# [Engineering Mathematics] 

## [Partial Differential Equations]

## Partial Differential Equations

## Chapter 1

### 1.1 Introduction

A differential equation which involves partial derivatives is called partial differential equation (PDE). The order of a PDE is the order of highest partial derivative in the equation and the degree of PDE is the degree of highest order partial derivative occurring in the equation. Thus order and degree of the PDE $\left(\frac{\partial^{2} z}{\partial x \partial y}\right)^{3}+\left(\frac{\partial z}{\partial x}\right)^{4}-x \frac{\partial z}{\partial x}=0$ are respectively 2 and 3 .

If ' $z$ ' is a function of two independent variables ' $x$ ' and ' $y$ ', let us use the following notations for the partial derivatives of ' $z$ ':
$\frac{\partial z}{\partial x} \equiv p, \quad \frac{\partial z}{\partial y} \equiv q, \quad \frac{\partial^{2} z}{\partial x^{2}} \equiv r, \quad \frac{\partial^{2} z}{\partial x \partial y} \equiv s, \quad \frac{\partial^{2} z}{\partial y^{2}} \equiv t$

### 1.2 Linear Partial Differential Equations of $1^{\text {st }}$ Order

If in a $1^{\text {st }}$ order PDE, both ' $p$ ' and ' $q$ ' occur in $1^{\text {st }}$ degree only and are not multiplied together, then it is called a linear PDE of $1^{\text {st }}$ order, i.e. an equation of the form $P p+Q q=R ; P, Q, R$ are functions of $x, y, z$, is a linear PDE of $1^{\text {st }}$ order.

Langrange's Method to Solve a Linear PDE of $1^{\text {st }}$ Order (Working Rule) :

1. Form the auxiliary equations $\frac{d x}{P}=\frac{d y}{Q}=\frac{d z}{R}$
2. Solve the auxiliary equations by the method of grouping or the method of multipliers* or both to get two independent solutions: $u=a, v=b$; where $a$ and $b$ are arbitrary constants.
3. $\varphi(u, v)=0$ or $u=f(v)$ is the general solution of the equation $P p+Q q=R$.

| $*$ Method of multipliers : Consider a |
| :---: |
| fraction $\frac{1}{2}=\frac{2}{4}=\frac{3}{6}$ |
| Taking $1,2,3$ as multipliers, each |
| fraction $=\frac{1 \times 1}{2 \times 1}=\frac{2 \times 2}{4 \times 2}=\frac{3 \times 3}{6 \times 3}$ |
| 1 |

Example 1. Solve the $\operatorname{PDE}\left(z^{2}-2 y z-y^{2}\right) p+(x y+z x) q=x y-z x$
Solution: Comparing with general form $P \equiv\left(z^{2}-2 y z-y^{2}\right), Q \equiv(x y+z x)$, $R \equiv(x y-z x)$

Step 1.
Auxiliary equations are $\frac{d x}{\left(z^{2}-2 y z-y^{2}\right)}=\frac{d y}{(x y+z x)}=\frac{d z}{(x y-z x)}$
Step 2.
Taking $x, y, z$ as multipliers, each fraction $=\frac{x d x+y d y+z d z}{\left(x z^{2}-2 x y z-x y^{2}+x y^{2}+x y z+x y z-x z^{2}\right)}$

$$
=\frac{x d x+y d y+z d z}{0}
$$

$\Rightarrow x d x+y d y+z d z=0$
Integrating, we get
$\frac{x^{2}}{2}+\frac{y^{2}}{2}+\frac{z^{2}}{2}=a_{1}$
$\Rightarrow u=x^{2}+y^{2}+z^{2}=a$
This is $1^{\text {st }}$ independent solution.
Now for $2^{\text {nd }}$ independent solution, taking last two members of auxiliary equations :
$\frac{d y}{x(y+z)}=\frac{d z}{x(y-z)}$
$\Rightarrow(y-z) d y=(y+z) d z$
$\Rightarrow y d y-(z d y+y d z)-z d z=0$
$\Rightarrow y d y-d(y z)-z d z=0$
Integrating, we get
$\frac{y^{2}}{2}-y z-\frac{z^{2}}{2}=b_{1}$
$\Rightarrow \quad v=y^{2}-2 y z-z^{2}=b$
Which is $2^{\text {nd }}$ independent solution

From (1) and (2), general solution is :
$\varphi\left(x^{2}+y^{2}+z^{2}, y^{2}-2 y z-z^{2}\right)=0$

### 1.3 Homogenous Linear Equations with Constant Coefficients

An equation of the form
$k_{0} \frac{\partial^{n} z}{\partial x^{n}}+k_{1} \frac{\partial^{n} z}{\partial x^{n-1} \partial y}+-\cdots--+k_{n} \frac{\partial^{n} z}{\partial y^{n}}=F(x, y)$
where $k^{\prime} s$ are constant is called a homogeneous linear PDE of $\mathrm{n}^{\text {th }}$ order with constant coefficients. It is homogeneous because all the terms contain derivatives of the same order.
Putting $\frac{\partial}{\partial x} \equiv D$ and $\frac{\partial}{\partial y} \equiv D^{r}$, (3) may be written as:
$\left(k_{0} D^{n}+k_{1} D^{n-1} D^{\prime}+----+k_{n} D^{\prime n}\right) z=F(x, y)$
or $f\left(D, D^{\prime}\right) \mathrm{z}=F(x, y)$

### 1.3.1 Solving Homogenous Linear Equations with Constant Coefficients

Case 1: When $F(x, y)=0$
i.e. equation is of the form $k_{0} \frac{\partial^{2} z}{\partial x^{2}}+k_{1} \frac{\partial^{2} z}{\partial x \partial y}+k_{2} \frac{\partial^{2} z}{\partial y^{2}}=0$
or $k_{0} D^{2}+k_{1} D D^{t}+k_{2} D^{t_{2}}=0$
In this case $Z=C . F$.
Case 2: When $F(x, y) \neq 0$
i.e. equation is of the form $k_{0} \frac{\partial^{2} z}{\partial x^{2}}+k_{1} \frac{\partial^{2} z}{\partial x \partial y}+k_{2} \frac{\partial^{2} z}{\partial y^{2}}=F(x, y)$
or $k_{0} D^{2}+k_{1} D D^{t}+k_{2} D^{t_{2}}=F(x, y)$
In this case $Z=$ C.F. + P.I.
Where C.F. denotes complimentary function and P.I. denotes Particular Integral.

## Rules for finding C.F. (Complimentary Function)

Step 1: Put $D=m$ and $D^{\prime}=1$ in (4) or (5) as the case may be
Then A.E. (Auxiliary Equation)is : $k_{0} m^{2}+k_{1} m+k_{2}=0$
Step 2: Solve the A.E. ( Auxiliary Equation):
i. If the roots of A.E. are real and different say $m_{1}$ and $m_{2}$, then

$$
\text { C.F. }=f_{1}\left(y+m_{1} x\right)+f_{2}\left(y+m_{2} x\right)
$$

ii. If the roots of A.E. are equal say $m$, then

$$
\text { C.F. }=f_{1}(y+m x)+x f_{2}(y+m x)
$$

Example 1.2 Solve $\frac{\partial^{2} z}{\partial x^{2}}-4 \frac{\partial^{2} z}{\partial x \partial y}-5 \frac{\partial^{2} z}{\partial y^{2}}=0$

Solution: $\Rightarrow\left(D^{2}-4 D D^{t}-5 D^{t}\right) z=0$

Auxiliary equation is: $m^{2}-4 m-5=0$
$\Rightarrow(m-5)(m+1)=0$
$\Rightarrow m=5,-1$
C.F. $=f_{1}(y+5 x)+f_{2}(y-x)$
$\Rightarrow \mathrm{Z}=f_{1}(y+5 x)+f_{2}(y-x)$
Example 1.3 Solve $\frac{\partial^{2} z}{\partial x^{2}}+2 \frac{\partial^{2} z}{\partial x \partial y}+\frac{\partial^{2} z}{\partial y^{2}}=0$

