

Introduction to Rheology

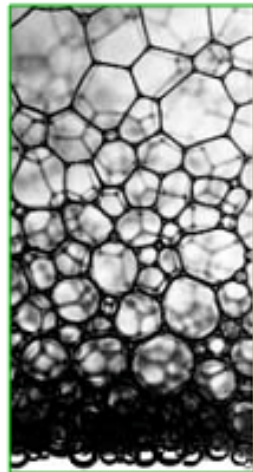
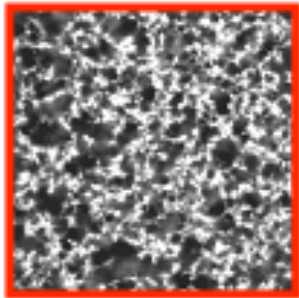
D. Vader, H.Wyss

Weitzlab group meeting tutorial

What is rheology?

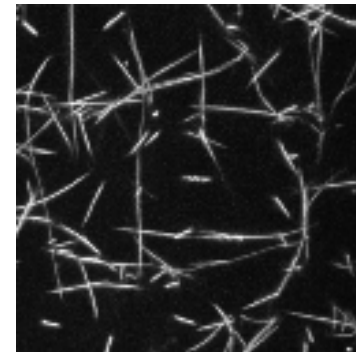
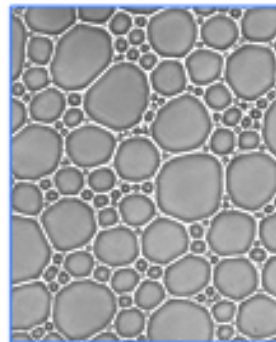
- Rheology is the study of the flow of matter: mainly liquids but also soft solids or solids under conditions in which they flow rather than deform elastically. It applies to substances which have a complex structure, including muds, sludges, suspensions, polymers, many foods, bodily fluids, and other biological materials.

Cheese



Foams

Emulsions

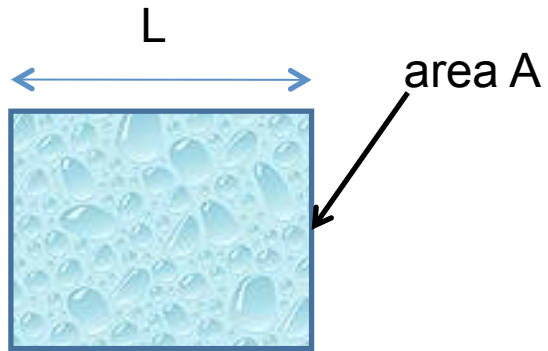


Biopolymers

What is rheology?

- The term rheology was coined in 1920s, and was inspired by a Greek quotation, "panta rei", "everything flows".
- In practice, rheology is principally concerned with extending the "classical" disciplines of elasticity and (Newtonian) fluid mechanics to materials whose mechanical behavior cannot be described with the classical theories.

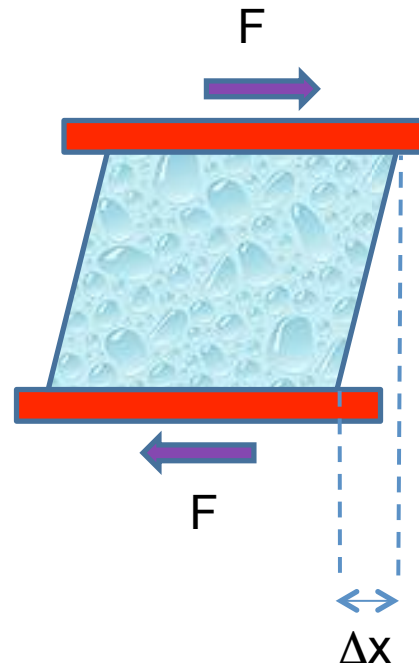
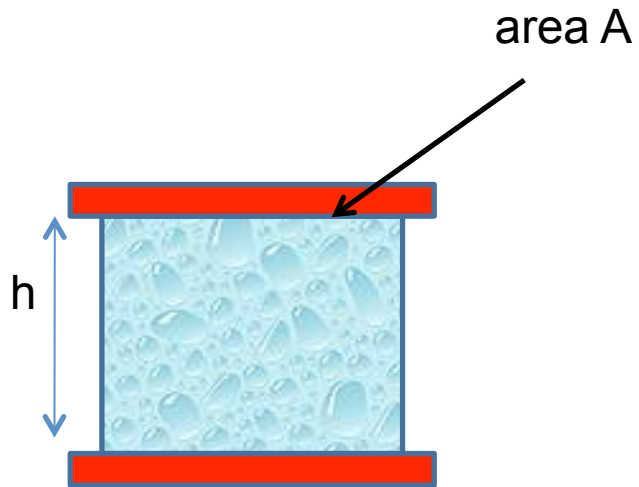
Basic concepts



