

Gauged Supergravity and Holographic RG Flows

Stonybrook, December, 2001

Based upon:

- [hep-th/9904017](#); [hep-th/9906194](#); [hep-th/004063](#)
Freedman, S. Gubser, K. Pilch and NPW
- Recent work with R. Corrado, M. Gunaydin and M. Zagermann
to appear

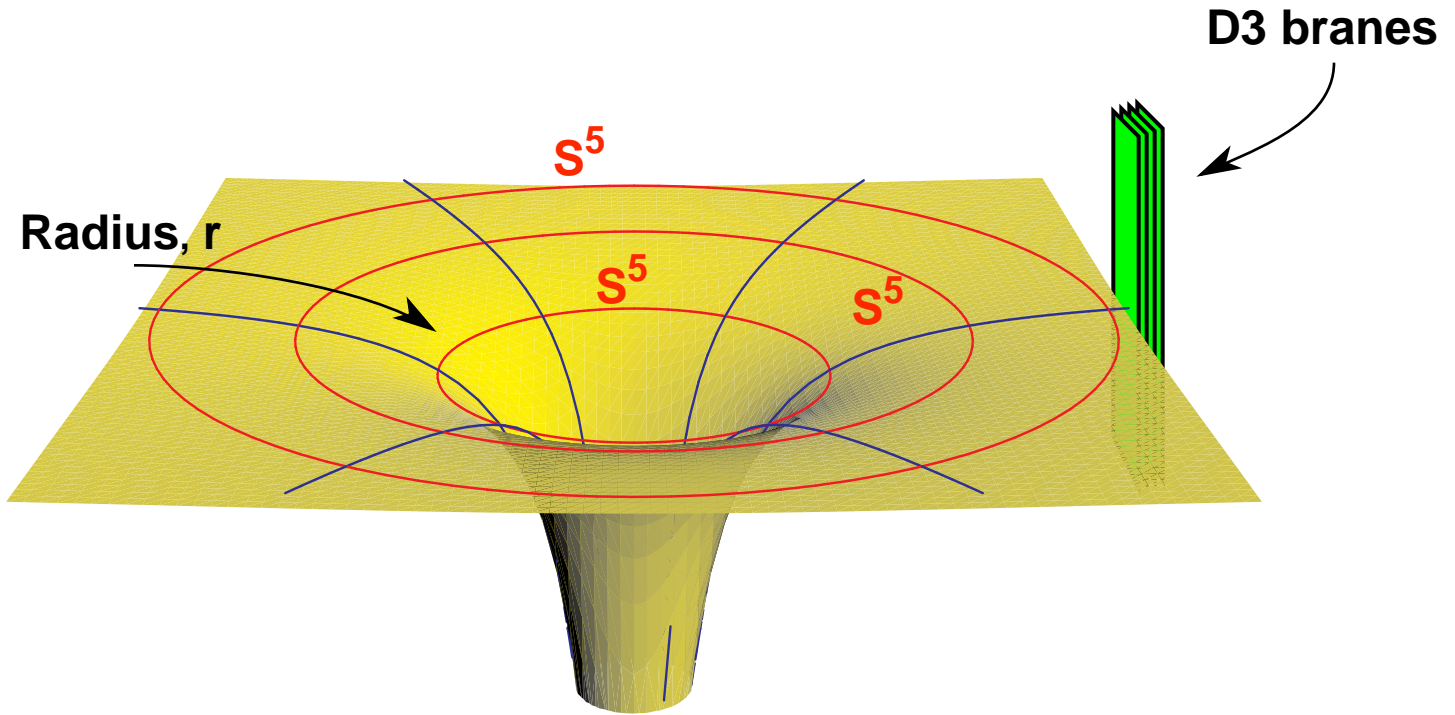
Overview

- Teaching an old dog new tricks: The AdS/CFT Correspondence
- Holographic Description of RG flows
- Open problems in holographic field theory and supergravity
 - Half-maximal supersymmetric models
 $\mathcal{N} = 2$ Quiver Gauge theories \longleftrightarrow $\mathcal{N} = 4$ supegravity + tensor multiplets
 - Half-maximal supersymmetric flows
 $\mathcal{N} = 2$ Seiberg Witten flows
 - Supersymmetric AdS geometries
- Other issues