Solution: $\Rightarrow \quad\left(D^{2}+2 D D^{\prime}+D^{\prime 2}\right) z=0$

Auxiliary equation is: $m^{2}+2 m+1=0$
$\Rightarrow(m+1)^{2}=0$
$\Rightarrow m=-1,-1$
C.F. $=f_{1}(y-x)+x f_{2}(y-x)$
$\Rightarrow \mathrm{Z}=f_{1}(y-x)+x f_{2}(y-x)$
$\underline{\text { Rules for finding P.I. (Particular Integral) }}$

$$
\text { * Applicable only if } F(x, y) \neq 0
$$

Let the given PDE be $f\left(D, D^{\prime}\right) \mathrm{z}=F(x, y)$

$$
\text { P.I }=\frac{F(x, y)}{f\left(D, D^{\prime}\right)}
$$

Case I : When $F(x, y)=e^{a x+b y}$

$$
\text { Put } D=a \text { and } D^{\prime}=b
$$

If $f(a, b) \neq 0$, P.I $=\frac{e^{a x+b y}}{f(a, b)}$
and if $f(a, b)=0$, P.I $=\frac{x e^{a x+b y}}{\frac{d}{d D} f\left(D, D^{\prime}\right)}$

Now put $D=a$ and $D^{\prime}=b$
Example 1.4 Solve $\frac{\partial^{2} z}{\partial x^{2}}-5 \frac{\partial^{2} z}{\partial x \partial y}+6 \frac{\partial^{2} z}{\partial y^{2}}=e^{x+y}$

Solution: $\Rightarrow\left(D^{2}-5 D D^{\prime}+6 D^{\prime 2}\right) z=e^{x+y}$
Auxiliary equation is: $m^{2}-5 m+6=0$
$\Rightarrow(m-2)(m-3)=0$
$\Rightarrow m=2,3$
C.F. $=f_{1}(y+2 x)+f_{2}(y+3 x)$
P. $\mathrm{I}=\frac{e^{x+y}}{D^{2}-5 D D^{\prime}+6 D^{\prime 2}}$

Put $D=1, D^{\prime}=1$
P.I $=\frac{e^{x+y}}{1-5+6}=\frac{e^{x+y}}{2}$
$\Rightarrow \mathrm{Z}=f_{1}(y+2 x)+f_{2}(y+3 x)+\frac{e^{x+y}}{2}$

Example 1.5 Solve $r-4 s+4 t=e^{2 x+y}$

Solution: $\Rightarrow \frac{\partial^{2} z}{\partial x^{2}}-4 \frac{\partial^{2} z}{\partial x \partial y}+4 \frac{\partial^{2} z}{\partial y^{2}}=e^{2 x+y}$

$$
\Rightarrow\left(D^{2}-4 D D^{\prime}+4 D^{\prime 2}\right) z=e^{2 x+y}
$$

Auxiliary equation is: $m^{2}-4 m+4=0$
$\Rightarrow(m-2)^{2}=0$
$\Rightarrow m=2,2$
C.F. $=f_{1}(y+2 x)+x f_{2}(y+2 x)$
P.I $=\frac{e^{2 x+y}}{\left(D-2 D^{\prime}\right)^{2}}$

Putting $D=2, D^{\prime}=1$, denominator $=0$
P.I $=\frac{x e^{2 x+y}}{\frac{d}{d D}\left(D-2 D^{\prime}\right)^{2}}=\frac{x e^{2 x+y}}{2\left(D-2 D^{\prime}\right)}$

Putting $D=2, D^{I}=1$, again denominator $=0$
P.I. $=\frac{x^{2} e^{2 x+y}}{\frac{d}{d D} 2\left(D-2 D^{\prime}\right)}$
$\Rightarrow$ P.I. $=\frac{x^{2} e^{2 x+y}}{2}$

Complete solution is $\mathrm{Z}=\mathrm{C} . \mathrm{F} .+$ P.I.
$\Rightarrow \mathrm{Z}=f_{1}(y+2 x)+x f_{2}(y+2 x)+\frac{x^{2} e^{2 x+y}}{2}$

Case II : When $F(x, y)=e^{a x+b y} \phi(x, y), \quad \phi(x, y)$ is a trigonometric function of sine or cosine.

$$
\mathrm{P} . \mathrm{I}=e^{a x+b y} \frac{1}{f\left(D+a, D^{\prime}+b\right)} \phi(x, y)
$$

Example 1.6 Solve $\left(D^{3}+D^{2} D^{\prime}-D D^{\prime 2}-D^{\prime^{3}}\right) z=e^{x} \sin 2 y$

Solution: Auxiliary equation is: $m^{3}+m^{2}-m-1=0$
$\Rightarrow m=1,-1,-1$
C.F. $=f_{1}(y+x)+f_{2}(y-x)+x f_{3}(y-x)$

$$
\begin{aligned}
& \text { P.I }=\frac{e^{x} \sin 2 y}{D^{3}+D^{2} D^{\prime}-D D^{\prime 2}-D^{\prime 3}} \\
& =\frac{e^{x} \sin 2 y}{D^{2}\left(D+D^{\prime}\right)-D^{\prime 2}\left(D+D^{\prime}\right)}=\frac{e^{x} \sin 2 y}{\left(D+D^{\prime}\right)^{2}\left(D-D^{\prime}\right)}
\end{aligned}
$$

$$
=e^{x} \frac{1}{f\left(D+1, D^{r}\right)} \sin 2 y \quad(\because a=1, b=0)
$$

$$
=e^{x} \frac{1}{\left(D+1+D^{\prime}\right)^{2}\left(D+1-D^{\prime}\right)} \sin 2 y
$$

$$
=e^{x} \frac{1}{\left(D^{2}+1+D^{\prime 2}+2 D+2 D^{\prime}+2 D D^{\prime}\right)\left(D+1-D^{\prime}\right)} \sin 2 y
$$

Put $\quad D^{2}=0, D D^{\prime}=0, D^{\prime 2}=-4$

$$
\begin{aligned}
& \text { P.I. }=e^{x} \frac{1}{\left(0+1-4+2 D+2 D^{\prime}+0\right)\left(D+1-D^{\prime}\right)} \sin 2 y \\
& =e^{x} \frac{1}{\left(-3+2 D+2 D^{\prime}\right)\left(D+1-D^{\prime}\right)} \sin 2 y \\
& =e^{x} \frac{1}{\left(-3 D+2 D^{2}+2 D D^{\prime}-3+2 D+2 D^{\prime}+3 D^{\prime}-2 D D^{\prime}-2 D^{\prime 2}\right)} \sin 2 y \\
& =e^{x} \frac{1}{\left(-D+0+0-3+5 D^{\prime}-0+8\right)} \sin 2 y \\
& =e^{x} \frac{1}{\left(5 D^{\prime}-D\right)+5} \sin 2 y \\
& =e^{x} \frac{\left(5 D^{\prime}-D\right)-5}{\left(5 D^{\prime}-D\right)^{2}-25} \sin 2 y \\
& =e^{x} \frac{\left(5 D^{\prime}-D-5\right) \sin 2 y}{25 D^{\prime 2}+D^{2}-10 D D^{\prime}-25} \\
& =e^{x} \frac{(10 \cos 2 y-0-5 \sin 2 y)}{25(-4)+0+0-25} \\
& =\frac{e^{x}}{-125}(10 \cos 2 y-5 \sin 2 y) \\
& =\frac{e^{x}}{25}(\sin 2 y-2 \cos 2 y) \\
& \Rightarrow \mathrm{Z}=f_{1}(y+x)+f_{2}(y-x)+x f_{3}(y-x)+\frac{e^{x}}{25}(\sin 2 y-2 \cos 2 y)
\end{aligned}
$$

Case III: When $F(x, y)=\operatorname{Sin}(a x+b y)$ or $\operatorname{Cos}(a x+b y)$
P.I $=\frac{\operatorname{Sin}(a x+b y) \text { or } \operatorname{Cos}(a x+b y)}{f\left(D^{2}, D D^{\prime}, D^{\prime 2}\right)}$

Put $D^{2}=-a^{2}, D^{\prime}=-a b, D^{\prime 2}=-b^{2}$
Hence P.I $=\frac{\operatorname{Sin}(a x+b y) o r \operatorname{Cos}(a x+b y)}{f\left(-a^{2},-a b,-b^{2}\right)}$

Example 1.7 Solve $\frac{\partial^{2} z}{\partial x^{2}}-\frac{\partial^{2} z}{\partial x \partial y}=\operatorname{Sin} x \operatorname{Cos} 2 y$
Solution: $\Rightarrow\left(D^{2}-D D^{\prime}\right) z=\operatorname{Sin} x \operatorname{Cos} 2 y$

