

FINITE ELEMENT MODELLING

Course code:AME014 VI- semester Regulation: IARE R-16

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COs	COURSE OUTCOMES
CO1	Describe the concept of FEM and difference between the FEM with other methods and problems based on 1-D bar elements and shape functions.
CO2	Derive elemental properties and shape functions for truss and beam elements and related problems.
CO3	Understand the concept deriving the elemental matrix and solving the basic problems of CST and axi-symmetric solids
CO4	Explore the concept of steady state heat transfer in fin and composite slab
CO5	Understand the concept of consistent and lumped mass models and solve the dynamic analysis of all types of elements.



UNIT-I INTRODUCTION TO FEM

Introduction



Introduction to FEM:

- Stiffness equations for a axial bar element in local co-ordinates using Potential Energy approach and Virtual energy principle.
- Finite element analysis of uniform, stepped and tapered bars subjected to mechanical and thermal loads.
- Assembly of Global stiffness matrix and load vector.
- Quadratic shape functions.
- Properties of stiffness matrix

Axially Loaded Bar – Governing Equations and Boundary Conditions

• Differential Equation

$$\frac{d}{dx} \begin{bmatrix} EA(x) & \frac{du}{dx} \end{bmatrix}_{+}^{+} f(x) = 0 \qquad 0 < x < L$$

- Boundary Condition Types
 - Prescribed displacement (essential BC)
 - Prescribed force/derivative of displacement (natural BC)

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Axially Loaded Bar – Boundary Conditions



Examples



□ Simple support



Potential Energy



Elastic Potential Energy (PE)

• Spring case





• Axially loaded bar

Unreformed: PE = 0deformed: $PE = \frac{1}{2} \int_{0}^{L} \sigma \varepsilon A dx$ • Elastic body $PE = \frac{1}{2} \int_{V}^{T} \sigma \varepsilon dv$

Potential Energy



• Work Potential (WE)



- f: distributed force over a line
- P: point force
- u: displacement

• Total Potential Energy

$$\Pi = \frac{1}{2} \int \sigma \mathcal{E} A dx - \int u \cdot f dx - P \cdot u_B$$

• Principle of Minimum Potential Energy

For conservative systems, of all the kinematically admissible displacement fields, those corresponding to equilibrium extremize the total potential energy. If the extremum condition is a minimum, the equilibrium state is stable.

Potential Energy + Rayleigh-Ritz Approach



Step 1: assume a displacement field

$$u = \sum_{i} a_i \phi_i(x) \quad i = 1 \text{ to } n$$

- f is shape function / basis function
- *n* is the order of approximation
- Step 2: calculate total potential energy

Potential Energy + Rayleigh-Ritz Approach

• Example:



• Step 3: select a_i so that the total potential energy is minimum



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Galerkin's Method



• Example:



 In the Galerkin's method, the weight function is chosen to be the same as the shape function.

Galerkin's Method



 FEM Formulation of Axially Loaded Bar – Governing Equations

Oifferential Equation

$$\frac{d}{dx} \begin{bmatrix} EA(x) & du \\ dx \end{bmatrix} + f(x) = 0 \qquad 0 < x < L$$

• Weighted-Integral Formulation

$$\int_{0}^{L} w \left(\frac{d}{dx} \left[\frac{du}{dx} \right] + f(x) \right) dx = 0$$

• Weak Form

$$0 = \int_{0}^{L} \left[\frac{dw}{dx} \left(EA(x) \frac{du}{dx} \right) - wf(x) \right] dx - w \left[EA(x) \frac{du}{dx} \right]_{0}^{L}$$



• Example:



• Step 2: Weak form of one element

$$\int_{x_1}^{x_2} \left[\frac{dw}{dx} \left(EA(x) \frac{du}{dx} \right) - w(x) f(x) \right] dx - w(x) \left(EA(x) \frac{du}{dx} \right) \right]_{x_1}^{x_2} = 0$$

$$\int_{x_1}^{x_2} \left[\frac{dw}{dx} \left(EA(x) \frac{du}{dx} \right) - w(x) f(x) \right] dx - w x_2 P_2 - w x_1 P_1 = 0$$



 x_2







• Example (*cont*):



• Step 3: Choosing shape functions- linear shape functions $u = \phi_1 u_1 + \phi_2 u_2$



• Example (*cont*):



• Step 4: Forming element equation

• Let
$$w = \phi$$
, weak form becomes

$$\int_{x_1}^{x_1} -\frac{1}{l} \left(EA \cdot \frac{u-u}{2} \right) |dx - \int_{x_1}^{x_2} \phi_1 f dx - \phi_1 P_2 - \phi_1 P_1 = 0 \longrightarrow \frac{EA}{l} - \frac{EA}{l} u_1 - \frac{EA}{l} u_2 = \int_{x_1}^{x_2} \phi_1 f dx + P_1$$
• Let $w = \phi$, weak form becomes

$$\int_{x_1}^{x_1} \frac{1}{l} \left(EA \cdot \frac{u-u}{l} \right) |dx - \int_{x_1}^{x_2} \phi_2 f dx - \phi_2 P_2 - \phi_2 P_1 = 0 \longrightarrow -\frac{EA}{l} u_1 + \frac{EA}{l} u_2 = \int_{x_1}^{x_2} \phi_2 f dx + P_2$$

$$EA \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \frac{1}{l} \left(u_1 \right) |dx - \int_{x_1}^{x_2} \phi_1 f dx | dx | dx + \int_{x_1}^{x_2} \phi_1 f dx | dx | dx + \int_{x_1}^{x_2} \phi_2 f dx - \phi_2 P_2 - \phi_2 P_1 = 0 \longrightarrow -\frac{EA}{l} u_1 + \frac{EA}{l} u_2 = \int_{x_1}^{x_2} \phi_2 f dx + P_2$$

$$EA \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} \frac{1}{l} \left(u_1 \right) |dx - \int_{x_1}^{x_2} \phi_1 f dx | dx | dx + \int_{x_1}^{x_2} \phi_1 f dx | dx | dx + \int_{x_1}^{x_2} \phi_1 f dx | dx | dx + \int_{x_1}^{x_2} \phi_2 f dx + P_2$$

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• Example (*cont*):



• Step 5: Assembling to form system equation Approach 1: $\begin{bmatrix} 1 & -1 & 0 & 0 \end{bmatrix} \begin{bmatrix} uf \\ fi' \end{bmatrix}$

Element 1: Element 2: Element 3:

• Example (*cont*):

 Step 5: Assembling to form system equation Assembled System:







Approximation Methods



• Example (*cont*):



Step 5: Assembling to form system equation
 Approach 2: Element connectivity table

$$k_{ij}^e \rightarrow K_{IJ}$$

	Element 1	Element 2	Element 3
1	1	2	3
2	2	3	4
↓ ocal noo (i,j)	de glob	al node inde	ex

Approximation Methods







• Step 6: Imposing boundary conditions and forming condense

Condensed system:

$$\begin{pmatrix} E^{I}A^{I} \\ \underline{l^{I}} \\ -\frac{E^{II}A^{II}}{l^{II}} + \frac{E^{II}A^{II}}{l^{II}} & \underline{E^{II}A^{II}} \\ -\frac{E^{II}A^{II}}{l^{II}} & \underline{E^{II}A^{II}} \\ 0 & -\frac{E^{III}A^{III}}{l^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} \\ 0 & -\frac{E^{III}A^{III}}{l^{III}} & \underline{E^{III}A^{III}} \\ 0 & -\frac{E^{III}A^{III}}{l^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} \\ 0 & -\frac{E^{III}A^{III}}{l^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{IIII}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{IIII}A^{III}} & \underline{E^{III}A^{III}} & \underline{E^{III}A^{I$$

Approximation Methods



• Example (*cont*):



- Step 7: solution
- Step 8: post calculation

$$u = u_1 \phi_1 + u_2 \phi_2 \quad \Longrightarrow \quad \mathcal{E} = \frac{du}{dx} = u_1 \frac{d\phi_1}{dx} + u_2 \frac{d\phi_2}{dx} \quad \Longrightarrow \quad O = E\mathcal{E} = Eu_1 \frac{d\phi_1}{dx} + Eu_2 \frac{d\phi_2}{dx}$$

Summary - Major Steps in FEM

- Discretization
- Derivation of element equation
 - weak form
 - construct form of approximation solution over oneelement
 - derive finite element model
- Assembling putting elements together
- Imposing boundary conditions
- Solving equations
- Post computation

Linear Formulation for Bar Element





Higher Order Formulation for Bar Element



 $u(x) = u_1 \phi_1(x) + u_2 \phi_2(x) + u_3 \phi_3(x)$







 $u(x) = u_1 \phi_1(x) + u_2 \phi_2(x) + u_3 \phi_3(x) + u_4 \phi_4(x) + \bullet \bullet \bullet \bullet \bullet + u_n \phi_n(x)$

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Natural Coordinates and Interpolation Functions



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Quadratic Formulation for Bar Element

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$$\begin{cases} P_{I} \\ P_$$



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Quadratic Formulation for Bar Element

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$$u(\xi) = u \phi_{1}(\xi) + u \phi_{2}(\xi) + u \phi_{3}(\xi) = u \frac{\xi(\xi-1)}{2} - u_{2}(\xi+1)(\xi-1) + u \frac{\xi(\xi+1)\xi}{2}$$



