

# Finite element programming by FreeFem++ – intermediate course

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## Numerical simulation with finite element method

- ▶ mathematical modeling
- ▶ discretization of time for evolution problem
- ▶ discretization scheme for the space
  - ▶ mesh generation / adaptive mesh refinement
  - ▶ stiffness matrix from finite elements and variational formulation
  - ▶ linear solver  $\leftarrow$  CG, GMRES, direct solver: UMFPACK, MUMPS

FreeFem++ provides vast amounts of tools

- ▶ nonlinear solver
- ▶ optimization solver

parallel computation is another topic.

distributed parallelization by MPI needs to be described by FreeFem++ script.

## Outline I

### Basics of FEM by examples from the Poisson equation

- Poisson equation with mixed boundary conditions
- error estimate by theory and FreeFem++ implementation

### Mixed formulation for the Stokes equations

- Stokes equations with inhomogenous Dirichlet conditions
- mixed formulation and inf-sup conditions
- finite element pair satisfying inf-sup conditions

### Nonlinear finite element problem by Newton method

- stationary Navier-Stokes equations
- differential calculus of nonlinear operator and Newton iteration

### Time-dependent Navier-Stokes equations around a cylinder

- boundary conditions of incompressible flow around a cylinder
- characteristic Galerkin method
- stream line for visualization

## Outline II

thermal convection problem by Rayleigh-Bénard eqs.

governing equations by Boussinesq approximation

time-dependent solution by characteristic Galerkin method

stationary solution by Newton method

### Conjugate Gradient solver in FreeFem++

basic CG method with preconditioning

CG method with orthogonal projection onto the image

CG method in Uzawa method

### Syntax and Data types of FreeFem++

syntax for loop and procedure

Data types

### Compilation from the source on Unix system

configure script with option and BLAS library

### References

## Poisson equation with mixed B.C. and a weak formulation: 1/2

$$\Omega \subset \mathbb{R}^2, \partial\Omega = \Gamma_D \cup \Gamma_N$$

$$-\Delta u = f \text{ in } \Omega,$$

$$u = g \text{ on } \Gamma_D,$$

$$\frac{\partial u}{\partial n} = h \text{ on } \Gamma_N.$$

weak formulation

$V$  : function space,  $V(g) = \{u \in V ; u = g \text{ on } \Gamma_D\}$ .

$V = C^1(\Omega) \cap C^0(\bar{\Omega})$  ?

Find  $u \in V(g)$  s.t.

$$\int_{\Omega} -\Delta u v dx = \int_{\Omega} f v dx \quad \forall v \in V(0)$$

Lemma (Gauss-Green's formula)

$u, v \in V, n = \begin{bmatrix} n_1 \\ n_2 \end{bmatrix}$  : outer normal to  $\partial\Omega$

$$\int_{\Omega} (\partial_i u) v dx = - \int_{\Omega} u \partial_i v dx + \int_{\partial\Omega} u n_i v ds .$$

## Poisson equation with mixed B.C. and a weak formulation: 2/2

$$\begin{aligned}\int_{\Omega} (-\partial_1^2 - \partial_2^2) u v \, dx &= \int_{\Omega} (\partial_1 u \partial_1 v + \partial_2 u \partial_2 v) \, dx - \int_{\partial\Omega} (\partial_1 u n_1 + \partial_2 u n_2) v \, ds \\ &= \int_{\Omega} \nabla u \cdot \nabla v \, dx - \int_{\Gamma_D \cup \Gamma_N} \nabla u \cdot n v \, ds \\ v = 0 \text{ on } \Gamma_D &\Rightarrow = \int_{\Omega} \nabla u \cdot \nabla v \, dx - \int_{\Gamma_N} h v \, ds\end{aligned}$$

Find  $u \in V(g)$  s.t.

$$\int_{\Omega} \nabla u \cdot \nabla v \, dx = \int_{\Omega} f v \, dx + \int_{\Gamma_N} h v \, ds \quad \forall v \in V(0)$$

- ▶  $a(\cdot, \cdot) : V \times V \rightarrow \mathbb{R}$  : bilinear form
- ▶  $F(\cdot) : V \rightarrow \mathbb{R}$  : functional

Find  $u \in V(g)$  s.t.

$$a(u, v) = F(v) \quad \forall v \in V(0)$$

## FreeFem++ script to solve Poisson equation

finite element basis,  $\text{span}[\varphi_1, \dots, \varphi_N] = V_h \subset V$

$u_h \in V_h \Rightarrow u_h = \sum_{1 \leq i \leq N} u_i \varphi_i$

Dirichlet data :  $u(P_j) = g(P_j) \quad P_j \in \Gamma_D$

Find  $u_h \in V_h(g)$  s.t.

$$\int_{\Omega} \nabla u_h \cdot \nabla v_h dx = \int_{\Omega} f v_h dx + \int_{\Gamma_N} h v_h ds \quad \forall v_h \in V_h(0).$$

```
mesh Th=square(20,20);
```

▶ example1.edp

▶ varf+matrix

```
fespace Vh(Th,P1);
```

```
Vh uh,vh;
```

```
func f = 5.0/4.0*pi*pi*sin(pi*x)*sin(pi*y/2.0);
```

```
func g = sin(pi*x)*sin(pi*y/2.0);
```

```
func h = (-pi)/2.0 * sin(pi * x);
```

```
solve poisson(uh,vh)=
```

```
  int2d(Th) ( dx(uh)*dx(vh)+dy(uh)*dy(vh) )
```

```
  - int2d(Th) ( f*vh ) - int1d(Th,1) ( h *vh )
```

```
  + on(2,3,4,uh=g);           // boudary 1 : (x,0)
```

```
plot(uh);
```

## discretization and matrix formulation : 1/2

finite element basis,  $\text{span}[\varphi_1, \dots, \varphi_N] = V_h \subset V$

$$u_h \in V_h \Rightarrow u_h = \sum_{1 \leq i \leq N} u_i \varphi_i$$

finite element nodes  $\{P_j\}_{j=1}^N$ ,  $\varphi_i(P_j) = \delta_{ij}$  Lagrange element

$\Lambda_D \subset \Lambda = \{1, \dots, N\}$  : index of node on the Dirichlet boundary

$$V_h(g) = \{u_h \in V_h; u_h = \sum u_i \varphi_i, u_k = g_k (k \in \Lambda_D)\}$$

Find  $u_h \in V_h(g)$  s.t.

$$a(u_h, v_h) = F(v_h) \quad \forall v_h \in V_h(0).$$

Find  $\{u_j\}$ ,  $u_k = g_k (k \in \Lambda_D)$  s.t.

$$a\left(\sum_j u_j \varphi_j, \sum_i v_i \varphi_i\right) = F\left(\sum_i v_i \varphi_i\right) \quad \forall \{v_i\}, v_k = 0 (k \in \Lambda_D)$$

Find  $\{u_j\}_{j \in \Lambda}$  s.t.

$$\begin{aligned} \sum_j a(\varphi_j, \varphi_i) u_j &= F(\varphi_i) & \forall i \in \Lambda \setminus \Lambda_D \\ u_k &= g_k & \forall k \in \Lambda_D \end{aligned}$$



## discretization and matrix formulation : 2/2

Find  $\{u_j\}_{j \in \Lambda \setminus \Lambda_D}$  s.t.

$$\sum_{j \in \Lambda \setminus \Lambda_D} a(\varphi_j, \varphi_i) u_j = F(\varphi_i) - \sum_{k \in \Lambda_D} a(\varphi_k, \varphi_i) g_k \quad \forall i \in \Lambda \setminus \Lambda_D$$

$A = \{a(\varphi_j, \varphi_i)\}_{i,j \in \Lambda \setminus \Lambda_D}$  : symmetric.

$A \in \mathbb{R}^{n \times n}$ ,  $f \in \mathbb{R}^n$ ,  $n = \#(\Lambda \setminus \Lambda_D)$

### Lemma

$A$  : (*symmetric*) *positive definite*  $\Leftrightarrow (Au, u) > 0 \quad \forall u \neq 0$

$\Rightarrow Au = f$  has a unique solution.

$A$  : bijective

► injective:  $Au = 0$ ,  $0 = (Au, u) > 0 \Rightarrow u = 0$ .

► surjective:

$\mathbb{R}^n = \text{Im}A \oplus (\text{Im}A)^\perp$ ,  $u \in (\text{Im}A)^\perp \Rightarrow (Av, u) = 0 \quad \forall v \in \mathbb{R}^n$

by putting  $v = u$ ,  $0 = (Au, u) \Rightarrow u = 0$

$(\text{Im}A)^\perp = \{0\} \Rightarrow \text{Im}A = \mathbb{R}^n$ .

$A$  : S.P.D.  $\Rightarrow$  solution by  $LDL^T$ -factorization, CG method

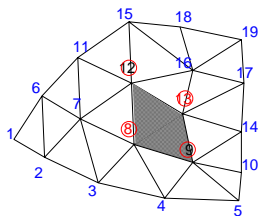
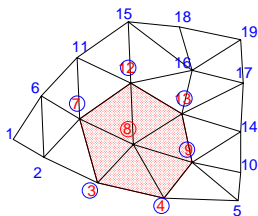
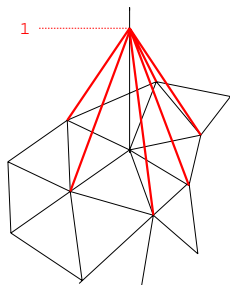
## P1 finite element and sparse matrix

$\mathcal{T}_h$  : triangulation of a domain  $\Omega$ , triangular element  $K \in \mathcal{T}_h$

piecewise linear element :  $\varphi_i|_K(x_1, x_2) = a_0 + a_1x_1 + a_2x_2$

$\varphi_i|_K(P_j) = \delta_{ij}$

$$[A]_{ij} = a(\varphi_j, \varphi_i) = \int_{\Omega} \nabla \varphi_j \cdot \nabla \varphi_i dx = \sum_{K \in \mathcal{T}_h} \int_K \nabla \varphi_j \cdot \nabla \varphi_i dx.$$



$A$  : sparse matrix, CRS (Compressed Row Storage) format to store

## FreeFem++ script to solve Poisson eq. using matrix

Find  $u_h \in V_h(g)$  s.t.  $a(u_h, v_h) = F(v_h) \forall v_h \in V_h(0)$ .

▶ example2.edp    ▶ solve

```
Vh u, v;  
varf aa(u, v) = int2d(Th) ( dx(u) * dx(v) + dy(u) * dy(v) )  
                + on(2, 3, 4, u=g);  
varf external(u, v) = int2d(Th) ( f*v ) + int1d(Th, 1) ( h*v )  
                + on(2, 3, 4, u=g);  
real tgv = 1.0e+30;  
matrix A = aa(Vh, Vh, tgv=tgv, solver=CG);  
real[int] ff = external(0, Vh, tgv=tgv);  
u[] = A^-1 * ff;      // u : fem unknown, u[] : vector  
plot(u);
```

useful liner solver; solver=

CG            iterative solver for SPD matrix

GMRES        iterative solver for nonsingular matrix

UMFPACK      direct solver for nonsingular matrix

sparsesolver other solvers called by dynamic link

## penalty method to solve inhomogeneous Dirichlet problem

modification of diagonal entries of  $A$  where index  $k \in \Lambda_D$

penalization parameter  $\tau = 1/\varepsilon$ ;  $\tau g_k$

$$[A]_{ij} = a(\varphi_j, \varphi_i)$$
$$\begin{matrix} u_k \\ \vdots \\ u_i \\ \vdots \end{matrix} = \begin{matrix} \tau g_k, k \in \Lambda_D \\ \vdots \\ f_i \\ \vdots \end{matrix}$$

$$\tau u_k + \sum_{j \neq k} a_{kj} u_j = \tau g_k \Leftrightarrow u_k - g_k = \varepsilon \left( - \sum_{j \neq k} a_{kj} u_j \right),$$
$$\sum_j a_{ij} u_j = f_i \quad \forall i \in \{1, \dots, N\} \setminus \Lambda_D.$$

keeping symmetry of the matrix without changing index numbering.