Auxiliary equation is: $m^{2}-m=0$
$\Rightarrow m(m-1)=0$
$\Rightarrow m=0,1$
C.F. $=f_{1}(y)+f_{2}(y+x)$

$$
\begin{aligned}
\text { P.I }=\frac{\operatorname{Sin} x \operatorname{Cos} 2 y}{D^{2}-D D^{\prime}} & =\frac{1}{D^{2}-D D^{\prime}} \frac{1}{2}[\operatorname{Sin}(x+2 y)+\operatorname{Sin}(x-2 y)] \\
& =\frac{1}{2} \frac{1}{D^{2}-D D^{\prime}} \operatorname{Sin}(x+2 y)+\frac{1}{2} \frac{1}{D^{2}-D D^{\prime}} \operatorname{Sin}(x-2 y)
\end{aligned}
$$

Putting $D^{2}=-1, D D^{\prime}=-2$ in the $1^{\text {st }}$ term, $D^{2}=-1, D D^{\prime}=2$ in the $2^{\text {nd }}$ term

$$
=\frac{1}{2} \frac{\operatorname{Sin}(x+2 y)}{-1-(-2)}+\frac{1}{2} \frac{\operatorname{Sin}(x-2 y)}{-1-(2)}
$$

$\Rightarrow$ P.I. $=\frac{1}{2} \operatorname{Sin}(x+2 y)-\frac{1}{6} \operatorname{Sin}(x-2 y)$
Complete solution is $\mathrm{Z}=\mathrm{C} . \mathrm{F}+\mathrm{P} . \mathrm{I}$
$\Rightarrow \mathrm{Z}=f_{1}(y)+f_{2}(y+x)+\frac{1}{2} \operatorname{Sin}(x+2 y)-\frac{1}{6} \operatorname{Sin}(x-2 y)$

Case IV: When $F(x, y)=x^{m} y^{n}$

$$
\begin{aligned}
\text { P.I } & =\frac{x^{m} y^{n}}{f\left(D, D^{\prime}\right)} \\
& =\left[f\left(D, D^{\prime}\right)\right]^{-1} \cdot x^{m} y^{n}
\end{aligned}
$$

Expand $\left[f\left(D, D^{\prime}\right)\right]^{-1}$ in ascending powers of $D$ or $D^{\prime}$ and operate on $x^{m} y^{n}$ term by term.
Example 1.8 Solve $\frac{\partial^{2} z}{\partial x^{2}}+\frac{\partial^{2} z}{\partial x \partial y}-6 \frac{\partial^{2} z}{\partial y^{2}}=x+y$

Solution: $\Rightarrow\left(D^{2}+D D^{\prime}-6 D^{\prime 2}\right) z=x+y$

Auxiliary equation is: $m^{2}+m-6=0$
$\Rightarrow(m+3)(m-2)=0$
$\Rightarrow m=-3,2$
C.F. $=f_{1}(y-3 x)+f_{2}(y+2 x)$
P.I $=\frac{x+y}{D^{2}+D D^{\prime}-6 D^{\prime 2}}$
$=\frac{1}{D^{2}}\left[1+\frac{D^{\prime}}{D}-6 \frac{D^{\prime 2}}{D^{2}}\right]^{-1}(x+y)$
$=\frac{1}{D^{2}}\left[1-\frac{D^{\prime}}{D}+-----\right](x+y)$
$\because(1+t)^{-1}=1-t+t^{2}-t^{3}+---$

$$
=\frac{1}{D^{2}}\left[(x+y)-\frac{1}{D}(0+1)\right]
$$

$$
=\frac{1}{D^{2}}[x+y-x]
$$

P.I. $=\frac{1}{D^{2}}[y]=\frac{y x^{2}}{2}$

Complete solution is $\mathrm{Z}=\mathrm{C} . \mathrm{F}+\mathrm{P} . \mathrm{I}$
$\Rightarrow \mathrm{Z}=f_{1}(y-3 x)+f_{2}(y+2 x)+\frac{y x^{2}}{2}$
Example 1.9 Solve $\left(D^{3}-3 D^{2} D^{\prime}\right) z=x^{2} y$

Solution: Auxiliary equation is: $m^{3}-3 m^{2}=0$

$$
\begin{aligned}
& \Rightarrow m^{2}(m-3)=0 \\
& \begin{aligned}
\Rightarrow m & =0,0,3 \\
\text { C.F. }= & f_{1}(y)+x f_{2}(y)+f_{3}(y+3 x) \\
\text { P.I }= & \frac{x^{2} y}{D^{3}-3 D^{2} D^{\prime}} \\
& =\frac{1}{D^{3}}\left[1-\frac{3 D^{\prime}}{D}\right]^{-1}\left(x^{2} y\right) \\
& =\frac{1}{D^{3}}\left[1+\frac{3 D^{\prime}}{D}\right]\left(x^{2} y\right) \\
\because(1- & -t)^{-1}=1+t+t^{2}+t^{3}+-- \\
& =\frac{1}{D^{3}}\left[\left(x^{2} y\right)+\frac{3}{D}\left(x^{2}\right)\right] \\
& =\frac{1}{D^{3}}\left[x^{2} y+x^{3}\right]
\end{aligned}
\end{aligned}
$$

$$
\begin{aligned}
& =\frac{1}{D^{2}}\left[\frac{x^{3} y}{3}+\frac{x^{4}}{4}\right] \\
& =\frac{1}{D}\left[\frac{x^{4} y}{12}+\frac{x^{5}}{20}\right] \\
\Rightarrow \text { P.I. } & =\left[\frac{x^{5} y}{60}+\frac{x^{6}}{120}\right]
\end{aligned}
$$

Complete solution is $\mathrm{Z}=\mathrm{C} . \mathrm{F}+\mathrm{P} . \mathrm{I}$
$\Rightarrow \mathrm{Z}=f_{1}(y)+x f_{2}(y)+f_{3}(y+3 x)+\left[\frac{x^{5} y}{60}+\frac{x^{6}}{120}\right]$
Case V: In case of any function of $F(x, y)$ or when solution fails for any case by above given methods

$$
\mathrm{P} . \mathrm{I}=\frac{F(x, y)}{f\left(D, D^{\prime}\right)}
$$

Resolve $\frac{1}{f\left(D_{s} D^{\prime}\right)}$ into partial fractions considering $f\left(D_{,} D^{\prime}\right)$ as function of $D$ alone.

$$
\mathrm{P} . \mathrm{I}=\frac{F(x, y)}{D-m D^{\prime}}=\int F(x, c-m x) d x
$$

where $C$ is replaced by $y+m x$ after integration.

Example 1.10 Solve $\left(D^{2}-D D^{\prime}-2 D^{\prime 2}\right) z=(y-1) e^{x}$

Solution: Auxiliary equation is: $m^{2}-m-2=0$
$\Rightarrow(m-2)(m+1)=0$
$\Rightarrow m=2,-1$
C.F. $=f_{1}(y+2 x)+f_{2}(y-x)$
P.I $=\frac{(y-1) e^{x}}{D^{2}-D D^{\prime}-2 D^{\prime 2}}=\frac{(y-1) e^{x}}{D^{2}-2 D D^{\prime}+D D^{\prime}-2 D^{\prime 2}}=\frac{(y-1) e^{x}}{\left(D-2 D^{\prime}\right)\left(D+D^{\prime}\right)}$

$$
=\frac{1}{\left(D-2 D^{\prime}\right)} \int(c+x-1) e^{x} d x
$$

Putting $y=c+x$ as $m=-1$

$$
\begin{aligned}
=\frac{1}{\left(D-2 D^{\prime}\right)}[(c & \left.+x-1) e^{x}-e^{x}\right] \\
& =\frac{1}{\left(D-2 D^{\prime}\right)}\left[(c+x) e^{x}-2 e^{x}\right] \\
& =\frac{1}{\left(D-2 D^{\prime}\right)}\left[(y-x+x) e^{x}-2 e^{x}\right]
\end{aligned}
$$

Putting $c=y-x$

$$
\begin{aligned}
& =\frac{1}{\left(D-2 D^{\prime}\right)}\left[y e^{x}-2 e^{x}\right] \\
& =\frac{1}{\left(D-2 D^{\prime}\right)}\left[(y-2) e^{x}\right] \\
& \quad=\int(c-2 x-2) e^{x} d x
\end{aligned}
$$