	$x - \frac{x_1 + x_2}{2}$	1	<u>d </u>	2
ξ =	$\frac{\lceil 2 \rceil}{l / 2}$	$\frac{1}{2} d\zeta = dx$	= dx	l

 $\frac{d\phi_1}{dx} = \frac{2}{l} \frac{d\phi_1}{d\xi} = \frac{2\xi - 1}{l}, \quad \frac{d\phi_2}{dx} = \frac{2}{l} \frac{d\phi_2}{d\xi} = -\frac{4\xi}{l}, \quad \frac{d\phi_3}{dx} = \frac{2}{l} \frac{d\phi_3}{d\xi} = \frac{2\xi + 1}{l}$



UNIT-II ANALYSIS OF TRUSSES AND BEAMS

INTRODUCTION

Finite Element Analysis of Trusses:

- Stiffness equations for a truss bar element oriented in 2D plane
- Finite Element Analysis of Trusses
- Plane Truss and Space Truss elements
- Methods of assembly

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Arbitrarily Oriented 1-D Bar

Element on 2-D Plane



Relationship Between Local Coordinates and Global Coordinates



$$\begin{cases} | \mathbf{u}_{1} \\ \mathbf{v} = 0 \\ | \mathbf{v} = 0 \\ | \mathbf{v} = 1 \\ | \mathbf{v} = 0 \\ | \mathbf{v} = 1 \\ | \mathbf{v} = 0 \\ | \mathbf$$

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Relationship Between Local Coordinates and Global Coordinates



Stiffness Matrix of 1-D Bar Element on 2-D Plane









 α , β , γ are the Direction Cosines of the bar in the x-y-z coordinate system

$$\begin{bmatrix} \underline{u} \\ 1 \\ 1 \\ \frac{1}{2} \end{bmatrix} \begin{bmatrix} \alpha & \beta & \gamma & 0 & 0 & 0 \\ x & \overline{x} & \overline{x} & x \\ w = 0 \end{bmatrix} \begin{bmatrix} \alpha & \beta & \gamma & 0 & 0 & 0 \\ \alpha & \beta & \gamma & 0 & 0 & 0 \\ w = 0 \end{bmatrix} \begin{bmatrix} \alpha & \beta & \gamma & 0 & 0 & 0 \\ \alpha & \beta & \gamma & 0 & 0 & 0 \\ \alpha & \beta & \gamma & 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} u \\ 1 \\ 1 \\ w \end{bmatrix}$$

$$\begin{bmatrix} \underline{P}_{1} \\ 0 \end{bmatrix} \begin{bmatrix} \alpha & \beta & \gamma & 0 & 0 & 0 \\ \alpha & \beta & \gamma & 0 & 0 & 0 \\ w \end{bmatrix} \begin{bmatrix} 1 \\ R_{1} \end{bmatrix}$$

$$\begin{bmatrix} 1 \\ R_{2} \end{bmatrix} \begin{bmatrix} 0 \\ 0 & 0 & 0 & 0 \\ \alpha & \beta & \gamma \end{bmatrix} \begin{bmatrix} 0 \\ R_{1} \end{bmatrix}$$

$$\begin{bmatrix} 1 \\ R_{2} \end{bmatrix} \begin{bmatrix} 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 1 \\ R_{2} \end{bmatrix}$$

$$\begin{bmatrix} 0 \\ 0 \end{bmatrix} \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} 0 \\$$



 $\begin{cases} P_{I} \\ Q = 0 \\ \hline R_{I} = 0 \\ \hline Q_{2} = 0 \\ \hline Q_{2} = 0 \\ \hline R_{2} = 0 \end{cases} = \frac{AE}{L} \begin{vmatrix} 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 & 0 \\ \hline 0 & 0 & 0 \\$ $\overline{P}_2, \overline{u}_2$ P_1, \overline{u}_1 $AE \begin{bmatrix} \alpha_{\overline{x}}^{2} & \alpha_{\overline{y}} & \alpha_{\overline{y}} & -\alpha_{\overline{x}}^{2} & -\alpha_{\overline{y}} & -\alpha_{\overline{x}} & x_{\overline{x}} \\ \alpha_{x}\beta_{x} & \beta_{x}^{2} & \beta_{x}\gamma_{x} & -\alpha_{\overline{x}}\beta_{\overline{x}} & -\beta_{\overline{x}}^{2} & -\beta_{\overline{x}}\gamma_{x} \\ \alpha_{x}\gamma_{x} & \beta_{x}\gamma_{x} & \gamma_{x} & -\alpha_{\overline{x}}\gamma_{x} & -\alpha_{\overline{x}}\beta_{\overline{x}} & -\beta_{\overline{x}}^{2} & -\beta_{\overline{x}}\gamma_{x} \\ \alpha_{x}\gamma_{x} & \beta_{x}\gamma_{x} & \gamma_{x} & \gamma_{x} & -\alpha_{\overline{x}}\gamma_{x} & \alpha_{\overline{x}}\gamma_{x} & -\beta_{\overline{x}}\gamma_{x} & \gamma_{\overline{y}} \end{bmatrix} \begin{bmatrix} u_{1} \\ v_{1} \\ v_{2} \end{bmatrix} \begin{bmatrix} AE \\ \alpha_{x}\gamma_{x} & \beta_{x}\gamma_{x} & -\beta_{x}\gamma_{x} & -\alpha_{x}\gamma_{x} & \alpha_{x}\gamma_{x} \\ -\alpha_{x}\gamma_{x} & \beta_{x}\gamma_{x} & -\beta_{\overline{x}}\gamma_{x} & -\alpha_{x}\gamma_{x} & \alpha_{x}\gamma_{x} \\ -\alpha_{x}\gamma_{x} & -\beta_{x}\gamma_{x} & -\beta_{x}\gamma_{x} & -\gamma_{x}^{2} & \alpha_{x}\gamma_{x} & \beta_{x}\gamma_{x} \\ -\alpha_{x}\gamma_{x} & -\beta_{x}\gamma_{x} & -\gamma_{x}^{2} & \alpha_{x}\gamma_{x} & \beta_{x}\gamma_{x} \\ -\alpha_{x}\gamma_{x} & -\beta_{x}\gamma_{x} & -\gamma_{x}^{2} & \alpha_{x}\gamma_{x} & \beta_{x}\gamma_{x} \\ -\gamma_{x}^{2} \end{bmatrix} \begin{bmatrix} v_{2} \\ v_{2} \\ v_{2} \end{bmatrix}$ P_1 $egin{array}{c|c} \mathcal{Q}_1 \ \mathcal{R}_1 \ \mathcal{R}_2 \ \mathcal{P}_2 \end{array}$

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Matrix Assembly of Multiple Bar Elements

• Element I

• Element II

Ilement III

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Matrix Assembly of Multiple Bar Elements

Iement I

Ilement II

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$$\begin{cases} P_{1} \\ Q_{1} \\ P_{2} \\ P_{3} \\ Q_{2} \\ P_{3} \\ Q_{3} \\ Q_{4} \\ Q_{2} \\ P_{3} \\ Q_{3} \\ Q_{3} \\ Q_{3} \\ Q_{4} \\ Q_{4} \\ Q_{2} \\ P_{3} \\ Q_{4} \\ Q_{4} \\ Q_{2} \\ P_{3} \\ Q_{3} \\ Q_{4} \\ Q_{2} \\ P_{3} \\ Q_{3} \\ Q_{4} \\ Q_{4} \\ Q_{4} \\ Q_{2} \\ P_{3} \\ Q_{4} \\ Q_{4} \\ Q_{4} \\ Q_{2} \\ P_{3} \\ Q_{4} \\ Q$$

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Matrix Assembly of Multiple Bar Elements



Apply known boundary conditions



Solution Procedures

$$\rightarrow$$
 $u_2 = 4FL/5AE, v_1 = 0$



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Recovery of axial forces

Element (I

Element

Element III

$$\begin{bmatrix} P_{1} \\ Q_{1} \\ P_{2} \\ Q_{2} \end{bmatrix} = \underbrace{AE} \begin{bmatrix} 1 & 0 & -1 & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} u_{1} = 0 \\ v_{1} = 0 \\ = \underbrace{AFL} \end{bmatrix} = F \begin{cases} -4 \\ 0 \end{bmatrix}$$

$$\begin{bmatrix} -4 \\ 0 \end{bmatrix} \begin{bmatrix} -4 \\ 0 \end{bmatrix} \begin{bmatrix} -4 \\ 0 \end{bmatrix} \\ = \underbrace{AFL} \end{bmatrix}$$

$$\begin{bmatrix} -4 \\ 0 \end{bmatrix} \begin{bmatrix} -4 \\ 0 \end{bmatrix} \begin{bmatrix} -4 \\ 0 \end{bmatrix} \end{bmatrix}$$

$$\begin{bmatrix} -4 \\ 0 \end{bmatrix} \begin{bmatrix} -4 \\ 0 \end{bmatrix} \begin{bmatrix} -4 \\ 0 \end{bmatrix} \end{bmatrix}$$

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$$\begin{bmatrix} -4 \\ 0 \end{bmatrix} \begin{bmatrix} -4 \\ 0 \end{bmatrix} \end{bmatrix}$$

$$\begin{bmatrix} -4 \\ 0 \end{bmatrix} \end{bmatrix}$$

 $\begin{cases} P_{1} \\ Q_{1} \\ P_{3} \\ P_{3} \\ |Q_{3}| \\$



Stresses inside members





Finite Element Analysis of Beams

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- Hermite shape functions
- Element stiffness matrix
- Load vector
- Problems

