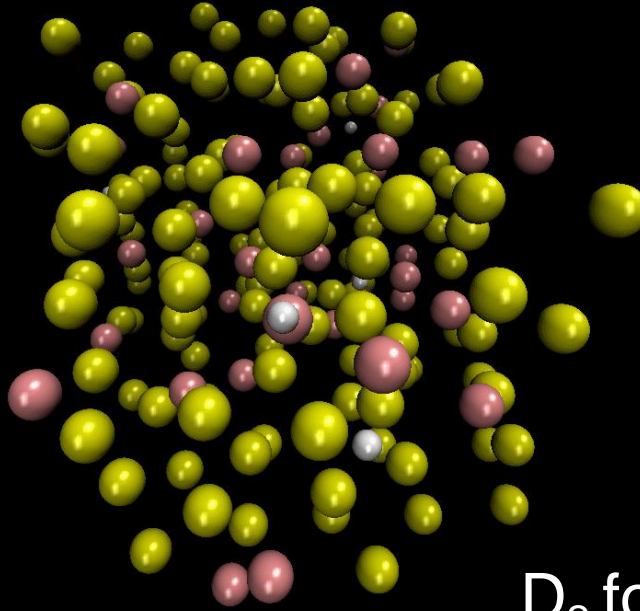
- 1、 ABACUS: LCAO-DZP
- 2、 BOMD-NVT;
- 4、 Time step: 0.2 fs, total time: 20 ps (100,000 steps);
- 5、 Li : Sn = 1:4, T=573 K,673 K and 873 K ,
deuterium for 10 and 100.





VideoMach unregistered

LiSn liquid



Pure Sn liquid

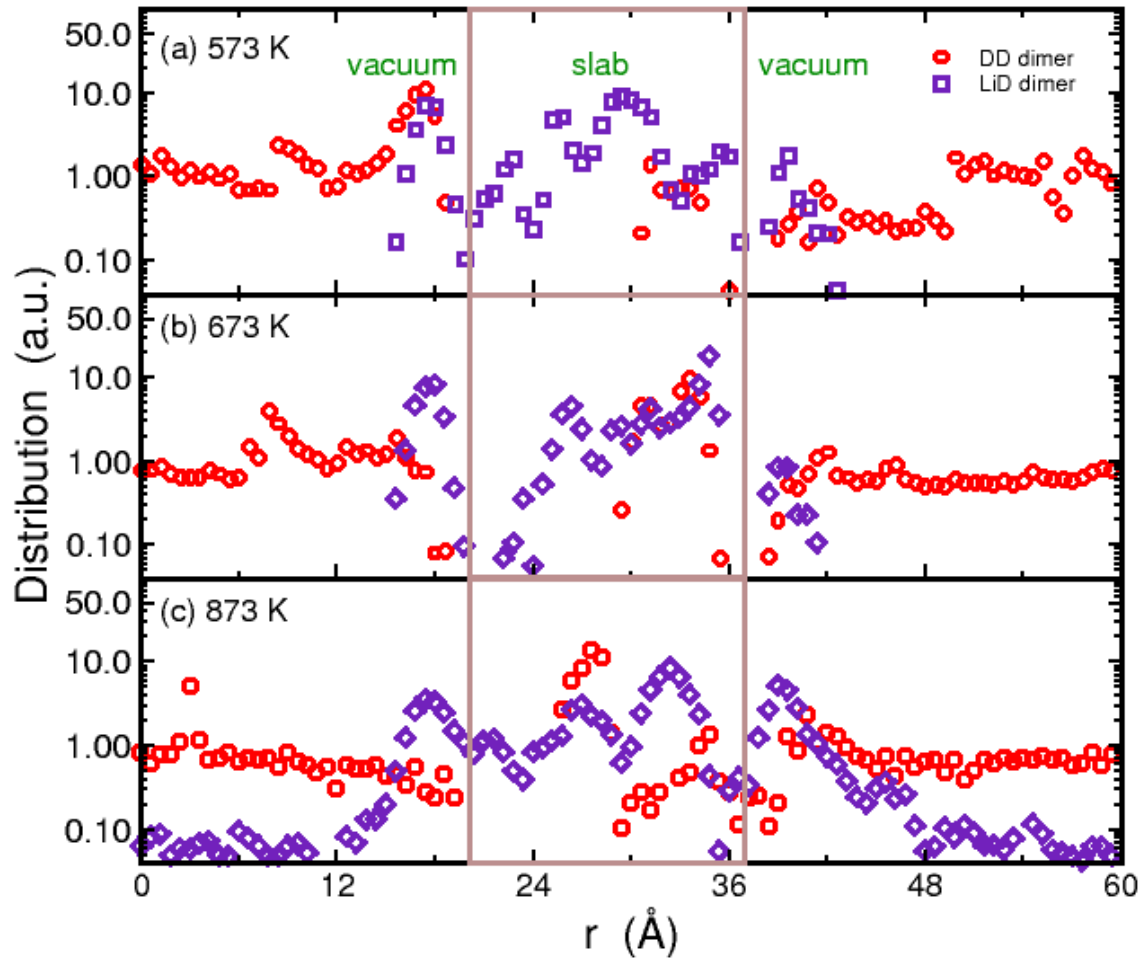
D_2 formation



(a) At 573 K and 673 K,
 D_2 dimer may escape
to the vacuum.

(a) At 873 K, small amount
of LiD dimer may also
escape to vacuum.

In preparation





Collaborators

1. Prof. Ren, Xinguo (USTC)
 2. Dr. Chen, Mohan (Temple U)
 3. Dr. Shen, Yu (USTC)
 4. Dr. Zhang, Wenshuai (USTC)
 5. Dr. Liu, Xiaohui (USTC)
 6. Mr. Li, Pengfei (USTC)
 7. Mr. Zheng, Daye (USTC)
 8. Mr. Lin, Peizhe (USTC)
 9. Mr. He, Fuxiang (USTC)
 10. Mr. Jin, Gan (USTC)
1. Prof. Meng, Sheng (IOP)
 2. Prof. Xia, Ke (BNU)
 3. Prof. Yuan, Zhe (BNU)



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建校三十周年誌慶

嚴濟慈
一九八八年五月

題