Putting $y=c-2 x$

$$
\begin{aligned}
& =(c-2 x-2) e^{x}-(-2) e^{x} \\
& =(c-2 x) e^{x} \\
& =(y+2 x-2 x) e^{x}
\end{aligned}
$$

Putting $c=y+2 x$
$\Rightarrow$ P.I. $=y e^{x}$

Complete solution is $\mathrm{Z}=\mathrm{C} . \mathrm{F}+\mathrm{P} . \mathrm{I}$
$\Rightarrow \mathrm{Z}=f_{1}(y+2 x)+f_{2}(y-x)+y e^{x}$
Example 1.11 Solve $2 \frac{\partial^{2} z}{\partial x^{2}}-5 \frac{\partial^{2} z}{\partial x \partial y}+2 \frac{\partial^{2} z}{\partial y^{2}}=5 \sin (2 x+y)$

Solution: $\Rightarrow\left(2 D^{2}-5 D D^{\prime}+2 D^{\prime 2}\right) z=5 \sin (2 x+y)$

Auxiliary equation is: $2 m^{2}-5 m+2=0$
$\Rightarrow(2 m-1)(m-2)=0$
$\Rightarrow m=\frac{1}{2}, 2$
C.F. $=f_{1}\left(y+\frac{x}{2}\right)+f_{2}(y+2 x)$

$$
\text { P.I }=\frac{5 \sin (2 x+y)}{2 D^{2}-5 D D^{\prime}+2 D^{\prime 2}}
$$

Putting $D^{2}=-4, D D^{\prime}=-2, D^{\prime 2}=-1$, denominator $=0$
$\therefore$ solution fails as per case II, resolving denominator into partial fractions

$$
\text { P.I }=\frac{5 \sin (2 x+y)}{\left(2 D-D^{\prime}\right)\left(D-2 D^{r}\right)}
$$

$=\frac{5}{\left(2 D-D^{\prime}\right)} \int \sin (2 x+(c-2 x)) d x$

Putting $y=c-2 x$

$$
\begin{aligned}
&=\frac{5}{2} \frac{1}{\left(D-\frac{D^{\prime}}{2}\right)} \int \sin c d x \\
&=\frac{5}{2} \frac{x \sin c}{\left(D-\frac{D^{\prime}}{2}\right)} \\
&=\frac{5}{2} \frac{x \sin (y+2 x)}{\left(D-\frac{D^{\prime}}{2}\right)} \\
&=\frac{5}{2} \int x \sin \left[\left(c-\frac{x}{2}\right)+2 x\right] d x
\end{aligned}
$$

Putting $y=c-\frac{x}{2}$

$$
\begin{aligned}
& =\frac{5}{2} \int x \sin \left(c+\frac{3}{2} x\right) d x \\
& =\frac{5}{2}\left[(x) \frac{\left(-\cos \left(c+\frac{3}{2} x\right)\right.}{\frac{3}{2}}-(1) \frac{\left(-\sin \left(c+\frac{3}{2} x\right)\right.}{\frac{9}{4}}\right] \\
& =-\frac{5}{3} x \cos \left(c+\frac{3}{2} x\right)+\frac{10}{9} \sin \left(c+\frac{3}{2} x\right) \\
& \quad=-\frac{5}{3} x \cos \left(y+\frac{x}{2}+\frac{3}{2} x\right)+\frac{10}{9} \sin \left(y+\frac{x}{2}+\frac{3}{2} x\right)
\end{aligned}
$$

Putting $c=y+\frac{x}{2}$
$\Rightarrow$ P.I. $=-\frac{5}{3} x \cos (y+2 x)+\frac{10}{9} \sin (y+2 x)$

Complete solution is $\mathrm{Z}=\mathrm{C} . \mathrm{F}+\mathrm{P} . \mathrm{I}$
$\Rightarrow \mathrm{Z}=f_{1}\left(y+\frac{x}{2}\right)+f_{2}(y+2 x)-\frac{5}{3} x \cos (y+2 x)+\frac{10}{9} \sin (y+2 x)$

### 1.3.2 Non Homogeneous Linear Equations

If in the equation $\left(D, D^{\prime}\right) z=F(x, y)$, the polynomial $f\left(D, D^{\prime}\right)$ in $D, D^{\prime}$ is not homogeneous, then it is called a non-homogeneous partial differential equation. Working Rule to Solve a Non Homogeneous Linear Equation

Step 1: Resolve $f\left(D, D^{\prime}\right)$ into linear factors of the form
$\left(D-m_{1} D^{\prime}-a_{1}\right)\left(D-m_{2} D^{\prime}-a_{2}\right) \ldots \ldots \ldots .\left(D-m_{n} D^{\prime}-a_{n}\right)$
Step2: Auxiliary equation is
$\left(D-m_{1} D^{\prime}-a_{1}\right)\left(D-m_{2} D^{\prime}-a_{2}\right) \ldots \ldots \ldots .\left(D-m_{n} D^{\prime}-a_{n}\right)=0$
Step3: C.F. $=e^{a_{1} x} f_{1}\left(y+m_{1} x\right)+e^{a_{2} x} f_{2}\left(y+m_{2} x\right)+\cdots \ldots \ldots \ldots+e^{a_{n} x} f_{n}\left(y+m_{n} x\right)$
In case of two repeated factors
C.F. $=e^{a x} f_{1}(y+m x)+x e^{a x} f_{2}(y+m x)$

Step4: Find P.I. by using usual methods of homogeneous PDE.
Step5: Complete solution is $\mathrm{Z}=$ C.F. + P.I.
Note: If the Auxiliary equation is of the form
$\left(D^{\prime}-m_{1} D-a_{1}\right)\left(D^{\prime}-m_{2} D-a_{2}\right) \ldots \ldots \ldots .\left(D^{\prime}-m_{n} D-a_{n}\right)=0$
Then C.F. $=e^{a_{1} y} f_{1}\left(x+m_{1} y\right)+e^{a_{2} y} f_{2}\left(x+m_{2} y\right)+\cdots \ldots \ldots \ldots+e^{a_{n} y} f_{n}\left(x+m_{n} y\right)$

Example 1.12 Solve $\left(D^{2}-D^{\prime 2}-3 D+3 D^{\prime}\right) z=e^{x+2 y}$

Solution: Auxiliary equation is: $\left(D^{2}-D^{\prime 2}-3 D+3 D^{\prime}\right)=0$

Clearly $D=D^{\prime}$ is satisfying the equation, $\therefore\left(D-D^{\prime}\right)$ is a factor.

Dividing by $\left(D-D^{\prime}\right)$, we get
$\left(D-D^{\prime}\right)\left(D+D^{\prime}-3\right)=0$
$\Rightarrow\left(D-D^{\prime}-0\right)\left(D+D^{\prime}-3\right)=0$
$\Rightarrow$ C.F. $=f_{1}(y+x)+e^{3 x} f_{2}(y-x)$

$$
\text { P.I }=\frac{e^{x+2 y}}{\left(D^{2}-D^{\prime 2}-3 D+3 D^{\prime}\right)}
$$

Putting $D=1, D^{\prime}=2, f(a, b)=0$
$\Rightarrow$ P.I $=\frac{x e^{x+2 y}}{\frac{d}{d D}\left(D^{2}-D^{\prime 2}-3 D+3 D^{\prime}\right)}=\frac{x e^{x+2 y}}{2 D-3}$

Putting D $=1$
P.I. $=-x e^{x+2 y}$

Complete solution is $\mathrm{Z}=$ C.F. + P.I.
$\Rightarrow \mathrm{Z}=f_{1}(y+x)+e^{3 x} f_{2}(y-x)-x e^{x+2 y}$
Example 1.13 Solve $\left(2 D D^{\prime}+D^{\prime 2}-3 D^{\prime}\right) z=3 \cos (3 x-2 y)$

Solution: Auxiliary equation is: $\left(2 D D^{\prime}+D^{\prime 2}-3 D^{\prime}\right)=0$

$$
\begin{array}{r}
\Rightarrow D^{\prime}\left(D^{\prime}+2 D-3\right)=0 \\
\Rightarrow \text { C.F. }=f_{1}(x)+e^{3 y} f_{2}(x-2 y) \\
\text { P.I }=\frac{3 \cos (3 x-2 y)}{\left(2 D D^{\prime}+D^{\prime 2}-3 D^{\prime}\right)}
\end{array}
$$