Bending Beam





Pure bending problems

- > Normal strain: $\varepsilon_x = -\frac{y}{\rho}$ > Normal stress: $o_x = -\frac{Ey}{\rho}$
- > Normal stress with bending moment: $\int -\sigma_x y dA = M$

 $O_x = -\frac{My}{I} \qquad I = \int y^2 dA$

 $\frac{1}{\rho} = \frac{M}{EI} \qquad \longrightarrow \qquad M = EI \frac{1}{\rho} \approx EI \frac{d^2 y}{dx^2}$

- Moment-curvature relationship:
- \succ Flexure formula:

Bending Beam





Relationship between shear force, bending moment and transverse load:



Governing Equation and Boundary Condition

Governing Equation $\frac{d^2}{dx^2} \left(EI \frac{d^2 v(x)}{dx^2} \right) - q(x) = 0,$ 0<x<L

Boundary Conditions

$$v = ? \& \frac{dv}{dx} = ? \& EI \frac{d v}{dx^2} = ? \& \frac{d}{dx} \left(EI \frac{d^2 v}{dx^2} \right) = ?, \quad \text{at } x = 0$$

$$v = ? \& \frac{dv}{dx} = ? \& EI \frac{d^2v}{dx^2} = ? \& \frac{d}{dx} \left(EI \frac{d^2v}{dx^2} \right) \models ?, \quad \text{at } x = L \qquad \frac{dv}{dx}$$

Essential BCs – if v or is specified at the boundary. Natural BCs – if v or is specified at the boundary. $EI \frac{d^2v}{dx^2}$ $\begin{pmatrix} d \\ EI \\ dx^2 \end{pmatrix} \begin{vmatrix} d \\ dx \end{vmatrix} \begin{pmatrix} d^2v \\ dx^2 \end{pmatrix}$



Weak Form from Integration-by-Parts ----- (2nd time)

$$0 = \int_{x_1}^{x_2} \left[\frac{d^2 w}{dx^2} \left(\frac{d^2 v}{EI \, dx^2} \right) - wq \left| \frac{dx}{dx} + w \frac{d}{dx} \left(\frac{d^2 v}{EI \, dx^2} \right) \right|_{x_1}^{x_2} - \frac{dw}{dx} \left(\frac{d^2 v}{EI \, dx^2} \right) \right|_{x_1}^{x_2}$$



$$0 = \int_{x_1}^{x_2} \left[\frac{d^2 w}{dx^2} \left(EI \frac{d^2 v}{dx^2} \right) - wq \right] dx + \left[wV - \frac{dw}{dx} M \right]_{x_1}^{x_2}$$



$$0 = \int_{x_1}^{x_2} \begin{bmatrix} d^2 w \begin{pmatrix} d^2 v \\ dx^2 \end{pmatrix} \\ dx^2 \end{bmatrix} = Wq \begin{bmatrix} dw \\ dx + \begin{bmatrix} wV - \frac{dw}{dx} \end{bmatrix}_{x_1}^{x^2}$$



$$Q_1 = V(x_1), \quad Q_2 = -M(x_1), \quad Q_3 = -V(x_2), \quad Q_4 = M(x_2)$$

$$\int_{x_{1}}^{x_{2}} \left[\frac{d^{2}w}{dx^{2}} \left(\frac{d^{2}v}{EI_{-}dx^{2}} \right) - wq \right] dx = w(x_{1})Q_{1} + w(x_{2})Q_{3} + \frac{dw}{dx} \left| \frac{Q_{2} + \frac{dw}{dx}}{dx} \right|_{2} Q_{4}$$

Ritz Method for Approximation









where
$$K = \int_{x_1}^{x_2} E \left[\frac{d^2 \phi}{d \phi} \frac{d^2 \phi}{d \phi} \right]_{x_1} dx$$
 and $q = x^2 \phi q dx$
 $i \int_{x_1}^{x_2} dx^2 dx^2 dx^2 dx^2$

Ritz Method for Approximation



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Selection of Shape Function

$$\begin{bmatrix} \left(1 \mid x\right) \begin{pmatrix} d\phi_1 \\ dx_1 \end{pmatrix} \\ \left(\frac{d\phi_2}{dx_1} \right) \begin{pmatrix} d\phi_2 \\ dx_2 \end{pmatrix} \\ \left(\frac{d\phi_3}{dx_1} \right) \begin{pmatrix} d\phi_3 \\ dx_2 \end{pmatrix} \\ \left(\frac{d\phi_3}{dx_1} \right) \begin{pmatrix} d\phi_3 \\ dx_2 \end{pmatrix} \\ \left(\frac{d\phi_4}{dx_1} \right) \begin{pmatrix} d\phi_4 \\ dx_2 \end{pmatrix} \\ \left(\frac{d\phi_4}{dx_1} \right) \begin{pmatrix} d\phi_4 \\ dx_2 \end{pmatrix} \\ \left(\frac{d\phi_4}{dx_1} \right) \end{pmatrix}$$

$$\begin{bmatrix} Q_{1} \\ Q_{2} \\ Q_{2} \\ Q_{3} \\ Q_{3} \\ Q_{4} \end{bmatrix} = \begin{bmatrix} K_{11} & K_{12} & K_{13} & K_{14} \end{bmatrix} \begin{bmatrix} u_{1} \\ u_{1} \\ W_{2} \\ K_{12} & K_{22} & K_{23} & K_{24} \\ K_{13} & K_{23} & K_{33} & K_{34} \\ K_{14} & K_{24} & K_{34} & K_{44} \end{bmatrix} \begin{bmatrix} u_{2} \\ u_{2} \\ u_{2} \\ u_{3} \\ u_{3} \\ u_{4} \end{bmatrix} = \begin{bmatrix} Q_{1} \\ Q_{2} \\ Q_{2} \\ Q_{3} \\ Q_{4} \end{bmatrix}$$

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Derivation of Shape Function for Beam Element Local Coordinates

$$v\left(\xi\right) = \tilde{u}_{1}\phi_{1} + \tilde{u}_{2}\phi_{2} + \tilde{u}_{3}\phi_{3} + \tilde{u}_{4}\phi_{4}$$

and
$$\frac{dv(\xi)}{d\xi} = \tilde{u}_{1}\frac{d\phi_{1}}{d\xi} + \tilde{u}_{2}\frac{d\phi_{2}}{d\xi} + \tilde{u}_{3}\frac{d\phi_{3}}{d\xi} + \tilde{u}_{4}\frac{d\phi_{4}}{d\xi}$$

where
$$\tilde{u}_{1} = v_{1}\tilde{u}_{2} = \frac{dv_{1}}{d\xi}\tilde{u}_{3} = v_{2}\tilde{u}_{4} = \frac{dv_{2}}{d\xi}$$

Let
$$\phi_{i} = a_{i} + b_{i}\xi + c_{i}\xi^{2} + d_{i}\xi^{3}$$

Find coefficients to satisfy the interpolation properties.

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In the global coordinates:

$$v(x) = v \phi_{1 \ 1}(x) + \frac{l}{2} \frac{dv_1}{dx} \phi_2(x) + v \phi_{2 \ 3}(x) + \frac{l}{2} \frac{dv_2}{dx} \phi_4(x)$$

$$\begin{cases} \left| 1 - \frac{3}{2} \frac{x}{1} + \frac{2}{2} \right| \left| \frac{x - x}{x_2 - x_2} \right|^3 \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_2} + \frac{x - x}{x_2 - x_2} \right| \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{x - x}{x_2 - x_2} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x_1} \\ \frac{2}{x_2 - x_1} + \frac{2}{x_2 - x$$

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Element Equations of 4th Order 1-D Model





Element Equations of 4th Order 1-D Model





where $q_i = \int_{x_1}^{x_2} \phi_i q dx$

Finite Element Analysis of 1-D Problems - Applications Example 1. L L L F

Governing equation:

$$\frac{d^2}{dx^2} \left(\frac{EI}{dx^2} dx^2 \right) - q(x) = 0 \quad 0 < x < L$$

Weak form for one element

where
$$\int_{x_1}^{x_2} \left(EI \frac{d^2 w}{dx^2} \frac{d^2 v}{dx^2} - wq \right) dx - w(x_1) Q_1 - \frac{dw}{dx} \Big|_{x_1} Q_2 - w(x_2) Q_3 - \frac{dw}{dx} \Big|_{x_2} Q_4 = 0$$
$$Q_1 = V(x_1) \qquad Q_2 = -M(x_1) \qquad Q_3 = -V(x_2) \qquad Q_4 = M(x_2)$$

Finite Element Analysis of 1-D Problems



• Approximation function: $v(x) = v \phi(x) + \frac{l}{dv_1} \frac{dv_1}{\phi(x) + v \phi(x)} + \frac{dv_2}{\phi(x)} \phi(x)$

$$\begin{array}{c} 11 & \overline{2} \ \overline{dx} \ 2 & 23 & \overline{2} \ \overline{dx} \ 4 \\ \left(\begin{array}{c} \left(\begin{array}{c} x - x \end{array}\right)^2 \\ 1 - 3 \end{array}\right)^2 \\ \left(\begin{array}{c} x - x \end{array}\right)^2 \\ \left(\begin{array}{c} x - x \end{array}\right)^2 \\ \left(\begin{array}{c} x - x \end{array}\right)^3 \\ \left(\begin{array}{c} x - x \end{array}\right)^3 \\ \left(\begin{array}{c} x - x \end{array}\right)^2 \\ \left(\begin{array}{c} x - x \end{array}\right)^3 \\ \left(\begin{array}{c} x - x \end{array}\right)^2 \\ \left(\begin{array}{c} x - x \end{array}\right)^3 \\ \left(\begin{array}{c} x$$