## abstract framework

$V$ : Hilbert space with inner product  $(\cdot, \cdot)$  and norm  $\|\cdot\|$ .  
bilinear form  $a(\cdot, \cdot) : V \times V \rightarrow \mathbb{R}$

- ▶ coercive :  $\exists \alpha > 0 \quad a(u, u) \geq \alpha \|u\|^2 \quad \forall u \in V$ .
- ▶ continuous :  $\exists \gamma > 0 \quad |a(u, v)| \leq \gamma \|u\| \|v\| \quad \forall u, v \in V$ .

functional  $F(\cdot) : V \rightarrow \mathbb{R}$ .

find  $u \in V$  s.t.  $a(u, v) = F(v) \quad \forall v \in V$

has a unique solution : Lax-Milgram's theorem

inf-sup conditions + continuity of  $a(\cdot, \cdot)$

- ▶  $\exists \alpha_1 > 0 \quad \sup_{v \in V, v \neq 0} \frac{a(u, v)}{\|v\|} \geq \alpha_1 \|u\| \quad \forall u \in V$ .
- ▶  $\exists \alpha_2 > 0 \quad \sup_{u \in V, u \neq 0} \frac{a(u, v)}{\|u\|} \geq \alpha_2 \|v\| \quad \forall v \in V$ .

find  $u \in V$  s.t.  $a(u, v) = F(v) \quad \forall v \in V$  has a unique solution.

## error estimate : theory 1 /2

$V$  : Hilbert space,  $V_h \subset V$  : finite element space.

▶  $u \in V$ ,  $a(u, v) = F(v) \quad \forall v \in V$ .

▶  $u_h \in V_h$ ,  $a(u_h, v_h) = F(v_h) \quad \forall v_h \in V_h \subset V$ .

$$a(u, v_h) = F(v_h) \quad \forall v_h \in V_h \subset V.$$

**Lemma (Galerkin orthogonality)**

$$a(u - u_h, v_h) = 0 \quad \forall v_h \in V_h.$$

assuming coercivity and continuity of  $a(\cdot, \cdot)$ .

**Lemma (Céa)**

$$\|u - u_h\| \leq \frac{\gamma}{\alpha} \inf_{v_h \in V_h} \|u - v_h\|.$$

*proof:*  $\|u - u_h\| \leq \|u - v_h\| + \|v_h - u_h\|$

$$\begin{aligned} \alpha \|u_h - v_h\|^2 &\leq a(u_h - v_h, u_h - v_h) \\ &= a(u_h, u_h - v_h) - a(v_h, u_h - v_h) \\ &= a(u, u_h - v_h) - a(v_h, u_h - v_h) \\ &= a(u - v_h, u_h - v_h) \leq \gamma \|u - v_h\| \|u_h - v_h\|. \end{aligned}$$

## Sobolev space : 1/2

P1 element space does not belong to  $C^1(\Omega)$ .

$V = H^1(\Omega)$ ,  $(u, v) = \int_{\Omega} u v + \nabla u \cdot \nabla v$ ,  $\|u\|_1^2 = (u, u) < +\infty$ .

$\|u\|_0^2 = \int_{\Omega} u u$ ,  $|u|_1^2 = \int_{\Omega} \nabla u \cdot \nabla u$ .

$H_0^1 = \{u \in H^1(\Omega); u = 0 \text{ on } \partial\Omega\}$ .

**Lemma (Poincaré's inequality)**

$\exists C(\Omega) u \in H_0^1 \Rightarrow \|u\|_0 \leq C(\Omega)|u|_1$ .

*proof:*

$\Omega \subset B = (0, s) \times (0, s)$ .  $v \in C_0^\infty(\Omega)$ ,  $\tilde{v}(x) = 0$  ( $x \in B \setminus \bar{\Omega}$ ).

$$v(x_1, x_2) = v(0, x_2) + \int_0^{x_1} \partial_1 v(t, x_2) dt$$

$$|v(x_1, x_2)|^2 \leq \int_0^{x_1} 1^2 dt \int_0^{x_1} |\partial_1 v(t, x_2)|^2 dt \leq s \int_0^s |\partial_1 v(t, x_2)|^2 dt$$

$$\int_0^s |v(x_1, x_2)|^2 dx_1 \leq s^2 \int_0^s |\partial_1 v(x)|^2 dx_1$$

$$\int_{\Omega} |v|^2 = \int_B |v|^2 dx_1 dx_2 \leq s^2 \int_B |\partial_1 u|^2 dx_1 dx_2 = s^2 \int_{\Omega} |\partial_1 u|^2.$$

## Sobolev space : 2/2

$$a(u, v) = \int_{\Omega} \nabla u \cdot \nabla v, \quad u, v \in H^1(\Omega).$$

▶  $a(\cdot, \cdot)$  is coercive on  $H_0^1(\Omega)$ .

▶  $a(\cdot, \cdot)$  is coercive on  $H^1(\Omega)/\mathbb{R}$ .

full-Neumann boundary problem

$$-\Delta u = f \text{ in } \Omega,$$

$$\partial_n u = h \text{ on } \partial\Omega.$$

▶  $(F, v) = F(v) = \int_{\Omega} f v + \int_{\partial\Omega} h v$

▶ compatibility condition :  $(F, 1) = \int_{\Omega} f + \int_{\partial\Omega} h = 0$

(N) Find  $u \in H^1(\Omega)$  s.t.

$$a(u, v) = F(v) \quad \forall v \in H^1(\Omega)$$

$u$ : solution of (N)  $\Rightarrow u + 1$ : solution of (N)

$[A]_{ij} = a(\varphi_j, \varphi_i)$ .  $A$  : singular ,  $\text{Ker} A = \vec{1}$ .

(N) has a unique solution in  $H^1(\Omega)/\mathbb{R} \simeq \{u \in H^1(\Omega) ; \int_{\Omega} u = 0\}$ .



## error estimate : theory 2 /2

$\Pi_h : C(\bar{\Omega}) \rightarrow V_h, \quad \Pi_h u = \sum_i u(P_i)\varphi_i,$   
 $\{\varphi_i\} : P_k$  finite element basis,  $\text{span}\{\{\varphi_i\}\} = V_h.$

**Theorem (interpolation error by polynomial)**

$K \in \mathcal{T}_h, P_k(K) \subset H^l(K), v \in H^{k+1}(\Omega)$   
 $\Rightarrow \exists c > 0 \quad |v - \Pi_h v|_{s,K} \leq C h_K^{k+1-s} |v|_{k+1,K},$   
 $0 \leq s \leq \min\{k+1, l\}.$

**Theorem (finite element error)**

$u \in H^{k+1}, u_h : \text{finite element solution by } P_k \text{ element.}$

$\Rightarrow \exists c > 0 \quad \|u - u_h\|_{1,\Omega} \leq C h^k |u|_{k+1,\Omega}$

*proof:*  $\|u - u_h\|_{1,\Omega} \leq C \inf_{v_h \in V_h} \|u - v_h\|_{1,\Omega}$   
 $\leq C \|u - \Pi_h u\|_{1,\Omega}$   
 $\leq C \sum_{K \in \mathcal{T}_h} (h_K^k + h_K^{(k+1)}) |u|_{k+1,K}$   
 $\leq C h^k |u|_{k+1,\Omega}$

$\mathcal{T}_h : \text{finite element mesh, } h_K = \text{diam}(K), h = \max_{K \in \mathcal{T}_h} h_K.$

## numerical integration

Numerical quadrature:

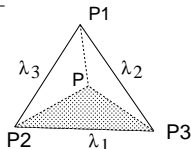
$\{P_i\}_{i \leq i \leq m}$  : integration points in  $K$ ,  $\{\omega_i\}_{i \leq i \leq m}$  : weights

$$|u - u_h|_{0,\Omega}^2 = \sum_{K \in \mathcal{T}_h} \int_K |u - u_h|^2 dx \sim \sum_{K \in \mathcal{T}_h} \sum_{i=1}^m |(u - u_h)(P_i)|^2 \omega_i$$

formula : degree 5, 7 points, **qf5pT**,

P.C. Hammer, O.J. Marlowe, A.H. Stroud [1956]

area coordinates $\{\lambda_i\}_{i=1}^3$	weight	
$(\frac{1}{3}, \frac{1}{3}, \frac{1}{3})$	$\frac{9}{40} K $	$\times 1$
$(\frac{6-\sqrt{15}}{21}, \frac{6-\sqrt{15}}{21}, \frac{9+2\sqrt{15}}{21})$	$\frac{155-\sqrt{15}}{1200} K $	$\times 3$
$(\frac{6+\sqrt{15}}{21}, \frac{6+\sqrt{15}}{21}, \frac{9-2\sqrt{15}}{21})$	$\frac{155+\sqrt{15}}{1200} K $	$\times 3$



### Remark

it is not good idea to use interpolation of continuous function to finite element space, for verification of convergence order.