Putting $D D^{\prime}=6, D^{\prime 2}=-4$
$\Rightarrow$ P.I $=\frac{3 \cos (3 x-2 y)}{\left(12-4-3 D^{\prime}\right)}=\frac{3 \cos (3 x-2 y)}{\left(8-3 D^{\prime}\right)}$

$$
\begin{aligned}
& =\frac{3\left(8+3 D^{\prime}\right) \cos (3 x-2 y)}{\left(8-3 D^{\prime}\right)\left(8+3 D^{\prime}\right)} \\
& =\frac{3\left(8+3 D^{\prime}\right) \cos (3 x-2 y)}{\left(64-9 D^{\prime 2}\right)} \\
& =\frac{3\left(8+3 D^{\prime}\right) \cos (3 x-2 y)}{(64+36)} \\
& =\frac{3}{100}\left(8+3 D^{\prime}\right) \cos (3 x-2 y) \\
& =\frac{3}{100}[8 \cos (3 x-2 y)+6 \sin (3 x-2 y)] \\
& =\frac{3}{50}[4 \cos (3 x-2 y)+3 \sin (3 x-2 y)]
\end{aligned}
$$

Complete solution is $\mathrm{Z}=\mathrm{C} . \mathrm{F} .+$ P.I.
$\Rightarrow \mathrm{Z}=f_{1}(x)+e^{3 y} f_{2}(x-2 y)+\frac{3}{50}[4 \cos (3 x-2 y)+3 \sin (3 x-2 y)]$

### 1.4 Applications of PDEs (Partial Differential Equations)

In this Section we shall discuss some of the most important PDEs that arise in various branches of science and engineering. Method of separation of variables is the most important tool, we will be using to solve basic PDEs that involve wave equation, heat flow equation and laplace equation.
Wave equation (vibrating string) $: \frac{\partial^{2} u}{\partial t^{2}}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}$
One- dimensional heat flow (in a rod) : $: \frac{\partial u}{\partial t}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}$
Two- dimensional heat flow in steady state (in a rectangular plate ) : $\frac{\partial^{2} u}{\partial t^{2}}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}$
Note: Two dimension heat flow equation in steady state is also known as laplace equation.

## Working Rule for Method of Separation of Variables

Let $u$ be a function of two independent variables $x$ and $t$.
Step1. Assume the solution to be the product of two functions each of which involves only one variable.

Step 2. Calculate the respective partial derivative and substitute in the given PDE.
Step 3. Arrange the equation in the variable separable form and put $\mathrm{LHS}=\mathrm{RHS}=\mathrm{K}$ (as both $x$ and $t$ are independent variables)

Step 4. Solve these two ordinary differential equations to find the two functions of $x$ and $t$ alone.

Example 1.14 Solve the equation $\frac{\partial u}{\partial x}=4 \frac{\partial u}{\partial t}+u$ given that $u(x, 0)=5 e^{-2 x}$
Solution: Step1.
Let $u=X T$.
where $X$ is a function of $x$ alone and $T$ be a function of $t$ alone.
Step 2.
$\frac{\partial u}{\partial x}=X^{\prime} T, \frac{\partial u}{\partial t}=X T^{\prime}$
Substituting these values of $\frac{\partial u}{\partial x}, \frac{\partial u}{\partial t}$ in the given equation
$X^{\prime} T=4 X T^{\prime}+X T \quad \Rightarrow X^{\prime} T=X\left(4 T^{I}+T\right)$
Step 3.
$\Rightarrow \frac{X^{I}}{X}=\frac{4 T^{I}}{T}+1$
Putting LHS $=$ RHS $=\mathrm{K}$
Step 4.

$$
\begin{gathered}
\text { i.e. } \quad \frac{X^{\prime}}{X}=K \\
\Rightarrow \log X=K x+\log C_{1} \\
\Rightarrow \log \frac{X}{C_{1}}=K x \\
\Rightarrow X=C_{1} e^{K x} \ldots \ldots \ldots
\end{gathered}
$$

$$
\begin{gather*}
\frac{4 T^{I}}{T}+1=K \\
\frac{T^{\prime}}{T}=\frac{1}{4}(K-1) \\
\Rightarrow \log T=\frac{1}{4}(K-1) t+\log C_{2} \\
\Rightarrow T=C_{2} e^{\frac{1}{4}(K-1) t} \ldots \ldots . . \tag{3}
\end{gather*}
$$

Using (2) and (3) in (1)
$u(x, t)=C_{1} C_{2} e^{K x} e^{\frac{1}{4}(K-1) t}$
$\Rightarrow u(x, 0)=C_{1} C_{2} e^{K x} \ldots \ldots \ldots$. (5)
Given that $u(x, 0)=5 e^{-2 x}$, using in (5)
$\Rightarrow 5 e^{-2 x}=C_{1} C_{2} e^{K x}$
$\Rightarrow C_{1} C_{2}=5, K=-2$
Using (6)in (4)
$u(x, t)=5 e^{-\left(2 x+\frac{3 t}{4}\right)}$

### 1.4.1 Solution of wave equation using method of separation of variables

Wave equation is given by $\frac{\partial^{2} u}{\partial t^{2}}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}$
whereugives displacement at distance $x$ from origin at any time $t$. To solve wave equation using method of separation of variables,
Let $u=X T$ $\qquad$
where $X$ is a function of $x$ alone and $T$ be a function of $t$ alone.
$\therefore \frac{\partial^{2} u}{\partial t^{2}}=X T^{\prime \prime}, \frac{\partial^{2} u}{\partial x^{2}}=X^{\prime \prime} T$
Substituting these values of $\frac{\partial^{2} u}{\partial t^{2}}, \frac{\partial^{2} u}{\partial x^{2}}$ in the wave equation given by (1)
$X T^{\prime \prime}=c^{2} X^{\prime \prime} \mathrm{T}$
Arranging in variable separable form $\Rightarrow \frac{X^{\prime \prime}}{X}=\frac{1}{c^{2}} \frac{T^{\prime \prime}}{T}$
Equating LHS $=$ RHS $=\mathrm{K} \quad(\because X$ and $T$ are independent $)$
$\Rightarrow \frac{X^{\prime \prime}}{X}=K$ and $\frac{1}{c^{2}} \frac{T^{\prime \prime}}{T}=K$
$\Rightarrow X^{I I}-K X=0$ and $T^{\prime \prime}-K C^{2} T=0$ $\qquad$
Solving ordinary differential equations given in (3), three cases arise
(i) When $K$ is +ve and $=p^{2}$ say

$$
X=c_{1} e^{p x}+c_{2} e^{-p x}, T=c_{3} e^{c p t}+c_{4} e^{-c p t}
$$

(ii) When $K$ is -ve and $=-p^{2}$ say

$$
X=c_{1} \cos p x+c_{2} \sin p x, T=c_{3} \cos c p t+c_{4} \sin c p t
$$

(iii)When $K=0$

$$
X=c_{1} x+c_{2}, T=c_{3} t+c_{4}
$$

Again since we are dealing with wave equation, $u$ must be a periodic function of $x$ and $t, \therefore$ the solution must involve trigonometric terms. Hence the solution given by (ii) i.e. corresponding to $K=-p^{2}$, is the most plausible solution, substituting (ii) in equation (2)
$u(x, t)=\left(c_{1} \operatorname{cosp} x+c_{2} \sin p x\right)\left(c_{3} \operatorname{coscp} t+c_{4} \sin c p t\right) \ldots \ldots$. (4)
Which is the required solution of wave equation.
Again if we consider a string of length $l$ tied at both ends at $x=0$ and $x=l$, then displacement of the string at end points at any time $t$ is zero .

$$
\begin{align*}
& \Rightarrow u(0, t)=0 \ldots \ldots \ldots  \tag{5}\\
& \operatorname{and} u(l, t)=0 \ldots \ldots \ldots .
\end{align*}
$$

using (5)in (4) $\Rightarrow 0=C_{1}\left(c_{3} \operatorname{coscpt}+c_{4} \operatorname{sincpt}\right)$
$\Rightarrow C_{1}=0$ $\qquad$ (7)