Finite element model:

$$\begin{cases} Q_{1} \\ Q_{2} \\ Q_{3} \\ Q_{3} \\ Q_{3} \\ Q_{4} \\ Q_$$

Discretization:

$$P_{1}, v_{1} \qquad \boxed{1} \qquad P_{2}, v_{2} \qquad \boxed{1} \qquad P_{3}, v_{3} \qquad \boxed{1} \qquad P_{4}, v_{4}$$
$$M_{1}, \theta_{1} \qquad M_{2}, \theta_{2} \qquad M_{3}, \theta_{3} \qquad M_{4}, \theta_{4}$$

Matrix Assembly of Multiple Beam Elements

		$\begin{vmatrix} Q^I \\ Q^I \end{vmatrix}$	6	3 <i>L</i>	-6	3 <i>L</i>	0	0	0	$0 \left[\begin{pmatrix} v_1 \\ \theta \end{pmatrix} \right]$
Element	1	$\begin{vmatrix} Q_I^I \\ Q_I^I \end{vmatrix}$	$\begin{vmatrix} 3L \\ -6 \\ 3I \end{vmatrix}$	$2L^{2}$ -3L L^{2}	-3L 6	L^2 -3L $2L^2$	$\begin{pmatrix} 0\\ 0\\ 0 \end{pmatrix}$	$\begin{array}{c} 0\\ 0\\ 0\\ \end{array}$	$\begin{array}{c} 0\\ 0\\ 0 \end{array}$	$\begin{array}{c} 0 \\ 0 \\ v_2^1 \\ 0 \\ v_2^2 \\ \end{array}$
		$\begin{cases} \begin{array}{c} 4 \\ 0 \\ \end{array} \\ = \begin{array}{c} 2 \\ 1 \\ 1 \\ \end{array} \\ \begin{array}{c} 2 \\ 1 \\ 3 \\ \end{array} \\ \begin{array}{c} 2 \\ 1 \\ 3 \\ \end{array} \\ \begin{array}{c} 2 \\ 1 \\ 3 \\ \end{array} \\ \begin{array}{c} 2 \\ 1 \\ 1 \\ 1 \\ \end{array} \\ \begin{array}{c} 2 \\ 1 \\ 1 \\ 1 \\ 1 \\ \end{array} \\ \begin{array}{c} 2 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\$	$\left\{ \begin{array}{c} J \\ J \\ J \\ J \\ J \end{array} \right\}$	L	-3L	$\frac{2L}{0}$	0	0	0	$0 \theta $ $0 v_3 $
			0	0	0	0	0	0	0	$0 \begin{vmatrix} \theta \\ \theta \end{vmatrix} = \begin{bmatrix} \theta \\ \theta \end{bmatrix}$
			$\begin{vmatrix} 0\\0 \end{vmatrix}$	0 0	0 0	0 0	0 0	0 0	$\begin{array}{c} 0 \\ 0 \end{array}$	$\begin{bmatrix} 0 \\ 0 \end{bmatrix} \begin{bmatrix} v \\ \theta^{\dagger} \end{bmatrix}$
	:	Ĺ	L							
Element	II	$\begin{bmatrix} 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0\\0 \end{bmatrix}$	$\begin{array}{c} 0 \\ 0 \\ 0 \end{array} \begin{array}{c} 0 \\ 0 \end{array}$	$\begin{array}{c} 0\\ 0\end{array}$	0 0		$\begin{array}{c} 0\\ 0\end{array}$		$\stackrel{0}{0}_{0} \stackrel{ v_{1}}{\theta}$
		$\left Q_{1}^{II} ight $		0^{1}	3 <i>L</i>	-6		3L	0	0 v2
		$ \langle Q^{I} \rangle \geq 2EI$		0 3L	$2L^2$	-3L		L^2	0	$0\left \left \left<\right \boldsymbol{\theta}\right>\right $
		$\begin{bmatrix} Q^{3II} \\ Q^{II} \end{bmatrix} = \begin{bmatrix} -L^3 \\ 0 \end{bmatrix}$	$\begin{vmatrix} 3 - \\ 0 \\ 0 \end{vmatrix}$	$\begin{array}{c} 0 \stackrel{\text{\tiny def}}{=} 6 \\ 0 \stackrel{\text{\tiny def}}{=} 3L \end{array}$	-3L L^2	6 -3L		-3L $2L^2$		$\begin{array}{c c}0 & \nu_3\\0 & \theta\end{array}$
		$\begin{vmatrix} \mathbf{\hat{c}} \\ 0 \end{vmatrix}$	$\overline{0}$	0 0	0	0		0	0	$0 v_4 $
		[0]]	0	0 0	0	0		0		$0 \left \left[\theta_{4} \right] \right $

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Matrix Assembly of Multiple Beam Elements



2 0 0 0

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Solution Procedures



Apply known boundary conditions

$$\begin{bmatrix} P_{1}=? \\ M = ? \\ M = ? \\ M = 0 \\ P_{2}=? \\ P_{2}=? \\ P_{4}=-F \\ M_{4}=0 \end{bmatrix} = \begin{bmatrix} 6 & 3L & -6 & 3L & L^{2} \\ 3L & 2L^{2} & -3L & L^{2} \\ 0 & 0 & 4L^{2} \\ 0 & 0 & 4L^{2} \\ 0 & 0 & 0 \\ P_{4}=-F \\ M_{4}=0 \end{bmatrix} = \begin{bmatrix} 0 & 0 & 3L & L^{2} \\ 0 & 0 & 3L & L^{2} \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 \\ P_{4}=-F \\ M_{4}=0 \end{bmatrix} = \begin{bmatrix} 3L & L^{2} & 0 & 4L^{2} \\ 0 & 0 & 3L & L^{2} \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 \\ P_{4}=-F \\ M_{4}=0 \end{bmatrix} = \begin{bmatrix} 3L & L^{2} & 0 & 4L^{2} \\ 0 & 0 & 3L & L^{2} \\ 0 & 0 & 0 \\ P_{4}=-F \\ 0 & 0 & 0 \\ P_{4}=-F \\ P_{4}=-F$$

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Solution Procedures

$$\begin{bmatrix} M_{2}=0 \\ M=0 \end{bmatrix} \begin{bmatrix} 3L & L^{2} & 0 & -3L & 4L^{2} & L^{2} & 0 & 0 \\ 0 & 0 & 3L & 0 & L^{2} & 4L^{2} & -3L & L^{2} \\ M_{2}=-F \\ M=0 \end{bmatrix} \begin{bmatrix} 0 & 0 & 0 & -6 & 0 & -3L & 6 & -3L \\ 0 & 0 & 0 & 3L & 0 & L^{2} & -3L & 2L^{2} \\ 0 & 0 & 0 & 3L & 0 & L^{2} & -3L & 2L^{2} \\ M_{2}=2 \\ M_{3}=2 \end{bmatrix} \begin{bmatrix} \frac{L^{3}}{6} \end{bmatrix} \begin{bmatrix} \frac{d}{6} & 3L & -6 & 0 & 3L & 0 & 0 \\ -6 & 3L & 12 & -6 & 0 & 3L & 0 & 0 \\ -6 & -3L & 12 & -6 & 0 & 3L & 0 & 0 \\ 0 & 0 & -6 & -12 \\ -3L & 0 & -6 & 3L \end{bmatrix} \begin{bmatrix} \theta_{2}=2 \\ \theta_{3}=2 \\ \theta_{4}=2 \end{bmatrix}$$

$$\begin{cases} M_{2} = 0 \\ M_{3} = 0 \\ P_{4}^{3} = -F \\ M_{4}^{3} = -F \\ M_{4}^{3} = 0 \\ M_{4}^{3} = 0 \\ M_{4}^{3} = -F \\ M_{4}^{3} = -F$$

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Shear Resultant & Bending







UNIT-III 2D ANALYSISENTS

Introduction



- Computation of shape functions for constant straintriangle
- Properties of the shape functions
- Computation of strain-displacement matrix
- Computation of element stiffness matrix
- Computation of nodal loads due to body forces
- Computation of nodal loads due to traction
- Recommendations for use
- Example problems

 Divide the body into connected to each other through special points ("nodes")