Yang-Mills theory and AdS/CFT

D3-branes in IIB Supergravity

$\mathcal{N} = 4$ Yang-Mills in strongly coupled, Large N limit



D3 branes + r \longleftrightarrow AdS₅

Operators on the brane \longleftrightarrow Sources in Supergravity

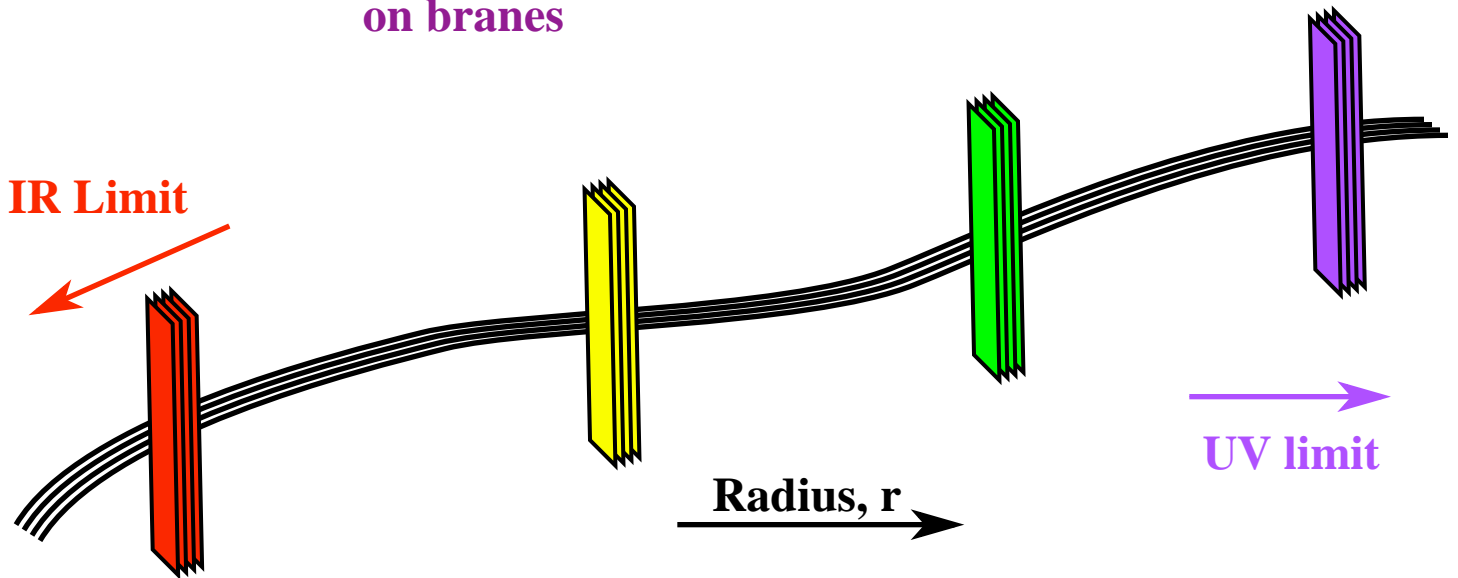
Correlators and states \longleftrightarrow Classical supergravity solutions

Five-dimensional perspective:

- View S^5 as an internal space
- Decompose theory into generalized Fourier modes on S^5
- Space-time = D3 branes + radius, r
- Interpret r as an RG scale for theory on brane

Interpreting the metric

$$ds_{10}^2 = \underbrace{\Omega^2(y)}_{\text{warp factor}} \left(\underbrace{e^{2A(r)}}_{\text{Physical scale on branes}} \underbrace{\eta_{\mu\nu} dx^\mu dx^\nu}_{\text{D3 brane}} + dr^2 \right) + \underbrace{ds_5^2(y)}_{\text{internal}}$$



- **Localize D3-branes at infinity** ↔ **Conformal, UV limit**
 Localize D3-branes finite r Explicit cut-off
- **Red shifts** ↔ **Wilsonian coarse graining**
- **Cosmological Entropy** ↔ **c-function / theorem**
- **No hair theorems** ↔ **Universality in IR limit**
- **Gauged supergravity potentials** ↔ **Phase diagrams / flows from relevant perturbations**
- **Consistent truncation** ↔ **Large N structure of OPE in E.M. supermultiplets**