Using (7) in (4), wave equation reduces to
$u(x, t)=c_{2} \operatorname{sinp} x\left(c_{3} \operatorname{coscp} t+c_{4} \operatorname{sincp} t\right)$

Now using (6) in (8) $\Rightarrow 0=c_{2} \operatorname{sinpl}\left(c_{3} \operatorname{coscpt}+c_{4} \operatorname{sincpt}\right)$
$\Rightarrow \sin p l=0 \quad \because c_{2} \neq 0$ and $\left(c_{3} \operatorname{coscp} t+c_{4} \operatorname{sinc} c t\right) \neq 0$
$\Rightarrow p l=n \pi \Rightarrow p=\frac{n \pi}{l}, n=1,2,3$.
using (9) in (8) $\Rightarrow u(x, t)=c_{2} \sin \frac{n \pi x}{l}\left(c_{3} \cos \frac{n \pi c t}{l}+c_{4} \sin \frac{n \pi c t}{l}\right)$
Adding up the solutions for different values of $n$, we get
$u(x, t)=\sum_{n=1}^{\infty}\left(a_{n} \cos \frac{n \pi c t}{l}+b_{n} \sin \frac{n \pi c t}{l}\right) \sin \frac{n \pi x}{l} \ldots \ldots \ldots(10$
(10) is also a solution of wave equation

Example 1.15 : A string is stretched and fastened to 2 points $l$ apart. Motion is started by displacing the string in the form $y=a \sin \frac{\pi x}{l}$ from which, it is released at time $t=0$. Show that the displacement at any point at a distance $x$ from one end at time $t$ is given by
$y(x, t)=a \sin \frac{\pi x}{l} \cos \frac{\pi c t}{l}$
Solution: Let the equation of vibrating string be given by $\frac{\partial^{2} y}{\partial t^{2}}=c^{2} \frac{\partial^{2} y}{\partial x^{2}}$
Boundary value conditions are given by :

$$
\begin{aligned}
& y(0, t)=0 \ldots \ldots \ldots .(2) \\
& y(l, t)=0 \ldots \ldots \ldots \ldots(3) \\
& \frac{\partial y}{\partial t}=0 \ldots \ldots \ldots .(4) \\
& y(x, 0)=a \sin \frac{\pi x}{l} \ldots \ldots \text { (5) }
\end{aligned}
$$

Let the solution of (1)be given by
$y(x, t)=\left(c_{1} \cos p x+c_{2} \sin p x\right)\left(c_{3} \operatorname{coscp} t+c_{4} \sin c p t\right) \ldots . .(6)$
Using (2) in (6)

$$
\begin{aligned}
& \Rightarrow 0=c_{1}\left(c_{3} \operatorname{coscpt}+c_{4} \sin c p t\right) \\
& \Rightarrow c_{1}=0 \ldots \ldots \ldots \ldots \ldots . .(7)
\end{aligned}
$$

Using (7) in (6), wave equation reduces to
$y(x, t)=c_{2} \sin p x\left(c_{3} \operatorname{coscp} t+c_{4} \operatorname{sincp} t\right)$
Now using (3) in(8) $\Rightarrow 0=c_{2} \sin p l\left(c_{3} \operatorname{coscpt}+c_{4} \operatorname{sincpt}\right)$
$\Rightarrow \sin p l=0 \quad \because c_{2} \neq 0$ and $\left(c_{3} \operatorname{coscp} t+c_{4} \operatorname{sincpt}\right) \neq 0$
$\Rightarrow p l=n \pi \Rightarrow p=\frac{n \pi}{l}$
using (9) in (8) $\Rightarrow y(x, t)=c_{2} \sin \frac{n \pi x}{l}\left(c_{3} \cos \frac{n \pi c t}{l}+c_{4} \sin \frac{n \pi c t}{l}\right) \ldots \ldots$.(10)
Now to use (4), differentiating (10) partially w.r.t. $t$
$\Rightarrow \frac{\partial y}{\partial t}=c_{2} \sin \frac{n \pi x}{l}\left(-c_{3} \frac{n \pi c}{l} \sin \frac{n \pi c t}{l}+c_{4} \frac{n \pi c}{l} \cos \frac{n \pi c t}{l}\right)$
Putting $t=0$

$$
\begin{align*}
& \frac{\partial y}{\partial t_{t=0}}=c_{2} \sin \frac{n \pi x}{l}\left(c_{4} \frac{n \pi c}{l}\right)=0 \text { using(4) } \\
& \Rightarrow c_{4}=0 \quad \because c_{2} \neq 0 \cdots \cdots \cdots \tag{11}
\end{align*}
$$

Using (11) in (10)
$\Rightarrow y(x, t)=c_{2} c_{3} \sin \frac{n \pi x}{l} \cos \frac{n \pi c t}{l} \ldots \ldots$. (12)
Using (5) in (12)
$\Rightarrow a \sin \frac{\pi x}{l}=c_{2} c_{3} \sin \frac{n \pi x}{l}$
$\Rightarrow c_{2} c_{3}=a, n=1$
using in (12)
$\Rightarrow y(x, t)=\operatorname{asin} \frac{\pi x}{l} \cos \frac{\pi c t}{l}$
Note : Above example can also be solved using solution of wave equation given by equation (10) in section 1.4.1
$y(x, t)=\sum_{n=1}^{\infty}\left(a_{n} \cos \frac{n \pi c t}{l}+b_{n} \sin \frac{n \pi c t}{l}\right) \sin \frac{n \pi x}{l}$

It is to be noted that boundary value conditions (2)and (3)have been already used in this solution.
Example 1.16 A tightly stretched string with fixed end points at $x=0$ and $x=l$ is initially at rest in its equilibrium position. If it is set vibrating by giving to each of its points a velocity $\lambda x(l-x)$, find the displacement of the string at any distance $x$ from one end at any time $t$.

Solution: Let the equation of vibrating string be given by $\frac{\partial^{2} u}{\partial t^{2}}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}$
Boundary value conditions are given by :

$$
\begin{align*}
& u(0, t)=0 \ldots \ldots \ldots .(2) \\
& u(l, t)=0 \ldots \ldots \ldots \ldots \text { (3) } \\
& u(x, 0)=0 \ldots \ldots \ldots \ldots \text { (4) }  \tag{3}\\
& \frac{\partial u}{\partial t}_{t=0}=\lambda x(l-x) \ldots \ldots \ldots \text { (5) } \tag{4}
\end{align*}
$$

solution of wave equation as given by equation (10) in section 1.4.1 is
$u(x, t)=\sum_{n=1}^{\infty}\left(a_{n} \cos \frac{n \pi c t}{l}+b_{n} \sin \frac{n \pi c t}{l}\right) \sin \frac{n \pi x}{l} \ldots . .(6)$
It is to be noted that boundary value conditions (2)and (3) have been already used in this solution.

Now using (4) in (6)
$0=\sum_{n=1}^{\infty} a_{n} \sin \frac{n \pi x}{l}$
$\Rightarrow a_{n}=0$
Using (7) in (6)
$u(x, t)=\sum_{n=1}^{\infty} b_{n} \sin \frac{n \pi c t}{l} \sin \frac{n \pi x}{l}$
Now to use (5), differentiating (8) partially w.r.t.t
$\frac{\partial u}{\partial t}=\sum_{n=1}^{\infty} b_{n} \frac{n \pi c}{l} \cos \frac{n \pi c t}{l} \sin \frac{n \pi x}{l}$