$$\varepsilon = \frac{\Delta L}{L}$$

$$\sigma = E\varepsilon$$

$$\sigma = \frac{F}{A}$$



$$\gamma = \frac{\Delta x}{h}$$

$$\sigma = G'\gamma$$

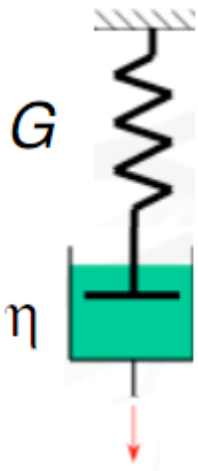
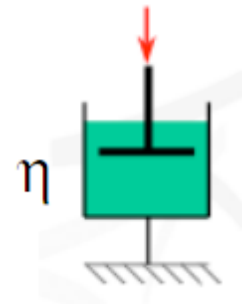
$$\sigma = G''\dot{\gamma}$$

Simple mechanical elements



Elastic solid: force (stress) proportional to strain

Viscous fluid: force (stress) proportional to strain rate



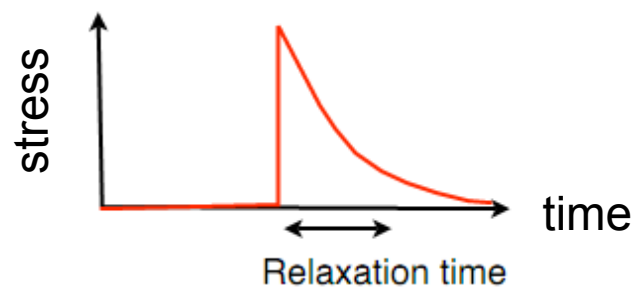
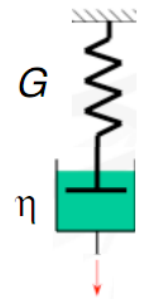
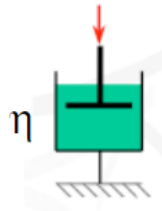
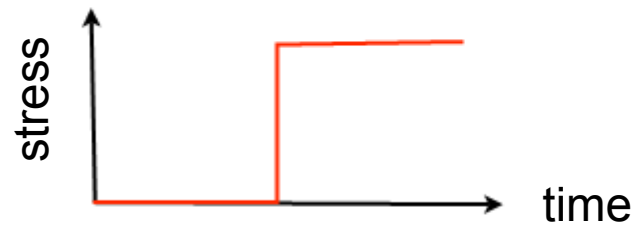
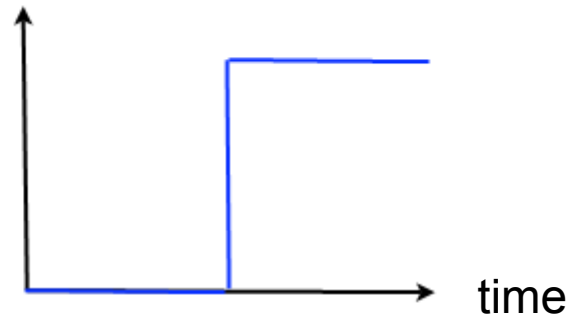
Viscoelastic material: time scales are important

