Fractals : Spectral properties Statistical physics

Course 1







6th Cornell Conference on Analysis, Probability, and Mathematical Physics on Fractals, June 13-17, 2017

Benefitted from discussions and collaborations with:

Technion:

Evgeni Gurevich (KLA-Tencor) Dor Gittelman Eli Levy (+ Rafael) Ariane Soret (ENS Cachan) Or Raz (HUJI, Maths) Omrie Ovdat Yaroslav Don

Rafael:

Assaf Barak Amnon Fisher

Elsewhere:

Gerald Dunne (UConn.) Alexander Teplyaev (UConn.) Jacqueline Bloch (LPN, Marcoussis) Dimitri Tanese (LPN, Marcoussis) Florent Baboux (LPN, Marcoussis) Alberto Amo (LPN, Marcoussis) Eva Andrei (Rutgers) Jinhai Mao (Rutgers) Arkady Poliakovsky (Maths. BGU)

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Plan of the 4 talks

- <u>Course 1</u>: Spectral properties of fractals Application in statistical physics
- <u>Talk</u>: quantum phase transition scale anomaly and fractals
- <u>Course 2</u>: topology and fractals measuring topological numbers with waves.
- <u>Elaboration</u> : Renormalisation group and Efimov physics

Program for today

- Introduction : spectral properties of self similar fractals.
- Heat kernel Asymptotic behaviour Weyl expansion - Spectral volume.
- Thermodynamics of the fractal blackbody.
- Summary Phase transitions.

Introduction : spectral properties of self similar fractals.

• attractive objects - Bear exotic names





Julia sets



Hofstadter butterfly





Sierpinski carpet

Sierpinski gasket



Diamond fractals

		1				_
	1/3					
1/9					_	_
1/27					_	_
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				•		

Triadic Cantor set

Convey the idea of highly symmetric objects yet with an unusual type of symmetry and a notion of extreme subdivision

Fractal : Iterative graph structure

Fractal : Iterative graph structure



Sierpinski gasket



Diamond fractals

As opposed to Euclidean spaces characterised by <u>translation</u> <u>symmetry</u>, fractals possess a <u>dilatation symmetry</u>.

Fractals are self-similar objects

Fractal ↔ Self-similar



Discrete scaling symmetry

- But not all fractals are obvious, good faith geometrical objects.
 - Sometimes, the fractal structure is not geometrical but it is hidden at a more abstract level.

- But generally, not all fractals are obvious, good faith geometrical objects.
 - Sometimes, the fractal structure is not geometrical but it is hidden at a more abstract level.

Exemple : Quasi-periodic chain of layers of 2 types A, B

Fibonacci sequence : $F_1 = B$; $F_2 = A$; $F_{j\geq 3} = \left[F_{j-2}F_{j-1}\right]$



Defines a cavity whose frequency spectrum is fractal.



Minicourse 2 - Tomorrow

Operators and fields on fractal manifolds

Operators are often expressed by local differential equations relating the space-time behaviour of a field

Ex. Wave equation

$$\frac{\partial^2 u}{\partial t^2} = \Delta u$$

Such local equations cannot be defined on a fractal

—	—	





But operators are essential quantities for physics!

- Quantum transport in fractal structures : e.g., networks, waveguides, ... electrons, photons
- Density of states
- Scattering matrix (transmission/reflection)

But operators are essential quantities for physics!

- Quantum fields on fractals, *e.g.*, fermions (spin 1/2), photons (spin 1) - canonical quantisation (Fourier modes) - path integral quantisation : path integrals, Brownian motion.
- "curved space QFT" or quantum gravity
- Scaling symmetry (renormalisation group) critical behaviour.



Michel Lapidus





Bob Strichartz

Jun Kigami

Recent new ideas >2000

Maths.



Intermezzo : heat and waves

From classical diffusion to wave propagation

Important relation between <u>classical diffusion</u> and <u>wave</u> <u>propagation</u> on a manifold.

Expresses the idea that it is possible to measure and characterise a manifold using waves (eigenvalue spectrum of the Laplace operator)



Use propagating waves/particles to probe :

- <u>spectral information</u>: density of states, transport, heat kernel, ...
- <u>geometric information</u>: dimension, volume, boundaries, shape, …

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Mathematical physics

<u>1910 Lorentz</u>: why is the Jeans radiation law only dependent on the volume ?

<u>1911 Weyl</u> : relation between asymptotic eigenvalues and dimension/volume.