Putting $t=0$,
$\Rightarrow \frac{\partial u}{\partial t}_{t=0}=\frac{\pi c}{l} \sum_{n=1}^{\infty} n b_{n} \sin \frac{n \pi x}{l}$.
Using (5) in (9)
$\lambda x(l-x)=\frac{\pi c}{l} \sum_{n=1}^{\infty} n b_{n} \sin \frac{n \pi x}{l}$
Multiplying both sides by $\sin \frac{n \pi x}{l}$ and integrating w.r.t. $x$ within the limits 0 to $l$
$\Rightarrow \int_{0}^{l} \lambda x(l-x) \sin \frac{n \pi x}{l} d x=\frac{\pi c}{l} n b_{n} \int_{0}^{l} \sin ^{2} \frac{n \pi x}{l} d x$
$=\frac{\pi c}{2 l} n b_{n} \int_{0}^{l}\left(1-\cos \frac{2 n \pi x}{l}\right) d x$
$=\frac{\pi c}{2} n b_{n}$
$\Rightarrow \pi c n b_{n}=2 \int_{0}^{l} \lambda x(l-x) \sin \frac{n \pi x}{l} d x$
$=2 \lambda \int_{0}^{l}\left(l x-x^{2}\right) \sin \frac{n \pi x}{l} d x$
$\Rightarrow \pi c n b_{n}=2 \lambda\left[\left(l x-x^{2}\right) \cdot\left(\frac{-l}{n \pi} \cos \frac{n \pi x}{l}\right)-(l-2 x)\left(-\frac{l^{2}}{n^{2} \pi^{2}} \sin \frac{n \pi x}{l}\right)\right.$
$\left.+(-2)\left(\frac{l^{3}}{n^{3} \pi^{3}} \cos \frac{n \pi x}{l}\right)\right]_{0}^{I}$
$\Rightarrow \pi c n b_{n}=2 \lambda\left[\frac{-2 l^{3}}{n^{3} \pi^{3}} \cos n \pi+\frac{2 l^{3}}{n^{3} \pi^{3}}\right]$
$\Rightarrow \pi c n b_{n}=\frac{4 \lambda l^{3}}{n^{3} \pi^{3}}\left[1-(-1)^{n}\right]$
$\Rightarrow \pi c n b_{n}=\left[\begin{array}{cc}\frac{8 \lambda l^{3}}{n^{3} \pi^{3}} & \text { if } n \text { is odd } \\ 0 & \text { if } n \text { is even }\end{array}\right]$
taking $n=2 m-1$

$$
b_{n}=\frac{8 \lambda l^{3}}{C \pi^{4}(2 m-1)^{4}} \ldots \ldots \ldots .(10)
$$

using(10) in 8), required solution is
$u(x, t)=\frac{8 \lambda l^{3}}{C \pi^{4}} \sum_{m=1}^{\infty} \frac{1}{(2 m-1)^{4}} \sin \frac{(2 m-1) \pi c t}{l} \sin \frac{(2 m-1) \pi x}{l}$

### 1.4.2 Solution of heat equation using method of separation of variables

One dimensional heat flow equation is given by $\frac{\partial u}{\partial t}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}$
where $u(x, t)$ is temperature function at distance $x$ from origin at any time $t$. To solve heat equation using method of separation of variables,

Let $u=X T$
where $X$ is a function of $x$ alone and $T$ be a function of $t$ alone.
$\therefore \frac{\partial u}{\partial t}=X T^{I}, \frac{\partial^{2} u}{\partial x^{2}}=X^{n} T$
Substituting these values of $\frac{\partial u}{\partial t}, \frac{\partial^{2} u}{\partial x^{2}}$ in the heat equation given by (1)
$X T^{\prime}=c^{2} X^{H} \mathrm{~T}$
Arranging in variable separable form $\Rightarrow \frac{X^{\prime \prime}}{X}=\frac{1}{c^{2}} \frac{T^{\prime}}{T}$
Equating LHS $=$ RHS $=\mathrm{K} \quad(\because X$ and $T$ are independent $)$
$\Rightarrow \frac{X^{\prime \prime}}{X}=K \operatorname{and} \frac{1}{c^{2}} \frac{T^{\prime}}{T}=K$
$\Rightarrow X^{\prime \prime}-K X=0$ and $T^{I}-K c^{2} T=0$ $\qquad$
Solving ordinary differential equations given in (3), three cases arise
(i) When $K$ is +ve and $=p^{2}$ say

$$
\begin{aligned}
X^{\prime \prime}-p^{2} X & =0 \text { and } T^{\prime}=p^{2} c^{2} T \\
& X=A e^{p x}+B e^{-p x}, T=C e^{p^{2} c^{2} t} \\
& \Rightarrow u(x, t)=\left(A e^{p x}+B e^{-p x}\right) C e^{p^{2} c^{2} t} \\
& \Rightarrow u(x, t)=\left(\left(c_{1} e^{p x}+c_{2} e^{-p x}\right) e^{p^{2} c^{2} t}\right.
\end{aligned}
$$

(ii) When $K$ is - ve and $=-p^{2}$ say

$$
\begin{aligned}
& X=A \cos p x+B \sin p x, T=C e^{-c^{2} p^{2} t} \\
& \Rightarrow u(x, t)=(A \cos p x+B \sin p x) C e^{-c^{2} p^{2} t} \\
& \Rightarrow u(x, t)=\left(c_{1} \cos p x+c_{2} \sin p x\right) e^{-c^{2} p^{2} t}
\end{aligned}
$$

(iii) When $K=0$

$$
\begin{aligned}
& X^{\prime \prime}=0 \text { and } T^{\prime}=0 \\
& \Rightarrow X=A x+B, T=C \\
& \Rightarrow u(x, t)=(A x+B) C \\
& \Rightarrow u(x, t)=c_{1} x+c_{2}
\end{aligned}
$$

The solution given by (ii) i.e. corresponding to $K=-p^{2}$, is the most plausible solution for steady state.

Special case: When the ends of a rod are kept at $0^{\circ} \mathrm{C}$
One dimensional heat equation in steady state is given by :
$u(x, t)=\left(c_{1} \cos p x+c_{2} \sin p x\right) e^{-c^{2} p^{2} t}$
Also since ends of a rod are kept at $0^{\circ} \mathrm{C}$

$$
\begin{align*}
& u(0, t)=0  \tag{5}\\
& u(l, t)=0 \tag{6}
\end{align*}
$$

$\qquad$

Using (5) in (4) $\Rightarrow 0=C_{1} e^{-c^{2} p^{2} t}$
$\Rightarrow C_{1}=0$ $\qquad$
Using (7) in (4), wave equation reduces to
$u(x, t)=c_{2} \sin p x e^{-c^{2} p^{2} t} \ldots \ldots$.
Now using (6) in (8) $\Rightarrow 0=c_{2} \operatorname{sinple} e^{-c^{2} p^{2} t} \Rightarrow \sin p l=0$

$$
\begin{equation*}
\because c_{2} \neq 0 \text { and } e^{-c^{2} p^{2} t} \neq 0 \tag{9}
\end{equation*}
$$

$\Rightarrow p l=n \pi \Rightarrow p=\frac{n \pi}{l}, n=1,2,3$.
using (9) in (8) $\Rightarrow u(x, t)=c_{2} \sin \frac{n \pi x}{l} e^{\frac{-c^{2} n^{2} \pi^{2} t}{l^{2}}}$
Adding up the solutions for different values of $n$, the most general solution is given by
$u(x, t)=\sum_{n=1}^{\infty} b_{n} \sin \frac{n \pi x}{l} e^{\frac{-c^{2} n^{2} \pi^{2} t}{l^{2}}} \ldots \ldots \ldots$.(10)
Example 1.17:A rod of length $l$ with insulated sides is initially at a uniform temperature $u_{0}$. Its ends are suddenly cooled to $0^{\circ} \mathrm{C}$ and are kept at that temperature, find the temperature formula $u(x, t)$.

Solution: One dimensional heat flow equation is given by $\frac{\partial u}{\partial t}=c^{2} \frac{\partial^{2} u}{\partial x^{2}}$
Solution of one dimensional heat equation in steady state is given by :

$$
\begin{gather*}
u(x, t)=\left(c_{1} \cos p x+c_{2} \sin p x\right) e^{-c^{2} p^{2} t} \\
u(0, t)=0 \ldots \ldots \ldots \ldots .(2) \\
u(l, t)=0 \ldots \ldots \ldots \ldots .(3) \tag{4}
\end{gather*}
$$