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 $u(x,y) \approx N_1(x,y)u_1 + N_2(x,y)u_2 + N_3(x,y)u_3 + N_4(x,y)u_4$ $v(x,y) \approx N_1(x,y)v_1 + N_2(x,y)v_2 + N_3(x,y)v_3 + N_4(x,y)v_4$ $\underline{\mathbf{u}} = \begin{cases} \mathbf{u} & (\mathbf{x}, \mathbf{y}) \\ \begin{pmatrix} \mathbf{u} \\ \mathbf{v} \\ \mathbf{x}, \mathbf{y} \end{pmatrix} = \begin{bmatrix} \mathbf{N}_{1} & \mathbf{0} & \mathbf{N}^{2} & \mathbf{0} & \mathbf{N}^{3} & \mathbf{0} & \mathbf{N}^{4} & \mathbf{0} \\ \\ \begin{pmatrix} \mathbf{v} & (\mathbf{x}, \mathbf{y}) \end{pmatrix} & \begin{bmatrix} \mathbf{0} & \mathbf{N}_{1} & \mathbf{0} & \mathbf{N}_{2} & \mathbf{0} & \mathbf{N}_{3} & \mathbf{0} & \mathbf{N}_{4} \end{bmatrix} \begin{cases} \mathbf{u}_{1} \\ \mathbf{u}_{2} \\ \mathbf{v}^{2} \\ \mathbf{u}_{3} \\ \mathbf{v}_{2} \end{cases}$ $= \mathbf{N} \mathbf{d}$ $\underline{\mathbf{u}} = \underline{\mathbf{N}} \underline{\mathbf{d}}$



TASK 2: APPROXIMATE THE STRAIN and STRESS WITHIN EACH ELEMENT

Approximation of the strain in element 'e'





$$\underline{\mathcal{E}} = \begin{cases} \left| \begin{array}{c} \mathcal{E}_{x} \\ \mathcal{E}_{y} \\ \\ \left| \begin{array}{c} \mathcal{K}_{y} \end{array} \right| \\ \\ \left| \begin{array}{c} \mathcal{K}_{y} \end{array} \right| \\ \\ \end{array} \right|$$



 $\underline{\varepsilon} = \underline{B} \underline{d}$

• Displacement approximation in terms of shape functions

 $\underline{u} = \underline{N} \underline{d}$

• Strain approximation in terms of strain-displacement

$$\underline{\varepsilon} = \underline{B} \underline{d}$$

- Stress approximation $O = \underline{DBd}$
- Element stiffness matrix matrix

$$\underline{k} = \int_{V^e} \underline{\mathbf{B}}^{\mathrm{T}} \underline{\mathbf{D}} \underline{\mathbf{B}} \mathrm{d} \mathbf{V}$$

 $f = \int_{\underbrace{e_V}^{e_V}} \underbrace{\underline{X}}_{V} dV + \int_{\underbrace{S_T}^{e}} \underbrace{\underline{N}}_{T} \underbrace{\underline{X}}_{S} dS$

• Element nodal load vector





Constant Strain Triangle (CST): Simplest 2D finite element



- In 3 nodes per element
- 2 dofs per node (each node can move in x- and y- directions)
- Hence 6 dofs per element



> The displacement approximation in terms of shape functionsis

$$\mathbf{u} \ (\mathbf{x}, \mathbf{y}) \approx N_{1} \mathbf{u}_{1} + N_{2} \mathbf{u}_{2} + N_{3} \mathbf{u}_{3}$$

$$\mathbf{v}(\mathbf{x}, \mathbf{y}) \approx N_{1} \mathbf{v}_{1} + N_{2} \mathbf{v}_{2} + N_{3} \mathbf{v}_{3}$$

$$\overset{\mathbf{u}}{=} \begin{cases} \mathbf{u} \ (\mathbf{x}, \mathbf{y}) \\ \mathbf{v} \ (\mathbf{x}, \mathbf{y}) \end{cases} = \begin{bmatrix} N_{1} & 0 & N_{2} & 0 & N_{3} & 0 \\ 0 & N_{1} & 0 & N_{2} & 0 & N_{3} \end{bmatrix} \begin{vmatrix} \mathbf{u}_{1} \\ \mathbf{v}_{2} \\ \mathbf{v}_{2} \\ \mathbf{u}_{3} \\ \mathbf{v}_{3} \end{vmatrix}$$

$$\mathbf{u}_{2 \times 1} = \mathbf{N}_{2 \times 6} \mathbf{d}_{6 \times 1}$$

$$\underbrace{\mathbf{N}}_{1} = \begin{bmatrix} N_{1} & 0 & N_{2} & 0 & N_{3} & 0 \\ 0 & N_{1} & 0 & N_{2} & 0 & N_{3} \end{bmatrix}$$




• The shape functions N_1 , N_2 and N_3 are linear functions of x and y



$$N_{i} = \begin{cases} 1 & at node 'i' \\ 0 & at other nodes \end{cases}$$

2 0 0 0

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• At every point in the domain

$$\sum_{i=1}^{3} N_{i} = 1$$

$$\sum_{i=1}^{3} N_{i} x_{i} = x$$

$$\sum_{i=1}^{3} N_{i} y_{i} = y$$



 Geometric interpretation of the shape functions, at any point P(x,y) that the shape functions are evaluated





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• Approximation of the strains

$$\underline{\mathcal{E}} = \left\{ \begin{bmatrix} \mathcal{E}_{x} \\ y \\ \gamma_{xy} \end{bmatrix} \right\} = \left\{ \begin{array}{c} \partial u \\ \partial \overline{\partial x} \\ \partial \overline{\partial y} \\ \partial \overline{y} \\ \partial \overline{y} \\ \partial \overline{y} \\ \partial \overline{y} + \partial \overline{x} \\ \partial \overline{y} \\ \partial$$

$$\underline{B} = \begin{bmatrix} \partial N_{1}(x,y) & 0 & \frac{\partial N_{2}(x,y)}{\partial x} & 0 & \frac{\partial N_{3}(x,y)}{\partial x} & 0 \\ \hline \partial P & \frac{\partial N_{1}(x,y)}{\partial y} & 0 & \frac{\partial P}{\partial Y} & \frac{\partial P}{\partial Y}(x,y) \\ \partial P & \frac{\partial P}{\partial Y} & \frac{\partial P}{\partial Y}(x,y) & \frac{\partial P}{\partial Y}(x,y) & \frac{\partial P}{\partial Y}(x,y) & \frac{\partial P}{\partial Y}(x,y) \\ \hline P & \frac{P}{\partial Y} & \frac{1}{\partial x} & \frac{2}{\partial y} & \frac{2}{\partial x} & \frac{3}{\partial y} & \frac{3}{\partial x} \end{bmatrix} \\
= \frac{1 \begin{bmatrix} b_{1} & 0 & b_{2} & 0 & b_{3} & 0 \end{bmatrix}}{2A \begin{bmatrix} 1 & 2 & 3 \\ 0 & c & 0 & c & 0 & c \end{bmatrix}} \\
= \frac{1 \begin{bmatrix} b_{1} & 0 & b_{2} & 0 & b_{3} & 0 \end{bmatrix}}{2A \begin{bmatrix} 1 & 2 & 3 \\ 0 & c & 0 & c & 0 & c \end{bmatrix}}$$



- Inside each element, all components of strain are constant: hence the name Constant Strain Triangle.
- Element stresses (constant inside each element).

$O = \underline{DB} \underline{d}$



IMPORTANT NOTE:

- > The displacement field is continuous across element boundaries
- The strains and stresses are NOT continuous across element

boundaries

Element stiffness matrix

$$\underline{k} = \int_{V^e} \underline{B}^{\mathrm{T}} \underline{D} \, \underline{B} \, \mathrm{dV}$$

Since \underline{B} is constant



$$\underline{k} = \underline{\mathbf{B}}^{\mathrm{T}} \underline{\mathbf{D}} \underline{\mathbf{B}} \int_{V^{e}} d\mathbf{V} = \underline{\mathbf{B}}^{\mathrm{T}} \underline{\mathbf{D}} \underline{\mathbf{B}} A t$$

t=thickness of the element A=surface area of the element

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Element nodal load vector

 $\underline{f} = \underbrace{\int_{e} \mathbf{N}^{T} \underline{X} \, dV}_{F_{b}} + \underbrace{\int_{S_{T}^{e}} \mathbf{N}^{T} T}_{S} \, dS}_{\underline{f}_{s}}$



Element nodal load vector due to body forces

$$\underline{f}_{b} = \int_{V^{e}} \underline{N}^{T} \underline{X} \quad dV = t \int_{A^{e}} \underline{N}^{T} \underline{X} \quad dA$$



$$f_{b} = \begin{cases} f_{b1x} \\ f_{b1y} \\ f_{b2x} \\ f_{b2x} \\ f_{b3x} \\ f_{b3y} \\ f_{b3y} \\ f_{b3y} \\ f_{b3y} \\ f_{b3y} \\ f_{b3y} \\ f_{b3x} \\ f_{b3y} \\ f_{b3x} \\ f_{b3y} \\ f_{b3x} \\ f_{b3y} \\ f_{b3x} \\ f_{b3y} \\ f_{by} \\$$



EXAMPLE: If X_a=1 and X_b=0

$$\mathcal{F}_{b} = \begin{cases} \left[f_{b1x} \right]_{f} & \left[t_{f} \right]_{e}^{N_{1}} X_{a} dA \right]_{A^{e}} N_{1} X_{b} dA \\ \left[t_{f} \right]_{e}^{N_{1}} X_{b} dA \\ f_{b2y} \\ f_{b2y} \\ f_{b3y} \\ f_{b3$$



Element nodal load vector due to traction

$$\underline{f}_{S} = \int_{S_T^e} \underline{\mathbf{N}}^T \underline{T}_S \, dS$$

EXAMPLE:



Element nodal load vector due to traction



Example

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$$f_{S} = t \int_{l_{2-3}^{e}} \underline{\mathbf{N}}^{T} \Big|_{along \, 2-3} \underline{T}_{S} \, dS$$

$$f_{S_{2x}} = t \int_{l_{2-3}^{e}} N_{2} \Big|_{along \, 2-3} \, (1) \, dy$$

$$= t \left(\frac{1}{2} \right) \times 2 \times 1 = t$$

— .