$|\Pi_h u - u_h|_{1,\Omega}$  may be smaller (in extreme cases, super convergence)

## numerical convergence order

for observing convergence order

$u \in H^2(\Omega)$  : manufactured solution

$u_h \in V_h(g)$  : finite element solution by  $P_k$  element.

$$\begin{aligned} \|u - u_h\|_{1,\Omega} &= c h^k, \\ \frac{\|u - u_{h_1}\|_{1,\Omega}}{\|u - u_{h_2}\|_{1,\Omega}} &= \frac{c h_1^k}{c h_2^k} = \left(\frac{h_1}{h_2}\right)^k \end{aligned}$$

numerical convergence order:

$$\kappa = \log\left(\frac{\|u - u_{h_1}\|_{1,\Omega}}{\|u - u_{h_2}\|_{1,\Omega}}\right) / \log\left(\frac{h_1}{h_2}\right).$$

## FreeFem++ script for error estimation

▶ example3.edp

```
real hh1,hh2,err1,err2;
func sol = sin(pi*x)*sin(pi*y/2.0);
func solx = pi*cos(pi*x)*sin(pi*y/2.0);
func soly = (pi/2.0)*sin(pi*x)*cos(pi*y/2.0);
mesh Th1=square(n1,n1);
mesh Th2=square(n2,n2);
fespace Vh1(Th1,P1);
...
solve poisson1(u1,v1) = ...
...
err1 = int2d(Th1) ((dx(u1)-solx)*(dx(u1)-solx) +
                  (dy(u1)-soly)*(dy(u1)-soly) +
                  (u1-sol)*(u1-sol));
err1 = sqrt(err1);
...
hh1 = 1.0/n1*sqrt(2.0);
hh2 = 1.0/n2*sqrt(2.0);
cout<<"O(h^2)="<<log(err1/err2)/log(hh1/hh2)<<endl;
```

## error estimate on unstructured mesh

unstructured mesh is generated by Delaunay triangulation

▶ example4.edp

```
n1 = 20;
border bottom(t=0,1) {x=t;y=0;   label=1;};
border right(t=0,1)  {x=1;y=t;   label=2;};
border top(t=0,1)    {x=1-t;y=1; label=3;};
border left(t=0,1)   {x=0;y=1-t; label=4;};
mesh Th1=buildmesh(bottom(n1)+right(n1)+top(n1)
                   +left(n1));

...
fespace Vh10(Th1,P0);
Vh10 h1 = hTriangle;
hh1 = h1[].max;

...
```

### Remark

$\min_K h_K$ ,  $\sum_K h_K / \#\mathcal{T}_h$ ,  $\max_K h_K$ , corresponding to mesh refinement are observed by following:

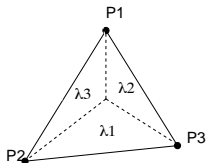
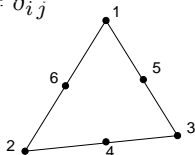
```
h1[].min;   hh1 = h1[].sum / h1[].n;   h1[].max;
```

## P2 finite element

$\mathcal{T}_h$  : triangulation of a domain  $\Omega$ , triangular element  $K \in \mathcal{T}_h$   
piecewise quadratic element : 6 DOF on element  $K$ .

$$\varphi_i|_K(x_1, x_2) = a_0 + a_1x_1 + a_2x_2 + a_3x_1^2 + a_4x_1x_2 + a_5x_2^2$$

$$\varphi_i|_K(P_j) = \delta_{ij}$$



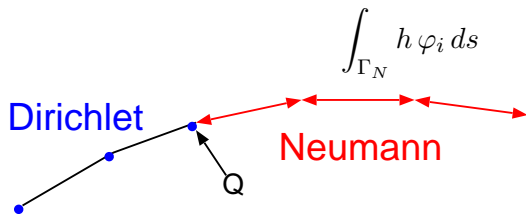
by using area coordinates  $\{\lambda_1, \lambda_2, \lambda_3\}$ ,  $\lambda_1 + \lambda_2 + \lambda_3 = 1$ .

$$\begin{pmatrix} \varphi_1 \\ \varphi_2 \\ \varphi_3 \\ \varphi_4 \\ \varphi_5 \\ \varphi_6 \end{pmatrix} = \begin{pmatrix} 1 & & & & & \\ & 1 & & & & \\ & & 1 & & & \\ & & & 1 & & \\ & & & & 4 & \\ & & & & & 4 \\ & & & & & & 4 \end{pmatrix} \begin{pmatrix} \lambda_1^2 \\ \lambda_2^2 \\ \lambda_3^2 \\ \lambda_2\lambda_3 \\ \lambda_3\lambda_1 \\ \lambda_1\lambda_2 \end{pmatrix} = \begin{pmatrix} \lambda_1(2\lambda_1 - 1) \\ \lambda_2(2\lambda_2 - 1) \\ \lambda_3(2\lambda_3 - 1) \\ 4\lambda_2\lambda_3 \\ 4\lambda_3\lambda_1 \\ 4\lambda_1\lambda_2 \end{pmatrix}$$

fespace  $V_h(\mathcal{T}_h, P2)$  ;

## treatment of Neumann data around mixed boundary

Neumann data is evaluated by line integral with FEM basis  $\varphi_i$ .



For given discrete Neumann data,  $h$  is interpolated in FEM space,  $h = \sum_j h_j \varphi_j|_{\Gamma_N}$ ,

$$\sum_j h_j \int_{\Gamma_N} \varphi_j \varphi_i ds.$$

On the node  $Q \in \bar{\Gamma}_D \cap \bar{\Gamma}_N$ , both Dirichlet and Neumann are necessary.

## advantages of finite element formulation

- ▶ weak formulation is obtained by integration by part with clear description on the boundary
- ▶ Dirichlet boundary condition is embedded in a functional space, called as essential boundary condition
- ▶ Neumann boundary condition is treated with surface/line integral by Gauss-Green's formula, called as natural boundary condition
- ▶ solvability of linear system is inherited from solvability of continuous weak formulation
- ▶ error of finite element solution is evaluated by approximation property of finite element space

## better to learn for efficient computation

- ▶ treatment of Dirichlet boundary conditions in FreeFem++ with explicit usage of matrix and linear solver



## Stokes equations and a weak formulation : 1/3

$$\Omega = (0, 1) \times (0, 1)$$

$$-2\nabla \cdot D(u) + \nabla p = f \text{ in } \Omega$$

$$\nabla \cdot u = 0 \text{ in } \Omega$$

$$u = g \text{ on } \partial\Omega$$

strain rate tensor :  $[D(u)]_{ij} = \frac{1}{2}(\partial_i u_j + \partial_j u_i)$ .

►  $V(g) = \{v \in H^1(\Omega)^2; v = g \text{ on } \partial\Omega\}$ ,  $V = V(0)$

►  $Q = L_0^2(\Omega) = \{p \in L^2(\Omega); \int_{\Omega} p \, dx = 0\}$

bilinear form and weak formulation :

$$a(u, v) = \int_{\Omega} 2D(u) : D(v) \, dx \quad u, v \in H^1(\Omega)^2$$

$$b(v, p) = - \int_{\Omega} \nabla \cdot v \, p \, dx \quad v \in H^1(\Omega)^2, p \in L^2(\Omega)$$

Find  $(u, p) \in V(g) \times Q$  s.t.

$$a(u, v) + b(v, p) = (f, v) \quad \forall v \in V,$$

$$b(u, q) = 0 \quad \forall q \in Q.$$

## Stokes equations and a weak formulation : 2/3

Lemma (Gauss-Green's formula)

$u, v \in H^1(\Omega)$ ,  $n$  : *outer normal to*  $\partial\Omega$

$$\int_{\Omega} (\partial_i u) v \, dx = - \int_{\Omega} u \partial_i v \, dx + \int_{\partial\Omega} u n_i v \, ds .$$

$$\begin{aligned} & -2 \int_{\Omega} (\nabla \cdot D(u)) \cdot v \, dx = \\ & -2 \int_{\Omega} \sum_i \sum_j \partial_j \frac{1}{2} (\partial_i u_j + \partial_j u_i) v_i \, dx = \int_{\Omega} \sum_{i,j} (\partial_i u_j + \partial_j u_i) \partial_j v_i \, dx \\ & \qquad \qquad \qquad - \int_{\partial\Omega} \sum_{i,j} (\partial_i u_j + \partial_j u_i) n_j v_i \, ds \\ & = \int_{\Omega} 2D(u) : D(v) \, dx - \int_{\partial\Omega} 2D(u) n \cdot v \, ds \end{aligned}$$

from the symmetry of  $D(u)$

$$\sum_{i,j} (\partial_i u_j + \partial_j u_i) \partial_j v_i = \sum_{i,j} (\partial_i u_j + \partial_j u_i) (\partial_j v_i + \partial_i v_j) / 2 = 2D(u) : D(v) .$$

## Stokes equations and a weak formulation : 3/3

$$\begin{aligned}\int_{\Omega} \sum_i (\partial_i p) v_i dx &= - \int_{\Omega} \sum_i p \partial_i v_i dx + \int_{\partial\Omega} \sum_i p n_i v_i \\ &= - \int_{\Omega} p \nabla \cdot v + \int_{\partial\Omega} p n \cdot v\end{aligned}$$

On the boundary  $\partial\Omega$ ,

$$\int_{\partial\Omega} (2D(u)n - np) \cdot v ds = 0 \quad v \in V \Rightarrow v = 0 \text{ on } \partial\Omega.$$

### Remark

compatibility condition on Dirichlet data :

$$0 = \int_{\Omega} \nabla \cdot u = - \int_{\Omega} u \cdot \nabla 1 + \int_{\partial\Omega} u \cdot n ds = \int_{\partial\Omega} g \cdot n ds.$$

### Remark

$$-2[\nabla \cdot D(u)]_i = - \sum_j \partial_j (\partial_i u_j + \partial_j u_i) = - \sum_j \partial_j^2 u_i = -[\Delta u]_i.$$

## existence of a solution of the Stokes equations

Find  $(u, p) \in V(g) \times Q$  s.t.

$$a(u, v) + b(v, p) = (f, v) \quad \forall v \in V,$$

$$b(u, q) = 0 \quad \forall q \in Q.$$

► coercivity :  $\exists \alpha_0 > 0 \quad a(u, u) \geq \alpha_0 \|u\|_1^2 \quad \forall u \in V.$

► inf-sup condition :

$$\exists \beta_0 > 0 \quad \sup_{v \in V, v \neq 0} \frac{b(v, q)}{\|v\|_1} \geq \beta_0 \|q\|_0 \quad \forall q \in Q.$$

bilinear form :  $A(u, p; v, q) = a(u, v) + b(v, p) + b(u, q)$

**Lemma**

$$\exists \alpha > 0 \quad \sup_{(u, p) \in V \times Q} \frac{A(u, p; v, q)}{\|(u, p)\|_{V \times Q}} \geq \alpha \|(v, q)\|_{V \times Q} \quad \forall (v, q) \in V \times Q.$$

**Here,**  $\|(u, p)\|_{V \times Q}^2 = \|u\|_1^2 + \|p\|_0^2.$

Find  $(u, p) \in V(g) \times Q$  s.t.

$$A(u, p; v, q) = (f, v) \quad \forall (v, q) \in V \times Q.$$

## mixed finite element method

$V_h \subset V$  : P2 finite element

$Q_h \subset Q$  : P1 finite element +  $\int_{\Omega} p_h dx = 0$ .

► coercivity :  $\exists \alpha_0 > 0 \quad a(u_h, u_h) \geq \alpha_0 \|u_h\|_1^2 \quad \forall u_h \in V_h$ .