<u>1966 Kac</u> : can one hear the shape of a drum ?

- Heat equation $\frac{\partial u}{\partial t} = \Delta u$
- Wave equation $\frac{\partial^2 u}{\partial t^2} = \Delta u$
 - Schr. equation. $i \frac{\partial u}{\partial t} = \Delta u$

- Heat equation $\frac{\partial u}{\partial t} = \Delta u$
- Wave equation $\frac{\partial^2 u}{\partial t^2} = \Delta u$ Schr. equation. $i\frac{\partial u}{\partial t} = \Delta u$

$$u(x,t) = \int d\mu(y) P_t(x,y) u(y,0)$$

• Heat equation
$$\frac{\partial u}{\partial t} = \Delta u$$

• Wave equation
$$\frac{\partial^2 u}{\partial t^2} = \Delta u$$

$$u(x,t) = \int d\mu(y) P_t(x,y) u(y,0)$$

Schr. equation.
$$i \frac{\partial u}{\partial t} = \Delta u$$

$$P_t(x,y) = \int_{x(0)=x, x(t)=y} \mathcal{D} x e^{-(i)\int_0^t \dot{x}^2 d\tau}$$

Brownian motion

$$P_t(x,y) \sim \frac{1}{t^{\frac{d}{2}}} \sum_n a_n(x,y) t^n$$

 $P_t(x,y) \sim \sum_{geodesics} (\#) e^{-(i)S_{classical}(x,y,t)}$

Heat kernel expansion

Gutzwiller - instantons



$$P_t(x,y) = \langle y | e^{-\Delta t} | x \rangle = \sum_{\lambda} \psi_{\lambda}^*(y) \psi_{\lambda}(x) e^{-\lambda t}$$

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$$Z(t) = Tr e^{-\Delta t} = \int dx \langle x | e^{-\Delta t} | x \rangle = \sum_{\lambda} e^{-\lambda t}$$

Heat kernel



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Heat kernel

$$\zeta_Z(s) \equiv \frac{1}{\Gamma(s)} \int_0^\infty dt t^{s-1} Z(t)$$

Mellin transform

$$\zeta_{Z}(s) = Tr \frac{1}{\Delta^{s}} = \sum_{\lambda} \frac{1}{\lambda^{s}}$$

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Small t behaviour of $Z(t) \iff \text{poles of } \zeta_Z(s)$

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Small t behaviour of $Z(t) \iff \text{poles of } \zeta_Z(s)$ Weyl expansion

The heat kernel is related to the density of states of the Laplacian

There are "Laplace transform" of each other:

$$Z(t) = \int_{0}^{\infty} d\omega \,\rho(\omega) \, e^{-t\,\omega}$$

From the Weyl expansion, it is possible to obtain the density of states.

How does it work ?



How does it work ?



whose spectral solution is
$$P_t(x,y) = \frac{1}{(4\pi Dt)^{1/2}} e^{-\frac{(x-y)^2}{4Dt}}$$

Probability of diffusing from x to y in a time t.
How does it work ?



Diffusion (heat) equation in d=1

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$$P_t(x,y) = \frac{1}{(4\pi Dt)^{1/2}} e^{-\frac{(x-y)^2}{4Dt}}$$

Probability of diffusing from x to y in a time t.

In d space dimensions:

$$P_t(x,y) = \frac{1}{(4\pi Dt)^{d/2}} e^{-\frac{(x-y)^2}{4Dt}}$$

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In d space dimensions:

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We can characterise the "spatial geometry" by watching how the heat flows. The heat kernel $Z_d(t)$ is

$$Z_{d}(t) = \int_{Vol.} d^{d}x P_{t}(x,x) = \frac{Volume}{(4\pi Dt)^{d/2}} \longrightarrow \text{volume of the manifold}}$$

 $\left(\right)$

Dirichlet:
$$\lambda_n = \left(\frac{n\pi}{L}\right)^2$$
, $n = 1, 2, ...$

 \bigcap

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Neumann: $\lambda_n = \left(\frac{n\pi}{L}\right)^2$, $n = 0, 1, 2, ...$

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$$Z_N(t) = \sum_{n=0}^{\infty} e^{-\left(\frac{n\pi}{L}\right)^2 t} = 1 + Z_D(t)$$

Dirichlet:
$$\lambda_n = \left(\frac{n\pi}{L}\right)^2$$
, $n = 1, 2, ...$
Neumann: $\lambda_n = \left(\frac{n\pi}{L}\right)^2$, $n = 0, 1, 2, ...$