Similarly, compute

$$f_{S_{y}} = 0$$
$$f_{S_{x}} = t$$
$$f_{S_{3y}} = 0$$



- Use in areas where strain gradients are small
- Use in mesh transition areas (fine mesh to coarse mesh)
- Avoid CST in critical areas of structures (e.g., stress concentrations, edges of holes, corners)
- In general CSTs are not recommended for general analysis purposes as a very large number of these elements are required

for reasonable accuracy.



Example



Thickness (t) = 0.5 in E= 30×10⁶ psi n=0.25

(a) Compute the unknown nodal displacements.

(b) Compute the stresses in the two elements.



Realize that this is a plane stress problem and therefore we need to use

$$\underline{D} = \underbrace{1 - v^2}_{0} \begin{bmatrix} 1 & v & 0 \\ v & 1 & 0 \\ 0 & 0 & \frac{1 - v}{2} \end{bmatrix} = \begin{bmatrix} 3.2 & 0.8 & 0 \\ 0.8 & 3.2 & 0 \\ 0 & 0 & 1.2 \end{bmatrix} \times 10^7 \ psi$$

Step 1: Node-element connectivity chart

ELEMEN T	Node 1	Node 2	Node 3	Area (sqin)
1	1	2	4	3
2	3	4	2	3

Node	x	У
1	3	0
2	3	2
3	0	2
4	0	0

Nodal coordinates

<u>Step 2</u>: Compute strain-displacement matrices for the elements





Step 3: Compute element stiffness matrices

$$\underline{k}^{(1)} = At \underline{B}^{(1)^{T}} \underline{D} \underline{B}^{(1)} = (3)(0.5) \underline{B}^{(1)^{T}} \underline{D} \underline{B}^{(1)}$$

$$= \begin{bmatrix} 0.9833 & -0.5 & -0.45 & 0.2 & -0.5333 & 0.3 \\ 1.4 & 0.3 & -1.2 & 0.2 & -0.2 \\ 0.45 & 0 & 0 & -0.3 \\ 1.2 & -0.2 & 0 \\ 0.5333 & 0 \end{bmatrix} \times 10^{7}$$

$$= \begin{bmatrix} u_{1} & v_{1} & u_{2} & v_{2} & u_{4} & v_{4} \end{bmatrix}$$



$$\underline{k}^{(2)} = At\underline{\mathbf{B}}^{(2)^{\mathrm{T}}} \underline{\mathbf{D}} \underline{\mathbf{B}}^{(2)} = (3)(0.5)\underline{\mathbf{B}}^{(2)^{\mathrm{T}}} \underline{\mathbf{D}} \underline{\mathbf{B}}^{(2)}$$

$$= \begin{bmatrix} 0.9833 & -0.5 & -0.45 & 0.2 & -0.5333 & 0.3 \\ 1.4 & 0.3 & -1.2 & 0.2 & -0.2 \\ 0.45 & 0 & 0 & -0.3 \\ 1.2 & -0.2 & 0 \\ 0.5333 & 0 \end{bmatrix} \times 10^7$$

$$= \begin{bmatrix} u_3 & v_3 & u_4 & v_4 & u_2 & v_2 \end{bmatrix}$$



Step 4: Assemble the global stiffness matrix corresponding to the $u_3 = v_3 = u_4 = v_4 = v_1 = 0$ nonzero degrees of freedom

• Hence we need to calculate only a small (3x3) stiffness matrix





$$\begin{array}{c} f \ = \ \left| \begin{pmatrix} f_{1x} \\ f_{2x} \\ \end{pmatrix} \right| = \left\{ \begin{array}{c} 0 \\ 0 \\ \\ f_{2y} \end{array} \right\} \\ \end{array}$$

$$f_{2y} = -1000 + f_{s_{2y}}$$

The consistent nodal load due to traction on the edge 3-2

$$f_{s}^{2y} = \int_{x=0}^{3} N_{3} \Big|_{3=2} (-300) t dx$$

= $(-300)(0.5) \int_{x=0}^{3} N_{3} \Big|_{3=2} dx$
= $-150 \int_{3}^{x} \frac{x}{2} dx$
= $-50 \left[\begin{array}{c} x^{2} \\ 2 \end{array} \right]_{0}^{3} = -50 \left(\begin{array}{c} 9 \\ 2 \end{array} \right) = -225 \, lb$





Hence

$$f_{2y} = -1000 + f_{S_{2y}}$$

= -1225 *lb*

Step 6: Solve the system equations to obtain the unknown nodalloads $\underline{Kd} = f$

 $10^{7} \times \begin{vmatrix} \Box 0.983 & -0.45 & 0.2 \\ 0.983 & 0 & 1.4 \end{vmatrix} \begin{pmatrix} u_{1} \\ u_{2} \\ u_{2} \\ v_{2} \end{vmatrix} = \begin{cases} 0.2337 \times 10^{-4} in \\ 0.1069 \times 10^{-4} in \\ -0.9084 \times 10^{-4} in \end{vmatrix}$



<u>Step 7</u>: Compute the stresses in the elements

In Element #1

$$\boldsymbol{O}^{(1)} = \underline{\mathbf{D}}\underline{\mathbf{B}}^{(1)}\,\underline{\mathbf{d}}^{(1)}$$

With

$$\underline{\mathbf{d}}^{(1)^{T}} = \begin{bmatrix} u_{1} & v_{1} & u_{2} & v_{2} & u_{4} & v_{4} \end{bmatrix}$$
$$= \begin{bmatrix} 0.2337 \times 10^{-4} & 0 & 0.1069 \times 10^{-4} & -0.9084 \times 10^{-4} & 0 & 0 \end{bmatrix}$$

Calculate

$$o^{(1)} = \begin{bmatrix} -114.1 \\ -1391.1 \end{bmatrix} psi$$
$$\lfloor -76.1 \rfloor$$



In Element #2

$$\underline{O}^{(2)} = \underline{\mathbf{DB}}^{(2)} \underline{\mathbf{d}}^{(2)}$$

With

$$\underline{\mathbf{d}}^{(2)^{T}} = \begin{bmatrix} u_{3} & v_{3} & u_{4} & v_{4} & u_{2} & v_{2} \end{bmatrix}$$
$$= \begin{bmatrix} 0 & 0 & 0 & 0 & 0.1069 \times 10^{-4} & -0.9084 \times 10^{-4} \end{bmatrix}$$

Calculate

$$O^{(2)} = \begin{bmatrix} 114.1 \\ 28.52 \\ -363.35 \end{bmatrix} psi$$

Notice that the stresses are constant in each element

Axi-symmetric Problems

Definition:

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A problem in which geometry, loadings, boundary conditions and materials are symmetric about one axis.

Axi-symmetric Analysis

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$$x = r\cos\theta; \ y = r\sin\theta; \ z = z$$

- quantities depend on *r* and *z* only
- 3-D problem
- 2-D problem



Axi-symmetric Analysis





Axi-symmetric Analysis – Single-Variable Problem



$$-\frac{1}{r}\frac{\partial \left(ra}{\partial r}\left(\frac{\partial u(r,z)}{\partial r}\right) - \frac{\partial \left(ra}{\partial z}\left(\frac{\partial u(r,z)}{\partial z}\right) + a u - f(r,z) = 0$$

Weak form: $0 = \int \left[\frac{\partial w}{\partial r} \left(a_{11} \frac{\partial u}{\partial r} \right) + \frac{\partial w}{\partial z} \left(a_{22} \frac{\partial u}{\partial z} \right) + a_{00} \frac{wu - v}{wf(r, z)} \right] r dr dz$ $-\oint_{\Gamma} wq_n ds$

 $q_{n} = a_{11} \frac{\partial u(r, z)}{\partial r} n_{r} + a_{22} \frac{\partial u(r, z)}{\partial z} n_{z}$ where

Finite Element Model – Single-Variable Problem



$$u = \sum_{j} u_{j} \phi_{j} \quad \text{where} \quad \phi_{j}(r, z) = \phi_{j}(x, y)$$

Ritz method:
Weak form

$$\sum_{j=1}^{e} K^{e} u^{e} = f^{e} + Q^{e}$$

$$\sum_{j=1}^{e} \left(a_{11} \partial \phi_{i} \partial \phi_{j} + a_{22} \partial \phi_{i} \partial \phi_{j} + a_{00 \ i \ j} \right) r dr dz$$
where

$$K^{e}_{ij} = \int_{\Omega_{e}} \left(a_{11} \partial \phi_{i} \partial \phi_{j} + a_{22} \partial \phi_{i} \partial \phi_{j} + a_{00 \ i \ j} \right) r dr dz$$

$$f^{e}_{i} = \int_{\Omega_{e}} \phi_{i} fr dr dz$$

$$Q^{e}_{i} = \int \phi_{i} q_{n} ds$$

 Γ_e

Single-Variable Problem – Heat Transfer

Heat Transfer:

$$-\frac{1}{r}\frac{\partial}{\partial r}\left(rk \frac{\partial T(r,z)}{\partial r}\right) - \frac{\partial}{\partial z}\left(k \frac{\partial T(r,z)}{\partial z}\right) - f(r,z) = 0$$

