## Chapter 8 -Virtual Work

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### 8.1 Introduction

### 8.1.1 General

Virtual Work is a fundamental theory in the mechanics of bodies. So fundamental in fact, that Newton's 3 equations of equilibrium can be derived from it. Virtual Work provides a basis upon which vectorial mechanics (i.e. Newton's laws) can be linked to the energy methods (i.e. Lagrangian methods) which are the basis for finite element analysis and advanced mechanics of materials.

Virtual Work allows us to solve determinate and indeterminate structures and to calculate their deflections. That is, it can achieve everything that all the other methods together can achieve.

Before starting into Virtual Work there are some background concepts and theories that need to be covered.

### 8.1.2 Background

## Strain Energy and Work Done

Strain energy is the amount of energy stored in a structural member due to deformation caused by an external load. For example, consider this simple spring:


We can see that as it is loaded by a gradually increasing force, $F$, it elongates. We can graph this as:


The line OA does not have to be straight, that is, the constitutive law of the spring's material does not have to be linear.

An increase in the force of a small amount, $\delta F$ results in a small increase in deflection, $\delta y$. The work done during this movement (force $\times$ displacement) is the average force during the course of the movement, times the displacement undergone. This is the same as the hatched trapezoidal area above. Thus, the increase in work associated with this movement is:

$$
\begin{align*}
\delta U & =\frac{F+(F+\delta F)}{2} \cdot \delta y \\
& =F \cdot \delta y+\frac{\delta F \cdot \delta y}{2}  \tag{1.1}\\
& \approx F \cdot \delta y
\end{align*}
$$

where we can neglect second-order quantities. As $\delta y \rightarrow 0$, we get:

$$
d U=F \cdot d y
$$

The total work done when a load is gradually applied from 0 up to a force of $F$ is the summation of all such small increases in work, i.e.:

$$
\begin{equation*}
U=\int_{0}^{y} F d y \tag{1.2}
\end{equation*}
$$

This is the dotted area underneath the load-deflection curve of earlier and represents the work done during the elongation of the spring. This work (or energy, as they are the same thing) is stored in the spring and is called strain energy and denoted $U$.

If the load-displacement curve is that of a linearly-elastic material then $F=k y$ where $k$ is the constant of proportionality (or the spring stiffness). In this case, the dotted area under the load-deflection curve is a triangle.


As we know that the work done is the area under this curve, then the work done by the load in moving through the displacement - the External Work Done, $W_{e}$ - is given by:

$$
\begin{equation*}
W_{e}=\frac{1}{2} F y \tag{1.3}
\end{equation*}
$$

We can also calculate the strain energy, or Internal Work Done, $W_{I}$, by:

$$
\begin{aligned}
U & =\int_{0}^{y} F d y \\
& =\int_{0}^{y} k y d y \\
& =\frac{1}{2} k y^{2}
\end{aligned}
$$

Also, since $F=k y$, we then have:

$$
W_{I}=U=\frac{1}{2}(k y) y=\frac{1}{2} F y
$$

But this is the external work done, $W_{e}$. Hence we have:

$$
\begin{equation*}
W_{e}=W_{I} \tag{1.4}
\end{equation*}
$$

Which we may have expected from the Law of Conservation of Energy. Thus:

The external work done by external forces moving through external displacements is equal to the strain energy stored in the material.

## Law of Conservation of Energy

For structural analysis this can be stated as:

Consider a structural system that is isolated such it neither gives nor receives energy; the total energy of this system remains constant.

The isolation of the structure is key: we can consider a structure isolated once we have identified and accounted for all sources of restraint and loading. For example, to neglect the self-weight of a beam is problematic as the beam receives gravitational energy not accounted for, possibly leading to collapse.

In the spring and force example, we have accounted for all restraints and loading (for example we have ignored gravity by having no mass). Thus the total potential energy of the system, $\Pi$, is constant both before and after the deformation.

In structural analysis the relevant forms of energy are the potential energy of the load and the strain energy of the material. We usually ignore heat and other energies.

## Potential Energy of the Load

Since after the deformation the spring has gained strain energy, the load must have lost potential energy, $V$. Hence, after deformation we have for the total potential energy:

$$
\begin{align*}
\Pi & =U+V \\
& =\frac{1}{2} k y^{2}-F y \tag{1.5}
\end{align*}
$$

In which the negative sign indicates a loss of potential energy for the load.

## Principle of Minimum Total Potential Energy

If we plot the total potential energy against $y$, equation (1.5), we get a quadratic curve similar to:


Consider a numerical example, with the following parameters, $k=10 \mathrm{kN} / \mathrm{m}$ and $F=10 \mathrm{kN}$ giving the equilibrium deflection as $y=F / k=1 \mathrm{~m}$. We can plot the following quantities:

- Internal Strain Energy, or Internal Work: $\quad U=W_{I}=\frac{1}{2} k y^{2}=\frac{1}{2} 10 y^{2}=5 y^{2}$
- Potential Energy:
$V=-F y=-10 y$
- Total Potential Energy:
$\Pi=U+V=5 y^{2}-10 y$
- External Work:

$$
W_{e}=\frac{1}{2} P y=\frac{1}{2} 10 y=5 y
$$

and we get the following plots (split into two for clarity):


From these graphs we can see that because $W_{I}$ increases quadratically with $y$, while the $W_{e}$ increases only linearly, $W_{I}$ always catches up with $W_{e}$, and there will always be a non-zero equilibrium point where $W_{e}=W_{I}$.

Admittedly, these plots are mathematical: the deflection of the spring will not take up any value; it takes that value which achieves equilibrium. At this point we consider a small variation in the total potential energy of the system. Considering $F$ and $k$ to be constant, we can only alter $y$. The effect of this small variation in $y$ is:

$$
\begin{align*}
\Pi(y+\delta y)-\Pi(y) & =\frac{1}{2} k(y+\delta y)^{2}-F(y+\delta y)-\frac{1}{2} k y^{2}+F y \\
& =\frac{1}{2} k(2 y \cdot \delta y)-F \cdot \delta y+\frac{1}{2} k(\delta y)^{2}  \tag{1.6}\\
& =(k y-F) \delta y+\frac{1}{2} k(\delta y)^{2}
\end{align*}
$$

Similarly to a first derivate, for $\Pi$ to be an extreme (either maximum or minimum), the first variation must vanish:

$$
\begin{equation*}
\delta^{(1)} \Pi=(k y-F) \delta y=0 \tag{1.7}
\end{equation*}
$$

Therefore:

$$
\begin{equation*}
k y-F=0 \tag{1.8}
\end{equation*}
$$

Which we recognize to be the $\sum F_{x}=0$. Thus equilibrium occurs when $\Pi$ is an extreme.

Before introducing more complicating maths, an example of the above variation in equilibrium position is the following. Think of a shopkeeper testing an old type of scales for balance - she slightly lifts one side, and if it returns to position, and no large rotations occur, she concludes the scales is in balance. She has imposed a
variation in displacement, and finds that since no further displacement occurs, the 'structure' was originally in equilibrium.

Examining the second variation (similar to a second derivate):

$$
\begin{equation*}
\delta^{(2)} \Pi=\frac{1}{2} k(\delta y)^{2} \geq 0 \tag{1.9}
\end{equation*}
$$

We can see it is always positive and hence the extreme found was a minimum. This is a particular proof of the general principle that structures take up deformations that minimize the total potential energy to achieve equilibrium. In short, nature is lazy!

To summarize our findings:

- Every isolated structure has a total potential energy;
- Equilibrium occurs when structures can minimise this energy;
- A small variation of the total potential energy vanishes when the structure is in equilibrium.

These concepts are brought together in the Principle of Virtual Work.

### 8.2 The Principle of Virtual Work

### 8.2.1 Definition

Based upon the Principle of Minimum Total Potential Energy, we can see that any small variation about equilibrium must do no work. Thus, the Principle of Virtual Work states that:

A body is in equilibrium if, and only if, the virtual work of all forces acting on the body is zero.

In this context, the word 'virtual' means 'having the effect of, but not the actual form of, what is specified'. Thus we can imagine ways in which to impose virtual work, without worrying about how it might be achieved in the physical world.

We need to remind ourselves of equilibrium between internal and external forces, as well as compatibility of displacement, between internal and external displacements for a very general structure:


Equilibrium


Compatibility

## The Two Principles of Virtual Work

There are two principles that arise from consideration of virtual work, and we use either as suited to the unknown quantity (force or displacement) of the problem under analysis.

## Principle of Virtual Displacements:

Virtual work is the work done by the actual forces acting on the body moving through a virtual displacement.

This means we solve an equilibrium problem through geometry, as shown:

## Principle of Virtual Displacements



## Virtual Displacements:

External compatible

displacements $\rightarrow$ Geometry $\longrightarrow$| Internal compatible |
| :---: |
| displacements |

## Principle of Virtual Forces:

Virtual work is the work done by a virtual force acting on the body moving through the actual displacements.

This means we solve a geometry problem through equilibrium as shown below:

## Principle of Virtual Forces



Internal Real 'Displacements':


### 8.2.2 Virtual Displacements

A virtual displacement is a displacement that is only imagined to occur. This is exactly what we did when we considered the vanishing of the first variation of $\Pi$; we found equilibrium. Thus the application of a virtual displacement is a means to find this first variation of $\Pi$.

So given any real force, $F$, acting on a body to which we apply a virtual displacement. If the virtual displacement at the location of and in the direction of $F$ is $\delta y$, then the force $F$ does virtual work $\delta W=F \cdot \delta y$.

There are requirements on what is permissible as a virtual displacement. For example, in the simple proof of the Principle of Virtual Work (to follow) it can be seen that it is assumed that the directions of the forces applied to $P$ remain unchanged. Thus:

- virtual displacements must be small enough such that the force directions are maintained.

The other very important requirement is that of compatibility:

- virtual displacements within a body must be geometrically compatible with the original structure. That is, geometrical constraints (i.e. supports) and member continuity must be maintained.

In summary, virtual displacements are not real, they can be physically impossible but they must be compatible with the geometry of the original structure and they must be small enough so that the original geometry is not significantly altered.

As the deflections usually encountered in structures do not change the overall geometry of the structure, this requirement does not cause problems.

### 8.2.3 Virtual Forces

So far we have only considered small virtual displacements and real forces. The virtual displacements are arbitrary: they have no relation to the forces in the system, or its actual deformations. Therefore virtual work applies to any set of forces in equilibrium and to any set of compatible displacements and we are not restricted to considering only real force systems and virtual displacements. Hence, we can use a virtual force system and real displacements. Obviously, in structural design it is these real displacements that are of interest and so virtual forces are used often.

A virtual force is a force imagined to be applied and is then moved through the actual deformations of the body, thus causing virtual work.

So if at a particular location of a structure, we have a deflection, $y$, and impose a virtual force at the same location and in the same direction of $\delta F$ we then have the virtual work $\delta W=y \cdot \delta F$.

Virtual forces must form an equilibrium set of their own. For example, if a virtual force is applied to the end of a spring there will be virtual stresses in the spring as well as a virtual reaction.

### 8.2.4 Simple Proof using Virtual Displacements

We can prove the Principle of Virtual Work quite simply, as follows. Consider a particle $P$ under the influence of a number of forces $F_{1}, \ldots, F_{n}$ which have a resultant force, $F_{R}$. Apply a virtual displacement of $\delta y$ to $P$, moving it to $P^{\prime}$, as shown:


The virtual work done by each of the forces is:

$$
\delta W=F_{1} \cdot \delta y_{1}+\ldots+F_{n} \cdot \delta y_{n}=F_{R} \cdot \delta y_{R}
$$

Where $\delta y_{1}$ is the virtual displacement along the line of action of $F_{1}$ and so on. Now if the particle $P$ is in equilibrium, then the forces $F_{1}, \ldots, F_{n}$ have no resultant. That is, there is no net force. Hence we have:

$$
\begin{equation*}
\delta W=0 \cdot \delta y_{R}=F_{1} \cdot \delta y_{1}+\ldots+F_{n} \cdot \delta y_{n}=0 \tag{2.1}
\end{equation*}
$$

Proving that when a particle is in equilibrium the virtual work of all the forces acting on it sum to zero. Conversely, a particle is only in equilibrium if the virtual work done during a virtual displacement is zero.

### 8.2.5 Internal and External Virtual Work

Consider the spring we started with, as shown again below. Firstly it is unloaded. Then a load, $F$, is applied, causing a deflection $y$. It is now in its equilibrium position and we apply a virtual displacement, $\delta y$ to the end of the spring, as shown:


A free body diagram of the end of the spring is:


Thus the virtual work done by the two forces acting on the end of the spring is:

$$
\delta W=F \cdot \delta y-k y \cdot \delta y
$$

If the spring is to be in equilibrium we must then have:

$$
\begin{aligned}
\delta W & =0 \\
F \cdot \delta y-k y \cdot \delta y & =0 \\
F \cdot \delta y & =k y \cdot \delta y \\
F & =k y
\end{aligned}
$$

That is, the force in the spring must equal the applied force, as we already know. However, if we define the following:

- External virtual work, $\delta W_{E}=F \cdot \delta y$;
- Internal virtual work, $\delta W_{I}=k y \cdot \delta y$;

We then have:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{I}-\delta W_{E} & =0
\end{aligned}
$$

Thus:

$$
\begin{equation*}
\delta W_{E}=\delta W_{I} \tag{2.2}
\end{equation*}
$$

which states that the external virtual work must equal the internal virtual work for a structure to be in equilibrium.

It is in this form that the Principle of Virtual Work finds most use.

Of significance also in the equating of internal and external virtual work, is that there are no requirements for the material to have any particular behaviour. That is, virtual work applies to all bodies, whether linearly-elastic, elastic, elasto-plastic, plastic etc. Thus the principle has more general application than most other methods of analysis.

Internal and external virtual work can arise from either virtual displacements or virtual forces.

### 8.3 Application of Virtual Displacements

### 8.3.1 Rigid Bodies

## Basis

Rigid bodies do not deform and so there is no internal virtual work done. Thus:

$$
\begin{align*}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I}  \tag{3.1}\\
\sum F_{i} \cdot \delta y_{i} & =0
\end{align*}
$$

A simple application is to find the reactions of statically determinate structures. However, to do so, we need to make use of the following principle:

## Principle of Substitution of Constraints

Having to keep the constraints in place is a limitation of virtual work. However, we can substitute the restraining force (i.e. the reaction) in place of the restraint itself. That is, we are turning a geometric constraint into a force constraint. This is the Principle of Substitution of Constraints. We can use this principle to calculate unknown reactions:

1. Replace the reaction with its appropriate force in the same direction (or sense);
2. Impose a virtual displacement on the structure;
3. Calculate the reaction, knowing $\delta W=0$.

## Reactions of Determinate and Indeterminate Structures

For statically determinate structures, removing a restraint will always render a mechanism (or rigid body) and so the reactions of statically determinate structures are easily obtained using virtual work. For indeterminate structures, removing a restraint does not leave a mechanism and hence the virtual displacements are harder to establish since the body is not rigid.

## Example 1

## Problem

Determine the reactions for the following beam:


## Solution

Following the Principle of Substitution of Constraints, we replace the geometric constraints (i.e. supports), by their force counterparts (i.e. reactions) to get the free-body-diagram of the whole beam:


Next, we impose a virtual displacement on the beam. Note that the displacement is completely arbitrary, and the beam remains rigid throughout:


In the above figure, we have imposed a virtual displacement of $\delta y_{A}$ at $A$ and then imposed a virtual rotation of $\delta \theta_{A}$ about $A$. The equation of virtual work is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum F_{i} \cdot \delta y_{i} & =0
\end{aligned}
$$

Hence:

$$
V_{A} \cdot \delta y_{A}-P \cdot \delta y_{C}+V_{B} \cdot \delta y_{B}=0
$$

Relating the virtual movements of points $B$ and $C$ to the virtual displacements gives:

$$
V_{A} \cdot \delta y_{A}-P\left(\delta y_{A}+a \cdot \delta \theta_{A}\right)+V_{B}\left(\delta y_{A}+L \cdot \delta \theta_{A}\right)=0
$$

And rearranging gives:

$$
\left(V_{A}+V_{B}-P\right) \delta y_{A}+\left(V_{B} L-P a\right) \delta \theta_{A}=0
$$

And here is the power of virtual work: since we are free to choose any value we want for the virtual displacements (ie. they are completely arbitrary), we can choose $\delta \theta_{A}=0$ and $\delta y_{A}=0$, which gives the following two equations:

$$
\begin{array}{rrr}
\left(V_{A}+V_{B}-P\right) \delta y_{A}=0 & \left(V_{B} L-P a\right) \delta \theta_{A}=0 \\
V_{A}+V_{B}-P=0 & V_{B} L-P a=0
\end{array}
$$

But the first equation is just $\sum F_{y}=0$ whilst the second is the same as $\sum M$ about $A=0$. Thus equilibrium equations occur naturally within the virtual work framework. Note also that the two choices made for the virtual displacements correspond to the following virtual displaced configurations:



Example 2
Problem
For the following truss, calculate the reaction $V_{C}$ :


Solution
Firstly, set up the free-body-diagram of the whole truss:


Next we release the constraint corresponding to reaction $V_{C}$ and replace it by the unknown force $V_{C}$ and we apply a virtual displacement to the truss to get:


Hence the virtual work done is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
-10 \cdot \frac{\delta y}{2}+V_{C} \cdot \delta y & =0 \\
V_{C} & =5 \mathrm{kN}
\end{aligned}
$$

Note that the reaction is an external force to the structure, and that no internal virtual work is done since the members do not undergo virtual deformation. The truss rotates as a rigid body about the support $A$.