Infra-red Fixed Points

I. Perturb the $\mathcal{N} = 4$ theory: $(\mathbf{A}_\mu, \lambda^1, \dots, \lambda^4, \mathbf{X}^1, \dots, \mathbf{X}^6)$

Freedman, Gubser, Pilch and Warner, hep-th/9901017

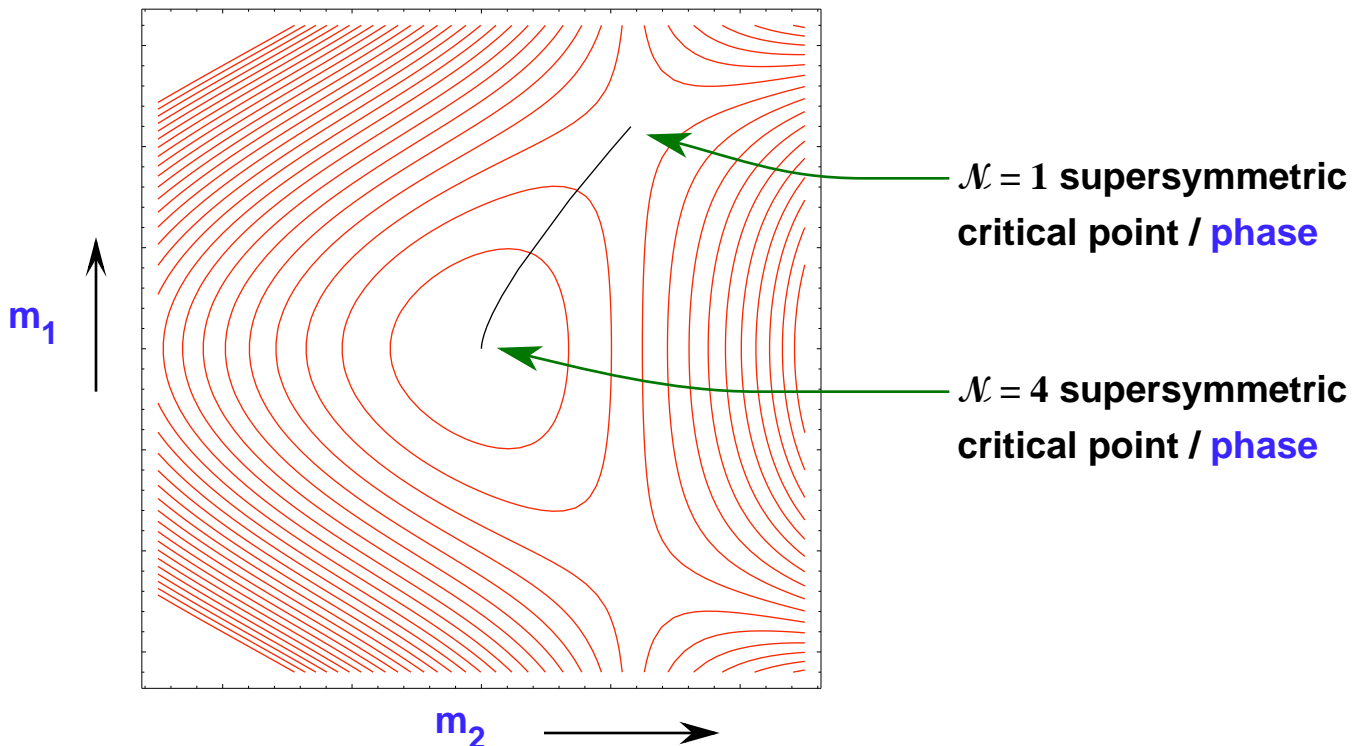
Give mass to an $\mathcal{N} = 1$ chiral multiplet: $\mathbf{W} = \mathbf{W}_0 + m \text{Tr}(\Phi^2)$

$$\Delta \mathcal{L} = m_1 \text{Tr}(\lambda^1 \lambda^1) + m_2 \text{Tr}((X^1)^2 + (X^2)^2); \quad m_2 = m_1^2 = m^2$$

Non-trivial $\mathcal{N} = 1$ fixed point
(Leigh and Strassler) for flow as $m \rightarrow \infty$

$$\frac{c_{\text{IR}}}{c_{\text{UV}}} = \frac{27}{32}$$

Supergravity: steepest descent on a supergravity superpotential, \mathcal{W} :