► uniform inf-sup condition :

$$\exists \beta_0 > 0 \quad \forall h > 0 \quad \sup_{v_h \in V_h, v_h \neq 0} \frac{b(v_h, q_h)}{\|v_h\|_1} \geq \beta_0 \|q_h\|_0 \quad \forall q_h \in Q_0.$$

### Lemma

$$\exists \alpha > 0 \quad \sup_{(u_h, p_h) \in V_h \times Q_h} \frac{A(u_h, p_h; v_h, q_h)}{\|(u_h, p_h)\|_{V \times Q}} \geq \alpha \|(v_h, q_h)\|_{V \times Q} \\ \forall (v_h, q_h) \in V_h \times Q_h.$$

Find  $(u_h, p_h) \in V_h(g) \times Q_h$  s.t.

$$A(u_h, p_h; v_h, q_h) = (f, v_h) \quad \forall (v_h, q_h) \in V_h \times Q_h.$$

### Lemma

$$\|u - u_h\|_1 + \|p - p_h\|_0 \leq C(\inf_{v_h \in V} \|u - v_h\|_1 + \inf_{q_h \in Q} \|p - q_h\|_0)$$

## FreeFem++ script to solve Stokes equations by P2/P1

Find  $(u, p) \in V_h(g) \times Q_h$  s.t.

$$a(u, v) + b(v, p) + b(u, q) - \epsilon \int_{\Omega} p q = (f, v) \quad \forall (v, q) \in V_h \times Q_h.$$

```
fespace Vh(Th, P2), Qh(Th, P1); ▶ example5.edp
func f1=5.0/8.0*pi*pi*sin(pi*x)*sin(pi*y/2.0)+2.0*x;
func f2=5.0/4.0*pi*pi*cos(pi*x)*cos(pi*y/2.0)+2.0*y;
func g1=sin(pi*x)*sin(pi*y/2.0)/2.0;
func g2=cos(pi*x)*cos(pi*y/2.0);
Vh u1,u2,v1,v2; Qh p,q;
macro d12(u1,u2) (dy(u1) + dx(u2))/2.0 //
real epsln=1.0e-6;
solve stokes(u1,u2,p1, v1,v2,q1) =
int2d(Th) ( 2.0*(dx(u1)*dx(v1)
+2.0*d12(u1,u2)*d12(v1,v2)+dy(u2)*dy(v2))
-p*dx(v1)-p*dy(v2)-dx(u1)*q-dy(u2)*q
-p*q*epsln ) // penalization
- int2d(Th) ( f1 * v1 + f2 * v1 )
+ on(1,2,3,4,u1=g1,u2=g2);
real meanp=int2d(Th) (p)/Th.area; //area=int2d(Th) (1.0)
p = p - meanp;
plot ([u1,u2],p,wait=1,value=true,coef=0.1);
```

## stabilized (penalty type) finite element method

$V_h \subset V$  : P1 finite element

$Q_h \subset Q$  : P1 finite element +  $\int_{\Omega} p_h dx = 0$ .

**Find**  $(u_h, p_h) \in V_h(g) \times Q_h$  **s.t.**

$$a(u_h, v_h) + b(v_h, p_h) = (f, v_h) \quad \forall v_h \in V_h,$$

$$b(u_h, q_h) - \delta d(p_h, q_h) = 0 \quad \forall q_h \in Q_h.$$

$\delta > 0$  : stability parameter,  $d(p_h, q_h) = \sum_{K \in \mathcal{T}} h_K^2 \int_K \nabla p_h \cdot \nabla q_h dx$ .

$|p_h|_h^2 = d(p_h, p_h)$  : mesh dependent norm on  $Q_h$ .

► uniform weak inf-sup condition : Franca-Stenberg [1991]

$$\exists \beta_0, \beta_1 > 0 \quad \forall h > 0 \quad \sup_{v_h \in V_h} \frac{b(v_h, q_h)}{\|v_h\|_1} \geq \beta_0 \|q_h\|_0 - \beta_1 |q_h|_h \quad \forall q_h \in Q_0.$$

**Lemma**

$$\exists \alpha > 0 \quad \sup_{(u_h, p_h) \in V_h \times Q_h} \frac{A(u_h, p_h; v_h, q_h)}{\|(u_h, p_h)\|_{V \times Q}} \geq \alpha \|(v_h, q_h)\|_{V \times Q}$$

$$\forall (v_h, q_h) \in V_h \times Q_h.$$

## FreeFem++ script to solve Stokes eqs. by P1/P1 stabilized

Find  $(u, p) \in V_h(g) \times Q_h$  s.t.

$$a(u, v) + b(v, p) + b(u, q) - \delta d(p, q) - \epsilon \int_{\Omega} p q = (f, v) \quad \forall (v, q) \in V_h \times Q_h.$$

```
fespace Vh(Th, P1), Qh(Th, P1);
```

▶ example6.edp

```
....
```

```
Vh u1, u2, v1, v2;
```

```
Qh p, q;
```

```
macro d12(u1, u2) (dy(u1) + dx(u2))/2.0 //
```

```
real delta=0.01;
```

```
real epsln=1.0e-6;
```

```
solve stokes(u1, u2, p1, v1, v2, q1) =
```

```
int2d(Th) ( 2.0*(dx(u1)*dx(v1)
```

```
+2.0*d12(u1, u2)*d12(v1, v2)+dy(u2)*dy(v2))
```

```
-p*dx(v1)-p*dy(v2)-dx(u1)*q-dy(u2)*q
```

```
-delta*hTriangle*hTriangle* // stabilization
```

```
(dx(p)*dx(q)+dy(p)*dy(q))
```

```
-p*q*epsln ) // penalization
```

```
- int2d(Th) ( f1 * v1 + f2 * v1 )
```

```
+ on(1, 2, 3, 4, u1=g1, u2=g2);
```

```
....
```



## matrix formulation of discretized form : homogeneous Dirichlet

Find  $(u_h, p_h) \in V_h \times Q_h$  s.t.

$$a(u_h, v_h) + b(v_h, p_h) = (f, v_h) \quad \forall v_h \in V_h,$$

$$b(u_h, q_h) = 0 \quad \forall q_h \in Q_h.$$

finite element bases,  $\text{span}\{\{\phi_i\}\} = V_h$ ,  $\text{span}\{\{\psi_\mu\}\} = S_h$ .

$$\begin{aligned} [A]_{ij} &= a(\phi_j, \phi_i) \\ [B]_{\mu j} &= b(\phi_j, \psi_\mu) \end{aligned} \quad K \begin{bmatrix} \vec{u} \\ \vec{p} \end{bmatrix} = \begin{bmatrix} A & B^T \\ B & 0 \end{bmatrix} \begin{bmatrix} \vec{u} \\ \vec{p} \end{bmatrix} = \begin{bmatrix} \vec{f} \\ \vec{0} \end{bmatrix}$$

$K \in \mathbb{R}^{(N_V+N_S) \times (N_V+N_S)}$  : symmetric, indefinite,  $\text{Ker}K = \begin{bmatrix} \vec{0} \\ \vec{1} \end{bmatrix}$ .

$B \in \mathbb{R}^{N_X \times N_S}$  : on the whole FE nodes of velocity/pressure

$$\begin{aligned} [B^T \vec{1}]_i &= \sum_\mu b(\phi_i, \psi_\mu) = b(\phi_i, \sum_\mu \psi_\mu) \\ &= b(\phi_i, 1) = - \int_\Omega \nabla \cdot \phi_i 1 = \int_\Omega \phi_i \cdot \nabla 1 - \int_{\partial\Omega} \phi_i \cdot n \, ds \\ &= 0 \quad \text{for } i \in \{1, \dots, N_X\} \setminus \Lambda_D. \end{aligned}$$

$b(\cdot, \cdot)$  satisfies inf-sup condition on  $V_h \times S_h \Leftrightarrow \text{Ker}B^T = \{\vec{1}\}$ .

## how to solve linear system of indefinite matrix

$\begin{bmatrix} A & B^T \\ B & 0 \end{bmatrix}$  : symmetric, indefinite, singular :

$\#\{\lambda > 0\} = N_V$ ,  $\#\{\lambda = 0\} = 1$ ,  $\#\{\lambda < 0\} = N_S - 1$ .

- ▶ penalization + direct factorization : **UMFPACK**

$\begin{bmatrix} A & B^T \\ B & -\epsilon M \end{bmatrix}$  : symmetric, indefinite, nonsingular :

$\#\{\lambda > 0\} = N_V$ ,  $\#\{\lambda < 0\} = N_S$ .

$[M]_{\mu\nu} = \int_{\Omega} \psi_{\nu} \psi_{\mu} dx$ ,  $\epsilon > 0$  : penalization parameter.

- ▶ preconditioned **CG** method with orthogonal projection  
Schur complement on pressure (aka Uzawa method)

$$-BA^{-1}B^T\vec{p} = -BA^{-1}\vec{f}$$

$BA^{-1}B^T$  : sym. positive definite on  $\{\vec{q} \in \mathbb{R}^{N_S} ; (\vec{q}, \vec{1}) = 0\}$ .

orthogonal projection  $P : \mathbb{R}^{N_S} \rightarrow \text{span}\{\vec{1}\}^{\perp}$ ,

$P\vec{q} = \vec{q} - (\vec{q}, \vec{1})/(\vec{1}, \vec{1})\vec{1}$        $[\vec{q}]_i = [\vec{q}]_i - \sum_{1 \leq j \leq n} [\vec{q}]_j / n$ .

preconditioner  $[M]_{\mu\nu} = \int_{\Omega} \psi_{\nu} \psi_{\mu} dx$ .

► Uzawa-CG

## FreeFem++ script to generate Stokes matrix

Find  $(u, p) \in V_h(g) \times Q_h$  s.t.

$$a(u, v) + b(v, p) + b(u, q) - \epsilon \int_{\Omega} p q = (f, v) \quad \forall (v, q) \in V_h \times Q_h.$$

```
fespace VQh(Th, [P2, P2, P1]);
... // func f1, f2, g1, g2 etc
Vh u1, u2, v1, v2; Qh p, q;
macro d12(u1, u2) (dy(u1) + dx(u2))/2.0 //
real epsln=1.0e-6;
varf stokes([u1, u2, p], [v1, v2, q]) =
  int2d(Th) ( 2.0*(dx(u1)*dx(v1)
    +2.0*d12(u1, u2)*d12(v1, v1)+dy(u2)*dy(v2))
    -p*dx(v1)-p*dy(v2)-dx(u1)*q-dy(u2)*q
    -p*q*epsln ) // penalization
  + on(1, 2, 3, 4, u1=1.0, u2=1.0);
varf external([u1, u2, p], [v1, v2, q])=
  int2d(Th) (f1 * v1 + f2 *v2)
  + on(1, 2, 3, 4, u1=g1, u2=g2); // Dirichlet data here
matrix A = stokes(VQh, VQh, solver=UMFPACK);
real[int] bc = stokes(0, VQh);
real[int] ff = external(0, VQh);
u1[] = A^-1 * ff;
```

▶ example7.edp

## FreeFem++ script to solve Stokes matrix by Dissection

Find  $(u, p) \in V_h(g) \times Q_h$  s.t.

$$a(u, v) + b(v, p) + b(u, q) = (f, v) \quad \forall (v, q) \in V_h \times Q_h.$$

```
load "Dissection"; // loading dynamic module
defaulttoDissection(); // sparsesolver=Dissection
fespace VQh(Th, [P2,P2,P1]);
Vh u1,u2,v1,v2; Qh p,q;
macro d12(u1,u2) (dy(u1) + dx(u2))/2.0 //
real epsln=1.0e-6;
varf stokes([u1,u2,p], [v1,v2,q]) =
  int2d(Th) ( 2.0*(dx(u1)*dx(v1)
    +2.0*d12(u1,u2)*d12(v1,v2)+dy(u2)*dy(v2))
    -p*dx(v1)-p*dy(v2)-dx(u1)*q-dy(u2)*q)
  + on(1,2,3,4,u1=1.0,u2=1.0); // no penalty term
varf external([u1,u2,p], [v1,v2,q])=p
  int2d(Th) (f1 * v1 + f2 *v2)
  + on(1,2,3,4,u1=g1,u2=g2); // Dirichlet data here
matrix A=stokes(VQh,VQh,solver=sparsesolver,
  strategy=2,tolpivot=1.0e-2); //new parameters
real[int] bc = stokes(0, VQh);
real[int] ff = external(0, VQh);
u1[] = A^-1 * ff;
```