Excercise: Find the horizontal reactions of the truss.

### 8.3.2 Deformable Bodies

## Basis

For a virtual displacement we have:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum F_{i} \cdot \delta y_{i} & =\int P_{i} \cdot \delta e_{i}
\end{aligned}
$$

For the external virtual work, $F_{i}$ represents an externally applied force (or moment) and $\delta y_{i}$ is its corresponding virtual displacement. For the internal virtual work, $P_{i}$ is the internal force in member $i$ and $\delta e_{i}$ is its virtual deformation. Different stress resultants have different forms of internal work, and we will examine these. The summations reflect the fact that all work done must be accounted for. Remember in the above, each the displacements must be compatible and the forces must be in equilibrium, summarized as:


These displacements are completely arbitrary (i.e. we can choose them to suit our purpose) and bear no relation to the forces above.

## Example 3

## Problem

For the simple truss, show the forces are as shown using the principle of virtual displacements.


Loading and dimensions


Equilibrium forces and actual deformation

## Solution

An arbitrary set of permissible compatible displacements are shown:


Compatible set of displacements

To solve the problem though, we need to be thoughtful about the choice of arbitrary displacement: we only want the unknown we are currently calculating to do work, otherwise there will be multiple unknowns in the equation of virtual work. Consider this possible compatible set of displacements:


Assuming both forces are tensile, and noting that the load does no work since the virtual external displacement is not along its line of action, the virtual work done is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
40 \times 0 & =\delta e_{1} \cdot F_{1}+\delta e_{2} \cdot F_{2}
\end{aligned}
$$

Since the displacements are small and compatible, we know by geometry that:

$$
\delta y=\delta e_{1}=\frac{3}{5} \delta e_{2}
$$

Giving:

$$
\begin{aligned}
0 & =\delta y \cdot F_{1}+\frac{3}{5} \delta y \cdot F_{2} \\
F_{1} & =-\frac{3}{5} F_{2}
\end{aligned}
$$

which tells us that one force is in tension and the other in compression. Unfortunately, this displaced configuration was not sufficient as we were left with two unknowns. Consider the following one though:


$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
-40 \cdot \delta y & =-\delta e_{2} \cdot F_{2}
\end{aligned}
$$

The negative signs appear because the 40 kN load moves against its direction, and the (assumed) tension member $A B$ is made shorter which is against its 'tendency' to elongate. Also, by geometry:

$$
\delta y=\frac{4}{5} \delta e_{2}
$$

And so:

$$
\begin{aligned}
40 \cdot \delta y & =\frac{4}{5} \delta y \cdot F_{2} \\
F_{2} & =50 \mathrm{kN}
\end{aligned}
$$

The positive result indicates that the member is in tension as was assumed. Now we can return to the previous equation to get:

$$
F_{1}=-\frac{3}{5}(+50)=-30 \mathrm{kN}
$$

And so the member is in compression, since the negative tells us that it is in the opposite direction to that assumed, which was tension.

## Internal Virtual Work by Axial Force

Members subject to axial force may have the following:

- real force by a virtual displacement:

$$
\delta W_{I}=P \cdot \delta e
$$

- virtual force by a real displacement:

$$
\delta W_{I}=e \cdot \delta P
$$

We have avoided integrals over the length of the member since we will only consider prismatic members.

## Internal Virtual Work in Bending

The internal virtual work done in bending is one of:

- real moment by a virtual curvature:

$$
\delta W_{I}=M \cdot \delta \kappa
$$

- virtual moment by a real curvature:

$$
\delta W_{I}=\kappa \cdot \delta M
$$

The above expressions are valid at a single position in a beam.

When virtual rotations are required along the length of the beam, the easiest way to do this is by applying virtual loads that in turn cause virtual moments and hence virtual curvatures. We must sum all of the real moments by virtual curvatures along the length of the beam, hence we have:

$$
\begin{aligned}
\delta W_{I} & =\int_{0}^{L} M_{x} \cdot \delta \kappa_{x} d x=\int_{0}^{L} \kappa_{x} \cdot \delta M_{x} d x \\
& =\int_{0}^{L} M_{x} \cdot \frac{\delta M_{x}}{E I_{x}} d x
\end{aligned}
$$

## Internal Virtual Work in Shear

At a single point in a beam, the shear strain, $\gamma$, is given by:

$$
\gamma=\frac{V}{G A_{v}}
$$

where $V$ is the applied shear force, $G$ is the shear modulus and $A_{v}$ is the cross-section area effective in shear which is explained below. The internal virtual work done in shear is thus:

- real shear force by a virtual shear strain:

$$
\delta W_{I}=V \cdot \delta \gamma=V \cdot \frac{\delta V}{G A_{v}}
$$

- virtual shear force by a real shear strain:

$$
\delta W_{I}=\gamma \cdot \delta V=\frac{V}{G A_{v}} \cdot \delta V
$$

These expressions are valid at a single position in a beam and must be integrated along the length of the member as was done for moments and rotations.

The area of cross section effective in shear arises because the shear stress (and hence strain) is not constant over a cross section. The stress $V / A_{v}$ is an average stress, equivalent in work terms to the actual uneven stress over the full cross section, $A$. We say $A_{v}=A / k$ where $k$ is the shear factor. Some values of $k$ are: 1.2 for rectangular sections; 1.1 for circular sections; and 2.0 for thin-walled circular sections.

## Internal Virtual Work in Torsion

At a single point in a member, the twist, $\phi$, is given by:

$$
\phi=\frac{T}{G J}
$$

where $T$ is the applied torque, $G$ is the shear modulus, $J$ is the polar moment of inertia. The internal virtual work done in torsion is thus:

- real torque by a virtual twist:

$$
\delta W_{I}=T \cdot \delta \phi=T \cdot \frac{\delta T}{G J}
$$

- virtual torque by a real twist:

$$
\delta W_{I}=\delta T \cdot \phi=\delta T \cdot \frac{T}{G J}
$$

Once again, the above expressions are valid at a single position in a beam and must be integrated along the length of the member as was done for moments and rotations.

Note the similarity between the expressions for the four internal virtual works.

## Example 4

## Problem

For the beam of Example 1 (shown again), find the bending moment at $C$.


## Solution

To solve this, we want to impose a virtual displacement configuration that only allows the unknown of interest, ie. $M_{C}$, to do any work. Thus choose the following:


Since portions $A C$ and $C B$ remain straight (or unbent) no internal virtual work is done in these sections. Thus the only internal work is done at $C$ by the beam moving through the rotation at $C$. Thus:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
P \cdot \delta y_{C} & =M_{C} \cdot \delta \theta_{C}
\end{aligned}
$$

Bu the rotation at $C$ is made up as:

$$
\begin{aligned}
\delta \theta_{C} & =\delta \theta_{C A}+\delta \theta_{C B} \\
& =\frac{\delta y_{C}}{a}+\frac{\delta y_{C}}{b} \\
& =\left(\frac{a+b}{a b}\right) \delta y_{C}
\end{aligned}
$$

But $a+b=L$, hence:

$$
\begin{aligned}
P \cdot \delta y_{C} & =M_{C} \cdot\left(\frac{L}{a b}\right) \delta y_{C} \\
M_{C} & =\frac{P a b}{L}
\end{aligned}
$$

We can verify this from the reactions found previously: $M_{C}=V_{B} \cdot b=(P a / L) b$.

## Sign Convention for Rotations

When imposing a virtual rotation, if the side that is already in tension due to the real force system elongates, then positive virtual work is done.


Tension side
gets langer

## Virtual Work done by Distributed Load

The virtual work done by an arbitrary distributed load moving through a virtual displacement field is got by summing the infinitesimal components:


Considering the special (but very common) case of a uniformly distributed load:

$$
\delta W_{e}=w \int_{x_{1}}^{x_{2}} \delta y(x) d x=w \times[\text { Area of } \delta y \text { diagram }]
$$

If we denote the length of the load as $L=x_{2}-x_{1}$ then we can say that the area of the displacement diagram is the average displacement, $\overline{\delta y}$, times the length, $L$. Thus:

$$
\delta W_{e}=w L \cdot \overline{\delta y}=[\text { Total load }] \times[\text { Average displacement }]
$$

## Example 5

## Problem

Calculate the force $F_{1}$ in the truss shown:


## Solution

To do this, we introduce a virtual displacement along the length of member 1 . We do this so that member 2 does not change length during this virtual displacement and so does no virtual work. Note also that compatibility of the truss is maintained. For example, the members still meet at the same joint, and the support conditions are met.


The virtual work done is then:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
-10 \cdot \frac{\delta y}{\sqrt{2}} & =-F_{1} \cdot \delta y \\
F_{1} & =\frac{10}{\sqrt{2}}=5 \sqrt{2} \mathrm{kN}
\end{aligned}
$$

Note some points on the signs used:

1. Negative external work is done because the 10 kN load moves upwards; i.e. the reverse direction to its action.
2. We assumed member 1 to be in compression but then applied a virtual elongation which lengthened the member thus reducing its internal virtual work. Hence negative internal work is done.
3. We initially assumed $F_{1}$ to be in compression and we obtained a positive answer confirming our assumption.

Exercise: Investigate the vertical and horizontal equilibrium of the loaded joint by considering vertical and horizontal virtual displacements separately.
8.3.3 Problems

1. For the truss shown, calculate the vertical reaction at $C$ and the forces in the members, using virtual work.

2. For the truss shown, find the forces in the members, using virtual work:

3. Using virtual work, calculate the reactions for the beams shown, and the bending moments at salient points.


Beam 1


Beam 2


Beam 3

### 8.4 Application of Virtual Forces

### 8.4.1 Basis

When virtual forces are applied, we have:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\int e_{i} \cdot \delta P_{i}
\end{aligned}
$$

And again note that we have an equilibrium set of forces and a compatible set of displacements:

$$
\begin{aligned}
& \text { Set of compatible } \\
& \text { displacements } \\
& \sum y_{i} \cdot \delta F_{i}=\int e_{i} \cdot \delta P_{i} \\
& \text { Set of forces in } \\
& \text { equilibrium }
\end{aligned}
$$

In this case the displacements are the real displacements that occur when the structure is in equilibrium and the virtual forces are any set of arbitrary forces that are in equilibrium.

### 8.4.2 Deflection of Trusses

## Basis

In a truss, we know:

1. The forces in the members (got from virtual displacements or statics);
2. Hence we can calculate the member extensions, $e_{i}$ as:

$$
e_{i}=\left(\frac{P L}{E A}\right)_{i}
$$

3. The virtual work equation is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i}
\end{aligned}
$$

4. To find the deflection at a single joint on the truss we apply a unit virtual force at that joint, in the direction of the required displacement, giving:

$$
y=\sum\left(\frac{P L}{E A}\right)_{i} \cdot \delta P_{i}
$$

5. Since in this equation, $y$ is the only unknown, we can calculate the deflection of the truss at the joint being considered.

## Example 6

## Problem

For the truss of Example 3, given that $E=200 \mathrm{kN} / \mathrm{mm}^{2}$ and $A=100 \mathrm{~mm}^{2}$ for each member, calculate the vertical and horizontal deflection of joint $B$.


## Solution

In Example 3 we found the forces in the members using virtual work to be:


The actual (real) compatible displacements are shown:


Two possible sets of virtual forces are:


In these problems we will always choose $\delta F=1$. Hence we apply a unit virtual force to joint $B$. We apply the force in the direction of the deflection required. In this way no work is done along deflections that are not required. Hence we have:


For horizontal deflection

The forces and elongations of the truss members are:

- Member AB:

$$
\begin{aligned}
P_{A B} & =+50 \mathrm{kN} \\
e_{A B} & =\left(\frac{+50 \cdot 5000}{200 \cdot 100}\right) \\
& =+12.5 \mathrm{~mm}
\end{aligned}
$$

- Member BC:

$$
\begin{aligned}
P_{B C} & =-30 \mathrm{kN} \\
e_{B C} & =\left(\frac{-30 \cdot 3000}{200 \cdot 100}\right) \\
& =-4.5 \mathrm{~mm}
\end{aligned}
$$

Note that by taking tension to be positive, elongations are positive and contractions are negative.

## Horizontal Deflection:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i} \\
y \cdot 1 & =(12.5 \cdot 0)_{A B}+(-4.5 \cdot+1)_{B C} \\
y & =-4.5 \mathrm{~mm}
\end{aligned}
$$

## Vertical Deflection:

$$
\begin{aligned}
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i} \\
y \cdot 1 & =\left(12.5 \cdot \frac{5}{4}\right)_{A B}+\left(-4.5 \cdot-\frac{3}{4}\right)_{B C} \\
y & =+18.4 \mathrm{~mm}
\end{aligned}
$$

Note that the sign of the result indicates whether the deflection occurs in the same direction as the applied force. Hence, joint $B$ moves 4.5 mm to the left.

## Example 7

## Problem

Determine the vertical and horizontal deflection of joint $D$ of the truss shown. Take $E$ $=200 \mathrm{kN} / \mathrm{mm}^{2}$ and member areas, $A=1000 \mathrm{~mm}^{2}$ for all members except $A E$ and $B D$ where $A=1000 \sqrt{ } 2 \mathrm{~mm}^{2}$.


## Solution

The elements of the virtual work equation are:

- Compatible deformations: The actual displacements that the truss undergoes;
- Equilibrium set: The external virtual force applied at the location of the required deflection and the resulting internal member virtual forces.

Firstly we analyse the truss to determine the member forces in order to calculate the actual deformations of each member:


Next, we apply a unit virtual force in the vertical direction at joint $D$. However, by linear superposition, we know that the internal forces due to a unit load will be $1 / 150$ times those of the 150 kN load.

For the horizontal deflection at $D$, we apply a unit horizontal virtual force as shown:


## Equations of Virtual Work

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i} \\
y_{D V} \cdot 1 & =\sum\left(\frac{P^{0} L}{E A}\right)_{i} \cdot \delta P_{i}^{1} \\
y_{D H} \cdot 1 & =\sum\left(\frac{P^{0} L}{E A}\right)_{i} \cdot \delta P_{i}^{2}
\end{aligned}
$$

In which:

- $P^{0}$ are the forces due to the 150 kN load;
- $\delta P^{1}$ are the virtual forces due to the unit virtual force applied in the vertical direction at $D$ :

$$
\delta P^{1}=\frac{P^{0}}{150}
$$

- $\delta P^{2}$ are the virtual forces due to the unit virtual force in the horizontal direction at $D$.

Using a table is easiest because of the larger number of members:

| Member | $L$ | $A$ | $P^{0}$ | $\delta P^{1}=\frac{P^{0}}{150}$ | $\delta P^{2}$ | $\left(\frac{P^{0} L}{A}\right) \cdot \delta P^{1}$ | $\left(\frac{P^{0} L}{A}\right) \cdot \delta P^{2}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $(\mathrm{~mm})$ | $\left(\mathrm{mm}^{2}\right)$ | $(\mathrm{kN})$ | $(\mathrm{kN})$ | $(\mathrm{kN})$ | $(\mathrm{kN} / \mathrm{mm}) \times \mathrm{kN}$ | $(\mathrm{kN} / \mathrm{mm}) \times \mathrm{kN}$ |
| AB | 2000 | 1000 | +150 | +1 | 0 | +300 | 0 |
| AE | $2000 \sqrt{ } 2$ | $1000 \sqrt{ } 2$ | $+150 \sqrt{ } 2$ | $+1 \sqrt{ } 2$ | 0 | +600 | 0 |
| AF | 2000 | 1000 | 0 | 0 | 0 | 0 | 0 |
| BC | 2000 | 1000 | 0 | 0 | 0 | 0 | 0 |
| BD | $2000 \sqrt{ } 2$ | $1000 \sqrt{ } 2$ | $+150 \sqrt{ } 2$ | $+1 \sqrt{ } 2$ | 0 | +600 | 0 |
| BE | 2000 | 1000 | -150 | -1 | 0 | +300 | 0 |
| CD | 2000 | 1000 | 0 | 0 | 0 | 0 | 0 |
| DE | 2000 | 1000 | -150 | -1 | +1 | +300 | -300 |
| EF | 2000 | 1000 | -300 | -2 | +1 | +1200 | -600 |
|  |  |  |  |  |  |  |  |

$E$ is left out because it is common. Returning to the equations, we now have:

$$
\begin{aligned}
y_{D V} \cdot 1 & =\frac{1}{E} \sum\left(\frac{P^{0} L}{A}\right)_{i} \cdot \delta P_{i}^{1} \\
y_{D V} & =\frac{+3300}{200}=+16.5 \mathrm{~mm}
\end{aligned}
$$

Which indicates a downwards deflection and for the horizontal deflection:

$$
\begin{aligned}
y_{D H} \cdot 1 & =\frac{1}{E} \sum\left(\frac{P^{0} L}{A}\right)_{i} \cdot \delta P_{i}^{2} \\
y_{D H} & =\frac{-900}{200}=-4.5 \mathrm{~mm}
\end{aligned}
$$

The sign indicates that it is deflecting to the left.

### 8.4.3 Deflection of Beams and Frames

## Example 8

## Problem

Using virtual work, calculate the deflection at the centre of the beam shown, given that $E I$ is constant.