II. Perturb an $\mathcal{N} = 2$ superconformal quiver theory:

Klebanov and Witten hep-th/9807080

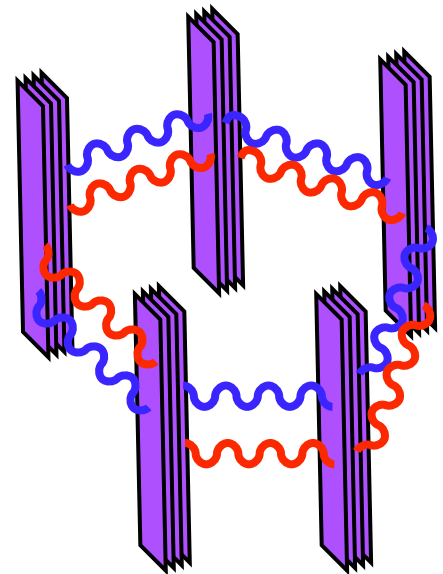
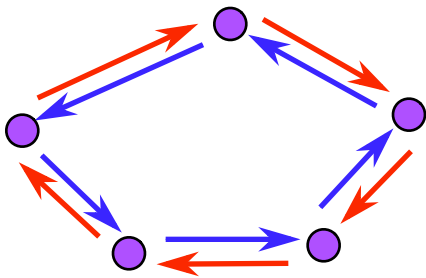
Gubser, Nekrasov and Shatashvili hep-th/9811230

$$\frac{c_{\text{IR}}}{c_{\text{UV}}} = \frac{27}{32}$$

Quiver Theories and Orbifolds

$\mathcal{N} = 2$, $(\text{SU}(N))^p$ Yang-Mills
+ bi-fundamental matter

\mathbb{Z}_p Orbifold of $p \times N$ D3-branes



= $\text{SU}(N)$ =

→ = (N, \bar{N}) =

← = (\bar{N}, N) =

Superconformal with R-symmetry = $\text{SU}(2) \times \text{U}(1)$

Untwisted sector: Vector multiplets

$$\left(\mathbf{A}^{(k)}_{\mu}; \lambda_1^{(k)}, \lambda_2^{(k)}; \phi^{(k)} = \chi_1^{(k)} + i \chi_2^{(k)} \right) \quad k=1, \dots, p$$

Twisted sector: Hypermultiplets $\left(\psi_1^{(k)}, \psi_2^{(k)}; \mathbf{A}^{(k)}, \mathbf{B}^{(k)} \right)$

Mass terms / flow: $\sum_k m_{(k)} (\phi^{(k)})^2$

Klebanov and Witten: $m \left((\phi^{(1)})^2 - (\phi^{(2)})^2 \right) \quad p=2, \mathbb{Z}_2 \text{ odd}$

Freedman, Gubser, Pilch and Warner: $m \left((\phi^{(1)})^2 + (\phi^{(2)})^2 \right) \quad \mathbb{Z}_2 \text{ even}$



Leigh-Strassler fixed points

Five-dimensional, gauged supergravity and quiver flows

Gauged $\mathcal{N} = 4$ supergravity + **Vector** or **Tensor** multiplets:

Dall'Agata, Herrmann and Zagermann, hep-th/0103106

R-symmetry on the brane = **Gauge group:** **SU(2) x U(1)**
(supergravity)

$$m_{(k)}^{ab} \lambda_a^{(k)} \lambda_b^{(k)} + \text{c.c.} \longleftrightarrow 3 (+2) + 3 (-2)$$

$$m_{(k)} (\phi^{(k)})^2 + \text{c.c.} \longleftrightarrow 1 (+4) + 1 (-4)$$

$$T_{(k)} = \frac{4\pi i}{g_{(k)}^2} + \frac{\theta_{(k)}}{2\pi} \longleftrightarrow 1 (0) + 1 (0)$$

Scalar content of a pair of five-dimensional, **charged tensor multiplets**

Gauged $\mathcal{N} = 4$ supergravity + **2p Tensor** multiplets: charges (+2) and (-2)

Scalar Manifold (Supergravity) \longleftrightarrow Couplings in Gauge theory

$$SO(1,1) \times \frac{SO(5, 2p)}{SO(5) \times SO(2p)} \supset SO(3) \times SO(2, 2p)$$

$$SU(2)_{\mathcal{R}} \times U(1)_{\mathcal{R}} \times SU(p, 1)$$

Invariance of supergravity potential

$$\frac{SU(p, 1)}{SU(p) \times U(1)}$$

= Non-compact CP_p parametrizing couplings: $T_{(k)}$

Generalizes:

$SU(1,1)/U(1)$ of IIB supergravity



$\mathcal{N} = 4$ Yang-Mills
Coupling, T

Unexpected Consequences: $SU(p)$ Symmetry

Untwisted scalar sector in $\mathcal{N} = 8$ supergravity + $\mathcal{N} = 4$ supergravity + two charged tensor multiplets $\Rightarrow \frac{SO(5, 2)}{SO(5) \times SO(2)}$

Twisted sector scalars from $2(p-1)$ additional tensor multiplets = $(p-1) \times (5, 2)$ of $SO(5) \times SO(2)$ $\Rightarrow \frac{SO(5, 2p)}{SO(5) \times SO(2p)}$

Flow of FGPW: $m ((\phi^{(1)})^2 + (\phi^{(2)})^2)$ lies entirely in scalars of **untwisted** sector

Flow of KW: $m ((\phi^{(1)})^2 - (\phi^{(2)})^2)$ lies entirely in scalars of **twisted** sector

$U(p)$ symmetry of $\mathcal{N} = 4$ supergravity potential

$SU(p)$ symmetry \Rightarrow

Critical point of FGPW extends to a $(p-1)$ -dimensional critical surface

Field Theory:

Flows driven by $\sum_k m_{(k)} (\phi^{(k)})^2$ go to a smooth, $SU(p)$ symmetric family of IR fixed points parametrized by $\tau_{(k)}^2 / m_{(k)}$ considered as homogeneous coordinates on CP_{p-1} .

IIB Supergravity: Bizarre

- FGPW flow: topologically trivial flux on S^5 / Z_2
 - KW flow: topologically **non-trivial** flux on **blow-up** of A_1 singularity on S^5 / Z_2
- $U(p)$
-

Comments:

- Family of **equivalent** fixed points with **same central charge**: parametrized by initial “velocities”, $\mathbf{m}_{(k)}$, of flow.
- This is a “**large N result:**” the $U(p)$ symmetry is broken by anomalies at finite N.

Residual Symmetry $U(p) \supset \mathcal{G} \supset (Z^{2N})^p$

finite ↙

- Discrete symmetry at finite N between FGPW and KW flow?
- Singular behavior at finite N when one or more $\mathbf{m}_{(k)}$ vanish?
- Does consistent truncation work for $\mathcal{N} = 4$ supergravity?
 - Unknown territory...
 - We have solutions in both five and ten dimensions.
 - Solutions coincide at linearized level
 - Same initial velocities
 - Same supersymmetry
 - Same \mathcal{R} -symmetry
- Try to find the $SU(2)$ family of solutions in IIB supergravity that interpolates between a topologically trivial flux on S^5/Z_2 and a topologically **non-trivial** flux on the **blow-up** of an A_1 singularity on S^5/Z_2 .
Corrado, Pilch and Warner: work in progress

Another duality that trades 3-form flux for Kähler moduli

Open Problems: Supersymmetric RR -Geometries

Geometry of supersymmetric solutions *with background RR -fluxes*?

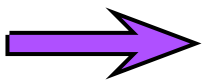
Ricci flat manifolds + $\mathcal{N} = 2$ supersymmetry: **Kähler**

Ricci flat manifolds + $\mathcal{N} = 4$ supersymmetry: **Hyperkähler**

What is the analogue of this for holographic RG flow solutions?

The issue for flow solutions in IIB supergravity:

- The flows generally start from AdS_5 , with a background 4-form RR -tensor gauge field
- Softly broken $\mathcal{N} = 4$ Yang-Mills: fermion masses are holographically dual to 2-form tensor gauge fields
- The dilaton and axion are dual to the Yang-Mills gauge couplings, and thus generically run.



Interesting supersymmetric flow solutions involve fluxes for all the background tensor gauge fields

Classification theorems for such solutions: very few, and none of them relevant to the important physical holographic flows.