Weak form

$$0 = \int \left[\left| \frac{\partial w}{\partial r} \left[k \frac{\partial T}{\partial r} \right] \right|_{+} \frac{\partial w}{\partial z} \left[\left[k \frac{\partial T}{\partial z} \right] \right]_{-\frac{\Omega_{e}}{\Gamma_{e}}} wq_{n} ds$$

where

$$q_{n} = k \frac{\partial T(r, z)}{\partial r} n_{r} + k \frac{\partial T(r, z)}{\partial z} n_{z}$$

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3-Node Axi-symmetric Element



$$T(r, z) = T_1 \phi_1 + T_2 \phi_2 + T_3 \phi_3$$





$$\phi_{2} = \underbrace{\left(\underline{1} \quad \underline{r} \quad \underline{z}\right)}_{2A_{e}} \begin{cases} r_{3}z_{1} - rz_{3} \\ z - z \\ \\ & r - r^{1} \\ \\ & 1 & 3 \end{cases}$$

$$\phi_{3} = \left(\frac{1 r z}{2A_{e}} \right) \begin{cases} r_{1}z_{2} - r_{z_{1}} \\ z - z \end{cases} \begin{cases} r_{1}z_{2} - r_{z_{1}} \\ r - r^{2} \\ z - z \end{cases}$$

4-Node Axi-symmetric Element



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Single-Variable Problem – Example





Step 1: Discretization

Step 2: Element equation

$$K_{ij}^{e} = \int_{\Omega_{e}} \left(\kappa \partial \phi \partial \phi_{j} + \kappa \partial \phi_{j} \partial \phi_{j} - \partial \phi_{j} \partial \phi_{j} - \partial \phi_{j} \partial \phi_{j} \right) r dr dz$$

$$f_i^{\ e} = \int_{\Omega_e} \phi_i \, fr dr dz \qquad Q_i^e = \int_{\Gamma_e} \phi_i q_n ds$$

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Constant Strain Triangle (CST) - easiest and simplest finite element

- Displacement field in terms of generalized coordinates
- $u=\beta_1+\beta_2x+\beta_3y$
- $\upsilon = \beta_4 + \beta_5 x + \beta_6 y$
 - Resulting strain field is

$$\varepsilon_x = \beta_2$$
 $\varepsilon_y = \beta_6$ $\gamma_{xy} = \beta_3 + \beta_5$

- Strains do not vary within the element. Hence, the name constant strain triangle (CST)
 - •Other elements are not so lucky.

•Can also be called linear triangle because displacement field is linear in x and y - sides remain straight.

Constant Strain Triangle



$$\begin{cases} \varepsilon_{x} \\ \varepsilon_{y} \\ \gamma_{xy} \end{cases} = \frac{1}{2A} \begin{bmatrix} y_{23} & 0 & y_{31} & 0 & y_{12} & 0 \\ 0 & x_{32} & 0 & x_{13} & 0 & x_{21} \\ x_{32} & y_{23} & x_{13} & y_{31} & x_{21} & y_{12} \end{bmatrix} \begin{cases} u_{1} \\ v_{1} \\ u_{2} \\ v_{2} \\ u_{3} \\ v_{3} \end{cases}$$

- Where, x_i and y_i are nodal coordinates (i=1, 2, 3)

$$- x_{ij} = x_i - x_j$$
 and $y_{ij} = y_i - y_j$

- 2A is twice the area of the triangle, $2A = x_{21}y_{31}-x_{31}y_{21}$
- Node numbering is arbitrary except that the sequence 123 must go clockwise around the element if A is to be positive.

2 0 0

Constant Strain Triangle



> Stiffness matrix for element $k = B^T E B t A$

- The CST gives good results in regions of the FE model where there is little strain gradient
 - Otherwise it does not work well.



Changes the shape functions and results in quadratic displacement distributions and linear strain distributions within the element.

$$\begin{aligned} &\lim_{x \to a} \sum_{y \to b_{10}} \sum_{x \to$$

$$\varepsilon_{x} = \beta_{2} + 2\beta_{4}x + \beta_{5}y$$

$$\varepsilon_{y} = \beta_{9} + \beta_{11}x + 2\beta_{12}y$$

$$\gamma_{xy} = (\beta_{3} + \beta_{8}) + (\beta_{5} + 2\beta_{10})x + (2\beta_{6} + \beta_{11})y$$
Example Problem



Consider the problem we were looking at:





UNIT-IV STEADY STATE HEAT TRANSFER ANALYSIS



CLOS	Course Learning Outcomes
CLO 1	Understand the concepts of steady state heat transfer analysis
	for one dimensional slab, fin and thin plate.
CLO 2	Derive the stiffness matrix for fin element.
CLO 3	Solve the steady state heat transfer problems for fin and
	composite slab.

Thermal Convection



Newton's Law of Cooling

$$q = h(T_s - T_\infty)$$

h: convective heat transfer coefficient ($W m^2 \cdot C^o$)

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Thermal Conduction in 1-D



Boundary conditions:

- Dirichlet BC
- Natural BC
- > Mixed BC

Weak Formulation of 1-D Heat Conduction (Steady State Analysis)



Governing Equation of 1-D Heat Conduction

$$-\frac{d}{dx}\left(\frac{\kappa(x)A(x)}{x}\frac{dT(x)}{dx}\right) - AQ(x) = 0 \qquad 0 < x < L$$

Weighted Integral Formulation

$$0 = \int_{0}^{L} w(x) \left[-\frac{d}{dx} \left(\kappa(x)A(x) - \frac{dT(x)}{dx} \right) - AQ(x) \right] dx$$

Weak Form from Integration-by-Parts

$$0 = \int_{0}^{L} \left[\frac{dw}{dx} \begin{pmatrix} dT \\ \kappa A \\ dx \end{pmatrix} \right] - wAQ \left[dx - w \begin{bmatrix} \kappa & dT \\ \kappa & dx \end{bmatrix} \right]_{0}^{L}$$

Formulation for 1-D Linear Element



 $T(\mathbf{x}) = T_1 \phi_1(\mathbf{x}) + T_2 \phi_2(\mathbf{x})$





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Formulation for 1-D Linear Element

• Let w(x) = f_i(x), i = 1, 2

$$0 = \sum_{j=1}^{T} \int_{-1}^{x_{2}} \kappa A \left(\frac{d\phi}{dx} \frac{d\phi}{dx} \right) dx \Big|_{-1}^{x} (\phi AQ) dx - [\phi(x)f + \phi(x)f] \\i = 2 - 2 - i - 1 - 1$$

$$= \sum_{j=1}^{2} K_{ij}T_{j} - Q_{i} - [\phi_{i}(x_{2})f_{2} + \phi_{i}(x_{1})f_{1}] \\\begin{cases} f_{1} \\ f_{2} \end{cases} + \begin{cases} Q_{1} \\ Q_{2} \end{cases} = \begin{bmatrix} K_{11} & K_{12} \\ K_{12} & K_{22} \end{bmatrix} \begin{bmatrix} T_{1} \\ T_{2} \end{bmatrix} \\K_{12} \end{bmatrix} \\k_{12} \end{bmatrix} \\where K_{ij} = \int_{x_{1}}^{x_{2}} \kappa A \left(\frac{d\phi}{dx} \frac{d\phi}{dx} \right) dx, Q_{i} = \int_{x_{1}}^{x_{1}} (\phi AQ) dx, f_{1} = -\kappa A \frac{dT}{dx} \Big|_{x_{1}}, f_{2} = \kappa A \frac{dT}{dx} \Big|_{x_{1}}$$

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Element Equations of 1-D Linear Element





where
$$\mathbf{Q}_i = \int_{x_1}^{x_2} (\phi AQ) dx$$
, $\mathbf{f}_1 = -\kappa A \frac{dT}{dx} \Big|_{x=x_1}$, $\mathbf{f}_2 = \kappa A \frac{dT}{dx} \Big|_{x=x_2}$

A composite wall consists of three materials, as shown in the figure below. The inside wall temperature is 200°C and the outside air temperature is 50°C with a convection coefficient of $h = 10 W(m^2.K)$. Find the temperature along the composite wall.