## stationary Navier-Stokes equations and a weak formulation

$$\Omega = (0, 1) \times (0, 1)$$

$$-2\nu \nabla \cdot D(u) + u \cdot \nabla u + \nabla p = f \text{ in } \Omega$$

$$\nabla \cdot u = 0 \text{ in } \Omega$$

$$u = g \text{ on } \partial\Omega$$

►  $V(g) = \{v \in H^1(\Omega)^2; v = g \text{ on } \partial\Omega\}$ ,  $V = V(0)$

►  $Q = L_0^2(\Omega) = \{p \in L^2(\Omega); \int_{\Omega} p \, dx = 0\}$

► outflow

bi/tri-linear forms and weak formulation :

$$a(u, v) = \int_{\Omega} 2\nu D(u) : D(v) \, dx \quad u, v \in H^1(\Omega)^2$$

$$a_1(u, v, w) = \frac{1}{2} \left( \int_{\Omega} (u \cdot \nabla v) \cdot w - (u \cdot \nabla w) \cdot v \, dx \right) \quad u, v, w \in H^1(\Omega)^2$$

$$b(v, p) = - \int_{\Omega} \nabla \cdot v \, p \, dx \quad v \in H^1(\Omega)^2, p \in L^2(\Omega)$$

Find  $(u, p) \in V(g) \times Q$  s.t.

$$a(u, v) + a_1(u, u, v) + b(v, p) = (f, v) \quad \forall v \in V,$$

$$b(u, q) = 0 \quad \forall q \in Q.$$

## trilinear form for the nonlinear term (Temam's trick)

$\nabla \cdot u = 0$ ,  $w \in H_0^1(\Omega)$  or  $u \cdot n = 0$  on  $\partial\Omega \Rightarrow$

$$a_1(u, v, w) = \int_{\Omega} (u \cdot \nabla v) \cdot w \, dx = \frac{1}{2} \left( \int_{\Omega} (u \cdot \nabla v) \cdot w - (u \cdot \nabla w) \cdot v \, dx \right).$$

$$\begin{aligned} \int_{\Omega} (u \cdot \nabla) v \cdot w \, dx &= \int_{\Omega} \sum_i \sum_j u_j (\partial_j v_i) w_i \, dx \\ &= - \int_{\Omega} \sum_{i,j} v_i \partial_j (u_j w_i) \, dx + \int_{\partial\Omega} \sum_{i,j} v_i n_j u_j w_i \, ds \\ &= - \int_{\Omega} \sum_{i,j} v_i (\partial_j u_j) w_i \, dx - \int_{\Omega} \sum_{i,j} v_i u_j \partial_j w_i \, dx \\ &= - \int_{\Omega} \sum_{i,j} u_j (\partial_j w_i) v_i \, dx \\ &= - \int_{\Omega} (u \cdot \nabla) w \cdot v \, dx. \end{aligned}$$

$a_1(u, u, u) = 0 \Rightarrow$  **corecivity** :  $a(u, u) + a_1(u, u, u) \geq \alpha \|u\|^2$ .

## nonlinear system of the stationary solution

$$A(u, p; v, q) = a(u, v) + a_1(u, u, v) + b(v, p) + b(u, q)$$

nonlinear problem:

**Find**  $(u, p) \in V(g) \times Q$  **s.t.**  $A(u, p; v, q) = (f, v) \quad \forall (v, q) \in V \times Q$ .

$a_1(\cdot, \cdot, \cdot)$  : trilinear form,

$$\begin{aligned} a_1(u + \delta u, u + \delta u, v) &= a_1(u, u + \delta u, v) + a_1(\delta u, u + \delta u, v) \\ &= a_1(u, u, v) + a_1(u, \delta u, v) + a_1(\delta u, u, v) + a_1(\delta u, \delta u, v) \end{aligned}$$

$$A(u + \delta u, p + \delta p; v, q) - A(u, p; v, q)$$

$$= a(u + \delta u, v) - a(u, v)$$

$$+ b(v, p + \delta p) - b(v, p) + b(u + \delta u, q) - b(u, q)$$

$$+ a_1(u + \delta u, u + \delta u, v) - a_1(u, u, v)$$

$$= a(\delta u, v) + b(v, \delta p) + b(\delta u, q) + a_1(\delta u, u, v) + a_1(u, \delta u, v) + O(\|\delta u\|^2)$$

**Find**  $(\delta u, \delta p) \in V \times Q$  **s.t.**

$$a(\delta u, v) + b(v, \delta p) + b(\delta u, q) + a_1(\delta u, u, v) + a_1(u, \delta u, v) =$$

$$- A(u, p; v, q) \quad \forall (v, q) \in V \times Q$$

## Newton iteration

$$(u_0, p_0) \in V(g) \times Q$$

loop  $n = 0, 1 \dots$

Find  $(\delta u, \delta p) \in V \times Q$  s.t.

$$a(\delta u, v) + b(v, \delta p) + b(\delta u, q) + a_1(\delta u, u_n, v) + a_1(u_n, \delta u, v) = A(u_n, p_n; v, q) \quad \forall (v, q) \in V \times Q$$

if  $\|(\delta u, \delta p)\|_{V \times Q} \leq \varepsilon$  then break

$$u_{n+1} = u_n - \delta u,$$

$$p_{n+1} = p_n - \delta p.$$

loop end.

► example8.edp

$(u^{(0)}, p^{(0)}) \in V(g) \times Q$  : solution of the Stokes eqs.,  $\nu = 1$ .

while  $(\nu > \nu_{\min})$

Newton iteration  $(u^{(k+1)}, p^{(k+1)}) \in V(g) \times Q$  from  $(u^{(k)}, p^{(k)})$ .

$$\nu = \nu/2, k++.$$

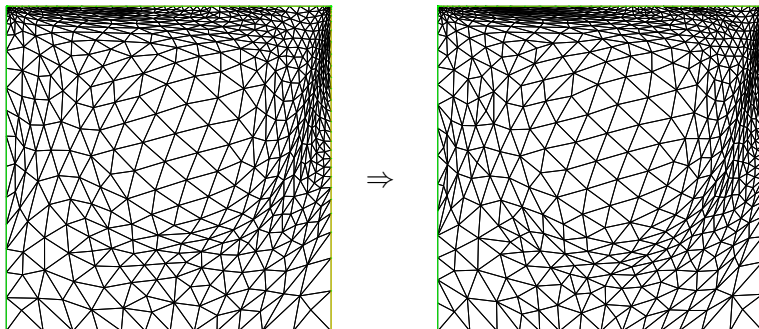
while end.

initial guess from the stationary state of lower Reynolds number



## mesh adaptation

```
fespace XXMh(Th, [P2,P2,P1]);  
XXMh [u1,u2,p];  
real lerr=0.01;  
Th=adaptmesh(Th, [u1,u2], p, err=lerr, nbvx=100000);  
[u1,u2,p]=[u1,u2,p]; // interpolation on the new mesh
```



`err` :  $P_1$  interpolation error level

`nbvx` : maximum number of vertexes to be generated.

## stream line for visualization of 2D flow : 1/2

$$\text{stream function } \psi : \Omega \rightarrow \mathbb{R}, u = \begin{bmatrix} \partial_2 \psi \\ -\partial_1 \psi \end{bmatrix} \Leftrightarrow u \perp \nabla \psi.$$

$$-\nabla^2 \psi = \nabla \times u = \partial_1 u_2 - \partial_2 u_1 \quad \text{in } \Omega$$

boundary conditions for the stream line:

$$0 = \int_0^x u_2(t, 0) dt = \int_0^x -\partial_1 \psi(t, 0) dt = -\psi(x, 0) + \psi(0, 0)$$

$$\psi(x, 0) = \psi(0, 0).$$

$$0 = \int_0^y u_1(t, 0) dt = \int_0^y \partial_2 \psi(0, t) dt = \psi(0, y) + \psi(0, 0)$$

$$\psi(0, y) = \psi(0, 0).$$

$$0 = \int_0^x u_2(t, 1) dt = \int_0^x -\partial_1 \psi(t, 1) dt = -\psi(x, 1) + \psi(0, 1)$$

$$\psi(x, 1) = \psi(0, 1) = \psi(0, 0).$$

$$\Rightarrow \psi = 0 \text{ on } \partial\Omega.$$

## stream line for visualization of 2D flow : 2/2

$$u = \begin{bmatrix} u_1 \\ u_2 \end{bmatrix} \in \text{P2 finite element space} \quad \Rightarrow \quad (\partial_1 u_2 - \partial_2 u_1) \in \text{P1}.$$

```
fespace Xh(Th,P2);
```

```
fespace Mh(Th,P1);
```

```
Xh u1, u2;          // computed from Navier-Stokes solver
```

```
Mh psi, phi;       // dy(u1), dx(u2) are polynomials of 1st
```

```
solve streamlines(psi, phi, solver=UMFPACK) =
```

```
    int2d(Th) ( dx(psi)*dx(phi) + dy(psi)*dy(phi) )
```

```
    + int2d(Th) ( (dx(u2)-dy(u1))*phi )
```

```
    + on(1,2,3,4,psi=0);
```

```
plot(psi, nbiso=30);
```

## Application of finite element method to fluid problems

### Time-dependent Navier-Stokes equations

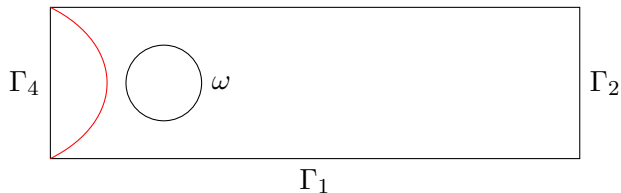
- ▶ material derivative is approximated by Characteristic Galerkin method
- ▶ functional space of pressure depends on boundary conditions of flow, e.g., inflow, non-slip, slip, and outflow.

### Thermal convection problem by Rayleigh-Bénard equations

- ▶ time-dependent problems for fluid and temperature by convection are solved by Characteristic Galerkin method
- ▶ stationary solution is obtained by Newton iteration using an initial value obtained from time-evolutionary solution

## incompressible flow around a cylinder : boundary conditions

$$\Omega = (-1, 9) \times (-1, 1) \quad \Gamma_3$$



$$\frac{\partial u}{\partial t} + u \cdot \nabla u - 2\nu \nabla \cdot D(u) + \nabla p = 0 \text{ in } \Omega$$

$$\nabla \cdot u = 0 \text{ in } \Omega$$

$$u = g \text{ on } \partial\Omega$$

boundary conditions:

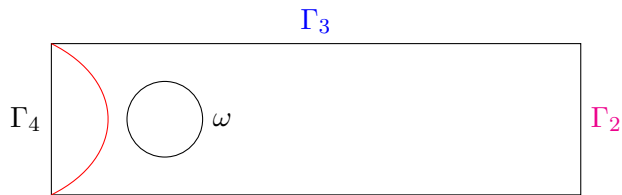
Poiseuille flow on  $\Gamma_4$  :  $u = (1 - y^2, 0)$ .

slip boundary condition on  $\Gamma_1 \cup \Gamma_3$  :  $\begin{cases} u \cdot n = 0 \\ (2\nu D(u)n - np) \cdot t = 0 \end{cases}$

no-slip boundary condition on  $\omega$  :  $u = 0$

outflow boundary condition on  $\Gamma_2$  :  $2\nu D(u)n - np = 0$

## slip boundary conditions and function space



slip boundary condition on  $\Gamma_1 \cup \Gamma_3$  :  $\begin{cases} u \cdot n = 0 \\ (2\nu D(u)n - np) \cdot t = 0 \end{cases}$

▶  $V(g) = \{v \in H^1(\Omega)^2; v = g \text{ on } \Gamma_4 \cup \omega, v \cdot n = 0 \text{ on } \Gamma_1 \cup \Gamma_3\}$ ,

▶  $Q = L^2(\Omega)$ .

