

ECE5463: Introduction to Robotics

Lecture Note 12: Dynamics of Open Chains: Lagrangian Formulation

Prof. Wei Zhang

Department of Electrical and Computer Engineering
Ohio State University
Columbus, Ohio, USA

Spring 2018

Outline

- Introduction
- Euler-Lagrange Equations
- Lagrangian Formulation of Open-Chain Dynamics

From Single Rigid Body to Open Chains

- Recall Newton-Euler Equation for a single rigid body:

$$\mathcal{F}_b = \mathcal{G}_b \dot{\mathcal{V}}_b - [\text{ad}_{\mathcal{V}_b}]^T (\mathcal{G}_b \mathcal{V}_b)$$

- Open chains consist of multiple rigid links connected through joints
- Dynamics of adjacent links are coupled.
- We are concerned with modeling multi-body dynamics subject to constraints.

Preview of Open-Chain Dynamics

- Equations of Motion are a set of 2nd-order differential equations:

$$\tau = M(\theta)\ddot{\theta} + h(\theta, \dot{\theta})$$

- $\theta \in \mathbb{R}^n$: vector of joint variables; $\tau \in \mathbb{R}^n$: vector of joint forces/torques
- $M(\theta) \in \mathbb{R}^{n \times n}$: mass matrix
- $h(\theta, \dot{\theta}) \in \mathbb{R}^n$: forces that lump together centripetal, Coriolis, gravity, and friction terms that depend on θ and $\dot{\theta}$
- **Forward dynamics:** Determine acceleration $\ddot{\theta}$ given the state $(\theta, \dot{\theta})$ and the joint forces/torques:

$$\ddot{\theta} = M^{-1}(\theta)(\tau - h(\theta, \dot{\theta}))$$

- **Inverse dynamics:** Finding torques/forces given state $(\theta, \dot{\theta})$ and desired acceleration $\ddot{\theta}$

$$\tau = M(\theta)\ddot{\theta} + h(\theta, \dot{\theta})$$

Lagrangian vs. Newton-Euler Methods

- There are typically two ways to derive the equation of motion for an open-chain robot: Lagrangian method and Newton-Euler method

Lagrangian Formulation

- Energy-based method
- Dynamic equations in closed form
- Often used for study of dynamic properties and analysis of control methods

Newton-Euler Formulation

- Balance of forces/torques
- Dynamic equations in numeric/recursive form
- Often used for numerical solution of forward/inverse dynamics

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Generalized Coordinates and Forces

- Consider k particles. Let f_i be the force acting on the i th particle, \bar{m}_i be its mass, p_i be its position. Newton's law: $f_i = \bar{m}_i \ddot{p}_i$, $i = 1, \dots, k$
- Now consider the case in which some particles are rigidly connected, imposing constraints on their positions

$$\alpha_j(p_1, \dots, p_k) = 0, \quad j = 1, \dots, n_c$$

- k particles in \mathbb{R}^3 under n_c constraints $\Rightarrow 3k - n_c$ degree of freedom
- Dynamics of this constrained k -particle system can be represented by $n \triangleq 3k - n_c$ independent variables q_i 's, called the **generalized coordinates**

$$\begin{cases} \alpha_j(p_1, \dots, p_k) = 0 \\ j = 1, \dots, n_c \end{cases} \Leftrightarrow \begin{cases} p_i = \gamma_i(q_1, \dots, q_n) \\ i = 1, \dots, k \end{cases}$$

Generalized Coordinates and Forces

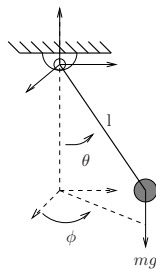
- To describe equation of motion in terms of generalized coordinates, we also need to express external forces applied to the system in terms components along generalized coordinates. These “forces” are called **generalized forces**.
- Generalized force f_i and coordinate rate \dot{q}_i are dual to each other in the sense that $f^T \dot{q}$ corresponds to power
- The equation of motion of the k -particle system can thus be described in terms of $3k - n_c$ independent variables instead of the $3k$ position variables subject to n_c constraints.
- This idea of handling constraints can be extended to interconnected rigid bodies (open chains).