Fast deformation: solid-like
Slow deformation: fluid-like

$$\tau = \frac{\eta}{G}$$

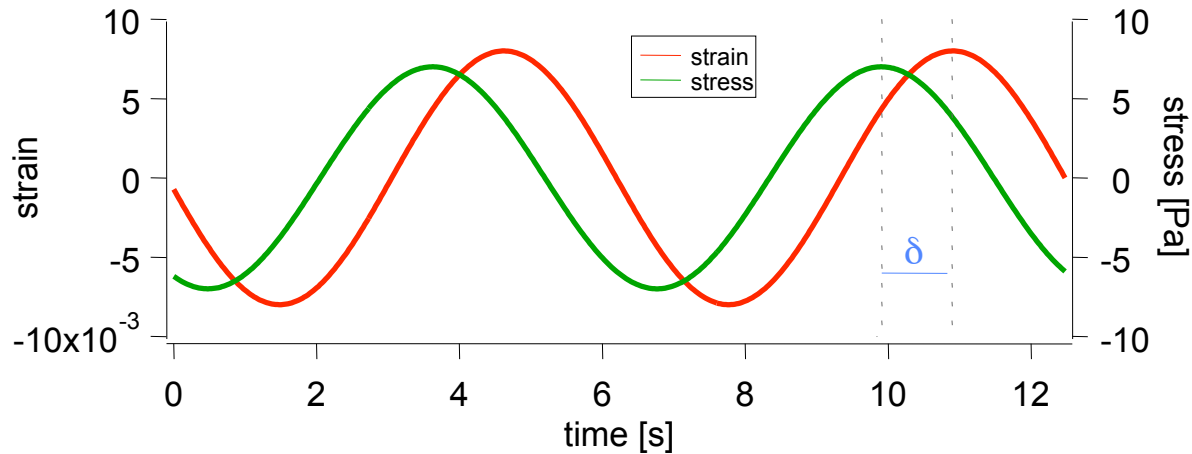
Response to deformation

Step strain



$$\tau = \frac{\eta}{G}$$

Oscillatory rheology



Elastic solid: $\sigma = G\gamma$

Stress and strain are in phase

Viscous fluid: $\sigma = \eta \dot{\gamma}$

Stress and strain are out of phase