## Solution

To calculate the deflection at $C$, we will be using virtual forces. Therefore the two relevant sets are:

- Compatibility set: the actual deflection at $C$ and the rotations that occur along the length of the beam;
- Equilibrium set: a unit virtual force applied at $C$ which is in equilibrium with the internal virtual moments it causes.


## Compatibility Set:

The external deflection at $C$ is what is of interest to us. To calculate the rotations along the length of the beam we have:

$$
\kappa_{x}=\frac{M_{x}}{E I_{x}}
$$

Hence we need to establish the bending moments along the beam:


For $A C$ the bending moment is given by (and similarly for $B$ to $C$ ):

$$
M_{x}=\frac{P}{2} x
$$

## Equilibrium Set:

As we choose the value for $\delta F=1$, we are only left to calculate the virtual moments:


For $A C$ the internal virtual moments are given by:

$$
\delta M_{x}=\frac{1}{2} x
$$

## Virtual Work Equation

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum \kappa_{i} \cdot \delta M_{i}
\end{aligned}
$$

Substitute in the values we have for the real rotations and the virtual moments, and use the fact that the bending moment diagrams are symmetrical:

$$
\begin{aligned}
y \cdot 1 & =2 \int_{0}^{L / 2}\left[\frac{M_{x}}{E I}\right] \cdot \delta M_{x} d x \\
y & =\frac{2}{E I} \int_{0}^{L / 2}\left(\frac{P}{2} x\right) \cdot\left(\frac{1}{2} x\right) d x \\
& =\frac{2 P}{4 E I} \int_{0}^{L / 2} x^{2} d x \\
& =\frac{P}{2 E I}\left[\frac{x^{3}}{3}\right]_{0}^{L / 2} \\
& =\frac{P}{6 E I} \cdot \frac{L^{3}}{8} \\
& =\frac{P L^{3}}{48 E I}
\end{aligned}
$$

Which is a result we expected.

## Example 9

## Problem

Find the vertical and horizontal displacement of the prismatic curved bar shown:


## Solution

Even though it is curved, from the statics, the bending moment at any point of course remains force by distance, so from the following diagram:


At any angle $\phi$, we therefore have:

$$
\begin{aligned}
M_{\phi} & =P(R-R \cos \phi) \\
& =P R(1-\cos \phi)
\end{aligned}
$$

To find the displacements, we follow our usual procedure and place a unit load at the location of, and in the direction of, the required displacement.

## Vertical Displacement

The virtual bending moment induced by the vertical unit load shown, is related to that for $P$ and is thus:

$$
\delta M_{\phi}=R(1-\cos \phi)
$$



Thus our virtual work equations are:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
1 \cdot \delta_{B V} & =\int \kappa \cdot \delta M \cdot d s \\
& =\int_{0}^{\pi / 2} \frac{M_{\phi}}{E I} \cdot \delta M_{\phi} \cdot R d \phi
\end{aligned}
$$

In which we have used the relation $d s=R d \phi$ to change the integration from along the length of the member to around the angle $\phi$. Next we introduce our equations for the real and virtual bending moments:

$$
\begin{aligned}
1 \cdot \delta_{B V} & =\int_{0}^{\pi / 2} \frac{P R(1-\cos \phi)}{E I} \cdot R(1-\cos \phi) \cdot R d \phi \\
& =\frac{P R^{3}}{E I} \int_{0}^{\pi / 2}(1-\cos \phi)^{2} \cdot d \phi \\
& =\frac{P R^{3}}{E I}(3 \pi-4) \approx 5.42 \frac{P R^{3}}{E I}
\end{aligned}
$$

## Horizontal Displacement

In this case, the virtual bending moment is:

$$
\delta M_{\phi}=R \sin \phi
$$



Thus the virtual work equations give:

$$
\begin{aligned}
1 \cdot \delta_{B H} & =\int_{0}^{\pi / 2} \frac{P R(1-\cos \phi)}{E I} \cdot R \sin \phi \cdot R d \phi \\
& =\frac{P R^{3}}{E I} \int_{0}^{\pi / 2}(\sin \phi-\sin \phi \cos \phi) \cdot d \phi \\
& =\frac{P R^{3}}{E I} \int_{0}^{\pi / 2}\left(\sin \phi-\frac{\sin 2 \phi}{2}\right) \cdot d \phi \\
& =\frac{P R^{3}}{E I}\left(\frac{1}{2}\right)=\frac{P R^{3}}{2 E I}
\end{aligned}
$$

### 8.4.4 Integration of Bending Moments

## General Volume Integral Expression

We are often faced with the integration of being moment diagrams when using virtual work to calculate the deflections of bending members. And as bending moment diagrams only have a limited number of shapes, a table of 'volume' integrals is used:

## Derivation

The general equation for the volume integral is derived consider the real bending moment diagram as parabolic, and the virtual bending moment diagram as linear:


In terms of the left and right end ( $y_{0}$ and $y_{L}$ respectively) and midpoint value $\left(y_{m}\right)$, a second degree parabola can be expressed as:

$$
y(x)=2\left(y_{L}-2 y_{m}+y_{0}\right)\left(\frac{x}{L}\right)^{2}+\left(4 y_{m}-3 y_{0}-y_{L}\right)\left(\frac{x}{L}\right)+y_{0}
$$

For the case linear case, $y_{m}=0.5\left(y_{L}-y_{0}\right)$, and the expression reduces to that of straight line. For the real bending moment (linear or parabolic) diagram, we therefore have:

$$
M(x)=2\left(M_{L}-2 M_{m}+M_{0}\right)\left(\frac{x}{L}\right)^{2}+\left(4 M_{m}-3 M_{0}-M_{L}\right)\left(\frac{x}{L}\right)+M_{0}
$$

And for the linear virtual moment diagram we have:

$$
\delta M(x)=\left(\delta M_{L}-\delta M_{0}\right)\left(\frac{x}{L}\right)+\delta M_{0}
$$

The volume integral is given by:

$$
V=\int_{0}^{L} M \cdot \delta M d x
$$

Which using Simpsons Rule, reduces to:

$$
\int_{0}^{L} M \cdot \delta M d x=\frac{L}{6}\left(\delta M_{0} \cdot M_{0}+4 \delta M_{m} \cdot M_{m}+\delta M_{L} \cdot M_{L}\right)
$$

This expression is valid for the integration of all linear by linear or parabolic diagrams.

## Example

As an example, if we have a rectangle by a triangle:

$$
\begin{gathered}
M(x)=k \\
\delta M(x)=j \frac{x}{L}
\end{gathered}
$$

We get:

$$
\int_{0}^{l} M \cdot \delta M d x=\frac{l}{6}\left(4 \frac{j}{2} k+j k\right)=\frac{1}{2} j k l
$$

And this entry is given in the table as follows:

|  |  | j |  |  |
| :---: | :---: | :---: | :---: | :---: |
|  | $\frac{1}{3} j k l$ | $\frac{1}{6} j k l$ | $\frac{1}{6}\left(j_{1}+2 j_{2}\right) k l$ | $\frac{1}{2} j k l$ |
| k | $\frac{1}{6} j k l$ | $\frac{1}{3} j k l$ | $\frac{1}{6}\left(2 j_{1}+j_{2}\right) k l$ | $\frac{1}{2} j k l$ |

This table is at the back page of these notes for ease of reference.

## Example 10

## Problem

Using the table of volume integrals, verify the answer to Example 8.

## Solution

In this case, the virtual work equation becomes:

$$
\begin{aligned}
y \cdot 1 & =\int\left[\frac{M}{E I}\right] \cdot \delta M d x \\
y & =\frac{1}{E I}\left[M_{x} \text { shape }\right] \times\left[\delta M_{x} \text { shape }\right]
\end{aligned}
$$



M


SM


M

$$
y=\frac{2}{E I}\left[\frac{1}{3} \cdot \frac{P L}{4} \cdot \frac{L}{4} \cdot \frac{L}{2}\right]=\frac{P L^{3}}{48 E I}
$$

In which the formula $\frac{1}{3} j k l$ is used from the table.

## Example 11

## Problem

Determine the influence of shear deformation on the deflection of a rectangular prismatic mild steel cantilever subjected to point load at its tip.

## Solution

Since we are calculating deflections, and must include the shear contribution, we apply a virtual unit point load at the tip and determine the real and virtual bending moment and shear force diagrams:


The virtual work done has contributions by both the bending and shear terms:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
1 \cdot \Delta & =\int \kappa \cdot \delta M \cdot d s+\int \gamma \cdot \delta V \cdot d s \\
& =\int \frac{M}{E I} \cdot \delta M \cdot d s+\int \frac{V}{G A_{v}} \cdot \delta V \cdot d s
\end{aligned}
$$

Using the Table of Volume Integrals, we have:

$$
\begin{aligned}
\Delta & =\frac{1}{E I}\left[\frac{1}{3} \cdot P L \cdot L \cdot L\right]+\frac{1}{G A_{v}}[P \cdot 1 \cdot L] \\
& =\frac{P L^{3}}{3 E I}+\frac{P L}{G A_{v}}
\end{aligned}
$$

The bending term and shear term contributions are readily apparent.

Since mild steel is an isotropic material, we will take the following properties for shear modulus and Poisson's ratio:

$$
G=\frac{E}{2(1+v)} \quad v=0.3
$$

We now have:

$$
\Delta=\frac{P L}{E}\left[\frac{L^{2}}{3 I}+\frac{2(1+v)}{A_{v}}\right]
$$

This is a general solution, independent of cross section. Now we will introduce the cross-section properties for a rectangle:

$$
I=\frac{b h^{3}}{12} \quad A_{v}=k A=k b h
$$

Where $k=0.67$ is a shear factor (proportion of area effective in shear), giving:

$$
\begin{aligned}
\Delta & =\frac{P L}{E}\left[\frac{12 L^{2}}{3 b h^{3}}+\frac{2(1+v)}{k b h}\right] \\
& =\frac{P L}{E b h}\left[4\left(\frac{L}{h}\right)^{2}+\frac{2(1+v)}{k}\right]
\end{aligned}
$$

For a large span-depth ratio of 10 , we have:

$$
\begin{aligned}
\frac{E b h}{P L} \Delta & =4(10)^{2}+\frac{2(1+0.3)}{0.67} \\
& =400+3.9
\end{aligned}
$$

And so the shear deflection is a very small percentage $(3.9 / 403.9=1 \%)$ of the total deflection. For a small span-depth ratio of say 3:

$$
\begin{aligned}
\frac{E b h}{P L} \Delta & =4(3)^{2}+\frac{2(1+0.3)}{0.67} \\
& =36+3.9
\end{aligned}
$$

And the shear deformation is $9.1 \%$ of the deflection - too small to ignore.


## Example 12

## Problem

Determine the static deflection a vehicle experiences as it crosses a simply-supported bridge. Model the vehicle as a point load.

## Solution

Since the vehicle moves across the bridge we need to calculate its deflection at some arbitrary location along the bridge. Let's say that at some point in time it is at a fraction $\xi$ along the length of the bridge:


To determine the deflection at $C$ at this point in time, we apply a unit virtual force at $C$. Thus, we have both real and virtual bending moment diagrams:


The equation for the real bending moment at $C$ is:

$$
M_{C}=\frac{P \cdot \xi L \cdot(1-\xi) L}{L}=P L \xi(1-\xi)
$$

And similarly:

$$
d M_{C}=L \xi(1-\xi)
$$

Thus:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
1 \cdot \Delta_{C} & =\int \kappa \cdot \delta M \cdot d s \\
& =\int \frac{M}{E I} \cdot \delta M \cdot d s \\
& =\frac{P}{E I}\left\{\begin{aligned}
& {\left[\frac{1}{3} \cdot \xi(1-\xi) L \cdot \xi(1-\xi) L \cdot \xi L\right]_{A C} } \\
&+\left[\frac{1}{3} \cdot \xi(1-\xi) L \cdot \xi(1-\xi) L \cdot(1-\xi) L\right]_{C B}
\end{aligned}\right\}
\end{aligned}
$$

Which after some algebra gives the rather delightful:

$$
\Delta_{C}=\frac{P L^{3}}{3 E I}[\xi(1-\xi)]^{2}
$$

For $\xi=0.5$ this reduces to the familiar:

$$
\Delta_{C}=\frac{P L^{3}}{3 E I}[0.5(1-0.5)]^{2}=\frac{P L^{3}}{48 E I}
$$

If the vehicle is travelling at a constant speed $v$, then its position at time $t$ is $x=v t$. Hence is fraction along the length of the bridge is:

$$
\xi=\frac{x}{L}=\frac{v t}{L}
$$

The proportion of deflection related to the midspan deflection is thus:

$$
\frac{\Delta}{\Delta_{\xi=0.5}}=\frac{\frac{P L^{3}}{3 E I}[\xi(1-\xi)]^{2}}{\frac{P L^{3}}{48 E I}}=\frac{48}{3}[\xi(1-\xi)]^{2}
$$

And so:

$$
\frac{\Delta}{\Delta_{\xi=0.5}}=\frac{48}{3}[\xi(1-\xi)]^{2}
$$



## Example 13 - Summer '07 Part (a)

## Problem

For the frame shown, determine the horizontal deflection of joint C. Neglect axial effects in the members and take $E I=36 \times 10^{3} \mathrm{kNm}^{2}$.


## Solution

Firstly we establish the real bending moment diagram:


Next, as usual, we place a unit load at the location of, and in the direction of, the required displacement:

band - SM

Now we have the following for the virtual work equation:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
1 \cdot \delta_{C H} & =\int \kappa \cdot \delta M \cdot d s \\
& =\int \frac{M}{E I} \cdot \delta M \cdot d s
\end{aligned}
$$

Next, using the table of volume integrals, we have:

$$
\begin{aligned}
\int \frac{M}{E I} \cdot \delta M \cdot d s & =\frac{1}{E I}\left\{\left[\frac{1}{3}(440)(2)(4 \sqrt{2})\right]_{A B}+\left[\frac{1}{6}(2)(160+2 \cdot 440)(4)\right]_{B C}\right\} \\
& =\frac{1659.3}{E I}+\frac{1386.7}{E I}=\frac{3046}{E I}
\end{aligned}
$$

Hence:

$$
1 \cdot \delta_{C H}=\frac{3046}{E I}=\frac{3046}{36 \times 10^{3}} \times 10^{3}=84.6 \mathrm{~mm}
$$

### 8.4.5 Problems

1. Determine the vertical deflection of joint $E$ of the truss shown. Take $E=200$ $\mathrm{kN} / \mathrm{mm}^{2}$ and member areas, $A=1000 \mathrm{~mm}^{2}$ for all members except $A C$ where $A=$ $1000 \sqrt{ } 2 \mathrm{~mm}^{2}$. (Ans. 12.2 mm ).

2. Determine the vertical and horizontal deflection of joint $C$ of the truss shown. Take $E=10 \mathrm{kN} / \mathrm{mm}^{2}$ and member areas, $A=1000 \mathrm{~mm}^{2}$ for all members except $A C$ where $A=1000 \sqrt{ } 2 \mathrm{~mm}^{2}$ and $C E$ where $A=2500 \mathrm{~mm}^{2}$.
(Ans. $32 / 7 \mathrm{~mm}$ horizontal, $74 / 7 \mathrm{~mm}$ vertical).

3. Determine the horizontal deflection of joint $\underline{A}$ and the vertical deflection of joint B of the truss shown. Take $E=200 \mathrm{kN} / \mathrm{mm}^{2}$ and member areas, $\mathrm{A}=1000 \mathrm{~mm}^{2}$ for all members except $B D$ where $A=1000 \sqrt{ } 2 \mathrm{~mm} 2$ and $A B$ where $A=2500$ $\mathrm{mm}^{2}$. (Ans. $15.3 \mathrm{~mm} ; 0 \mathrm{~mm}$ )

4. Determine the vertical and horizontal deflection of joint $C$ of the truss shown. Take $E=200 \mathrm{kN} / \mathrm{mm}^{2}$ and member areas, $A=150 \mathrm{~mm}^{2}$ for all members except $A C$ where $A=150 \sqrt{ } 2 \mathrm{~mm} 2$ and $C D$ where $A=250 \mathrm{~mm}^{2}$. $($ Ans. $8.78 \mathrm{~mm} ; 0.2 \mathrm{~mm})$

5. Determine the vertical deflection of joint $D$ of the truss shown. Take $E=200$ $\mathrm{kN} / \mathrm{mm}^{2}$ and member areas, $A=1000 \mathrm{~mm}^{2}$ for all members except $B C$ where $A=$ $1000 \sqrt{ } 2 \mathrm{~mm}^{2} .($ Ans. 24 mm$)$

6. Verify that the deflection at the centre of a simply-supported beam under a uniformly distributed load is given by:

$$
\delta_{C}=\frac{5 w L^{4}}{384 E I}
$$

7. Show that the flexural deflection at the tip of a cantilever, when it is subjected to a point load on the end is:

$$
\delta_{B}=\frac{P L^{3}}{3 E I}
$$

8. Show that the rotation that occurs at the supports of a simply supported beam with a point load in the middle is:

$$
\theta_{C}=\frac{P L^{2}}{16 E I}
$$

9. Show that the vertical and horizontal deflections at $B$ of the following frame are:

$$
\delta_{B y}=\frac{P R^{3}}{2 E I} ; \quad \delta_{B x}=\frac{P R^{3} \pi}{4 E I}
$$


10. For the frame of Example 13 - Summer '07 Part (a), using virtual work, verify the following displacements in which the following directions are positive: vertical upwards; horizontal to the right, and; clockwise:

- Rotation at $A$ is $-\frac{1176.3}{E I}$;
- Vertical deflection at $B$ is $-\frac{3046}{E I}$;
- Horizontal deflection at $D$ is $\frac{8758.7}{E I}$;
- Rotation at $D$ is $\frac{1481.5}{E I}$.