Euler-Lagrange Equation

- Now let $q \in \mathbb{R}^n$ be the generalized coordinates and $f \in \mathbb{R}^n$ be the generalized forces of some constrained dynamical system.
- **Lagrangian function:** $\mathcal{L}(q, \dot{q}) = \mathcal{K}(q, \dot{q}) - \mathcal{P}(q)$
 - $\mathcal{K}(q, \dot{q})$: kinetic energy of system
 - $\mathcal{P}(q)$: potential energy
- **Euler-Lagrange Equations:**

$$f = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{q}} - \frac{\partial \mathcal{L}}{\partial q} \quad (1)$$

Example: Spherical Pendulum



Example: Spherical Pendulum (Continued)

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Lagrangian Formulation of Open Chains

- For open chains with n joints, it is convenient and always possible to choose the joint angles $\theta = (\theta_1, \dots, \theta_n)$ and the joint torques $\tau = (\tau_1, \dots, \tau_n)$ as the generalized coordinates and generalized forces, respectively.
 - If joint i is revolute: θ_i joint angle and τ_i is joint torque
 - If joint i is prismatic: θ_i joint position and τ_i is joint force

- Lagrangian function: $\mathcal{L}(\theta, \dot{\theta}) = \mathcal{K}(\theta, \dot{\theta}) - \mathcal{P}(\theta, \dot{\theta})$

- Dynamic Equations:

$$\tau_i = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}_i} - \frac{\partial \mathcal{L}}{\partial \theta_i}$$

- To obtain the Lagrangian dynamics, we need to derive the kinetic and potential energies of the robot in terms of joint angles θ and torques τ .

Some Notations

For each link $i = 1, \dots, n$, Frame $\{i\}$ is attached to the center of mass of link i . All the following quantities are expressed in frame $\{i\}$

- \mathcal{V}_i : Twist of link $\{i\}$
- \bar{m}_i : mass; \mathcal{I}_i : rotational inertia matrix;
- $G_i = \begin{bmatrix} \mathcal{I}_i & 0 \\ 0 & \bar{m}_i I \end{bmatrix}$: Spatial inertia matrix
- Kinetic energy of link i : $\mathcal{K}_i = \frac{1}{2} \mathcal{V}_i^T \mathcal{G}_i \mathcal{V}_i$
- $J_b^{(i)} \in \mathbb{R}^{6 \times i}$: body Jacobian of link i

$$J_b^{(i)} = \begin{bmatrix} J_{b,1}^{(i)} & \dots & J_{b,i}^{(i)} \end{bmatrix}$$

where $J_{b,j}^{(i)} = \left[\text{Ad}_{e^{-[\mathcal{B}_i]\theta_i} \dots e^{-[\mathcal{B}_{j+1}]\theta_{j+1}}} \right] \mathcal{B}_j$, $j < i$ and $J_{b,i}^{(i)} = \mathcal{B}_i$

Kinetic and Potential Energies of Open Chains

- $J_{ib} = [J_b^{(i)} \quad 0] \in \mathbb{R}^{6 \times n}$

- Total Kinetic Energy:

$$\mathcal{K}(\theta, \dot{\theta}) = \frac{1}{2} \sum_{i=1}^n \mathbf{v}_i^T \mathcal{G}_i \mathbf{v}_i = \frac{1}{2} \dot{\theta}^T \left(\sum_{i=1}^n (J_{ib}^T(\theta) \mathcal{G}_i J_{ib}(\theta)) \right) \dot{\theta} \triangleq \frac{1}{2} \dot{\theta}^T M(\theta) \dot{\theta}$$

- Potential Energy:

$$\mathcal{P}(\theta) = \sum_{i=1}^n \bar{m}_i g h_i(\theta)$$

- $h_i(\theta)$: height of CoM of link i

Lagrangian Dynamic Equations of Open Chains

- Lagrangian: $\mathcal{L}(\theta, \dot{\theta}) = \mathcal{K}(\theta, \dot{\theta}) - \mathcal{P}(\theta)$

- $\tau_i = \frac{d}{dt} \frac{\partial \mathcal{L}}{\partial \dot{\theta}_i} - \frac{\partial \mathcal{L}}{\partial \theta_i} \Rightarrow$

$$\tau_i = \sum_{j=1}^n M_{ij}(\theta) \ddot{\theta}_j + \sum_{j=1}^n \sum_{k=1}^n \Gamma_{ijk}(\theta) \dot{\theta}_j \dot{\theta}_k + \frac{\partial \mathcal{P}}{\partial \theta_i}, \quad i = 1, \dots, n$$