▶ non-slip

$$\begin{aligned} \int_{\Gamma_1 \cup \Gamma_3} (2\nu D(u)n - np) \cdot v ds &= \int_{\Gamma_1 \cup \Gamma_3} (2\nu D(u)n - np) \cdot (v_n n + v_t t) ds \\ &= \int_{\Gamma_1 \cup \Gamma_3} (2\nu D(u)n - np) \cdot (v \cdot n) n ds \\ &\quad + \int_{\Gamma_1 \cup \Gamma_3} (2\nu D(u)n - np) \cdot t v_t ds = 0 \end{aligned}$$

## characteristic line and material derivative

$u(x_1, x_2, t) : \Omega \times (0, T] \rightarrow \mathbb{R}^2$ , given velocity field.

$\phi(x_1, x_2, t) : \Omega \times (0, T] \rightarrow \mathbb{R}$ .

$X(t) : (0, T] \rightarrow \mathbb{R}^2$ , characteristic curve :

$$\frac{dX}{dt}(t) = u(X(t), t), X(0) = X_0$$

$$\begin{aligned}\frac{d}{dt}\phi(X(t), t) &= \nabla\phi(X(t), t) \cdot \frac{d}{dt}X(t) + \frac{\partial}{\partial t}\phi(X(t), t) \\ &= \nabla\phi(X(t), t) \cdot u(X(t), t) + \frac{\partial}{\partial t}\phi(X(t), t)\end{aligned}$$

material derivative :  $\frac{D\phi}{Dt} = \frac{\partial}{\partial t}\phi + u \cdot \nabla\phi$ .

approximation by difference

$$\frac{D\phi(X(t), t)}{Dt} \sim \frac{\phi(X(t), t) - \phi(X(t - \Delta t), t - \Delta t)}{\Delta t}$$

## characteristic Galerkin method to discretize material derivative

approximation by Euler method :

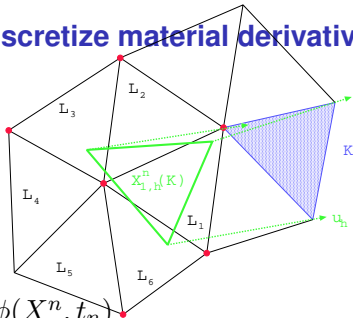
$$t_n < t_{n+1}, t_{n+1} = \Delta t + t_n.$$

$$X(t_{n+1}) = x$$

$$X(t_n) = X^n(x) + O(\Delta t^2)$$

$$X^n(x) = x - u(x, t_n)$$

$$\begin{aligned} \frac{D\phi(X(t_{n+1}), t_{n+1})}{Dt} &= \frac{\phi(x, t_{n+1}) - \phi(X^n, t_n)}{\Delta t} + O(\Delta t) \\ &\sim \frac{\phi^{n+1} - \phi^n \circ X^n}{\Delta t}. \end{aligned}$$



$u^n$  : obtained in the previous time step.

Find  $(u^{n+1}, p^{n+1}) \in V(g) \times Q$  s.t.

$$\left( \frac{u^{n+1} - u^n \circ X^n}{\Delta t}, v \right) + a(u^{n+1}, v) + b(v, p^{n+1}) = 0 \quad \forall v \in V,$$

$$b(u^{n+1}, q) = 0 \quad \forall q \in Q.$$



## FreeFem++ script using characteristic Galerkin method

FreeFem++ provides `convect` to compute  $(u^n \circ X^n, \cdot)$ .

```
real nu=1.0/Re;
real alpha=1.0/dt;
int i;
problem NS([u1,u2,p],[v1,v2,q],solver=UMFPACK,init=i) =
  int2d(Th)(alpha*(u1*v1 + u2*v2)
    +2.0*nu*(dx(u1)*dx(v1)+2.0*d12(u1,u2)*d12(v1,v2)
    +dy(u2)*dy(v2))
    - p * div(v1, v2) - q * div(u1, u2))
- int2d(Th)(alpha*( convect([up1,up2],-dt,up1)*v1
    +convect([up1,up2],-dt,up2)*v2) )
+ on(1,3,u2=0)+on(4,u1=1.0-y*y,u2=0)+on(5,u1=0,u2=0);

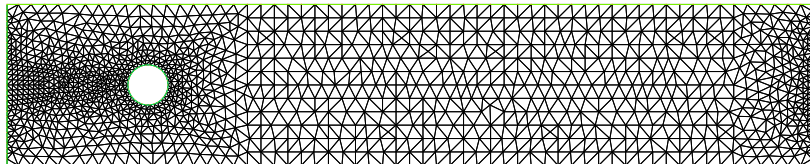
for (i = 0; i <= timestepmax; i++) {
  up1 = u1; up2 = u2; pp = p;
  NS; // factorization is called when i=0
  plot([up1,up2],pp,wait=0,value=true,coef=0.1);
}
```

▶ example9.edp

## FreeFem++ script for mesh generation around a cylinder

Delaunay triangulation from nodes given on the boundary  
boundary segments are oriented and should be connected.

```
int n1 = 30;
int n2 = 60;
border ba(t=0,1.0){x=t*10.0-1.0;y=-1.0;label=1;};
border bb(t=0,1.0){x=9.0;y=2.0*t-1.0;label=2;};
border bc(t=0,1.0){x=9.0-10.0*t;y=1.0;label=3;};
border bd(t=0,1.0){x=-1.0;y=1.0-2.0*t;label=4;};
border cc(t=0,2*pi){x=cos(t)*0.25+0.75;
                    y=sin(t)*0.25;label=5;};
mesh Th=buildmesh(ba(n2)+bb(n1)+bc(n2)+bd(n1)+cc(-n1));
plot(Th);
```



## stream line for visualization of flow around a cylinder : 1/2

stream function  $\psi : \Omega \rightarrow \mathbb{R}$ ,  $u = \begin{bmatrix} \partial_2 \psi \\ -\partial_1 \psi \end{bmatrix}$ .

boudnary conditions for the stream line:

$$\text{inlet: } y - \frac{y^3}{3} = \int_0^y u_1(x_1, t) dt = \int_0^y \partial_2 \psi(x_1, t) dt = \psi(x_1, y) - \psi(x_1, 0)$$

$$\psi(x_1, y) = \psi(x_1, 0) + y - \frac{y^3}{3} = y - \frac{y^3}{3}.$$

$$\text{slip: } 0 = \int_{x_1}^x u_2(t, \pm 1) dt = \int_{x_1}^x -\partial_1 \psi(t, \pm 1) dt = \psi(x_1, \pm 1) - \psi(x, \pm 1)$$

$$\psi(x, \pm 1) = \psi(x_1, \pm 1) = \psi(x_1, 0) \pm 2/3 = \pm 2/3.$$

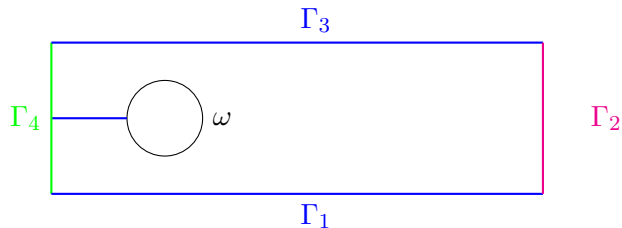
$$\text{cylinder: } 0 = \int_{-\pi}^{\theta} u \cdot n d\theta = \int_{-\pi}^{\theta} -\partial_1 \psi r \sin \theta + \partial_2 \psi r \cos \theta$$

$$= \int_{-\pi}^{\theta} \frac{\partial}{\partial \theta} \psi(r, \theta) d\theta.$$

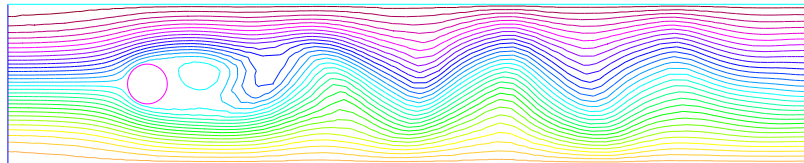
center from inlet:  $u_2 = 0 \Rightarrow$  same as slip wall ,

$$\psi|_{\omega} = \psi(x_1, 0) = 0.$$

## stream line for visualization of flow around a cylinder : 2/2



slip boundary condition on  $\Gamma_1 \cup \Gamma_3$ , outflow on  $\Gamma_2$ .



## thermal convection in a box : 1/2

$$\Gamma_3 : \theta = \theta_0, u_2 = 0$$

$$\Gamma_4 : \partial_n \theta = 0, u_1 = 0$$

$$\Gamma_2 : \partial_n \theta = 0, u_1 = 0$$

$$\Gamma_1 : \theta = \theta_0 + \Delta\theta, u_2 = 0$$

Rayleigh-Bénard equations

$$\rho \left( \frac{\partial u}{\partial t} + u \cdot \nabla u \right) - 2\nabla \cdot \mu_0 D(u) + \nabla p = -\rho g \vec{e}_2 \text{ in } \Omega,$$

$$\nabla \cdot u = 0 \text{ in } \Omega,$$

$$\frac{\partial \theta}{\partial t} + u \cdot \nabla \theta - \nabla \cdot (\kappa \theta) = 0 \text{ in } \Omega.$$

$\vec{e}_2$  : unit normal of  $y$ -direction

$d$  : height of the box,  $g$  : gravity acceleration,

$\kappa$  : thermal diffusivity,  $\mu_0$  : viscosity

## thermal convection in a box : 2/2

Boussinesq approximation :  $\rho = \rho_0\{1 - \alpha(\theta - \theta_0)\}$ ,  $\theta_0 = 0$ .

$\rho_0$ : representative density,  $\alpha$  : thermal expansion coefficient.

non-dimensional Rayleigh-Bénard equations

$$\frac{1}{Pr} \left( \frac{\partial u}{\partial t} + u \cdot \nabla u \right) - 2\nabla \cdot D(u) + \nabla p = Ra\theta\vec{e}_2 \text{ in } \Omega,$$

$$\nabla \cdot u = 0 \text{ in } \Omega,$$

$$\frac{\partial \theta}{\partial t} + u \cdot \nabla \theta - \Delta \theta = 0 \text{ in } \Omega$$

$$u \cdot n = 0 \text{ on } \partial\Omega,$$

$$\theta = 1 \text{ on } \Gamma_1,$$

$$\theta = 0 \text{ on } \Gamma_3,$$

$$\partial_n \theta = 0 \text{ on } \Gamma_2 \cup \Gamma_4.$$

- ▶  $Pr = \frac{\mu_0}{\kappa\rho_0}$ : Prandtl number,
- ▶  $Ra = \frac{\rho_0 g \alpha \Delta \theta d^3}{\kappa \mu_0}$ : Rayleigh number.