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$$Z_{d=2}(t) \sim \frac{Vol.}{4\pi t} - \frac{L}{4}\frac{1}{\sqrt{4\pi t}} + \frac{1}{6} + \dots$$

Dirichlet:
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bulk

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$$Z_{d=2}(t) \sim \frac{Vol.}{4\pi t} - \frac{L}{4\sqrt{4\pi t}} + \frac{1}{6} + \dots$$

sensitive to boundary
bulk

Dirichlet:
$$\lambda_n = \left(\frac{n\pi}{L}\right)^2$$
, $n = 1, 2, ...$
Neumann: $\lambda_n = \left(\frac{n\pi}{L}\right)^2$, $n = 0, 1, 2, ...$



$$Z_{d=2}(t) \sim \frac{Vol.}{4\pi t} - \frac{L}{4\sqrt{4\pi t}} + \frac{1}{6} + \dots$$

integral of bound.
sensitive to boundary curvature

$$\zeta$$
-function $\zeta_z(s) = Tr \frac{1}{\Delta^s} = \sum_{\lambda} \frac{1}{\lambda^s}$

Dirichlet :
$$\lambda_n = \frac{n^2 \pi^2}{L^2}$$
, $n = 1, 2, ...$

$$\zeta(s) = \sum_{n=0}^{\infty} \left(\frac{L^2}{n^2 \pi^2}\right)^s = \frac{L^{2s}}{\pi^{2s}} \sum_{n=0}^{\infty} \frac{1}{n^{2s}} \equiv \frac{L^{2s}}{\pi^{2s}} \zeta_R(2s)$$

$$\zeta_R(2s)$$
 has a simple pole at $s = \frac{1}{2} \left(s = \frac{d}{2}\right)$ so that,

$$Z(t) = \frac{1}{2i\pi} \int_{a-i\infty}^{a+i\infty} ds t^{-s} \Gamma(s) \zeta(s) \sim \frac{L}{2\pi} t^{-\frac{1}{2}} \Gamma(\frac{1}{2}) + \dots$$
$$= \frac{L}{\sqrt{4\pi t}} + \dots$$

How does it work on a fractal?

How does it work on a fractal?

Differently...

No simple access to the eigenvalue spectrum but we know how to calculate the heat kernel.

$$Z(t) = Tr e^{-\Delta t} = \int dx \langle x | e^{-\Delta t} | x \rangle = \sum_{\lambda} e^{-\lambda t}$$

and thus, the density of states,

$$Z(t) = \int_{0}^{\infty} d\omega \,\rho(\omega) \, e^{-t\,\omega}$$

More precisely,

$$Z(t) = \sum_{k=1}^{\infty} e^{-k^2 \pi^2 t} + B \sum_{n=0}^{\infty} L_n^{d_h} \sum_{k=1}^{\infty} e^{-k^2 \pi^2 t L_n^{d_w}}$$

 $L_n = a^n$ is the total length upon iteration of the elementary step

$$\zeta(s) = \frac{\zeta_R(2s)}{\pi^{2s}} + \sum_{n=0}^{\infty} L_n^{d_h} \sum_{k=1}^{\infty} \left(\frac{1}{k^2 \pi^2 L_n^{d_w}}\right)^s$$
$$= \frac{\zeta_R(2s)}{\pi^{2s}} \left(1 + \sum_{n=0}^{\infty} L_n^{d_h - d_w s}\right)$$
$$= \frac{\zeta_R(2s)}{\pi^{2s}} \left(\frac{2 - a^{d_h - d_w s}}{1 - a^{d_h - d_w s}}\right) \text{ which has poles at } a^{d_h - d_w s} = 1$$

$$a^{d_h - d_w s} = 1 \qquad \Longleftrightarrow \qquad s_n = \frac{d_s}{2} + \frac{2i\pi n}{d_w \ln a}$$

Infinite number of complex poles : complex fractal dimensions. They control the behaviour of the heat kernel which exhibits oscillations.