11. For the cantilever beam below, show that the vertical deflection of $B$, allowing for both flexural and shear deformation, is $\frac{P L^{3}}{3 E I}+\frac{P L}{G A_{v}}$.

12. For the beam below, show that the vertical deflection of $D$ is zero and that that of $C$ is $\frac{1850 / 3}{E I}$ :


Problems 13 to 17 are considered Genius Level!
13. Show that the vertical deflection at the centre of a simply supported beam subject to UDL, allowing for both flexural and shear deformation, is $\frac{5 w L^{4}}{384 E I}+\frac{w L^{2}}{8 G A_{v}}$.
14. For the frame shown, show that the vertical deflection of point $D$ is $\frac{4594}{E I} \downarrow$. Neglect axial deformation and take $E I=120 \times 10^{3} \mathrm{kNm}^{2}$.

15. For the frame shown, determine the vertical deflection of $C$ and the rotation at $D$ about the global $x$ axis. Take the section and material properties as follows:
$A=6600 \mathrm{~mm}^{2} \quad E=205 \mathrm{kN} / \mathrm{mm}^{2} \quad G=100 \mathrm{kN} / \mathrm{mm}^{2}$
$I=36 \times 10^{6} \mathrm{~mm}^{4} \quad J=74 \times 10^{6} \mathrm{~mm}^{4}$.


(Ans. $12.23 \mathrm{~mm} \uparrow ; 0.0176$ rads ACW)
16. Determine the horizontal and vertical deflections at $B$ in the following frame. Neglect shear and axial effects.

(Ans. $\left.\delta_{B x}=\frac{2 w R^{4}}{E I} \rightarrow ; \delta_{B y}=\frac{3 \pi w R^{4}}{2 E I} \uparrow\right)$
17. For the frame shown, neglecting shear effects, determine the vertical deflection of points $B$ and $A$.

$\left(A n s . \delta_{B y}=\frac{(P+w a) b^{3}}{3 E I} \downarrow ; \delta_{A y}=\frac{1}{E I}\left[\frac{P a^{3}}{3}+\frac{w a^{4}}{8}+\frac{(P+w a) b^{3}}{3}\right]+\left(P a+\frac{w a^{2}}{2}\right) \frac{a b}{G J} \downarrow\right)$

### 8.5 Virtual Work for Indeterminate Structures

### 8.5.1 General Approach

Using the concept of compatibility of displacement, any indeterminate structure can be split up into a primary and reactant structures. In the case of a one-degree indeterminate structure, we have:


At this point the primary structure can be analysed. However, we further break up the reactant structure, using linear superposition:


We summarize this process as:

$$
M=M^{0}+\alpha M^{1}
$$

- $M$ is the force system in the original structure (in this case moments);
- $M^{0}$ is the primary structure force system;
- $M^{1}$ is the unit reactant structure force system.

For a truss, the procedure is the same:



The final system forces are:

$$
P=P^{0}+\alpha P^{1}
$$

The primary structure can be analysed, as can the unit reactant structure. Therefore, the only unknown is the multiplier, $\alpha$.

We use virtual work to calculate the multiplier $\alpha$.

### 8.5.2 Using Virtual Work to Find the Multiplier

We must identify the two sets for use:

- Displacement set: We use the actual displacements that occur in the real structure;
- Equilibrium set: We use the unit reactant structure's set of forces as the equilibrium set. We do this, as the unit reactant is always a determinate structure and has a configuration similar to that of the displacement set.

A typical example of these sets for an indeterminate structure is shown:


Compatibility Set
Internal and external real
displacements linked through constitutive relations


## Equilibrium Set

Internal (bending moments) and external virtual forces are linked through equilibrium

Note that the Compatibility Set has been chosen as the real displacements of the structure. The virtual force (equilibrium) set chosen is a statically determine substructure of the original structure. This makes it easy to use statics to determine the internal virtual 'forces' from the external virtual force. Lastly, note that the external virtual force does no virtual work since the external real displacement at $B$ is zero. This makes the right hand side of the equation zero, allowing us to solve for the multiplier.

The virtual work equation (written for trusses) gives:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i} \\
0 \cdot 1 & =\sum\left(\frac{P L}{E A}\right)_{i} \cdot \delta P_{i}^{1}
\end{aligned}
$$

There is zero external virtual work. This is because the only virtual force applied is internal; no external virtual force applied. Also note that the real deformations that occur in the members are in terms of $P$, the unknown final forces. Hence, substituting $P=P^{0}+\alpha \cdot \delta P^{1}$ (where $\delta$ is now used to indicate virtual nature):

$$
\begin{aligned}
0 & =\sum\left(\frac{\left(P^{0}+\alpha \cdot \delta P^{1}\right) L}{E A}\right)_{i} \cdot \delta P_{i}^{1} \\
& =\sum\left(\frac{P^{0} L}{E A}\right)_{i} \cdot \delta P_{i}^{1}+\alpha \cdot \sum\left(\frac{\delta P^{1} L}{E A}\right)_{i} \cdot \delta P_{i}^{1} \\
0 & =\sum \frac{P^{0} \cdot \delta P_{i}^{1} \cdot L_{i}}{E A_{i}}+\alpha \cdot \sum \frac{\left(\delta P_{i}^{1}\right)^{2} L_{i}}{E A_{i}}
\end{aligned}
$$

For beams and frames, the same development is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\lceil ] \kappa_{i} \cdot \delta M_{i} \\
0 \cdot 1 & =\left[\frac{M_{i}}{E I} \cdot \delta M_{i}\right.
\end{aligned}
$$

Where again there is no external displacement of the virtual force.

Note that we use the contour integral symbol $\|[f]$ simply to indicate that we integrate around the structure, accounting for all members in the beam/frame (i.e. integrate along the length of each member separately, and then sum the results).

Also, substitute $M=M^{0}+\alpha \cdot \delta M^{1}$ to get:

$$
\begin{aligned}
& 0=\mathfrak{j} \frac{\left(M^{0}+\alpha \cdot \delta M^{1}\right)}{E I_{i}} \delta M_{i}^{1} d x \\
& 0=\left\{\left[\frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x+\alpha \cdot \mathfrak{j} \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x\right.\right.
\end{aligned}
$$

Thus in both bases we have a single equation with only one unknown, $\alpha$. We can establish values for the other two terms and then solve for $\alpha$ and the structure as a whole.

## Example 14

## Problem

Calculate the bending moment diagram for the following prismatic propped cantilever. Find also the deflection under the point load.


## Solution

To illustrate the methodology of selecting a redundant and analysing the remaining determinate structure for both the applied loads ( $M^{0}$ ) and the virtual unit load ( $\delta M$ ), we will analyse the simple case of the propped cantilever using two different redundants. The final solution will be the same (as it should), and this emphasis that it does not matter which redundant is taken.

Even though the result will be the same, regardless of the redundant chosen, the difficulty of the calculation may not be the same. As a result, a good choice of redundant can often be made so that the ensuing calculations are made easier. A quick qualitative assessment of the likely real and virtual bending moment diagrams for a postulated redundant should be made, to see how the integrations will look.

## Redundant: Vertical Reaction at $B$

Selecting the prop at $B$ as the redundant gives a primary structure of a cantilever. We determine the primary and (unit) virtual bending moment diagrams as follows:


The virtual work equation is:

$$
\left.0=\emptyset \frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x+\alpha \cdot \emptyset\right\rceil \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x
$$

The terms of the virtual work equation are:

$$
\begin{gathered}
E I \int \frac{M^{0} \cdot \delta M}{E I} d x=\left[\frac{1}{6}\left(\frac{P L}{2}\right)\left(-\frac{L}{2}+2 \cdot-L\right)\left(\frac{L}{2}\right)\right]_{A C} \\
=\frac{-5 P L^{3}}{48} \\
E I \int \frac{(\delta M)^{2}}{E I} d x=\frac{1}{3}(L)(L)(L)=\frac{L^{3}}{3}
\end{gathered}
$$

Giving:

$$
0=\frac{1}{E I}\left(-\frac{5 P L^{3}}{48}\right)+\alpha \cdot \frac{1}{E I}\left(\frac{L^{3}}{3}\right)
$$

And so:

$$
\alpha=\left(\frac{3}{L^{3}}\right)\left(\frac{5 P L^{3}}{48}\right)=\frac{5 P}{16}
$$

And this is the value of the reaction at $B$.

## Redundant: Moment Reaction at $A$

Choosing the moment restraint at $A$ as the redundant gives a simply-support beam as the primary structure. The bending moments diagrams and terms of are thus:

$$
\begin{aligned}
& E I \int \frac{M^{0} \cdot \delta M}{E I} d x=\left[\frac{1}{6}\left(\frac{P L}{4}\right)\left(2 \cdot-\frac{1}{2}+-1\right)\left(\frac{L}{2}\right)\right]_{A C}+\left[\frac{1}{3}\left(\frac{P L}{4}\right)\left(-\frac{1}{2}\right)\left(\frac{L}{2}\right)\right]_{B C} \\
& =\frac{-P L^{2}}{48}+\frac{-2 P L^{2}}{48} \\
& =-\frac{P L^{2}}{16}
\end{aligned}
$$

$$
E I \int \frac{(\delta M)^{2}}{E I} d x=\frac{1}{3}(-1)(-1)(L)=\frac{L}{3}
$$

Giving:

$$
0=\frac{1}{E I}\left(-\frac{P L^{2}}{16}\right)+\alpha \cdot \frac{1}{E I}\left(\frac{L}{3}\right)
$$

And so:

$$
\alpha=\left(\frac{3}{L}\right)\left(\frac{P L^{2}}{16}\right)=\frac{3 P L}{16}
$$

And this is the value of the moment at $A$.

## Overall Solution

Both cases give the following overall solution, as they should:


### 8.5.3 Indeterminate Trusses

## Example 15

## Problem

Calculate the forces in the truss shown and find the horizontal deflection at $C$. Take $E A$ to be $10 \times 10^{4} \mathrm{kN}$ for all members.


## Solution

## Solve for the Truss Member Forces

Choose member $A C$ as the redundant:


Next analyse for the $P^{0}$ and $P^{1}$ force systems.


Using a table:

| Member | $L$ | $P^{0}$ | $\delta P^{1}$ | $P^{0} \cdot \delta P^{1} \cdot L$ | $\left(\delta P^{1}\right)^{2} \cdot L$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  | $(\mathrm{~mm})$ | $(\mathrm{kN})$ | $(\mathrm{kN})$ | $\times 10^{4}$ | $\times 10^{4}$ |
|  | 4000 | +40 | $-4 / 5$ | -12.8 | 0.265 |
| BC | 3000 | 0 | $-3 / 5$ | 0 | 0.108 |
| CD | 4000 | 0 | $-4 / 5$ | 0 | 0.256 |
| AC | 5000 | 0 | 1 | 0 | 0.5 |
| BD | 5000 | -50 | 1 | -25 | 0.5 |

Hence:

$$
\begin{aligned}
0 & =\sum \frac{P^{0} \cdot \delta P_{i}^{1} \cdot L_{i}}{E A_{i}}+\alpha \cdot \sum \frac{\left(\delta P_{i}^{1}\right)^{2} L_{i}}{E A_{i}} \\
& =\frac{-37.8 \times 10^{4}}{E A}+\alpha \cdot \frac{1.62 \times 10^{4}}{E A}
\end{aligned}
$$

And so

$$
\alpha=-\frac{-37.8}{1.62}=23.33
$$

The remaining forces are obtained from the compatibility equation:

| Member | $P^{0}$ | $\delta P^{1}$ | $P=P^{0}+\alpha \cdot \delta P^{1}$ |
| :---: | :---: | :---: | :---: |
|  | $(\mathrm{kN})$ | $(\mathrm{kN})$ | $(\mathrm{kN})$ |
|  | +40 | $-4 / 5$ | 21.36 |
| BC | 0 | $-3 / 5$ | -14 |
| CD | 0 | $-4 / 5$ | -18.67 |
| AC | 0 | 1 | 23.33 |
| BD | -50 | 1 | -26.67 |

Note that the redundant always has a force the same as the multiplier.

## Calculate the Horizontal Deflection at C

To calculate the horizontal deflection at $C$, using virtual work, the two relevant sets are:

- Compatibility set: the actual deflection at $C$ and the real deformations that occur in the actual structure;
- Equilibrium set: a horizontal unit virtual force applied at $C$ to a portion of the actual structure, yet ensuring equilibrium.

We do not have to apply the virtual force to the full structure. Remembering that the only requirement on the virtual force system is that it is in equilibrium; choose the force systems as follows:


Thus we have:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i} \\
y_{C H} \cdot 1 & =\sum\left(\frac{P L}{E A}\right)_{i} \cdot \delta P_{i} \\
& =\left(\frac{23.33 \times 5000}{10 \times 10^{4}}\right) \cdot \frac{5}{3}+\left(\frac{-18.67 \times 4000}{10 \times 10^{4}}\right) \cdot \frac{-4}{3} \\
y_{C H} & =2.94 \mathrm{~mm}
\end{aligned}
$$

Because we have chosen only two members for our virtual force system, only these members do work and the calculation is greatly simplified.

### 8.5.4 Indeterminate Frames

## Example 16

## Problem

For the prismatic frame shown, calculate the reactions and draw the bending moment diagram. Determine the vertical deflection at joint $A$.


## Solution

Choosing the horizontal reaction at $C$ as the redundant gives the primary structure as:


And the unit redundant structure as:


The terms of the virtual work equation are:

$$
\begin{aligned}
E I \int \frac{M^{0} \cdot \delta M}{E I} d x & =\left[\frac{1}{3}(80)(-3)(5)\right]_{B C}=-400 \\
E I \int \frac{(\delta M)^{2}}{E I} d x & =\left[\frac{1}{3}(-3)(-3)(5)\right]_{B C}+\left[\frac{1}{3}(-3)(-3)(3)\right]_{B D}=24
\end{aligned}
$$

Giving:

$$
\begin{aligned}
0 & =\int \frac{M^{0} \cdot \delta M}{E I} d x+\alpha \cdot \int \frac{(\delta M)^{2}}{E I} d x \\
& =\frac{-400}{E I}+\alpha \cdot \frac{24}{E I}
\end{aligned}
$$

And so the multiplier on the unit redundant is:

$$
\alpha=\frac{400}{24}=16.7
$$

As a result, using superposition of the primary and unit redundant times the multiplier structures, we have:


To determine the deflection at $A$, we must apply a unit vertical virtual force at $A$ to obtain the second virtual force system, $\delta M^{2}$. However, to ease the computation, and since we only require an equilibrium system, we apply the unit vertical force to the primary structure subset of the overall structure as follows:


The deflection is then got from the virtual work equation as follows:

$$
1 \times \Delta=\int \frac{M \cdot \delta M^{2}}{E I}
$$

$$
\begin{aligned}
\Delta & =\frac{1}{E I}\left\{\left[\frac{1}{3}(80)(2)(2)\right]_{A B}+\left[\frac{1}{3}(30)(2)(5)\right]_{B C}\right\} \\
& =\frac{620 / 3}{E I}
\end{aligned}
$$



## Example 17

## Problem

For the prismatic frame shown, calculate the reactions and draw the bending moment diagram. Determine the horizontal deflection at joint $C$. Take $E I=40 \times 10^{3} \mathrm{kNm}^{2}$.


## Solution

Break the frame up into its reactant and primary structures:


Establish the $M^{0}$ and $M^{1}$ force systems:


Apply the virtual work equation:

$$
0=\oint \frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x+\alpha \cdot \mathfrak{f j} \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x
$$

We will be using the table of volume integrals to quicken calculations. Therefore we can only consider lengths of members for which the correct shape of bending moment diagram is available. Also, we must choose sign convention: we consider tension on the outside of the frame to be positive.

As each term has several components we consider them separately:

Term 1 - $\left\{\left[\frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x:\right.\right.$
We show the volume integrals beside the real and virtual bending moment diagrams for each member of the frame:

## Member

BMDs
Volume Integral


Hence:

$$
\int \frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x=\frac{-23400}{E I}
$$

## Term 2:

$$
\begin{aligned}
\int f \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x & =\frac{1}{E I}\left\{[j k l]_{A B}+\left[\frac{1}{3} j k l\right]_{B C}\right\} \\
& =\frac{1}{E I}\left\{[(-6)(-6) 8]_{A B}+\left[\frac{1}{3}(-6)(-6) 6\right]_{B C}\right\} \\
& =\frac{360}{E I}
\end{aligned}
$$

Note that Term 2 is always easier to calculate as it is only ever made up of straight line bending moment diagrams.

Thus we have:

$$
\begin{aligned}
0 & =\left\{\frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x+\alpha \cdot \oint \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x\right. \\
& =\frac{-23400}{E I}+\alpha \cdot \frac{360}{E I}
\end{aligned}
$$

And so:

$$
\alpha=\frac{23400}{360}=65.0
$$

Thus the vertical reaction at $C$ is 65.0 kN upwards. With this information we can solve for the moments using $M=M^{0}+\alpha M^{1}$ (or by just using statics) and the remaining reactions to get:


To calculate the horizontal deflection at $C$ using virtual work, the two relevant sets are:

- Compatibility set: the actual deflection at $C$ and the real deformations (rotations) that occur in the actual structure;
- Equilibrium set: a horizontal unit virtual force applied at $C$ to a determinate portion of the actual structure.