Important Example: (unsolved)

Find the Holographic Dual of the $\mathcal{N} = 2$ Seiberg-Witten effective action

Holography and the $\mathcal{N} = 2$ Seiberg-Witten effective action

Simplest version:

$\mathcal{N} = 4$ Yang-Mills \longrightarrow $\mathcal{N} = 2$ Yang-Mills + massive hypermultiplet

Coulomb branch: $\mathbf{u}_n = \text{Tr}(\phi^n)$

\downarrow
D3-brane distribution Source: $\rho(\mathbf{y}) = \sum \mathbf{u}_n \mathbf{y}^n$

Complex scalar, ϕ \longrightarrow Complex coordinate, \mathbf{y}

Problem: Find the general $\mathcal{N} = 4$ supersymmetric flow solution determined by an arbitrary source function, $\rho(\mathbf{y})$, of two variables

What is known: the solution for one point in the moduli space

$$\rho(\mathbf{y}) = (a^2 - |\mathbf{y}|^2)^{1/2}$$

A uniform disk-like distribution of D3-branes

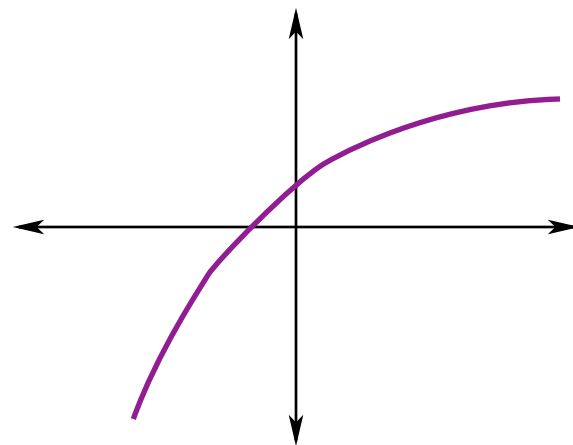
Pilch and Warner; hep-th/0006066
Buchel, Peet and Polchinski; hep-th/0008076
Evans, Johnson and Petrini; hep-th/0008081

The general $\mathcal{N} = 8$ supersymmetric flow solution (the Coulomb branch of $\mathcal{N} = 4$ Yang-Mills) is well-known, and is determined by an arbitrary harmonic source function, $\mathbf{H}(\mathbf{y})$, of six variables. The half-maximal supersymmetric solutions shouldn't be much more difficult....

Harder problem: Find the general $\mathcal{N} = 4$ supersymmetric flow solution for the quiver gauge theories

Other Issues

- $\mathcal{N} = 1$ supersymmetric flows: generalized Kähler structure in the presence of RR fluxes?
 - Brane probe results find the Kähler structure on the moduli space of the probe.
Johnson, Lovis, C. Page, hep-th/0107261
- Non-compact gaugings: Holographic interpretation???
 - Vacua: De-Sitter space \longrightarrow All supersymmetry broken
 - Domain wall solutions can be (half-maximally) supersymmetric
Hull, hep-th/0110048; Gibbons and Hull hep-th/0111072
- Potentials without critical points:
 - Domain wall solutions...
 - \downarrow ???
 - Gauged Supergravity description + Generalizations and variations on Duality Cascade of Klebanov and Strassler? hep-th/0007191
- Most of the exactly known flow solutions have been obtained by lifting solutions of gauged supergravity
 - Systematics of lifting
 - General methods of construction in 10 or 11 dimensions



Conclusions

- Gauged supergravity is a very valuable tool in the study of holographic RG flows
- $\mathcal{N} = 4$ gauged supergravity coupled to tensor multiplets can be used to study a class of $\mathcal{N} = 1$ supersymmetric flows in large N , $\mathcal{N} = 2$ quiver gauge theories
- The large N , \mathbb{Z}_p quiver models have a $(p-1)$ -dimensional surface, $\mathbb{C}P_{p-1}$, of $\mathcal{N} = 1$ supersymmetric fixed points, and all these fixed point theories are equivalent under the action of an $SU(p)$.
- In ten dimensions, this $SU(p)$ must act as a “generalized duality symmetry”, mapping compactifications with fluxes on S^5/\mathbb{Z}_p to compactifications with blow-ups of the A_{p-1} singularity:

Flux \longleftrightarrow Kähler Moduli

- Important open question in holographic RG flows:

Geometry of supersymmetric compactifications with RR fluxes???

- Holographic description of Seiberg-Witten actions
- Holographic interpretation of the many supersymmetric domain-wall solutions?