1.00.8 $Z_{diamond}(t) 0.6$ 1.005 1.000 0.40.995 0.990 0.985 0.2 0.980 $\overline{2.10}^{\underline{t}_3}$ 10-3 2.10^{-4} $\overline{0.25}^{t}$ 0.15 0.20 0.05 0.10 $0\overline{.00}$

A new fractal dimension : **<u>spectral dimension</u>** d_s

Notion of spectral volume

Consider for simplicity
$$n = 1$$
, namely $s_1 = \frac{d_s}{2} + \frac{2i\pi}{d_w \ln a} \equiv \frac{d_s}{2} + i\delta$

$$Z(t) = \operatorname{Re}\left(\frac{V_{S}}{\frac{d_{s}}{2} + i\delta}\right)$$

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$$Z(t) = \operatorname{Re}\left(\frac{V_{S}}{\frac{d_{s}}{2} + i\delta}\right) \text{ so tha}$$

It

$$Z(t) \sim \frac{V_s}{t^{\frac{d_s}{2}}} \cos\left(\frac{2\pi}{d_w \ln a} \ln t\right)$$

Consider for simplicity
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to compare with

$$Z_d(t) = \int_{Vol.} d^d x P_t(x,x) = \frac{Volume}{(4\pi Dt)^{d/2}}$$

Spectral volume ?



Geometric volume described by the Hausdorff dimension is large (infinite)

Spectral volume ?



Geometric volume described by the Hausdorff dimension is large (infinite)



Spectral volume V_s is the finite volume occupied by the modes

Numerical solution of Maxwell eqs. on the Sierpinski gasket

Physical application : Thermodynamics of photons on fractals



Electromagnetic field in a waveguide fractal structure.

How to measure the spectral volume ?

The radiating fractal

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and, integrating the inequality $P_{\Omega}(\vec{\rho} | \vec{\rho}; t) \leq 1/2\pi t$ over $\Omega \leq |\Omega|/2\pi t$.



Equation of state at thermodynamic equilibrium relating pressure, volume of the state and internal energy:

To George Eugene Uhlenbech

Before I explain the title and ir In an enclosure with a perfectly stellecting surface there was analogous to tones of an organapipe a worshall confine lowill asks for the energy in the frequency interval ded...ines Ra problem to prove that the number of streamining method is independent of the shape of the sh Noting that $N(a)a^2 = |\Omega(a)|$ we record the fruits of our form of the inequality

$$|\Omega(a)| \frac{4}{a^2} \sum_{\substack{m,n \ \text{odd}}} \exp\left[-\frac{(m^2+n^2)\pi^2}{2a^2}t\right] \leq \sum_{n=1}^{\infty} e^{-\lambda_n t} \leq 1$$

From the fact (already noted above) that

$$\lim_{t \to 0} 2\pi t \frac{4}{a^2} \sum_{\substack{m,n \\ \text{odd}}} \exp\left[-\frac{(m^2 + n^2)\pi^2}{2a^2}t\right] = 1$$

we conclude easily that

$$\left| \Omega(a) \right| \leq \liminf_{t \to 0} 2\pi t \sum_{n=1}^{\infty} e^{-\lambda_n t} \leq \limsup_{t \to 0} 2\pi t \sum_{n=1}^{\infty} e^{-\lambda_n t}$$

and since, by choosing *a* sufficiently small, we can make $|\Omega(a to |\Omega|)$, we must have $\lim_{t\to 0} 2\pi t \sum_{n=1}^{\infty} e^{-\lambda_n t} = |\Omega|$ or, in other

$$\sum_{n=1}^{\infty} e^{-\lambda_n t} \sim \frac{\left| \Omega \right|}{2\pi t}, \qquad t \to 0.$$

9. Are we now through with rigor? Not quite. For while

$$P_{\Omega}(\vec{\rho} \mid \vec{r}; t) \leq \frac{\exp\left[-\frac{\|\vec{r} - \vec{\rho}\|^2}{2t}\right]}{2\pi t}$$

The radiating fractal blackbody

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and, integrating the inequality $P_{\Omega}(\vec{\rho} \mid \vec{\rho}; t) \leq 1/2\pi t$ over Ω

Noting that $N(a)a^2 = |\Omega(a)|$ we record the fruits of our

PV = U/



Equation of state at thermodynamic equilibrium relating pressure, volume of the state and internal energy:

 $\leq |\Omega|/2\pi t.$

To George Eugene Uhlenbecl

form of the inequality In an enclosure with a perfectly states there with $\left| \Omega(a) \right| \frac{4}{a^2} \sum_{m,n} \exp\left[-\frac{(m^2+n^2)\pi^2}{2a^2} t \right] \leq \sum_{n=1}^{\infty} e^{-\lambda_n t} \leq$ analogous to tones of an organapipteawershallecomfine Iouvila From the fact (already noted above) that asks for the energy in the frequency interval dv... It is he occasions, the luxury of stopping problem to prove that the number of sufficiently high over spent on streamlining mathemati $\lim_{t \to 0} 2\pi t \frac{4}{a^2} \sum_{m,n} \exp\left[-\frac{(m^2 + n^2)\pi^2}{2a^2}t\right] = 1$ we conclude easily that is independent of the shaper of these parts of the shaper of the shaper of the shaper of the second vector $D = -\frac{1}{\beta} \frac{\partial \ln Z}{\partial V}$ $U = -\frac{\partial U}{\partial V}$ $|\Omega(a)| \leq \liminf_{t \to 0} 2\pi t \sum_{n=1}^{\infty} e^{-\lambda_n t} \leq \limsup_{t \to 0} 2\pi t \sum_{n=1}^{\infty} e^{-\lambda_n t}$ and since, by choosing sufficiently small, we can make $|\Omega(a)|$ to $|\Omega|$, we must be $\lim_{t\to 0} 2\pi t \sum_{n=1}^{\infty} e^{-\lambda_n t} = |\Omega|$ or, in other $\sum_{n=1}^{\infty} e^{-\lambda_n t} \sim \frac{\left| \Omega \right|}{2\pi t}, \qquad t \to 0.$ **Spectral** w through with rigor? Not quite. For while volume ? $P_{\Omega}(\vec{\rho} \mid \vec{r}; t) \leq \frac{\exp\left[-\frac{||r-\rho||^2}{2t}\right]}{2\pi t}$

Usual approach : count modes in momentum space

Calculate the partition (generating)

function z(T,V) for a blackbody of PV = -Ularge volume V in dimension d



Mode decomposition of the field

Z

Τ

$$\omega = c \left| \vec{k} \right| = c V^{-\frac{1}{d}} 2\pi \left| \vec{n} \right|$$
$$\ln \mathcal{Z}(T, V)$$

$$\ln z(T,V) = Q\left(\frac{L_{\beta}}{V'^{d}}\right) \quad \text{with}$$
$$U = -\frac{\partial \beta}{\partial \beta} \quad \beta = \frac{1}{k_{B}T}$$

with
$$L_{\beta} \equiv \beta \hbar c$$

(photon thermal
wavelength)

Thermodynamics :

$$U = -\frac{\partial}{\partial \beta} \ln z(T, V) = -\left(\frac{dQ}{dx}\right) \hbar c V^{-\frac{1}{d}}$$
$$P = \frac{1}{\beta} \left(\frac{\partial}{\partial V} \ln z\right)_{T} = -\left(\frac{dQ}{dx}\right) \frac{\hbar c V^{-\frac{1}{d}}}{V d}$$

Thermodynamics :

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$$P = \frac{1}{\beta} \left(\frac{\partial}{\partial V} \ln z\right)_{T} = -\left(\frac{dQ}{dx}\right) \frac{\hbar c V^{-\frac{1}{d}}}{V d}$$

so that $PV = \frac{U}{d}$ (The exact expression of Q is unimportant)

Thermodynamics :

$$U = -\frac{\partial}{\partial \beta} \ln z(T, V) = -\left(\frac{dQ}{dx}\right) \hbar c V^{-\frac{1}{d}}$$
$$P = \frac{1}{\beta} \left(\frac{\partial}{\partial V} \ln z\right)_{T} = -\left(\frac{dQ}{dx}\right) \frac{\hbar c V^{-\frac{1}{d}}}{V d}$$

so that $PV = \frac{U}{d}$ (The exact expression of Q is unimportant) Stefan-Boltzmann $U \propto VT^{d+1}$ is a consequence of $\left(\frac{\partial U}{\partial V}\right)_T = T\left(\frac{\partial P}{\partial T}\right)_V - P$

Adiabatic expansion

 $VT^d = Cte$

Dimensions of momentum and position spaces are usually different : problem with the conventional formulation in terms of phase space cells.

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Volume of a fractal is usually infinite.

Nevertheless,

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Volume of a fractal is usually infinite.

Nevertheless,





Re-phrase the thermodynamic problem in terms of heat kernel and zeta function.

Partition function of equilibrium quantum radiation

$$\ln z(T,V) = -\frac{1}{2} \ln Det_{M \times V} \left(\frac{\partial^2}{\partial \tau^2} + c^2 \Delta \right)$$

Looks (almost) like a bona fide wave equation **but** proper time.

This expression does not rely on mode decomposition.