Choose the following force system as it is easily solved:


Thus we have:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\emptyset\left[\theta_{i} \cdot \delta M_{i}\right. \\
y_{C H} \cdot 1 & =\emptyset\left[\frac{M_{x}}{E I}\right] \cdot \delta M_{x} d x
\end{aligned}
$$

The real bending moment diagram, $M$, is awkward to use with the integral table. Remembering that $M=M^{0}+\alpha M^{1}$ simplifies the calculation by using:


To give:


And so using the table formulae, we have:

$$
\begin{aligned}
{\left[f\left[\frac{M_{x}}{E I}\right] \cdot \delta M_{x} d x\right.} & =\left[\frac{1}{2}(-4)(-6 \cdot \alpha) 4+\frac{1}{6}(2 \cdot 360+600)(-4) 4\right]_{A D}+ \\
& {\left[\frac{1}{2}(-4-8)(-6 \cdot \alpha) 4+\frac{1}{2}(-4-8)(360) 4\right]_{D B}+} \\
& {\left[\frac{1}{3}(-8)(-6 \cdot \alpha) 6+\frac{1}{4}(-8)(360) 6\right]_{B C} }
\end{aligned}
$$

Which gives us:

$$
\begin{aligned}
\delta_{C x} & =\left\lceil\left[\frac{M_{x}}{E I}\right] \cdot \delta M_{x} d x\right. \\
& =\frac{1}{E I}[288 \alpha+(-16480)] \\
& =\frac{2240}{E I} \\
& =0.056 \\
& =56 \mathrm{~mm}
\end{aligned}
$$

The answer is positive, indicating that the structure moves to the right at $C$ : the same direction in which the unit virtual force was applied.

We could have chosen any other statically determinate sub-structure and obtained the same result. Some sub-structures will make the analysis easier to perform.

Exercise: Verify that the same deflection is obtained by using a sub-structure obtained by removing the vertical support at $C$ and applying the virtual force. Does this sub-structure lead to an easier analysis for the deflection?

### 8.5.5 Multiply-Indeterminate Structures

## Basis

In this section we will introduce structures that are more than 1 degree statically indeterminate. We do so to show that virtual work is easily extensible to multiplyindeterminate structures, and also to give a method for such beams that is easily worked out, and put into a spreadsheet.

Consider the example 3 -span beam. It is 2 degrees indeterminate, and so we introduce 2 hinges (i.e. moment releases) at the support locations and unit reactant moments in their place, as shown:


Alongside these systems, we have their bending moment diagrams:

$M$


Using the idea of the multiplier and superposition again, we can see that:

$$
M=M^{0}+\alpha_{1} \cdot \delta M^{1}+\alpha_{2} \cdot \delta M^{2}
$$

The virtual work equation is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
0 \cdot 1 & =\left\lceil\int \kappa \cdot \delta M\right.
\end{aligned}
$$

There is no external virtual work done since the unit moments are applied internally. Since we have two virtual force systems in equilibrium and one real compatible system, we have two equations:

$$
0=\iint \frac{M}{E I} \cdot \delta M^{1} \quad \text { and } \quad 0=\int \mathfrak{f} \frac{M}{E I} \cdot \delta M^{2}
$$

For the first equation, expanding the expression for the real moment system, $M$ :

$$
\begin{aligned}
& \int \frac{\left(M^{0}+\alpha_{1} \cdot \delta M^{1}+\alpha_{2} \cdot \delta M^{2}\right)}{E I} \cdot \delta M^{1}=0 \\
& \int \frac{M^{0} \cdot \delta M^{1}}{E I}+\alpha_{1} \int \frac{\delta M^{1} \cdot \delta M^{1}}{E I}+\alpha_{2} \int \frac{\delta M^{2} \cdot \delta M^{1}}{E I}=0
\end{aligned}
$$

In which we've dropped the contour integral - it being understood that we sum for all members. Similarly for the second virtual moments, we have:

$$
\int \frac{M^{0} \cdot \delta M^{2}}{E I}+\alpha_{1} \int \frac{\delta M^{1} \cdot \delta M^{2}}{E I}+\alpha_{2} \int \frac{\delta M^{2} \cdot \delta M^{2}}{E I}=0
$$

Thus we have two equations and so we can solve for $\alpha_{1}$ and $\alpha_{2}$. Usually we write this as a matrix equation:

$$
\left\{\begin{array}{l}
\int \frac{M^{0} \cdot \delta M^{1}}{E I} \\
\int \frac{M^{0} \cdot \delta M^{2}}{E I}
\end{array}\right\}+\left[\begin{array}{ll}
\int \frac{\delta M^{1} \cdot \delta M^{1}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{1}}{E I} \\
\int \frac{\delta M^{1} \cdot \delta M^{2}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{2}}{E I}
\end{array}\right]\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0
\end{array}\right\}
$$

Each of the integral terms is easily found using the integral tables, and the equation solved.

Similarly the (simpler) equation for a 2 -span beam is:

$$
\int \frac{M^{0} \cdot \delta M^{1}}{E I}+\alpha_{1} \int \frac{\delta M^{1} \cdot \delta M^{1}}{E I}=0
$$

And the equation for a 4-span beam is:

$$
\left\{\begin{array}{l}
\int \frac{M^{0} \cdot \delta M^{1}}{E I} \\
\int \frac{M^{0} \cdot \delta M^{2}}{E I} \\
\int \frac{M^{0} \cdot \delta M^{3}}{E I}
\end{array}\right\}+\left[\begin{array}{ccc}
\int \frac{\delta M^{1} \cdot \delta M^{1}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{1}}{E I} & 0 \\
\int \frac{\delta M^{1} \cdot \delta M^{2}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{2}}{E I} & \int \frac{\delta M^{3} \cdot \delta M^{2}}{E I} \\
0 & \int \frac{\delta M^{2} \cdot \delta M^{3}}{E I} & \int \frac{\delta M^{3} \cdot \delta M^{3}}{E I}
\end{array}\right]\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2} \\
\alpha_{3}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0 \\
0
\end{array}\right\}
$$

The diagram shows why there are no terms involving $\delta M^{3} \cdot \delta M^{1}$, and why it is only adjacent spans that have non-zero integrals:


## Example 18

## Problem

Using virtual work, find the bending moment diagram for the following beam, and determine the deflection at midspan of span $B C$.


## Solution

Proceeding as described above, we introduce releases (hinges) at the support points, apply the unit virtual moments, and find the corresponding bending moment diagrams:





Next we need to evaluate each term in the matrix virtual work equation. We'll take the two 'hard' ones first:

$$
\begin{aligned}
\int \frac{M^{0} \cdot \delta M^{1}}{E I} & =\frac{1}{E I}\left[\frac{1}{6}(50)(-1)(4+2)\right]_{A B}+\frac{1}{2 E I}\left[\frac{1}{3}(45)(-1)(6)\right]_{B C} \\
& =\frac{1}{E I}(-50-45)=-\frac{95}{E I}
\end{aligned}
$$

Note that since $\delta M^{1}=0$ for span $C D$, there is no term for it above. Similarly, for the following evaluation, there will be no term for span $A B$. Note also the $2 E I$ term for member $B C$ - this could be easily overlooked.

$$
\begin{aligned}
\int \frac{M^{0} \cdot \delta M^{2}}{E I} & =\frac{1}{2 E I}\left[\frac{1}{3}(45)(-1)(6)\right]_{B C}+\frac{1}{E I}\left[\frac{1}{3}(93.75)(-1)(5)\right]_{C D} \\
& =\frac{1}{E I}(-45-156.25)=-\frac{201.25}{E I}
\end{aligned}
$$

The following integrals are more straightforward since they are all triangles:

- $\int \frac{\delta M^{1} \cdot \delta M^{1}}{E I}=\frac{1}{E I}\left[\frac{1}{3}(-1)(-1)(4)\right]_{A B}+\frac{1}{2 E I}\left[\frac{1}{3}(-1)(-1)(6)\right]_{B C}=\frac{2.333}{E I}$
- $\int \frac{\delta M^{2} \cdot \delta M^{1}}{E I}=\frac{1}{2 E I}\left[\frac{1}{6}(-1)(-1)(6)\right]_{B C}=\frac{0.5}{E I}$
- $\int \frac{\delta M^{1} \cdot \delta M^{2}}{E I}=\frac{0.5}{E I}$, since it is equal to $\int \frac{\delta M^{2} \cdot \delta M^{1}}{E I}$ by the commutative property of multiplication.

$$
\text { - } \int \frac{\delta M^{2} \cdot \delta M^{2}}{E I}=\frac{1}{2 E I}\left[\frac{1}{3}(-1)(-1)(6)\right]_{B C}+\frac{1}{E I}\left[\frac{1}{3}(-1)(-1)(5)\right]_{C D}=\frac{2.667}{E I}
$$

With all the terms evaluated, enter them into the matrix equation:

$$
\left\{\begin{array}{l}
\int \frac{M^{0} \cdot \delta M^{1}}{E I} \\
\int \frac{M^{0} \cdot \delta M^{2}}{E I}
\end{array}\right\}+\left[\begin{array}{ll}
\int \frac{\delta M^{1} \cdot \delta M^{1}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{1}}{E I} \\
\int \frac{\delta M^{1} \cdot \delta M^{2}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{2}}{E I}
\end{array}\right]\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0
\end{array}\right\}
$$

To give:

$$
\frac{1}{E I}\left\{\begin{array}{c}
-95 \\
-201.25
\end{array}\right\}+\frac{1}{E I}\left[\begin{array}{cc}
2.333 & 0.5 \\
0.5 & 2.667
\end{array}\right]\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0
\end{array}\right\}
$$

And solve, as follows:

$$
\begin{aligned}
{\left[\begin{array}{cc}
2.333 & 0.5 \\
0.5 & 2.667
\end{array}\right]\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2}
\end{array}\right\} } & =\left\{\begin{array}{c}
95 \\
201.25
\end{array}\right\} \\
\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2}
\end{array}\right\} & =\frac{1}{(2.333 \cdot 2.667-0.5 \cdot 0.5)}\left[\begin{array}{cc}
2.667 & -0.5 \\
-0.5 & 2.333
\end{array}\right]\left\{\begin{array}{c}
95 \\
201.25
\end{array}\right\} \\
& =\left\{\begin{array}{l}
25.57 \\
70.67
\end{array}\right\}
\end{aligned}
$$

Now using our superposition equation for moments, $M=M^{0}+\alpha_{1} \cdot \delta M^{1}+\alpha_{2} \cdot \delta M^{2}$, we can show that the multipliers are just the hogging support moments:

$$
\begin{aligned}
& M_{B}=0+25.57 \cdot 1+70.67 \cdot 0=25.57 \mathrm{kNm} \\
& M_{C}=0+25.57 \cdot 0+70.67 \cdot 1=70.67 \mathrm{kNm}
\end{aligned}
$$

From these we get the final BMD:


To determine the deflection we identify a statically determinate subset of the original structure as our equilibrium set. The compatibility set is the actual deflection we seek, along with the real curvatures caused by the real bending moment diagram:


Notice that only the bending moment diagrams for the span being examined are needed for the calculation. Next, we express the real bending moment diagram in terms of its constituent parts the primary and two reactant diagrams:


And so the volume integrals are determined by multiplying the following diagrams:


Thus we have:

$$
\begin{aligned}
\delta & =\int\left[\frac{M_{x}}{E I}\right] \cdot \delta M_{x} d x \\
& =\frac{1}{2 E I}\left[2\left(\frac{5}{12}\right)(1.5)(45)(3)+\alpha_{1}\left(\frac{1}{6}\right)(1.5)(-1)(6+3)+\alpha_{2}\left(\frac{1}{6}\right)(1.5)(-1)(6+3)\right] \\
& =\frac{84.4-1.125 \alpha_{1}-1.125 \alpha_{2}}{E I}=-\frac{23.9}{E I}
\end{aligned}
$$

The negative sign indicates movement opposite to the direction of the virtual unit load and so the deflection is upwards.

## Spreadsheet Solution

A simple spreadsheet for a 3-span beam with centre-span point load and UDL capabilities, showing Example 18, is:


## Example 19

## Problem

Analyze the following frame for the reactions and bending moment diagram. Draw the deflected shape of the frame.


## Solution

The frame is three degrees indeterminate. We must therefore choose three redundants. We will remove the fixed support at $D$ to obtain the primary structure and its bending moment diagram:


Next we apply a unit virtual force in lieu of each redundant, $V_{D}, H_{D}$, and $M_{D}$, and determine the corresponding virtual bending moment diagram:


The virtual work equation is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
0 \cdot 1 & =\llbracket \kappa \cdot \delta M
\end{aligned}
$$

And the real curvature is $\kappa=M / E I$ giving:

$$
0=\int \frac{M}{E I} \cdot \delta M
$$

There are three separate virtual bending moment diagrams, $\delta M$, and so we have three spate virtual work equations of the above form:

$$
\begin{aligned}
& 0=\iint \frac{M}{E I} \cdot \delta M^{1} \\
& 0=\iint \frac{M}{E I} \cdot \delta M^{2} \\
& 0=\int \frac{M}{E I} \cdot \delta M^{3}
\end{aligned}
$$

The final bending moment diagram is expressed in terms of the unit redundant bending moments and the primary bending moment diagram using superposition:

$$
M=M^{0}+\alpha_{1} \cdot \delta M^{1}+\alpha_{2} \cdot \delta M^{2}+\alpha_{3} \cdot \delta M^{3}
$$

Giving three equations of the form:

$$
0=\emptyset \int \frac{M^{0}}{E I} \delta M+\alpha_{1} \oint \frac{\delta M^{1}}{E I} \delta M+\alpha_{2} \oint \frac{\delta M^{2}}{E I} \delta M+\alpha_{3}\left\lceil\mathfrak{f M ^ { 3 }} \frac{E I}{E M}\right.
$$

Since there are three unknowns in this equation, and there are three such equations, the multipliers can be found. The three equations can be written in matrix form:

$$
\left\{\begin{array}{l}
\int \frac{M^{0} \cdot \delta M^{1}}{E I} \\
\int \frac{M^{0} \cdot \delta M^{2}}{E I} \\
\int \frac{M^{0} \cdot \delta M^{3}}{E I}
\end{array}\right\}+\left[\begin{array}{lll}
\int \frac{\delta M^{1} \cdot \delta M^{1}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{1}}{E I} & \int \frac{\delta M^{3} \cdot \delta M^{1}}{E I} \\
\int \frac{\delta M^{1} \cdot \delta M^{2}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{2}}{E I} & \int \frac{\delta M^{3} \cdot \delta M^{2}}{E I} \\
\int \frac{\delta M^{1} \cdot \delta M^{3}}{E I} & \int \frac{\delta M^{2} \cdot \delta M^{3}}{E I} & \int \frac{\delta M^{3} \cdot \delta M^{3}}{E I}
\end{array}\right]\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2} \\
\alpha_{3}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0 \\
0
\end{array}\right\}
$$

Virtual work problems can always be written like this. The general form of the matrix equation is:

$$
\left\{\mathbf{A}_{\mathbf{0}}\right\}+[\mathbf{\delta} \mathbf{A}]\{\boldsymbol{\alpha}\}=\{\mathbf{0}\}
$$

where $\left\{\mathbf{A}_{\mathbf{0}}\right\}$ is the free actions vector; $[\mathbf{\delta A}]$ is the virtual actions matrix; and $\{\boldsymbol{\alpha}\}$ is the vector of multipliers.