## a weak form to solve time-dependent Rayleigh-Bénard eqs.

- ▶ velocity :  $V = \{v \in H^1(\Omega)^2; v \cdot n = 0 \text{ on } \partial\Omega\}$ ,
- ▶ pressure :  $Q = L_0^2(\Omega) = \{p \in L^2(\Omega); \int_{\Omega} p \, dx = 0\}$ ,
- ▶ temperature :  $\Psi_D = \{\theta \in H^1(\Omega); \theta = 1 \text{ on } \Gamma_1, \theta = 0 \text{ on } \Gamma_0\}$ .

bilinear forms:

$$a_0(u, v) = \int_{\Omega} 2D(u) : D(v), \quad b(v, p) = - \int_{\Omega} \nabla \cdot v p,$$

$$c_0(\theta, \psi) = \int_{\Omega} \nabla \theta \cdot \nabla \psi.$$

using Characteristic Galerkin method:

▶ example10.edp

$(u^n, \theta^n) \in V \times \Psi_D$  : from previous time step

Find  $(u^{n+1}, p^{n+1}, \theta^{n+1}) \in V \times Q \times \Psi_D$  s.t.

$$\frac{1}{Pr} \left( \frac{u^{n+1} - u^n \circ X^n}{\Delta t}, v \right) + a_0(u^{n+1}, v) + b(v, p^{n+1}) = Ra(\theta^n \vec{e}_2, v) \quad \forall v \in V,$$

$$b(u^{n+1}, q) = 0 \quad \forall q \in Q,$$

$$\left( \frac{\theta^{n+1} - \theta^n \circ X^n}{\Delta t}, \psi \right) + c_0(\theta^{n+1}, \psi) = 0 \quad \forall \psi \in \Psi_0.$$

## a weak form to solve stationary Rayleigh-Bénard eqs.

trilinear forms and bilinear form for the Navier-Stokes eqs.

$$\blacktriangleright a_1(u, v, w) = \frac{1}{2Pr} (\int_{\Omega} (u \cdot \nabla v) \cdot w - (u \cdot \nabla w) \cdot v)$$

$$\blacktriangleright c_1(u, \theta, \psi) = \frac{1}{2} (\int_{\Omega} (u \cdot \nabla \theta) \cdot \psi - (u \cdot \nabla \psi) \cdot \theta)$$

$$\blacktriangleright A(u, p; v, q) = a_0(u, v) + a_1(u, u, v) + b(v, p) + b(u, q)$$

Newton iteration  $(u_0, p_0, \theta_0) \in V \times Q \times \Psi_D$

▶ example11.edp

loop  $n = 0, 1 \dots$

Find  $(\delta u, \delta p, \delta \theta) \in V \times Q \times \Psi_0$  s.t.

$$a_0(\delta u, v) + b(v, \delta p) + b(\delta u, q) + a_1(\delta u, u_n, v) + a_1(u_n, \delta u, v)$$

$$- Ra(\delta \theta \vec{e}_2, v) = A(u_n, p_n; v, q) - Ra(\theta_n \vec{e}_2, v) \quad \forall (v, q) \in V \times Q$$

$$c_0(\delta \theta, \psi) + c_1(u_n, \delta \theta, \psi) + c_1(\delta u, \theta_n, \psi) = c_0(\theta_n, \psi) + c_1(u_n, \theta_n, \psi)$$

$$\forall \psi \in \Psi_0$$

if  $\|(\delta u, \delta p, \delta \theta)\|_{V \times Q \times \Psi} \leq \varepsilon$  then break

$$u_{n+1} = u_n - \delta u, \quad p_{n+1} = p_n - \delta p, \quad \theta_{n+1} = \theta_n - \delta \theta.$$

loop end.

initial data  $\Leftarrow$  stationary solution by time-dependent problem.



## Details on iterative linear solver

FreeFem++ provides iterative solvers for symmetric positive definite matrix and general unsymmetric invertible matrix.

- ▶ Conjugate Gradient method `LinearCG`  
`src/femlib/MatriceCreuse.hpp::ConjuguedGradient2()`
- ▶ Generalized Minimal Residual method `LinearGMRES`  
`src/femlib/gmres.hpp::GMRES()`

They are useful to get solution with less memory consumption.

- ▶ treatment of boundary condition is somewhat different from penalization technique for direct solver
- ▶ definition of SpMV (Sparse matrix vector multiplication) operator:  $y = Ax$ , and preconditioning operator by  
`func real[int] SpMV(real[int] &x);`

To use good preconditioner is very important for faster convergence.

- ▶ preconditioner for time-dependent generalized Stokes equations

## conjugate gradient method

$$A_\tau \vec{u} = \vec{f}_\tau, \quad [A_\tau]_{kk} = \tau, [f_\tau]_k = \tau g_k \text{ for } k \in \Lambda_D.$$

preconditioner  $Q \sim A_\tau^{-1}$

Krylov subsp. :

$$K_n(Q\vec{r}^0, QA_\tau) = \text{span}[Q\vec{r}^0, QA_\tau Q\vec{r}^0, \dots, (QA_\tau)^n Q\vec{r}^0]$$

Find  $\vec{u}^n \in K_n(Q\vec{r}^0, QA_\tau) + \vec{u}^0$  s.t.

$$(A\vec{u}^n - \vec{f}_\tau, \vec{v}) = 0 \quad \vec{v} \in K_n(Q\vec{r}^0, QA_\tau).$$

### Preconditioned CG method

$\vec{u}^0$  : initial step for CG.

$$\vec{r}^0 = \vec{f}_\tau - A_\tau \vec{u}^0$$

$$\vec{p}^0 = Q\vec{r}^0.$$

loop  $n = 0, 1, \dots$

$$\alpha_n = (Q\vec{r}^n, \vec{r}^n) / (A_\tau \vec{p}^n, \vec{p}^n),$$

$$\vec{u}^{n+1} = \vec{u}^n + \alpha_n \vec{p}^n,$$

$$\vec{r}^{n+1} = \vec{r}^n - \alpha_n A_\tau \vec{p}^n,$$

if  $\|\vec{r}^{n+1}\| < \epsilon$  exit loop.

$$\beta_n = (Q\vec{r}^{n+1}, \vec{r}^{n+1}) / (Q\vec{r}^n, \vec{r}^n),$$

$$\vec{p}^{n+1} = Q\vec{r}^{n+1} + \beta_n \vec{p}^n.$$

LinearCG(opA, u, f, precon=opQ, nbiter=100, eps=1.0e-10)

```
Vh u, v;
varf aa(u, v) = int2d(Th) ( dx(u) * dx(v) + dy(u) * dy(v) )
                + on(2, 3, 4, u = 1.0);
varf external(u, v) = int2d(Th) ( f * v ) + int1d(Th, 1) ( h * v )
real tgv = 1.0e+30;    matrix A;
real[int] bc = aa(0, Vh, tgv = tgv);
func real[int] opA(real[int] &pp) { // SpMV operation with
    pp = bc ? 0.0 : pp;                // homogeneous data
    real[int] qq = A * pp;
    pp = bc ? 0.0 : qq; return pp; } // qq: locally allocated
func real[int] opQ(real[int] &pp) {
    for (int i = 0; i < pp.n; i++)
        pp(i) = pp(i) / A(i, i);
    pp = bc ? 0.0 : pp; return pp; }
A = aa(Vh, Vh, tgv = tgv, solver = sparsesolver);
real[int] ff = external(0, Vh);
u = bc ? v[] : 0.0;    // v : Dirichlet data without tgv
v[] = A * u[]; ff -= v[]; // Dirichlet data goes to RHS
ff = bc ? 0.0 : ff;    // CG works in zero-Dirichlet
LinearCG(opA, u[], ff, precon = opQ, nbiter = 100, eps = 1.0e-10);
```

## conjugate gradient method on the image space

$$\vec{f} \in \text{Im}A, \quad \text{find } u \in \text{Im}A \quad A\vec{u} = \vec{f}.$$

preconditioner  $Q : \text{Im}A \rightarrow \text{Im}A$ ,  $Q \sim A|_{\text{Im}A}^{-1}$

orthogonal projection  $P : \mathbb{R}^n \rightarrow \text{Im} A$ .

### Preconditioned CG method

$\vec{u}^0$  : initial step for CG.

$$\vec{r}^0 = P(\vec{f} - A\vec{u}^0)$$

$$\vec{p}^0 = Q\vec{r}^0.$$

loop  $n = 0, 1, \dots$

$$\alpha_n = (Q\vec{r}^n, \vec{r}^n) / (A\vec{p}^n, \vec{p}^n),$$

$$\vec{u}^{n+1} = P(\vec{u}^n + \alpha_n \vec{p}^n),$$

$$\vec{r}^{n+1} = P(\vec{r}^n - \alpha_n A\vec{p}^n),$$

if  $\|\vec{r}^{n+1}\| < \epsilon$  exit loop.

$$\beta_n = (Q\vec{r}^{n+1}, \vec{r}^{n+1}) / (Q\vec{r}^n, \vec{r}^n),$$

$$\vec{p}^{n+1} = P(Q\vec{r}^{n+1} + \beta_n \vec{p}^n).$$

$P$  is used to avoid numerical round-off error which perturbs vectors from the image space

LinearCG cannot handle this safe operation.  $PQ$  and  $PA$  only.

## FreeFem++ script for full-Neumann problem

▶ example13.edp

```
Vh u, v;
varf aa(u, v) = int2d(Th) ( dx(u) * dx(v) + dy(u) * dy(v) );
varf external(u, v) = int2d(Th) ( f * v );
matrix A;
func real[int] opA(real[int] &pp) {
  real[in] qq = A * pp;
  pp = qq; pp -= pp.sum / pp.n; // projection
  return pp;
}
func real[int] opQ(real[int] &pp) {
  for (int i = 0; i < pp.n; i++)
    pp(i) = pp(i) / A(i, i);
  pp -= pp.sum / pp.n; // projection
  return pp;
}
A = aa(Vh, Vh, solver = CG); real[int] ff = external(0, Vh);
ff -= ff.sum / ff.n; // projection
u[] = 0.0; // initial step for CG
LinearCG(opA, u[], ff, precon = opQ, nbiter = 100, eps = 1.0e-10);
```

## conjugate gradient in Uzawa method for Stokes eqs.

► Stokes Solver

$$\begin{bmatrix} A_\tau & B^T \\ B & 0 \end{bmatrix} \begin{bmatrix} \vec{u} \\ \vec{p} \end{bmatrix} = \begin{bmatrix} \vec{f}_\tau \\ \vec{0} \end{bmatrix} \quad [A_\tau]_{kk} = \tau, \quad [\vec{f}_\tau]_k = \tau g_k \quad \text{for } k \in \Lambda_D.$$

orthogonal projection  $P : \mathbb{R}^{N_S} \rightarrow \text{span}\{[\vec{1}]\}^\perp$ , preconditioner  $Q$   
 $(BA^{-1}B^T)^{-1} \sim I_h^{-1} = Q$ : inverse of mass matrix.