Also initial condition is $u(x, 0)=u_{0}$ $\qquad$
The most general solution of heat equation(1) using (2)and(3) is given by
$u(x, t)=\sum_{n=1}^{\infty} b_{n} \sin \frac{n \pi x}{l} e^{\frac{-c^{2} n^{2} \pi^{2} t}{l^{2}}}$
Using (4)in (5) $\Rightarrow u_{0}=\sum_{n=1}^{\infty} b_{n} \sin \frac{n \pi x}{l}$
Multiplying both sides of (6) by $\sin \frac{n \pi x}{l}$, and integrating w.r.t. $x$ within the limits 0 to $l$ $\int_{0}^{l} u_{0} \sin \frac{n \pi x}{l} d x=b_{n} \int_{0}^{l} \sin ^{2} \frac{n \pi x}{l} d x$

$$
\begin{align*}
& \qquad b_{n}=\frac{2}{l} \int_{0}^{l} u_{0} \sin \frac{n \pi x}{l} d x \\
& =\frac{-2 u_{0}}{l}\left[\frac{l}{n \pi} \cos \frac{n \pi x}{l}\right]_{0}^{l} \\
& \qquad \Rightarrow b_{n}=\frac{2 u_{0}}{n \pi}\left[1-(-1)^{n}\right]=\left[\begin{array}{cc}
\frac{4 u_{0}}{n \pi}, & \text { if } n \text { is odd } \\
0 & \text { if } n \text { is even }
\end{array}\right] \\
& \text { Putting } n=2 m-1, \Rightarrow b_{n}=\frac{4 u_{0}}{(2 m-1) \pi} \ldots \ldots . .7
\end{align*}
$$

Using (7) in (5), the required temperature formula is
$u(x, t)=\frac{4 u_{0}}{\pi} \sum_{m=1}^{\infty} \frac{1}{2 m-1} \sin \frac{(2 m-1) \pi x}{l} e^{\frac{-c^{2}(2 m-1)^{2} \pi^{2} t}{l^{2}}}$

### 1.4.3 Solution of laplace equation (two dimensional heat flow) using method of

 separation of variablesConsider the heat flow in a uniform rectangular metal plate at any time $t$; if $u(x, y)$ be the temperature at time $t$, two dimensional heat flow equation is given by

$$
\frac{\partial u}{\partial t}=c^{2}\left(\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}\right)
$$

In steady state, $u$ doesn't change with $t$ and hence $\frac{\partial u}{\partial t}=0$
$\therefore$ Two dimensional heat flow equation in steady state is given by $\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}=0$
where $u(x, y)$ is temperature function at any point $(x, y)$ of the rectangular metal plate. This is called Laplace equation in two dimensions. To solve Laplace equation using method of separation of variables,
Let $u=X Y$
where $X$ is a function of $x$ alone and $Y$ be a function of $y$ alone.
$\therefore \frac{\partial^{2} u}{\partial x^{2}}=X^{n} Y, \frac{\partial^{2} u}{\partial y^{2}}=X Y^{n}$

Substituting these values of $\frac{\partial^{2} u}{\partial x^{2}}, \frac{\partial^{2} u}{\partial y^{2}}$ in the Laplace equation given by (1)
$X^{\prime \prime} Y+X Y^{\prime \prime}=0$
Arranging in variable separable form $\Rightarrow \frac{X^{\prime \prime}}{X}=-\frac{Y^{\prime \prime}}{Y}$
Equating LHS $=\mathrm{RHS}=\mathrm{K} \quad(\because x$ and $y$ are independent $)$
$\Rightarrow \frac{X^{\prime \prime}}{X}=K$ and $-\frac{Y^{\prime \prime}}{Y}=K$
$\Rightarrow X^{I I}-K X=0$ and $Y^{I I}+K Y=0$ $\qquad$
Solving ordinary differential equations given in (3), three cases arise
(i) When $K$ is +ve and $=p^{2}$ say

$$
\begin{aligned}
& X=c_{1} e^{p x}+c_{2} e^{-p x}, Y=c_{3} \cos p y+c_{4} \sin p y \\
& \Rightarrow u(x, y)=\left(c_{1} e^{p x}+c_{2} e^{-p x}\right)\left(c_{3} \cos p y+c_{4} \sin p y\right)
\end{aligned}
$$

(ii) When $K$ is -ve and $=-p^{2}$ say

$$
\begin{aligned}
& X=\left(c_{1} \cos p x+c_{2} \sin p x, Y=\left(c_{3} e^{p y}+c_{4} e^{-p y}\right)\right. \\
& \Rightarrow u(x, y)=\left(c_{1} \cos p x+c_{2} \sin p x\right)\left(c_{3} e^{p y}+c_{4} e^{-p y}\right)
\end{aligned}
$$

(iii)When $K=0$

$$
\begin{aligned}
& X^{\prime \prime}=0 \text { and } Y^{\prime \prime}=0 \\
& \Rightarrow X=c_{1} x+c_{2}, Y=c_{3} y+c_{4} \\
& \Rightarrow u(x, t)=\left(c_{1} x+c_{2}\right)\left(c_{3} y+c_{4}\right)
\end{aligned}
$$

The solution given by (ii) i.e. corresponding to $K=-p^{2}$, is the most plausible solution for steady state.

Example 1.18 An infinitely long rectangular uniform plate with breadth $\pi$ is bounded by two parallel edges maintained at $0^{\circ} \mathrm{C}$. Base of the plate is kept at a temperature $u_{0}$ at all points. Determine the temperature at any point of the plate in the steady state.

Solution : In steady state, two dimensional heat flow equation is given by $\frac{\partial^{2} u}{\partial x^{2}}+\frac{\partial^{2} u}{\partial y^{2}}=0 \ldots$ (1)
Boundary value conditions are

$$
\begin{align*}
& u(0, y)=0 \ldots \ldots \ldots \ldots(2) \\
& u(\pi, y)=0 \ldots \ldots \ldots \ldots \text { (3) } \tag{3}
\end{align*}
$$

$\lim _{y \rightarrow \infty} u(x, y)=0,0<x<\pi \cdots \cdots \cdots(4)$
$u(x, 0)=u_{0}, 0<x<\pi$
Solution of (1) is given by :
$u(x, y)=\left(c_{1} \cos p x+c_{2} \sin p x\right)\left(c_{3} e^{p y}+c_{4} e^{-p y}\right)$
Using (2) in (6)
$\Rightarrow u(0, y)=0=c_{1}\left(c_{3} e^{p y}+c_{4} e^{-p y}\right)$
$\Rightarrow c_{1}=0$ $\qquad$
Using (7) in (6)
$u(x, y)=c_{2} \sin p x\left(c_{3} e^{p y}+c_{4} e^{-p y}\right)$
Using (3) in (8)

$$
\begin{aligned}
& \begin{array}{l}
\Rightarrow u(\pi, y)=0=c_{2} \sin p \pi\left(c_{3} e^{p y}+c_{4} e^{-p y}\right) \\
\\
\quad \Rightarrow \sin p \pi=0 \\
\quad \Rightarrow p \pi=\mathrm{n} \pi
\end{array} \\
& \Rightarrow p=\mathrm{n} \ldots \ldots \ldots .9
\end{aligned}
$$

Using (9) in (8)
$\Rightarrow u(x, y)=c_{2} \sin n x\left(c_{3} e^{n y}+c_{4} e^{-n y}\right)$ $\qquad$
Using (4) in (10)

$$
\begin{align*}
& \lim _{y \rightarrow \infty} u(x, y)=0=c_{2} \sin n x \lim _{y \rightarrow \infty}\left(c_{3} e^{n y}+c_{4} e^{-n y}\right) \\
& \Rightarrow c_{3}=0 \ldots \ldots \ldots \tag{1}
\end{align*}
$$

Using (11) in (10)
$\Rightarrow u(x, y)=c_{2} c_{4} \sin n x e^{-n y}$
The most general solution of heat equation is given by
$u(x, t)=\sum_{n=1}^{\infty} b_{n} \sin n x . e^{-n y} \ldots \ldots$. (12) where $c_{2} c_{4}=b_{n}$

Using (5) in (10)
$\Rightarrow u(x, 0)=u_{0}=\sum_{n=1}^{\infty} b_{n} \sin n x$
Multiplying both sides bysin $n x$, and integrating w.r.t. $x$ within the limits 0 to $\pi$
$b_{n}=\frac{2}{\pi} \int_{0}^{\pi} u_{0} \sin n x d x$
$\Rightarrow b_{n}=\frac{2 u_{0}}{\pi}\left[\frac{-\cos n x}{n}\right]_{0}^{\pi}=\left[\begin{array}{cc}\frac{4 u_{0}}{n \pi}, & \text { if } n \text { is odd } \\ 0 & \text { if } n \text { is even }\end{array}\right] \ldots . .(13)$
Let $n=2 m-1$ as $n$ is odd
Using (13) in (12) the required temperature formula is:
$u(x, t)=\frac{4 u_{0}}{\pi} \sum_{n=1}^{\infty} \frac{1}{(2 m-1)} \sin (2 m-1) x . e^{-(2 m-1) y}$