Rescale by $L_{\beta} \equiv \beta \hbar c$
Partition function of equilibrium quantum radiation

$$\ln z(T,V) = -\frac{1}{2} \ln Det_{M \times V} \left(\frac{\partial^2}{\partial u^2} + L_{\beta}^2 \Delta \right)$$

le of radius $L_{\alpha} \equiv \beta \hbar c$

M: circ 'β -

Spatial manifold (fractal)

Thermal equilibrium of photons on a spatial manifold V at temperature T is described by the (scaled) wave equation on $M \times V$

$$\ln z(T,V) = -\frac{1}{2} \ln Det_{M \times V} \left(\frac{\partial^2}{\partial u^2} + L_{\beta}^2 \Delta \right)$$

$$\ln z(T,V) = -\frac{1}{2} \ln Det_{M \times V} \left(\frac{\partial^2}{\partial u^2} + L_{\beta}^2 \Delta \right)$$

$$\ln z(T,V) = \frac{1}{2} \int_{0}^{\infty} \frac{d\tau}{\tau} f(\tau) Tr_{V} e^{-\tau L_{\beta}^{2} \Delta}$$

$$\ln z(T,V) = -\frac{1}{2} \ln Det_{M \times V} \left(\frac{\partial^2}{\partial u^2} + L_{\beta}^2 \Delta \right)$$

$$\ln z(T,V) = \frac{1}{2} \int_{0}^{\infty} \frac{d\tau}{\tau} f(\tau) Tr_{V} e^{-\tau L_{\beta}^{2} \Delta}$$
$$f(\tau) = \sum_{n=-\infty}^{\infty} e^{-(2\pi n)^{2} \tau}$$

$$\ln z(T,V) = -\frac{1}{2} \ln Det_{M \times V} \left(\frac{\partial^2}{\partial u^2} + L_{\beta}^2 \Delta \right)$$



$$\ln z(T,V) = -\frac{1}{2} \ln Det_{M \times V} \left(\frac{\partial^2}{\partial u^2} + L_{\beta}^2 \Delta \right)$$



Large volume limit (a high temperature limit) $V \gg L_{\beta}^{d} \Leftrightarrow k_{B}T \gg \frac{\hbar c}{V/d}$

Weyl expansion:

$$Z(L_{\beta}^{2}\tau) \sim \frac{V}{\left(4\pi L_{\beta}^{2}\tau\right)^{d/2}}$$

$$\ln z(T,V) = \frac{1}{2} \int_{0}^{\infty} \frac{d\tau}{\tau} f(\tau) Tr_{V} e^{-\tau L_{\beta}^{2} \Delta}$$

+ Weyl expansion
$$\implies \ln z(T,V) \sim \frac{V}{L_{\beta}^{d}}$$

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Thermodynamics :

$$U = -\frac{\partial}{\partial\beta} \ln z(T,V) = -\left(\frac{dQ}{dx}\right) \hbar c V^{-\frac{1}{d}}$$

$$P = \frac{1}{\beta} \left(\frac{\partial}{\partial V} \ln z \right)_T = -\left(\frac{dQ}{dx} \right) \frac{\hbar c V^{-\frac{1}{d}}}{V d}$$

so that $PV = \frac{U}{d}$ (The exact expression of Q is unimportant)

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+ Weyl expansion
$$\implies \ln z(T,V) \sim \frac{V}{L_{\beta}^{d}}$$

 $PV = \frac{U}{d}$

Thermodynamics measures the spectral volume

On a fractal...

$$Z(L_{\beta}^{2}\tau) \sim \frac{V_{s}}{\left(4\pi L_{\beta}^{2}\tau\right)^{d_{s}/2}} f(\ln\tau)$$

On a fractal...







Thermodynamic equation of state for a fractal manifold

$$PV_s = \frac{U}{d_s}$$

Thermodynamics measures the spectral volume and the spectral dimension.

Summary

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- Thermodynamics is directly related to the heat kernel (partition function) - fractal blackbody importance of the spectral volume.
- Phase transitions on fractals : scaling/hyperscaling relations are modified on fractals (dependence on distinct fractal dimensions).

 Non gaussian fixed points (limit cycles) - Harris criterion : fractal geometry is a specific type of disorder similar to quasicrystals.

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- Off-diagonal long range order superfluidity (Mermin, Wagner, Coleman theorem) - Non diagonal Green's function.
- Applications to other problems : quantum phase transitions quantum Einstein gravity, ...

Thank you for your attention.