**Preconditioned CG method with projection**

$\vec{p}^0 = \vec{0}$ : initial step for CG.

$$\vec{g}^0 = PBA_\tau^{-1}\vec{f}_\tau,$$

$$\vec{w}^0 = PQ\vec{g}^0.$$

loop  $n = 0, 1, \dots$

$$\alpha_n = (PQ\vec{g}^n, \vec{g}^n) / (PBA_\tau^{-1}B^T\vec{w}^n, \vec{w}^n),$$

$$\vec{p}^{n+1} = \vec{p}^n + \alpha_n \vec{w}^n,$$

$$\vec{g}^{n+1} = \vec{g}^n - \alpha_n (BA_\tau^{-1}B^T)\vec{w}^n,$$

$$\beta_n = (PQ\vec{g}^{n+1}, \vec{g}^{n+1}) / (\vec{g}^n, \vec{g}^n),$$

$$\vec{w}^{n+1} = PQ\vec{g}^{n+1} + \beta_n \vec{w}^n.$$

$$\vec{u}^{n+1} = A_\tau^{-1}(\vec{f}_\tau - B^T\vec{p}^{n+1}).$$

$$A_\tau^{-1}\vec{f}_\tau \Leftrightarrow A_\tau\vec{u} = \vec{f}_\tau \quad \text{with } u_k = g_k, \quad k \in \Lambda_D$$

► penalty

$$A_\tau^{-1}B^T\vec{w} \Leftrightarrow A_\tau\vec{u} = B^T\vec{w} \quad \text{with } u_k = 0, \quad k \in \Lambda_D$$

## FreeFem++ script for CG with Uzawa 1/2

▶ example14.edp

```
fespace Vh(Th, [P2,P2]), Qh(Th, P1);
... // func f1, f2, g1, g2 etc
Vh [u1, u2], [v1, v2], [bcsol1, bcsol2];
Qh p, q;
macro d12(u1, u2) (dy(u1) + dx(u2))/2.0 //

varf a([u1, u2], [v1, v2]) =
  int2d(Th) ( 2.0*(dx(u1)*dx(v1)
    +2.0*d12(u1, u2)*d12(v1, v1)+dy(u2)*dy(v2))
  + on(1, 2, 3, 4, u1=g1, u2=g2));
varf b([u1, u2], [q])= int2d(Th) (- q*(dx(u1)+dy(u2)));
varf external([u1, u2], [v1, v2])=
  int2d(Th) (f1 * v1 + f2 *v2);
varf massp(p, q)= int2d(Th) (p * q);
matrix A = a(Vh, Vh, solver=UMFPACK, init=true);
matrix B = b(Vh, Qh);
matrix Mp = massp(Qh, Qh, solver=UMFPACK, init=true);
real[int] bc = a(0, Vh);
real[int] ff = external(0, Vh);
```

## FreeFem++ script for CG with Uzawa 2/2

```
func real[int] UzawaStokes(real[int] &pp) {
    real[int] b = B'*pp;
    real[int] uu = A^-1 * b;
    pp = B * uu;  pp -= pp.sum / pp.n;
    return pp;
}
func real[int] PreconMass(real[int] &pp) {
    real[int] ppp = Mp^-1 * pp;
    pp = ppp; pp -= pp.sum / pp.n;
    return pp;
}
p = 0.0;
ff += bc;           //bc keeps Dirichlet data with tgv
real[int] uu = A^-1 * ff;
q[] = B * uu;
LinearCG(UzawaStokes, p[], q[], precon=PreconMass,
         nbiter=100,eps=1.0e-10,verbosity=100);
ff = external(0, Vh); real[int] b = B'*p[];
ff -= b; ff += bc; //bc keeps Dirichlet data with tgv
u1[] = A^-1 * ff;   // to access [u1, u2]
```



descretized Navier-Stokes equations by characteristic Galerkin

$\Delta t$  : time step,  $\nu$  : Reynolds number

Find  $(u^{n+1}, p^{n+1}) \in V(g) \times Q$  s.t.

$$\left( \frac{u^{n+1}}{\Delta t}, v \right) + a(\nu; u^{n+1}, v) + b(v, p^{n+1}) = - \left( \frac{u^n \circ X^n}{\Delta t}, v \right) \quad \forall v \in V,$$

$$b(u^{n+1}, q) = 0 \quad \forall q \in Q.$$

- ▶  $I_v, I_p$  : mass matrix for velocity, pressure
- ▶  $A_p$  : stiffness matrix of Laplacian for pressure with B.C.

$$\begin{bmatrix} \frac{1}{\Delta t} I_v + \nu A & B^T \\ B & 0 \end{bmatrix} \begin{bmatrix} \vec{u} \\ \vec{p} \end{bmatrix} = \begin{bmatrix} \vec{f} \\ \vec{0} \end{bmatrix}$$

Preconditioner by Cahouet-Chabard [1988]

$$\left( B \left( \frac{1}{\Delta t} I_v + \nu A \right)^{-1} B^T \right)^{-1} \sim \frac{1}{\Delta t} A_p^{-1} + \nu I_p^{-1}$$

Uzawa with CG : to compute large problem with less memory

## syntax of FreeFem++ script

### loops

```
for (int i=0; i<10; i++) {
    ...
    if (err < 1.0e-6) break;
}

int i = 0;
while (i < 10) {
    ...
    if (err < 1.0e-6) break;
    i++;
}
```

### finite element space, variational form, and matrix

```
fespace Xh(Th,P1)
Xh u,v;           // finite element data
varf a(u,v)=int2d(Th) ( ... ) ;
matrix A = a(Xh,Xh,solver=UMFPACK);
real [int] v;    // array
v = A*u[];      // multiplication matrix to array
```

### procedure (function)

```
func real[int] ff(real[int] &pp) { // C++ reference
    ...
    return pp;                       // the same array
}
```

## array, vector, FEM data, sparse matrix, block data : 1/2

### fundamental data types

```
bool flag; // true or false
int i;
real w;
string st = "abc";
```

### array

```
real[int] v(10); // real array whose size is 10
real[int] u; // not yet allocated
u.resize(10); // same as C++ STL vector
real[int] vv = v; // allocated as same size of v.n
a(2)=0.0 ; // set value of 3rd index
a += b; // a(i) = a(i) + b(i)
a = b .* c ; // a(i) = b(i) * c(i); element-wise
a = b < c ? b : c // a(i) = min(b(i), c(i)); C-syntax
a.sum; // sum a(i);
a.n; // size of array
```

There are other operations such as  $\ell^1, \ell^2, \ell^\infty$ -norms, max, min.  
cf. Finite element analysis by mathematical programming  
language FreeFem++, Ohtsuka-Takaishi [2014].

## array, vector, FEM data, sparse matrix, block data : 2/2

### FEM data

```
func fnc = sin(pi*x)*cos(pi*y); // function with x,y
mesh Th = ...;
fespace Vh(Th,P2);           // P2 space on mesh Th
Vh f;                        // FEM data on Th with P2
f[];                          // access data of FEM DOF
f = fnc;                      // interpolation onto FEM space
fespace Vh(Th,[P2,P2]);     // 2 components P2 space
Vh [u1,u2];                  // u1[], u2[] is allocated
u1[]=0.0;                    // access all data of [u1,u2];
real[int] uu([u1[].n+u2[].n);
u1[] = uu;                   // u1[], u2[] copied from uu
[u1[], u2[]] = uu;          // using correct block data
```

### dense and sparse matrices

```
real[int,int] B(10,10); // 2D array
varf aa(u,v)=int2d(Th)(u*v); // L2-inner prod. for mass
matrix A=aa(Vh,Vh,solver=sparseovler); //sparse matrix
```

file I/O is same as C++, ofstream/ifstream

▶ example{10,11}.edp

## Compilation with configure : 1/2

- ▶ download the latest source from `http://www.freefem.org/ff++/`
- ▶ run `configure` script.  
% `./configure --enable-m64 CXXFLAGS=-std=c++11 --enable-download`  
this enables automatic downloading of all sources including MUMPS etc.
- ▶ run `make`.  
% `make`
- ▶ binaries will be created in `src/nw`

GNU bison and flex are necessary for FreeFem+++ language parser.

OpenGL compatible libraries are also necessary for `ffglut` graphics viewer.

Other options to minimize the capability,

`--disable-superlu --disable-scotch --without-mpi`

## Compilation with configure : 2/2

Fortran is mandatory for MUMPS linear solver.

Without Fortran compiler, by adding

```
--disable-fortran --disable-mumps
```

It is necessary to remove `mumps-seq` from `LIST_SOFT` of `download/Makefile` and

to remove `ffnewuoa.$(DYLIB_SUFFIX)` from

`LIST_COMPILE_PKG` of `example++-load/Makefile` when Fortran is disabled.

BLAS library is automatically detected by `configure` and information is written in

`examples++-load/WHERE_LIBRARY-config.`

`$(INTEL_MKL)` is described as appropriate directory:

```
lapack LD -L$(INTEL_MKL)/lib/intel64 -lmkl_rt \
```

```
    -lmkl_sequential -lmkl_core -liomp5 -lpthread
```

```
lapack INCLUDE -I$(INTEL_MKL)/include
```

```
mkl LD -L$(INTEL_MKL)/lib/intel64 -lmkl_rt \
```

```
    -lmkl_intel_thread -lmkl_core -liomp5 -lpthread
```

```
mkl INCLUDE -I$(INTEL_MKL)/include
```

```
blas LD -L$(INTEL_MKL)/mkl/lib/intel64 \
```

```
    -lmkl_rt -lmkl_sequential -lmkl_core -liomp5 -lpthread
```

## References : 1/2

### FreeFem++:

- ▶ F. Hecht, FreeFem++ manual, 3rd ed., 2015.
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### Finite element theory:

- ▶ D. Braess, Finite elements – Theory, fast solvers and application in solid mechanics, 3rd ed., Cambridge Univ. Press, 2007.
- ▶ A. Ern, J.-L. Guermond, Theory and practice of finite elements, Springer Verlag, New-York, 2004.
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## References : 2/2

specialized topics:

- ▶ T. A. Davis, Direct Methods for Sparse Linear Systems, SIAM, Philadelphia, 2006.
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- ▶ A.H. Stroud, Approximate calculation of multiple integrals, Prentice-Hall, 1971.
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## Appendix: Lagrange multiplier approach for full-Neumann problem

full-Neumann boundary problem

$$-\Delta u = f \text{ in } \Omega,$$

$$\partial_n u = h \text{ on } \partial\Omega.$$

- ▶ compatibility condition :  $\int_{\Omega} f + \int_{\partial\Omega} h = 0$
- ▶  $[A]_{ij} = a(\varphi_j, \varphi_i)$ .  $A$  : singular ,  $\text{Ker}A = \vec{1}$ .
- ▶  $[\vec{b}]_i = F(\varphi_i)$  : compatibility condition  $\Leftrightarrow \vec{b} \in \text{Im}A = (\text{Ker}A)^{\perp}$ .

solution in image of  $A$  : find  $\vec{u} \in \text{Im}A$   $A\vec{u} = \vec{b}$

Lagrange multiplier to deal with constraint  $(\vec{x}, \vec{1}) = 0$ .

$$\begin{bmatrix} A & \vec{1} \\ \vec{1}^T & 0 \end{bmatrix} \begin{bmatrix} \vec{u} \\ \lambda \end{bmatrix} = \begin{bmatrix} \vec{b} \\ 0 \end{bmatrix}$$

$\vec{u} \in \text{Im}A$  and  $\lambda = 0$ .

▶ example13b.edp

```
matrix A = aa(Vh, Vh, solver=sparsesolver);  
real[int] c(u[.n]); c = 1.0; // kernel of A  
matrix AA=[ [A, c], [c', 0] ]; // matrix with constraint  
set(AA, solver=UMFPACK);
```