For this problem, we establish the terms of the free actions vector as follows:

$$
\begin{aligned}
\int \frac{M^{0} \cdot \delta M^{1}}{E I} & =[(-8)(80)(6)]_{A B}+\left[\frac{1}{2}(80)(-6)(2)+\frac{1}{3}(-2)(80)(2)\right]_{B C} \\
& =-3840-1760 / 3 \\
& =-13280 / 3 \\
\int \frac{M^{0} \cdot \delta M^{2}}{E I} & =\left[\frac{1}{2}(6)(80)(6)\right]_{A B}+\left[\frac{1}{2}(6)(80)(2)\right]_{B C} \\
& =1440+480 \\
& =1920 \\
\int \frac{M^{0} \cdot \delta M^{3}}{E I} & =[(80)(-1)(6)]_{A B}+\left[\frac{1}{2}(-1)(80)(2)\right]_{B C} \\
& =-480-80 \\
& =-560
\end{aligned}
$$

Note that the volume integrals are commutative and so the virtual actions matrix is symmetrical. Thus we only calculate the upper triangle terms:

$$
\begin{aligned}
& \int \frac{\delta M^{1} \cdot \delta M^{1}}{E I}=\frac{1}{E I}[(-8)(-8)(6)]_{A B}+\frac{1}{E I}\left[\frac{1}{3}(-8)(-8)(8)\right]_{B C} \\
& =\frac{1}{E I}(384+512 / 3) \\
& =1664 / 3 E I \\
& \int \frac{\delta M^{2} \cdot \delta M^{1}}{E I}=\frac{1}{E I}\left[\frac{1}{2}(6)(-8)(6)\right]_{A B}+\frac{1}{E I}\left[\frac{1}{2}(6)(-8)(8)\right]_{B C} \\
& =\frac{1}{E I}(-144-192) \\
& =-336 / E I \\
& \int \frac{\delta M^{3} \cdot \delta M^{1}}{E I}=\frac{1}{E I}[(-1)(-8)(6)]_{A B}+\frac{1}{E I}\left[\frac{1}{2}(-1)(-8)(8)\right]_{B C} \\
& =\frac{1}{E I}(48+32) \\
& =80 / E I \\
& \int \frac{\delta M^{2} \cdot \delta M^{2}}{E I}=\frac{1}{E I}\left[\frac{1}{3}(6)(6)(6)\right]_{A B}+\frac{1}{E I}[(6)(6)(8)]_{B C}+\frac{1}{E I}\left[\frac{1}{3}(6)(6)(6)\right]_{C D} \\
& =\frac{1}{E I}(72+288+72) \\
& =432 / E I \\
& \int \frac{\delta M^{3} \cdot \delta M^{2}}{E I}=\frac{1}{E I}\left[\frac{1}{2}(-1)(6)(6)\right]_{A B}+\frac{1}{E I}[(-1)(6)(8)]_{B C}+\frac{1}{E I}\left[\frac{1}{2}(-1)(6)(6)\right]_{C D} \\
& =\frac{1}{E I}(-18-48-18) \\
& =-84 / E I \\
& \int \frac{\delta M^{3} \cdot \delta M^{3}}{E I}=\frac{1}{E I}[(-1)(-1)(6)]_{A B}+\frac{1}{E I}[(-1)(-1)(8)]_{B C}+\frac{1}{E I}[(-1)(-1)(6)]_{C D} \\
& =\frac{1}{E I}(6+8+6) \\
& =20 / E I
\end{aligned}
$$

Substituting the values in, we obtain:

$$
\frac{1}{E I}\left\{\begin{array}{c}
-13280 / 3 \\
1920 \\
-560
\end{array}\right\}+\frac{1}{E I}\left[\begin{array}{ccc}
1664 / 3 & -336 & 80 \\
-336 & 432 & -84 \\
80 & -84 & 20
\end{array}\right]\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2} \\
\alpha_{3}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0 \\
0
\end{array}\right\}
$$

Giving:

$$
\left\{\begin{array}{l}
\alpha_{1} \\
\alpha_{2} \\
\alpha_{3}
\end{array}\right\}=-\left[\begin{array}{ccc}
1664 / 3 & -336 & 80 \\
-336 & 432 & -84 \\
80 & -84 & 20
\end{array}\right]^{-1}\left\{\begin{array}{c}
-13280 / 3 \\
1920 \\
-560
\end{array}\right\}=\left\{\begin{array}{c}
9.32 \\
5.45 \\
13.6
\end{array}\right\} \begin{gathered}
\mathrm{kN} \\
\mathrm{kN} \\
\mathrm{kNm}
\end{gathered}
$$

This solution is easily found using software tools. For example MINVERSE and MMULT in MS Excel, or using the following script in MATLAB, writing A0 for $\left\{\mathbf{A}_{\mathbf{0}}\right\}$ and dA for $[\boldsymbol{\delta A}]$ :

```
A0 = [-13280/3; 1920; -560];
\(\mathrm{dA}=\)\begin{tabular}{cll}
{\([1664 / 3\)} & -336 & \(80 ;\) \\
-336 & 432 & \(-84 ;\) \\
80 & -84 & \(20] ;\)
\end{tabular}
```

$a=-i n v(d A) * A 0$;

Thus the complete solution for this problem is:


### 8.5.6 Problems

1. For the truss shown, calculate the vertical deflection at $C$ when member $E B$ is not present and when it is present. Members $A B, D C$ and $D E$ have $E A=60 \times 10^{3} \mathrm{kN}$; members $E B, B C$ and $A D$ have $E A=100 \times 10^{3} \mathrm{kN}$, and; member $B D$ has $E A=80$ $\times 10^{3} \mathrm{kN}$. (Ans. 15 mm and 11.28 mm )

2. Calculate the vertical deflection at $C$ when member $E B$ is not present and when it is present. Members $A B, D C$ and $D E$ have $E A=60 \times 10^{3} \mathrm{kN}$; members $E B, B C$ and $A D$ have $E A=100 \times 10^{3} \mathrm{kN}$, and; member $B D$ has $E A=80 \times 10^{3} \mathrm{kN}$.

3. Autumn 2007 For the truss shown, determine the force in each member and determine the horizontal deflection of joint B . Take $E A=200 \times 10^{3} \mathrm{kN}$ for all members. (Ans. Choosing BD: $\alpha=50 \sqrt{2} ; \sqrt{2} \mathrm{~mm}$ )

4. For the truss shown, calculate the horizontal deflection at $D$. All members have $E A=100 \times 10^{3} \mathrm{kN}$.

5. Calculate the member forces for the truss. All members have $E A=100 \times 10^{3} \mathrm{kN}$. Determine the horizontal deflection at $B$ and the vertical displacement of joint $C$.

6. For the truss shown, determine the member forces and calculate the vertical displacement of joint $B$. All members have $E A=100 \times 10^{3} \mathrm{kN}$.

7. For the beam of Example 18, show that the deflections at the centre of the following spans, taking downwards as positive, are:

- $\operatorname{Span} A B: \frac{41.1}{E I}$;
- Span $C D: \frac{133.72}{E I}$.

8. For the beam shown, determine the reactions and bending moment diagram. Determine the deflection at midspan.

9. For the beam shown, determine the reactions and bending moment diagram. Determine the deflection at $D$.

10. A prestressed post-tensioned beam is shown at transfer; determine the reactions, bending moment diagram, and deflection at midspan.

11. A prestressed post-tensioned beam is shown in service; determine the reactions, bending moment diagram, and deflection at midspan.

12. Using virtual work, analyse the following prismatic beam to show that the support moments are $M_{B}=51.3 \mathrm{kNm}$ and $M_{C}=31.1 \mathrm{kNm}$. Draw the bending moment diagram. Show that the deflections at middle of the spans moving left to right are 19.6/EI, 204/EI, and 71.6/EI respectively.

13. For the following prismatic frame, determine the bending moment diagram and the vertical deflection at $C$.

14.Summer 2007, Part (b): For the frame shown, draw the bending moment diagram and determine the horizontal deflection of joint C . Neglect axial effects in the members and take $E I=36 \times 10^{3} \mathrm{kNm}^{2}$. (Ans. $H_{D}=75.8 \mathrm{kN} ; 19.0 \mathrm{~mm}$ )

15.For the following prismatic frame, determine the bending moment diagram and the vertical deflection at $C$.

16.For the following prismatic frame, determine the bending moment diagram and the vertical deflection at $C$.

17.For the following prismatic frame, determine the bending moment diagram and the vertical deflection at $C$.

14. For the following prismatic portal frame, determine the bending moment diagram and the horizontal deflection at B and $D$.

19.For the frame shown, calculate the bending moment diagram, and verify the following deflections: $\delta_{B x}=\frac{856.2}{E I} \rightarrow$ and $\delta_{C y}=\frac{327.6}{E I} \downarrow .\left(A n s . M_{B}=30.1 \mathrm{kNm}\right)$

20.The defections for the portal frame of Q18 are found to be too large and an alternative scheme using fixed supports is proposed. Determine the bending moment diagram and the new horizontal deflections at B and $D$. Compare the results and discuss the pros and cons of the revised scheme.


### 8.6 Virtual Work for Self-Stressed Structures

### 8.6.1 Background

## Introduction

Self-stressed structures are structures that have stresses induced, not only by external loading, but also by any of the following:

- Temperature change of some of the members (e.g. solar gain);
- Lack of fit of members from fabrication:
- Error in the length of the member;
- Ends not square and so a rotational lack of fit;
- Incorrect support location from imperfect construction;
- Non-rigid (i.e. spring) supports due to imperfect construction.

Since any form of fabrication or construction is never perfect, it is very important for us to know the effect (in terms of bending moment, shear forces etc.) that such errors, even when small, can have on the structure.

Here we introduce these sources, and examine their effect on the virtual work equation. Note that many of these sources of error can exist concurrently. In such cases we add together the effects from each source.

## Temperature Change - Axial

The source of self-stressing in this case is that the temperature change causes a member to elongate:

$$
\Delta L_{T}=\alpha L(\Delta T)
$$

where $\alpha$ is the coefficient of linear thermal expansion (change in length, per unit length, per degree Celsius), $L$ is the original member length and $\Delta T$ is the temperature change.


Since temperature changes change the length of a member, the internal virtual work is affected. Assuming a truss member is being analysed, we now have changes in length due to force and temperature, so the total change in length of the member is:

$$
e=\frac{P L}{E A}+\alpha L(\Delta T)
$$

Hence the internal virtual work for this member is:

$$
\begin{aligned}
\delta W_{I} & =e \cdot \delta P \\
& =\left(\frac{P L}{E A}+\alpha L(\Delta T)\right) \cdot \delta P
\end{aligned}
$$

## Temperature Change - Curvature

When two sides of a flexural member are exposed to different temperature changes, both an axial elongation and curvature result. A typical example is a bridge deck: the road surface is exposed to the sun and heats up whilst the underside (soffit) is not exposed to directly sunlight and so heats less.


A short element of the member, $d x$, which has depth $h$, expands as shown. This deformation can be broken up into axial ( $\Delta T_{e}$ ) and flexural ( $\Delta T_{\kappa}$ ) components:


Hence the curvature at the section is:

$$
\kappa_{x}=\alpha \frac{2 \Delta T_{\kappa}}{h} d x
$$

In terms of the change in temperature at the top and bottom of the beam, we have:

$$
\Delta T_{\kappa}=\frac{\Delta T_{t}-\Delta T_{b}}{2}
$$

And so:

$$
\kappa_{x}=\alpha \frac{\Delta T_{t}-\Delta T_{b}}{h} d x
$$

Since the virtual work done by a virtual moment moving through this real curvature at the cross section is:

$$
\delta W_{I}=\kappa \cdot \delta M
$$

These expressions are valid for a point along the beam. Hence, along the full length of the member, the virtual work done is:

$$
\delta W_{I}=\int_{0}^{L} \kappa_{x} \cdot \delta M_{x} d x=\int_{0}^{L} \frac{\alpha}{h}\left(\Delta T_{t}-\Delta T_{b}\right) \cdot \delta M_{x} d x
$$

If we assume that the member is prismatic and the temperature changes are the same long the length of the beam, we have:

$$
\delta W_{I}=\frac{\alpha}{h}\left(\Delta T_{t}-\Delta T_{b}\right) \int_{0}^{L} \delta M_{x} d x
$$

Which is just $\alpha\left(\Delta T_{t}-\Delta T_{b}\right) / h$ times the area under the virtual moment diagram.

## Linear Lack of Fit

For a linear lack of fit, the member needs to be artificially elongated or shortened to fit it into place, thus introducing additional stresses. This is denoted:

$$
\lambda_{L}
$$

Considering a truss member subject to external loading, the total change in length will be the deformation due to loading and the linear lack of fit:

$$
e=\frac{P L}{E A}+\lambda_{L}
$$

Hence the internal virtual work for this member is:

$$
\begin{aligned}
\delta W_{I} & =e \cdot \delta P \\
& =\left(\frac{P L}{E A}+\lambda_{L}\right) \cdot \delta P
\end{aligned}
$$

## Rotational Lack of Fit

A rotational lack of fit, which applies to frames only, occurs when the end of a member is not square. Thus the member needs to be artificially rotated to get it into place, as shown below. This is denoted as:


Considering a frame member which has a lack of fit, $\lambda_{\theta}$ and a virtual moment $\delta M$ at the same point, then the internal virtual work done at this point is:

$$
\delta W_{I}=\lambda_{\theta} \cdot \delta M
$$

This must be added to the other forms of internal virtual work. Not also that the signs must be carefully chosen so that the virtual moment closes the gap - we will see this more clearly in an example.

## Errors in Support Location

The support can be misplaced horizontally and/or vertically. It is denoted:

$$
\lambda_{S}
$$

A misplaced support affects the external movements of a structure, and so contributes to the external virtual work. Denoting the virtual reaction at the support, in the direction of the misplacement as $\delta R$, then we have:

$$
\delta W_{e}=\lambda_{s} \cdot \delta R
$$

## Spring Supports

For spring supports we will know the spring constant for the support, denoted:

$$
k_{S}
$$

Since movements of a support are external, spring support movements affect the external virtual work. The real displacement $\Delta_{S}$ that occurs is:

$$
\Delta_{S}=R / k_{S}
$$

In which $R$ is the real support reaction in the direction of the spring. Further, since $R$ will be known in terms of the multiplier and virtual reaction, $\delta R$, we have:

$$
R=R^{0}+\alpha \cdot \delta R
$$

Hence:

$$
\Delta_{S}=\left(R^{0}+\alpha \cdot \delta R\right) / k_{S}
$$

And so the external work done is:

$$
\begin{aligned}
\delta W_{e} & =\Delta_{S} \cdot \delta R \\
& =\frac{R^{0}}{k_{S}} \cdot \delta R+\alpha \cdot \frac{(\delta R)^{2}}{k_{S}}
\end{aligned}
$$

The only unknown here is $\alpha$ which is solved for from the virtual work equation.

### 8.6.2 Trusses

## Example 20

## Problem

Here we take the truss of Example 7 and examine the effects of:

- Member $E D$ was found to be 5 mm too long upon arrival at site;
- Member $A B$ is subject to a temperature increase of $+100{ }^{\circ} \mathrm{C}$.

For this truss, $E=200 \mathrm{kN} / \mathrm{mm}^{2 ;}$ member areas, $A=1000 \mathrm{~mm}^{2}$ for all members except $A E$ and $B D$ where $A=1000 \sqrt{ } 2 \mathrm{~mm}^{2}$, and; coefficient of thermal expansion is $\alpha=10 \times 10^{-6}{ }^{\circ} \mathrm{C}^{-1}$.


## Solution

The change in length due to the temperature change is:

$$
\begin{aligned}
\Delta L_{T} & =\alpha L(\Delta T) \\
& =\left(10 \times 10^{-6}\right)(2000)(+100) \\
& =2 \mathrm{~mm}
\end{aligned}
$$

## Vertical Displacement of Joint D

The virtual work equation is now:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i} \\
y_{D V} \cdot 1 & =\sum\left(\frac{P^{0} L}{E A}\right)_{i} \cdot \delta P_{i}^{1}+5 \cdot \delta P_{E D}^{1}+2 \cdot \delta P_{A B}^{1}
\end{aligned}
$$

In Example 7 we established various values in this equation:

- $\sum\left(\frac{P^{0} L}{E A}\right)_{i} \cdot \delta P_{i}^{1}=16.5$
- $\delta P_{E D}^{1}=-1$
- $\delta P_{A B}^{1}=+1$

Hence we have:

$$
\begin{aligned}
y_{D V} \cdot 1 & =16.5+5(-1)+2(+1) \\
y_{D V} & =13.5 \mathrm{~mm}
\end{aligned}
$$

## Horizontal Displacement of Joint D:

Similarly, taking the values obtained in Example 7, we have:

$$
\begin{aligned}
y_{D H} \cdot 1 & =\sum\left(\frac{P^{0} L}{E A}\right)_{i} \cdot \delta P_{i}^{2}+5 \cdot \delta P_{E D}^{2}+2 \cdot \delta P_{A B}^{2} \\
& =-4.5+5(+1)+2(0) \\
y_{D H} & =+0.5 \mathrm{~mm} \text { to the right }
\end{aligned}
$$

## Example 21

## Problem

Here we use the truss of Example 15 and examine, separately, the effects of:

- Member $A C$ was found to be 3.6 mm too long upon arrival on site;
- Member $B C$ is subject to a temperature reduction of $-50^{\circ} \mathrm{C}$;
- Support $D$ is surveyed and found to sit 5 mm too far to the right.

Take $E A$ to be $10 \times 10^{4} \mathrm{kN}$ for all members and the coefficient of expansion to be $\alpha=10 \times 10^{-6}{ }^{\circ} \mathrm{C}^{-1}$.


## Solution

Error in Length:
To find the new multiplier, we include this effect in the virtual work equation:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\sum y_{i} \cdot \delta F_{i} & =\sum e_{i} \cdot \delta P_{i} \\
0 & =\sum\left(\frac{P L}{E A}\right)_{i} \cdot \delta P_{i}^{1}+\lambda_{L} \cdot \delta P_{A C}^{1}
\end{aligned}
$$

But $P=P^{0}+\alpha \cdot \delta P^{1}$, hence:

$$
\begin{aligned}
0 & =\sum\left(\frac{\left(P^{0}+\alpha \cdot \delta P^{1}\right) L}{E A}\right)_{i} \cdot \delta P_{i}^{1}+\lambda_{L} \cdot \delta P_{A C}^{1} \\
& =\sum\left(\frac{P^{0} L}{E A}\right)_{i} \cdot \delta P_{i}^{1}+\alpha \cdot \sum\left(\frac{\delta P^{1} L}{E A}\right)_{i} \cdot \delta P_{i}^{1}+\lambda_{L} \cdot \delta P_{A C}^{1} \\
0 & =\sum \frac{P^{0} \cdot \delta P_{i}^{1} \cdot L_{i}}{E A_{i}}+\alpha \cdot \sum \frac{\left(\delta P_{i}^{1}\right)^{2} L_{i}}{E A_{i}}+\lambda_{L} \cdot \delta P_{A C}^{1}
\end{aligned}
$$

As can be seen, the $\lambda_{L} \cdot \delta P_{A C}^{1}$ term is simply added to the usual VW equation. From before, we have the various values of the summations, and so have:

$$
\begin{gathered}
0=\frac{-37.8 \times 10^{4}}{E A}+\alpha \cdot \frac{1.62 \times 10^{4}}{E A}+(3.6)(+1) \\
\alpha=\frac{37.8 \times 10^{4}-3.6 E A}{1.62 \times 10^{4}}=+1.11
\end{gathered}
$$

Thus member $A C$ is 1.1 kN in tension. Note the change: without the error in fit it was 23.3 kN in tension and so the error in length has reduced the tension by 22.2 kN .