2 0 0

Thermal Conduction and Convection- Fin



dx

Objective: to enhance heat transfer

Governing equation for 1-D heat transfer in thin fin

$$\frac{\frac{d}{dx}\left(\kappa A \frac{dT}{c}\right)}{\left(\kappa A \frac{dT}{c}\right)} + A Q_{c} = 0$$

$$Q_{loss} = \frac{2h(T - T_{\infty}) \cdot dx \cdot w + 2h(T - T_{\infty}) \cdot dx \cdot t}{A_{c} \cdot dx} = \frac{2h(T - T_{\infty}) \cdot (w + t)}{A_{c}}$$

$$\frac{d}{dx} \left(\kappa A_{c} \frac{dT}{dx} \right) - Ph\left(T - T_{\infty}\right) + A_{c}Q = 0$$

where P = 2(w + t)

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Fin (Steady State Analysis)

Governing Equation of 1-D Heat Conduction

$$-\frac{d}{dx}\left(\kappa(x)A(x)\frac{dT(x)}{-dx}\right) + Ph\left(T - T_{\infty}\right) - AQ = 0 \qquad 0 < x < L$$

Weighted Integral Formulation

$$0 = \int_{0}^{L} w(x) \left[-\frac{d}{dx} \left(\kappa(x)A(x) - \frac{dT(x)}{dx} \right) + Ph(T - T_{\infty}) - AQ(x) \right] dx$$

Weak Form from Integration-by-Parts

$$0 = \int_{0}^{L} \left[\frac{dw}{dx} \left(\frac{dT}{dx} \right) + wPh(T - T_{\infty}) - wAQ \right] \left| \frac{dx}{dx} - \frac{dT}{dx} \right|_{0}^{L}$$

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Formulation for 1-D Linear Element

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Let W(x) = f_i(x),
$$i = 1, 2$$

$$0 = \sum_{j=1}^{T} \prod_{i=1}^{n} \left[\int_{1}^{x_{i}} \left(\kappa A \frac{d\phi}{dx} \frac{d\phi}{dx} \frac{d\phi}{dx} + Ph\phi\phi}{dx dx} \right) dx \right] = \int_{1}^{x_{i}} \phi \left(AQ + PhT \right) dx$$

$$- \left[\phi(x_{2}) f_{2} + \phi_{i}(x_{1}) f_{1} \right]$$

$$= \sum_{j=1}^{2} K_{ij} T_{j} - Q_{i} - \left[\phi(x_{2}) f_{2} + \phi(x_{1}) f_{1} \right]$$

$$\left\{ \begin{cases} f_{1} \\ f_{2} \end{cases} + \left\{ Q_{1} \\ Q_{2} \end{cases} = \left[K_{11} \quad K_{12} \\ K_{12} \quad K_{22} \end{cases} \right] \left\{ T_{1} \\ T_{2} \end{cases} \right\}$$
where $K_{ij} = \sum_{x_{1}}^{x} \left(\kappa A \frac{d\phi}{dx} \frac{d\phi}{dx} + Ph\phi\phi}{dx dx} \right) dx$, $Q_{i} = \int_{1}^{x} \phi \left(AQ + PhT \right) dx$,

$$f_{1} = -\kappa A \frac{dT}{dx} \Big|_{x=x_{1}}, f_{2} = \kappa A \frac{dT}{dx} \Big|_{x=x_{2}}$$

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Element Equations of 1-D Linear Element





where
$$Q_i = \int_{x_1}^{x_2} \phi_i \left(AQ + PhT_{\infty} \right) dx$$
, $f_1 = -\kappa A \left. \frac{dT}{dx} \right|_{x=x_1}$, $f_2 = \kappa A \left. \frac{dT}{dx} \right|_{x=x_2}$

Time-Dependent Problems



Two approaches:

 $u(x,t) = \sum u_i \phi_i(x,t)$

 $u(x,t) = \sum u_i(t)\phi_i(x)$

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$$c \frac{\partial u}{\partial t} = \frac{\partial}{\partial x} \left(a \frac{\partial u}{\partial x} \right) + f(x, t)$$

Weak form:

$$0 = \int_{x_1}^{x_2} \left(a \frac{\partial w}{\partial x} \frac{\partial u}{\partial x} + cw \frac{\partial u}{\partial t} - wf \right) dx - Q_1 w(x_1) - Q_2 w(x_2)$$

$$Q_{1} = -\begin{bmatrix} a & du \\ dx \end{bmatrix}_{x_{1}}; \qquad Q_{2} = \begin{bmatrix} a & du \\ dx \end{bmatrix}_{x_{2}}$$

Transient Heat Conduction

 $u(x,t) = \sum_{j=1}^{n} u_j(t) \phi_j(x) \quad \text{and} \quad w = \phi_i(x)$ let: $0 = \int_{1}^{x_2} \left(a \frac{\partial w \partial u}{\partial x \partial x} + cw \frac{\partial u}{\partial t} - wf \right) dx - Q w(x) - Q w(x)$ $\begin{bmatrix} K \\ u \\ + \end{bmatrix} \{ u \\ + \end{bmatrix} \{ u \\ dx \end{bmatrix} \{ u \\ = \{ F \}$ $M_{ij} = \int c \phi_i \phi_j dx$ $K_{ij} = \int_{0}^{x_2} a \frac{\partial \phi_i \partial \phi_j}{\partial x \partial x} dx$ x_1 $F_i = \int \phi_i f dx + Q_i$

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$$a\frac{du}{dt} + bu = f(t) \qquad 0 < t < T \qquad u(0) = u_0$$

Forward difference approximation – explicit

$$u_{k+1} = u_k + \frac{\Delta t}{a} \left[f_k - b u_k \right]$$

Backward difference approximation - implicit

$$u_{k+1} = u_k + \frac{\Delta}{a + b\Delta t} [f_k - bu_k]$$

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Stability Requirement



$$\Delta t \leq \Delta t_{cri} = \frac{2}{(1 - 2\alpha)\lambda_{max}}$$

where
$$([K] - \lambda[M]) \{u\} = \{Q\}$$

Note: One must use the same discretization for solvingthe eigenvalue problem.

Transient Heat Conduction - Example

$$\frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} = 0 \qquad \qquad 0 < x < 1$$

$$u(0,t) = 0 \qquad \qquad \frac{\partial u}{\partial t} (1,t) = 0 \qquad t > 0$$
$$\frac{\partial u}{\partial t} = 0 \qquad t > 0$$

$$u(x,0) = 1.0$$

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UNIT-V DYNAMIC ANALYSIS



For many structural system, the approximation of linear structural behavior is made in order to convert the physical equilibrium statement, to the following set of second order linear differential equation

$$M\ddot{u}(t)_a + C\dot{u}(t)_a + Ku(t)_a = \mathbf{P(t)}$$

Stiffness and flexibility stiffness matrix

Consider a uniform elastic spring subjected to a load P. This Structure obeys Hook's law. If a force P is applied to a spring fixed. At one end, to produce a displacement then the linear force displacement is u.

$$\begin{array}{c} k & P \\ & & \\ &$$

Stiffness and flexibility stiffeness matrix



- K is called the stiffeness of the spring
- f is called the flexibility of spring

Suppose the uniform elastic spring has nodal points and 2 at its ends, and that the forces at these points are P_1 and P_2 with corresponding displacements u_1 and u_2 .



> From equilibrium consid $\overrightarrow{P_1} = \sqrt{1 + 2}$

$$P_1 = k(u_1 - u_2)$$
$$P_2 = -P_1 = k(u_2 - u_1)$$

> It is convenient to show the above in matrix form as follows

$$\begin{cases} P_1 \\ P_2 \end{cases} = \begin{bmatrix} k & -k \\ -k & k \end{bmatrix} \begin{cases} u_1 \\ u_2 \end{cases}$$

Simple system consisting of just two springs



The system is in equilibrium

 $P_1 + P_2 + P_3 = 0$ $P_1 = k_1(u_1 - u_2)$ $P_2 = k_2(u_3 - u_2)$ $P_2 = -k_1u_1 + (k_1 + k_2) - k_2u_3$



actual distributed mass of the element.

> The element mass matrix is defined as

 $[M] = \int_{v} \rho[N]^{T} [N] \, \mathrm{dV}$

Dynamic equations



I ne torce equiprium of a multi degree of freedom lumped mass system

$$P(t)_i + P(t)_D + P(t)_s = P(t)$$

- Vector of inertia forces acting on the node masses $P(t)_i$
- Vector of viscous damping or energy dissipation forces $P(t)_D$
- A vector of internal forces carried by the structure $P(t)_s$
- Vector of externally applied loads P(t)
- For many structural systems the approximation of linear structural behavior is made in order to convert the physical equilibrium statement, to the following set of second order linear differential equation

$$M\ddot{u}(t)_a + C\dot{u}(t)_a + Ku(t)_a = P(t)$$

Vibration analysis



when loads are suddenly applied or when the loads are of a variable nature, the mass and acceleration effects come into the picture. If a solid such as an engineering structure is deformed elastically and suddenly released. It tends to vibrate about its equilibrium position. This periodic motion due to the restoring strain energy is called free vibration. The number of cycles per unit time is called frequency. The maximum displacement from the equilibrium position is the amplitude.

Equation for damped forced vibration

$$M\ddot{u}(t) + C\dot{u}(t) + Ku(t) = P(t)$$

If there is no damping the equation become

 $M\ddot{u}(t) + Ku(t) = P(t)$

• Free undamped vibration equation $M\ddot{u}(t) + Ku(t) = 0$



- The free undamped vibration equation is linear and homogeneous. Its general solution is a linear combination of exponentials. Under matrix definiteness conditions the exponentials can be expressed as a combination of trignometric functions: sines and cosines of argument ωt.
- A compact representation of such functions is obtained by using the exponential form $e^{j\omega t}$

$$u(t) = \sum v_i e^{j\omega t}$$

Replace
$$u(t) = v_i e^{j\omega t}$$

 $M\ddot{u}(t) + Ku(t) = 0$

The Vibration Eigen problem



I ne time dependence to the exponential is segregated

$$(-\omega^2 M + K)ve^{j\omega t} = 0$$

Since is not identically zero, it can be dropped leaving the algebraic condition

$$(-\omega^2 M + K)v = 0$$

> Because v cannot be the null vector this equation is an algebraic Eigen value problem in ω^2 . The Eigen values $\lambda_i = \omega_i^2$ are the roots of the characteristic polynomial be index by I

$$\det(K - \omega_i^2 M) = 0$$

Dropping the index I this Eigen problem is usually written as

$$Kv = \omega^2 Mv$$

hank you