- $\Gamma_{ijk}(\theta)$ is called the **Christoffel symbols of the first kind**

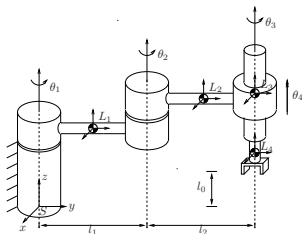
$$\Gamma_{ijk}(\theta) = \frac{1}{2} \left(\frac{\partial M_{ij}}{\partial \theta_k} + \frac{\partial M_{ik}}{\partial \theta_j} - \frac{\partial M_{jk}}{\partial \theta_i} \right)$$

Lagrangian Dynamic Equations of Open Chains

- Dynamic equation in vector form:

$$\tau = M(\theta)\ddot{\theta} + C(\theta, \dot{\theta})\dot{\theta} + g(\theta)$$

- $C_{ij}(\theta, \dot{\theta}) \triangleq \sum_{k=1}^n \Gamma_{ijk} \dot{\theta}_k$ is called the **Coriolis matrix**



$$\begin{aligned} M_{11} &= I_{y2}s_2^2 + I_{y3}s_{23}^2 + I_{z1} + I_{z2}c_2^2 + I_{z3}c_{23}^2 \\ &\quad + m_2r_1^2c_2^2 + m_3(l_1c_2 + r_2c_{23})^2 \\ M_{12} &= 0 \\ M_{13} &= 0 \\ M_{21} &= 0 \\ M_{22} &= I_{x2} + I_{x3} + m_3l_1^2 + m_2r_1^2 + m_3r_2^2 + 2m_3l_1r_2c_3 \\ M_{23} &= I_{x3} + m_3r_2^2 + m_3l_1r_2c_3 \\ M_{31} &= 0 \\ M_{32} &= I_{x3} + m_3r_2^2 + m_3l_1r_2c_3 \\ M_{33} &= I_{x3} + m_3r_2^2. \end{aligned}$$

$$\begin{aligned} \Gamma_{112} &= (I_{y2} - I_{z2} - m_2r_1^2)c_2s_2 + (I_{y3} - I_{z3})c_{23}s_{23} \\ &\quad - m_3(l_1c_2 + r_2c_{23})(l_1s_2 + r_2s_{23}) \\ \Gamma_{113} &= (I_{y3} - I_{z3})c_{23}s_{23} - m_3r_2s_{23}(l_1c_2 + r_2c_{23}) \\ \Gamma_{121} &= (I_{y2} - I_{z2} - m_2r_1^2)c_2s_2 + (I_{y3} - I_{z3})c_{23}s_{23} \\ &\quad - m_3(l_1c_2 + r_2c_{23})(l_1s_2 + r_2s_{23}) \end{aligned}$$

$$\begin{aligned} \Gamma_{131} &= (I_{y3} - I_{z3})c_{23}s_{23} - m_3r_2s_{23}(l_1c_2 + r_2c_{23}) \\ \Gamma_{211} &= (I_{z2} - I_{y2} + m_2r_1^2)c_2s_2 + (I_{z3} - I_{y3})c_{23}s_{23} \\ &\quad + m_3(l_1c_2 + r_2c_{23})(l_1s_2 + r_2s_{23}) \\ \Gamma_{223} &= -l_1m_3r_2s_3 \\ \Gamma_{232} &= -l_1m_3r_2s_3 \\ \Gamma_{233} &= -l_1m_3r_2s_3 \end{aligned}$$

$$\begin{aligned} \Gamma_{311} &= (I_{z3} - I_{y3})c_{23}s_{23} + m_3r_2s_{23}(l_1c_2 + r_2c_{23}) \\ \Gamma_{322} &= l_1m_3r_2s_3 \end{aligned}$$

$$\begin{bmatrix} 0 \\ -(m_2gr_1 + m_3gl_1)\cos\theta_2 - m_3r_2\cos(\theta_2 + \theta_3) \\ -m_3gr_2\cos(\theta_2 + \theta_3) \end{bmatrix}$$

More Discussions

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