## Temperature Change:

The same derivation from the VW equation gives us:

$$
0=\sum \frac{P^{0} \cdot \delta P_{i}^{1} \cdot L_{i}}{E A_{i}}+\alpha \cdot \sum \frac{\left(\delta P_{i}^{1}\right)^{2} L_{i}}{E A_{i}}+\Delta L_{T} \cdot \delta P_{B C}^{1}
$$

The change in length due to the temperature change is:

$$
\begin{aligned}
\Delta L_{T} & =\alpha L(\Delta T) \\
& =\left(10 \times 10^{-6}\right)(3000)(-50)=-1.5 \mathrm{~mm}
\end{aligned}
$$

Noting that the virtual force in member $B C$ is $\delta P_{B C}^{1}=-0.6$, we have:

$$
0=\frac{-37.8 \times 10^{4}}{E A}+\alpha \cdot \frac{1.62 \times 10^{4}}{E A}+(-1.5)(-0.6)
$$

Giving:

$$
\alpha=\frac{37.8 \times 10^{4}-0.9 E A}{1.62 \times 10^{4}}=+17.7
$$

Thus member $A C$ is 17.7 kN in tension, a decrease of 5.6 kN from 23.3 kN .

## Error in Support Location:

Modifying the VW equation gives:

$$
\lambda_{S} \cdot \delta H_{D}=\sum \frac{P^{0} \cdot \delta P_{i}^{1} \cdot L_{i}}{E A_{i}}+\alpha \cdot \sum \frac{\left(\delta P_{i}^{1}\right)^{2} L_{i}}{E A_{i}}
$$

The value of the virtual horizontal reaction is found from the virtual force system to be 0.6 kN to the right. Hence:

$$
+(5)(0.6)=\frac{-37.8 \times 10^{4}}{E A}+\alpha \cdot \frac{1.62 \times 10^{4}}{E A}
$$

Note the sign on the support displacement: since the real movement is along the same direction as the virtual force, it does positive virtual work. Solving:

$$
\alpha=\frac{-3 E A+37.8 \times 10^{4}}{1.62 \times 10^{4}}=+4.8
$$

Thus member $A C$ is 4.8 kN in tension; a reduction of 18.5 kN .

## All Effects Together:

In this case, the virtual work equation is:

$$
\lambda_{S} \cdot \delta H_{D}=\sum \frac{P^{0} \cdot \delta P_{i}^{1} \cdot L_{i}}{E A_{i}}+\alpha \cdot \sum \frac{\left(\delta P_{i}^{1}\right)^{2} L_{i}}{E A_{i}}+\lambda_{L} \cdot \delta P_{A C}^{1}+\Delta L_{T} \cdot \delta P_{B C}^{1}
$$

Substituting the various values in gives:

$$
+(5)(0.6)=\frac{-37.8 \times 10^{4}}{E A}+\alpha \cdot \frac{1.62 \times 10^{4}}{E A}+(3.6)(+1)+(-1.5)(-0.6)
$$

And solving:

$$
\alpha=\frac{(-3.6-0.9-3) E A+37.8 \times 10^{4}}{1.62 \times 10^{4}}=-22.9
$$

And so member $A C$ is 22.9 kN in compression. Thus should be the same as the original force of 23.3 kN plus all the changes induced by the errors:

$$
23.3-22.2-5.6-18.5=-23.0 \mathrm{kN}
$$

Exercise: Find the remaining forces in the truss and compare to the forces without the presence of self-stresses and calculate the horizontal deflection at $C$ without selfstresses.

### 8.6.3 Beams

## Example 22

## Problem

Determine the deflection at the centre of a simply-supported beam subject to a linear temperature increase over the section depth:


## Solution

We will assume that $T_{2}>T_{1}$ for developing the solution, and consider positive displacements to be upwards consequently. We identify the compatibility set as the real internal and external displacements:


Next, we identify the equilibrium set, which is the external virtual force and internal virtual bending moments:



SM

The curvature is constant, since the temperature changes are constant along the length of the beam. Thus the virtual work equation is:

$$
\begin{aligned}
1 \times \Delta & =\int \kappa_{x} \delta M d x \\
& =\frac{\alpha}{h}\left(T_{2}-T_{1}\right) \int_{0}^{L} \delta M_{x} d x
\end{aligned}
$$

The area under the virtual moment diagram is $0.5(L)(L / 4)=L^{2} / 8$. Hence we have the solution as:

$$
\Delta=\frac{\alpha L^{2}}{8 h}\left(T_{2}-T_{1}\right)
$$

## Numerical Example

If we take the following values:

- $T_{2}=+50^{\circ} \mathrm{C}$;
- $T_{1}=+20^{\circ} \mathrm{C}$;
- $L=10 \mathrm{~m}$;
- $\alpha=10 \times 10^{-6}{ }^{\circ} \mathrm{C}^{-1}$;
- $h=500 \mathrm{~mm}$

Then we have:

$$
\Delta=\frac{10 \times 10^{-6}\left(10 \times 10^{3}\right)^{2}}{8(500)}(50-20)=7.5 \mathrm{~mm}
$$

## Example 23

## Problem

Determine a general solution for the two-span beam subject to linear temperature change over its section depth, uniform along the length of the beam.


## Solution

As usual, we will assume that $T_{2}>T_{1}$ for developing the solution. Since this is a onedegree indeterminate structure, we choose one redundant: the reaction at $B$. Hence, the compatibility set is the real displacements, the internal curvatures along the beam, and the external displacement at $B$ :


As an equilibrium set, we choose the statically determinate subset given by an external unit load at $B$ and corresponding internal virtual bending moments:


Next we write the general virtual work equation, noting the external unit load moves through no virtual displacement:

$$
1 \times 0=\int \kappa_{x} \delta M d x
$$

And establish that the real curvatures are made up of both temperature-induced curvatures and load (or bending moment)-induced curvatures:

$$
\kappa_{x}=\kappa_{T}+\kappa_{M}
$$

But we know the temperature curvature to be constant, and causes tension on top so we have:

$$
\kappa_{T}=-\frac{\alpha}{h}\left(T_{2}-T_{1}\right)
$$

And the moment (or load)-induced curvature is found as usual:

$$
\kappa_{M}=\frac{M}{E I}=\frac{M_{0}+R \cdot \delta M}{E I}
$$

where we have used superposition to break the total bending moment at a point into the primary and reactant bending moments. Notice also that we have altered the notation for the virtual work multiplier (usually $\alpha$ ) to $R$, since we are now using that symbol for the coefficient of linear thermal expansion. Thus we have:

$$
\kappa_{x}=-\frac{\alpha}{h}\left(T_{2}-T_{1}\right)+\frac{M_{0}+R \cdot \delta M}{E I}
$$

And using this in the virtual work formula gives:

$$
0=\int\left[-\frac{\alpha}{h}\left(T_{2}-T_{1}\right)+\frac{M_{0}+R \cdot \delta M}{E I}\right] \delta M d x
$$

And so:

$$
0=-\frac{\alpha}{h}\left(T_{2}-T_{1}\right) \int \delta M d x+\int \frac{M_{0} \cdot \delta M}{E I} d x+R \int \frac{(\delta M)^{2}}{E I} d x
$$

Giving the redundant, $R$ as:

$$
R=\frac{\frac{\alpha}{h}\left(T_{2}-T_{1}\right) \int \delta M d x-\int \frac{M_{0} \cdot \delta M}{E I} d x}{\int \frac{(\delta M)^{2}}{E I} d x}
$$

This is a general expression for any structure, of which the usual solution is a special case when there is no temperature change. In the special case where there is no external applied load ( $M_{0}=0$ ), only the temperature change on the two-span beam, the area of the virtual moment diagram is $0.5(2 L)(L / 2)=L^{2} / 2$ and so we have:

$$
R=\frac{\frac{\alpha}{h}\left(T_{2}-T_{1}\right) \frac{L^{2}}{2}}{\frac{1}{E I}\left(2 \cdot \frac{1}{3} \cdot \frac{L}{2} \cdot \frac{L}{2} L\right)}=\frac{3 E I}{L} \cdot \frac{\alpha}{h}\left(T_{2}-T_{1}\right)
$$

In which we have used the volume integral for the triangular virtual bending moment diagram.

## Numerical Example

If we take the following values:

- $T_{2}=+50^{\circ} \mathrm{C}$;
- $T_{1}=+20^{\circ} \mathrm{C}$;
- $L=10 \mathrm{~m}$;
- $\alpha=10 \times 10^{-6}{ }^{\circ} \mathrm{C}^{-1}$;
- $h=500 \mathrm{~mm}$;
- $E I=40 \times 10^{3} \mathrm{kNm}^{2}$

Then we have:

$$
R=\frac{3\left(40 \times 10^{3}\right)}{10} \cdot \frac{10 \times 10^{-6}}{0.5}(50-20)=7.2 \mathrm{kN} \downarrow
$$

And there is a bending moment across the beam as follows:


An interesting part of the temperature-induced deformation is that the deflected shape diagram shows two points of contraflexure which are not evident from the bending moment diagram. This is because the curvature (and hence deflection) induced by the temperature increase does not have an associated bending moment. Only the restraint of the middle support causes bending moments (and associated curvatures). The final deflected shape occurs from the superposition of both curvatures:

$$
\begin{gathered}
\kappa_{x}=\kappa_{T}+\kappa_{M} \\
\kappa_{M}=\frac{M}{E I}=\frac{36}{40 \times 10^{3}}=0.9 \times 10^{-3} \mathrm{~m}^{-1} \text { at } B \\
\kappa_{T}=\frac{\alpha}{h}\left(T_{2}-T_{1}\right)=\frac{10 \times 10^{-6}}{0.5}(50-20)=0.6 \times 10^{-3} \mathrm{~m}^{-1} \text { constant }
\end{gathered}
$$

The curvature diagrams are therefore as follows:


## Example 24

## Problem

Consider the following 2 -span beam with central spring support. Determine an expression for the central support reaction in terms of its spring stiffness.


## Solution

To do this, we will consider the central support as the redundant:


The virtual work equation, accounting for the spring, is:

$$
1 \cdot \Delta_{s}=\left\{\int \frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x+\alpha \cdot \mathfrak{f} \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x\right.
$$

The deflection at the central support will be:

$$
\Delta_{S}=\frac{V_{B}}{k}=\frac{\alpha}{k}
$$

Since the reaction at $B$ is $1 \cdot \alpha$. Noting that the deflection of the spring will be opposite to the unit load, and using the volume integrals, we have:

$$
\begin{aligned}
-\frac{\alpha}{k} & =\frac{2}{E I}\left[\frac{5}{12}\left(-\frac{l}{2}\right)\left(\frac{w l^{2}}{2}\right)(l)\right]+\alpha \cdot \frac{2}{E I}\left[\frac{1}{3}\left(-\frac{l}{2}\right)\left(-\frac{l}{2}\right)(l)\right] \\
-\frac{\alpha}{k} & =\frac{-5 w l^{4}}{24 E I}+\alpha \cdot \frac{l^{3}}{6 E I} \\
\alpha\left(-\frac{E I}{k}-\frac{l^{3}}{6}\right) & =\frac{-5 w l^{4}}{24}
\end{aligned}
$$

And finally:

$$
\begin{aligned}
\alpha & =\frac{5 w l^{4}}{\frac{24 E I}{k}+4 l^{3}} \\
& =\frac{5 w l}{24 \frac{E I}{k l^{3}}+4}
\end{aligned}
$$

It is important we understand the implications of this result:

- For no support present, $k=0$ and so $24 E I / k \rightarrow \infty$ meaning $\alpha \rightarrow 0$ and there is no support reaction (as we might expect);
- For $k=\infty$ we have the perfectly rigid (usual) roller support and so $24 E I / k \rightarrow 0$ giving us $\alpha \rightarrow \frac{5}{4} w l$ - a result we established previously;
- For intermediate relative stiffnesses, the value of the reaction is between these extremes. Therefore, we as the designer have the ability to choose the reaction most favourable to us.

A plot of the reaction and relative stiffnesses is:


Exercise: Verify that for $k=0$ the deflection at the centre of the beam reduces to the the familiar expression for a simply-supported beam under UDL.

## Example 25

## Problem

A simply supported beam has a rotational lack of fit at midspan as shown. Determine the deflection at midspan.


## Solution

We apply a unit virtual load at midspan. The virtual work equation is:

$$
1 \cdot \Delta_{C}=\int \frac{M \cdot \delta M}{E I} d x+\lambda_{\theta} \cdot \delta M_{C}
$$

The first term returns the deflection due to the loading alone. The second term accounts for the lack of fit. Note that the work done is positive when closing a lack of fit. In this case, the applied virtual moment opens the lack of fit and so this term is negative overall. From the usual real and virtual bending moment diagrams, we have:

$$
\begin{aligned}
\Delta_{C} & =\frac{1}{E I}\left[2\left(\frac{5}{12}\right)\left(\frac{w L^{2}}{8}\right)\left(\frac{L}{4}\right)\left(\frac{L}{2}\right)\right]-\lambda_{\theta} \cdot\left(\frac{L}{4}\right) \\
& =\frac{5 w L^{4}}{384 E I}-\frac{\lambda_{\theta} L}{4}
\end{aligned}
$$

That the lack of fit improves the deflection can be understood by considering the movement required to close the lack of fit prior to application of the load as follows:


It is obvious from this diagram that the upward deflection from lack of fit only is given by $S=R \theta$, giving $\Delta_{C}=\left(\lambda_{\theta} / 2\right)(L / 2)=\lambda_{\theta} L / 4$.

## Numerical Example

Prior to gang-nailing a domestic joist together the builder notices the ends aren't square and sends the following sketch. Can he proceed?


The rotational lack of fit is (approximately):

$$
\lambda_{\theta}=\frac{5}{200}=0.025 \mathrm{rads}
$$

The upward displacement is thus:

$$
\Delta=\frac{(0.025)\left(4 \times 10^{3}\right)}{4}=25 \mathrm{~mm}
$$

Taking a deflection limit of span/200 for a domestic case, the limit is:

$$
\Delta_{\text {allow }}=\frac{4 \times 10^{3}}{200}=20 \mathrm{~mm}
$$

On this analysis the builder should not proceed. However, the clever engineer can make this work, allowing the builder to proceed - what can be taken into account to allow this?

### 8.6.4 Frames

## Example 26

## Problem

The following frame, in addition to its loading, is subject to:

- Support $A$ is located 10 mm too far to the left;
- End $C$ of member $B C$ is $1.2 \times 10^{-3}$ rads out of square, as shown;
- Member $C D$ is 12 mm too short.

Determine the bending moment diagram. Take $E I=36 \times 10^{3} \mathrm{kNm}^{2}$ for each member.


## Solution

We will choose the horizontal reaction at $A$ as the redundant. Since we are dealing with a linear lack of fit in member $C D$, we need to allow for the virtual work done by the axial forces in this member and so we solve for the axial force diagrams also.

For the primary structure:
$\sum F_{x}=0: \quad H_{D}^{0}-90=0 \quad \therefore H_{D}^{0}=90 \mathrm{kN}$
$\sum M$ about $A=0$ :
$24 \cdot \frac{6^{2}}{2}+90 \cdot 4+90 \cdot 2-V_{D}^{0} \cdot 6=0 \quad \therefore V_{D}^{0}=162 \mathrm{kN}$
$\sum F_{y}=0: \quad-V_{A}^{0}-24 \cdot 6+162=0$
$\therefore V_{A}^{0}=18 \mathrm{kN}$

Also, $M_{C}=90 \cdot 6=540 \mathrm{kNm}$ giving:


For the redundant structure:

$\sum M$ about $A=0:-1 \cdot 2+V_{D}^{1} \cdot 6=0 \quad \therefore V_{D}^{1}=\frac{1}{3} \mathrm{kN} \downarrow$
$\sum F_{y}=0: \quad-V_{A}^{1}+\frac{1}{3}=0 \quad \therefore V_{A}^{1}=\frac{1}{3} \mathrm{kN} \uparrow$

The virtual work equation, accounting for the relevant effects is:

$$
\begin{aligned}
\delta W & =0 \\
\delta W_{E} & =\delta W_{I} \\
\lambda_{S} \cdot H_{A}^{1} & =\emptyset\left[\frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x+\alpha \cdot \emptyset \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x+\lambda_{\theta} \cdot \delta M_{C}^{1}+\lambda_{L} \cdot \delta P_{C D}^{1}\right.
\end{aligned}
$$

We take each term in turn:
(a) $\lambda_{S} \cdot H_{A}^{1}$ : The applied unit load is in the same direction as the error in the support location, and keeping all units in metres:

$$
\lambda_{S} \cdot H_{A}^{1}=10 \times 10^{-3} \cdot+1=10 \times 10^{-3}(\mathrm{kN} \cdot \mathrm{~m})
$$

(b) $\oint \frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x$ : Using the volume integrals:

$$
\begin{aligned}
\int \frac{M^{0} \cdot \delta M_{i}^{1}}{E I_{i}} d x & =\frac{1}{E I}\left\{\left[\frac{1}{12}(4+3 \cdot 6)(-540) 6\right]_{B C}+\left[\frac{1}{3}(6)(-540) 6\right]_{C D}\right\} \\
& =-\frac{12420}{E I}
\end{aligned}
$$

(c) $\mathfrak{f} \frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x$ : Again using the integrals table:

$$
\begin{aligned}
& f\left(\frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x=\frac{1}{E I}\left[\frac{1}{3}(4)(4) 4\right]_{A B}+\frac{1}{E I}\left[\frac{1}{6}(4(2 \cdot 4+6)+6(4+2 \cdot 6)) 6\right]_{B C}\right. \\
&+\frac{1}{E I}\left[\frac{1}{3}(6)(6) 6\right]_{C D} \\
& \int\left[\frac{\left(\delta M_{i}^{1}\right)^{2}}{E I_{i}} d x\right.=\frac{245.33}{E I}
\end{aligned}
$$

(d) $\lambda_{\theta} \cdot \delta M_{C}^{1}$ : The virtual bending moment at $C$ is 6 kNm with tension on the inside of the frame. This is in the opposite direction to that needed to close the lack of fit and so the sign of this term is negative, as shown:

(e) $\lambda_{L} \cdot \delta P_{C D}^{1}:$ Since the member is too short, the lack of fit is negative, whilst the virtual force is in tension and so positive:

$$
\lambda_{L} \cdot \delta P_{C D}^{1}=\left(-12 \times 10^{-3}\right) \cdot\left(+\frac{1}{3}\right)=-4 \times 10^{-3}(\mathrm{kN} \cdot \mathrm{~m})
$$

Substituting these values into the equation, along with $E I=36 \times 10^{3} \mathrm{kNm}^{2}$, gives:

$$
10 \times 10^{-3}=-\frac{12420}{36 \times 10^{3}}+\alpha \cdot \frac{245.33}{36 \times 10^{3}}-7.2 \times 10^{-3}-4 \times 10^{-3}
$$

And so we can solve for $\alpha$ :

$$
\begin{aligned}
10 & =-345+\alpha \cdot 6.815-7.2-4 \\
\alpha & =\frac{10+345+7.2+4}{6.815} \\
& =53.73
\end{aligned}
$$

And so the horizontal reaction at $A$ is:

$$
H_{A}=H_{A}^{0}+\alpha \cdot H_{A}^{1}=0+53.73 \cdot 1=53.73 \mathrm{kN}
$$

and similarly for the other reactions. Also, using the superposition equation for moments, $M=M^{0}+\alpha \cdot M^{1}$, we have:


### 8.6.5 Problems

1. For the truss of Example 21, show that the force in member $A C$ is 22.2 kN in compression when there is no external loads present, and only the lack of fit of 3.6 mm of member $A C$.
2. For the truss of Example 21, determine the horizontal deflection of joint $C$ due to each of the errors separately, and then combined.
3. For the beam of Example 24, show that the stiffness of the spring support that optimizes the bending moments is $89.5 E I / l^{3}$, i.e. makes the sagging and hogging moments equal in magnitude.
4. For the following beam, show that the vertical deflection at $C$ can be given by:

$$
\delta_{C y}=P\left(\frac{2 L^{3}}{3 E I}+\frac{1}{k}\right)
$$


5. For the following beam, determine the bending moment diagram. Take $\alpha=10 \times 10^{-6}{ }^{\circ} \mathrm{C}^{-1}, h=500 \mathrm{~mm}$, and $E I=40 \times 10^{3} \mathrm{kNm}^{2}$.

6. For the following prismatic beam, determine the bending moment diagram, reactions and midspan deflections for the lack of fit scenario shown.

7. Solve the following prismatic beam, for the bending moment diagram. Show that the bending moment and deflection at $C$ approaches the limiting cases for a pinned-pinned beam and a fixed-fixed beam as $k$ approaches 0 and infinity respectively. You may use the symmetry of the problem in your solution.

8. For the following prismatic beam, investigate the behaviour for different ratios of $E I$ to the support stiffness, $k$. Is the reaction at $C$ always downwards?

9. For the frame of Example 26, show that the horizontal deflection of joint C due to the applied loads only is 77.5 mm . Find the deflection of joint C due to both the loads and errors given.
10.For the following frame, the support at D was found to yield horizontally by 0.08 $\mathrm{mm} / \mathrm{kN}$. Also, the end of member BC is not square, as shown. Draw the bending moment diagram and determine the horizontal deflection at D . Take EI $=100 \times 103$ kNm 2 . (Ans. 5.43 mm to the right)


### 8.7 Past Exam Questions

## Summer 1997

3. (a) Calculate the horizontal deflection of joint $C$ in the pin-jointed truss shown in Fig. Q3(a) when a load of 100 kN is applied as shown at $B$.
(b) If the truss were constructed with an additional member $B D$ as shown in Fig. Q3(b), calculate the horizontal deflection of joint $C$ when the load of 100 kn is applied, as before, at $B$.
(c) What "initial lack of fit" in member $B D$ will give zero horizontal deflection at $C$ when the load of 100 kN is applied, as before, at $B$ in the truss in Fig. Q2(b).

Assume $E=200 \mathrm{kN} / \mathrm{mm}^{2}$.
Assume the cross-sectional of $A B, B C$ and $C D$ is $100 \mathrm{~mm}^{2}$, and $A C$ and $B D$ is $141 \mathrm{~mm}^{2}$. ( 25 marks)


FIG Q3(a)


FIG Q3(b)

Answers:
(a) $\delta_{C x}=30.0 \mathrm{~mm} \rightarrow$;
(b) $\delta_{C x}=12.9 \mathrm{~mm} \rightarrow$;
(c) $\lambda_{B D}=21.1 \mathrm{~mm}$ too long.

## Summer 1998

3. (a) Calculate the horizontal deflection of joint $C$ in the pin-jointed truss loaded as shown in Fig. Q3 (a).

(b) . If the truss were propped with an additional member DF as shown in Fig. Q3 (b), calculate the new horizontal deflection of joint C when the truss is loaded, as before.

Assume $\mathrm{E}=200 \mathrm{kN} / \mathrm{mm}^{2}$.
Assume the cross-sectional area of $\mathrm{AB}, \mathrm{BD}, \mathrm{CD}$ and DE is $100 \mathrm{~mm}^{2}$, and $\mathrm{BC}, \mathrm{BE}$ and DF is $141 \mathrm{~mm}^{2}$.
( 25 marks)


Answers:
(a) $\delta_{C x}=45.0 \mathrm{~mm} \rightarrow$;
(b) $\delta_{C x}=26.2 \mathrm{~mm} \rightarrow$.

## Summer 1999



Answers:
(a) $V_{E}=42.5 \mathrm{kN} \uparrow$;
(b) $\delta_{C y}=28.6 \mathrm{~mm} \downarrow$.

## Summer 2000

3. Use the Flexibility method to determine, for the frame shown in Fig. Q3,
(a) the bending moment in the frame and
(b) the vertical deflection at C .

Sketch the bending moment diagram giving the value of the bending moment at all salient points.
Sketch the deflected shape of the frame.
Take constant $\mathrm{EI}=12 \times 10^{4} \mathrm{kNm}{ }^{2}$.
(25 marks)


FIG. Q3 $九 \ldots-4 \mathrm{~m}$

Answers:
(a) $V_{A}=2.78 \mathrm{kN} \downarrow$;
(b) $\delta_{C}=0.47 \mathrm{~mm} \downarrow$.

## Summer 2001

3. (a) On assembly of the pin-jointed truss shown in Fig. Q3 it was found that member $A B$ was 4 mm too long and support E was 10 mm too high.
The truss was then loaded as shown at $B$. Find the forces in the members.
(17 marks)
(b) If in addition to all the conditions in (a) above, (i.e. AB 4 mm too long, support E 100 mm too high, and the loading as shown), it is found that support A yields to the right for each 1 kN of reaction at support A , determine the forces in the members.
(8 marks)
Assume EA $=120 \times 10^{3} \mathrm{kN}$ for all members.


## FIG. Q3

## Answers:

(a) $P_{B D}=+44.4 \mathrm{kN}$;
(b) $P_{B D}=+300 \mathrm{kN}$.

## Summer 2002

3. (a) Using Virtual Work determine by what length member AC, in the truss loaded as shown in Fig. Q3 (a), must be adjusted to ensure that the vertical deflection at node C is zero. Assume $E A=200 \times 10^{3} \mathrm{kN}$ for all members.

(b) (i) Using Virtual Work, and neglecting axial deformation,determine the bending moments in the frame loaded as shown in Fig. Q3(b).
(ii) Sketch the bending moment diagram for the frame.
(iii) Sketch the deflected shape of the frame.


## Answers:

(a) $\lambda_{L}=1.35 \mathrm{~mm}$ shorter;
(b) $V_{C}=160 \mathrm{kN} \uparrow$.

## Summer 2004

3. (a) Calculate the vertical deflection of joint $A$ in the pin-jointed truss shown in Fig. Q3 when a load of 100 kN is applied at A as shown.
(20 marks)
(b) What initial lack of fit in member $A C$ will give zero vertical deflection at $A$ when the load of 100 kN is applied, as before, at $A$ in the truss in Fig. Q3.

Assume $E=200 \mathrm{kN} / \mathrm{mm}^{2}$.
Assume the cross-sectional area of $A B$ and $A D$ is $500 \mathrm{~mm}^{2}$ each, and of $A C$ is $400 \mathrm{~mm}^{2}$.


Answers:
(a) $\delta_{A y}=2.19 \mathrm{~mm} \downarrow$;
(b) $\lambda_{L}=5 \mathrm{~mm}$ shorter.

## Summer 2007

Using Virtual Work:
(i) For the frame in Fig. Q3(a), determine the horizontal deflection of joint $C$;
(ii) Draw the bending moment diagram for the frame in Fig. Q3(b);
(iii) Determine the horizontal deflection of joint $C$ for the frame in Fig. Q3(b).

## Note:

You may neglect axial effects in the members.
Take $E I=36 \times 10^{3} \mathrm{kNm}^{2}$ for all members.


FIG. Q3(a)
FIG. Q3(b)

Answers:
(a) $\delta_{C x}=84.6 \mathrm{~mm} \rightarrow$;
(b) $H_{D}=75.8 \mathrm{kN} \leftarrow$;
(c) $\delta_{C x}=19.02 \mathrm{~mm} \rightarrow$.

## Semester 2, 2008

For the frame shown in Fig. Q2, using Virtual Work:
(a) Determine the following:
reactions;
bending moment diagram.
(b) After construction, it is found that the support at $C$ is not rigid but a spring support of stiffness $k=2000 / 3 \mathrm{kN} / \mathrm{m}$. Taking this into account, determine the following:
vertical displacement of joint $C$;
bending moment diagram.

## Note:

You may neglect axial effects in the members.
Take $E I=36 \times 10^{3} \mathrm{kNm}^{2}$ for all members.


FIG. Q2

Answers:
(a) $V_{C}=35.63 \mathrm{kN} \uparrow$;
(b) $\delta_{C y}=45 \mathrm{~mm} \downarrow$.

## Semester 2, 2009

## QUESTION 2

Using Virtual Work, for the frame shown in Fig. Q2, do the following:
(i) Determine the reactions;
(ii) Draw the bending moment diagram;
(iii) Draw the deflected shape of the structure;
(iv) Determine the horizontal deflection of joint $D$.
(25 marks)

## Note:

- Take $E=200 \mathrm{kN} / \mathrm{mm}^{2}$ and $I=45 \times 10^{7} \mathrm{~mm}^{4}$ for all members.


FIG. Q2

Answers:
(a) $V_{B}=225 \mathrm{kN} \uparrow$;
(b) $\delta_{D x}=81.25 \mathrm{~mm} \leftarrow$.

## Semester 2, 2009

## QUESTION 3

For the structure shown in Fig. Q2, due to an error on site, it is required to make the moment at support $A$ zero. This is to be done by the introduction of a spring support at $B$, in place of the roller support, as shown in Fig. Q3. Using Virtual Work:
(i) Determine the spring stiffness required so that the moment at A is zero.
(ii) Drawn the revised bending moment diagram;
(iii) Draw the revised deflected shape of the structure;
(iv) Determine the vertical deflection of the structure at $B$.

## Note:

- Take $E=200 \mathrm{kN} / \mathrm{mm}^{2}$ and $I=45 \times 10^{7} \mathrm{~mm}^{4}$ for all members.
- You may use any relevant results from your workings for Q2, but in doing so acknowledge their source.


FIG. Q3

Answers:
(a) $k=20 \times 10^{3} \mathrm{kN} / \mathrm{m}$;
(b) $\delta_{B y}=7.5 \mathrm{~mm} \downarrow$.

## Semester 2, 2010

## QUESTION 2

Using Virtual Work, for the continuous beam shown in Fig. Q2, do the following:
(i) Draw the bending moment diagram, noting all important values;
(ii) Draw the deflected shape of the structure;
(iii) Determine the deflection at $E$.

## Note:

- Take $E=200 \mathrm{kN} / \mathrm{mm}^{2}$ and $I=20 \times 10^{7} \mathrm{~mm}^{4}$ for all members;
- The following expression may be of assistance:


FIG. Q2

## Answers:

(a) $M_{B}=174.5 \mathrm{kNm} ; M_{C}=97.6 \mathrm{kNm}$;
(b) $\delta_{E y}=16.9 \mathrm{~mm} \downarrow$.

## Semester 2, 2011

## QUESTION 3

Using Virtual Work, for the truss shown in Fig. Q3, do the following:
(i) Determine the member forces;
(ii) Draw the member forces using conventional arrow notation;
(v) Draw the deflected shape of the structure;
(iii) Determine the vertical deflection of joint $E$.

## Note:

- Take $E=200 \mathrm{kN} / \mathrm{mm}^{2}$;
- $A=1500 \mathrm{~mm}^{2}$ for members $B C, B F, C E$, and $E F$;
- $A=1500 \sqrt{2} \mathrm{~mm}^{2}$ for members $B E$ and $C F$;
- $A=2000 \mathrm{~mm}^{2}$ for members $A F$ and $D E$;
- $A=2500 \mathrm{~mm}^{2}$ for members $A B$ and $C D$.


FIG. Q3

Answers:
(a) $F_{B E}=15 \sqrt{2} \mathrm{kN}$ compression;
(b) $\delta_{E y}=4.25 \mathrm{~mm} \downarrow$.

### 8.8 References

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### 8.9 Appendix - Volume Integrals

|  |  | j |  | $\square_{1}{ }^{j}$ |
| :---: | :---: | :---: | :---: | :---: |
|  | $\frac{1}{3} j k l$ | $\frac{1}{6} j k l$ | $\frac{1}{6}\left(j_{1}+2 j_{2}\right) k l$ | $\frac{1}{2} j k l$ |
|  | $\frac{1}{6} j k l$ | $\frac{1}{3} j k l$ | $\frac{1}{6}\left(2 j_{1}+j_{2}\right) k l$ | $\frac{1}{2} j k l$ |
|  | $\frac{1}{6} j\left(k_{1}+2 k_{2}\right) l$ | $\frac{1}{6} j\left(2 k_{1}+k_{2}\right) l$ | $\begin{aligned} & \frac{1}{6}\left[j_{1}\left(2 k_{1}+k_{2}\right)+\right. \\ & \left.j_{2}\left(k_{1}+2 k_{2}\right)\right] l \end{aligned}$ | $\frac{1}{2} j\left(k_{1}+k_{2}\right) l$ |
| $k$ <br> I | $\frac{1}{2} j k l$ | $\frac{1}{2} j k l$ | $\frac{1}{2}\left(j_{1}+j_{2}\right) k l$ | $j k l$ |
|  | $\frac{1}{6} j k(l+a)$ | $\frac{1}{6} j k(l+b)$ | $\begin{aligned} & \frac{1}{6}\left[j_{1}(l+b)+\right. \\ & \left.\quad j_{2}(l+a)\right] k \end{aligned}$ | $\frac{1}{2} j k l$ |
|  | $\frac{5}{12} j k l$ | $\frac{1}{4} j k l$ | $\frac{1}{12}\left(3 j_{1}+5 j_{2}\right) k l$ | $\frac{2}{3} j k l$ |
|  | $\frac{1}{4} j k l$ | $\frac{5}{12} j k l$ | $\frac{1}{12}\left(5 j_{1}+3 j_{2}\right) k l$ | $\frac{2}{3} j k l$ |
|  | $\frac{1}{4} j k l$ | $\frac{1}{12} j k l$ | $\frac{1}{12}\left(j_{1}+3 j_{2}\right) k l$ | $\frac{1}{3} j k l$ |
| $k$ | $\frac{1}{12} j k l$ | $\frac{1}{4} j k l$ | $\frac{1}{12}\left(3 j_{1}+j_{2}\right) k l$ | $\frac{1}{3} j k l$ |
|  | $\frac{1}{3} j k l$ | $\frac{1}{3} j k l$ | $\frac{1}{3}\left(j_{1}+j_{2}\right) k l$ | $\frac{2}{3} j k l$ |