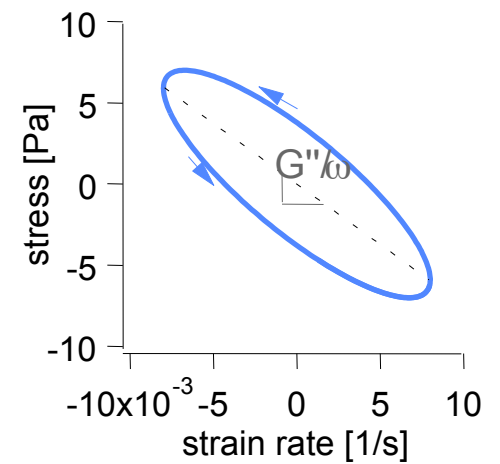
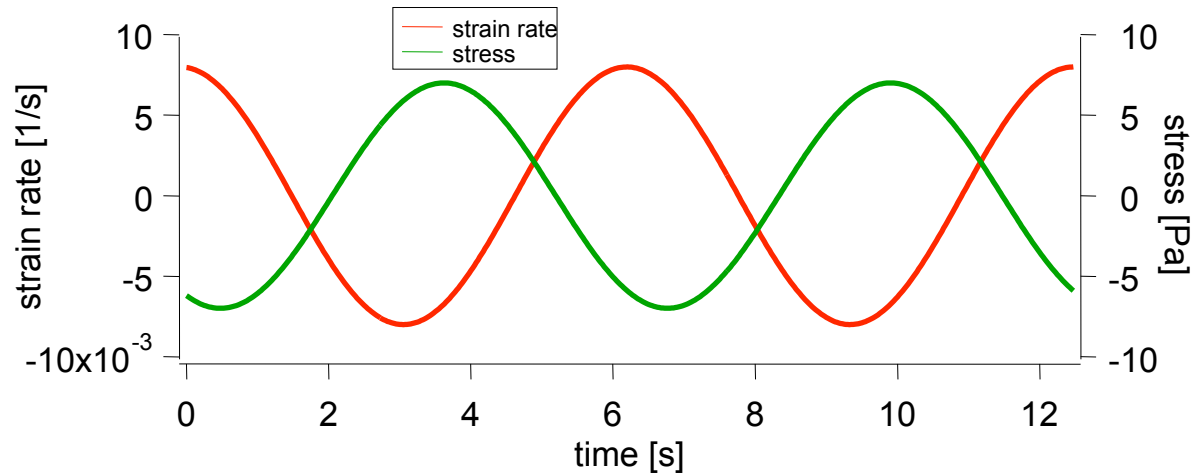
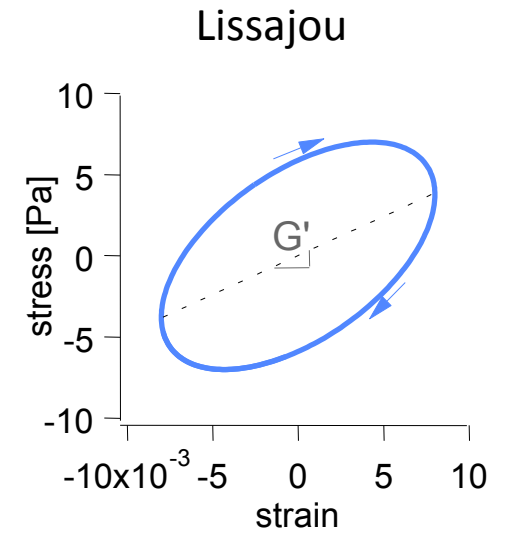
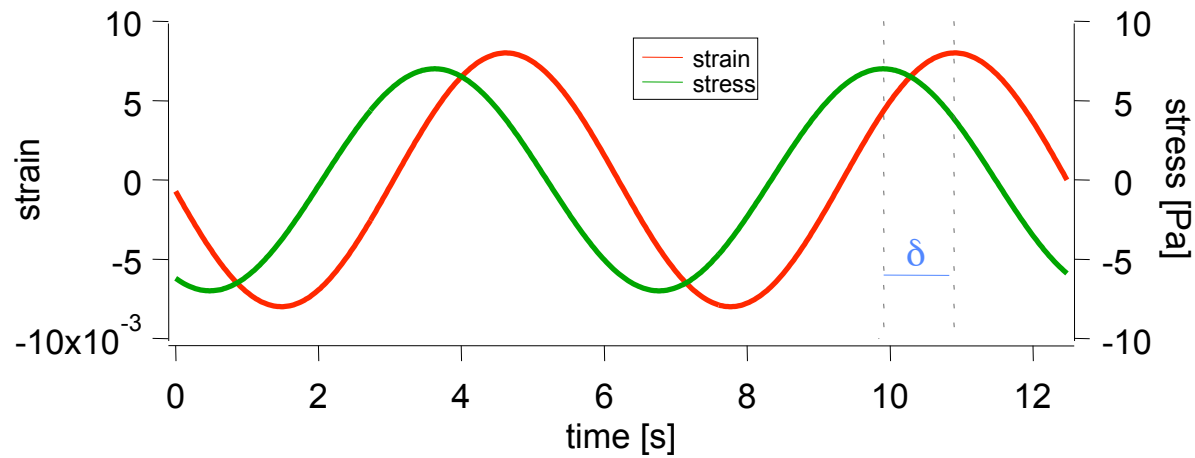
Viscoelastic material, use: $\gamma(\omega, t) = \gamma_0 \sin(\omega t)$

$$\sigma(\omega, t) = G' \gamma_0 \cdot \sin(\omega t) + G'' \gamma_0 \cdot \cos(\omega t)$$

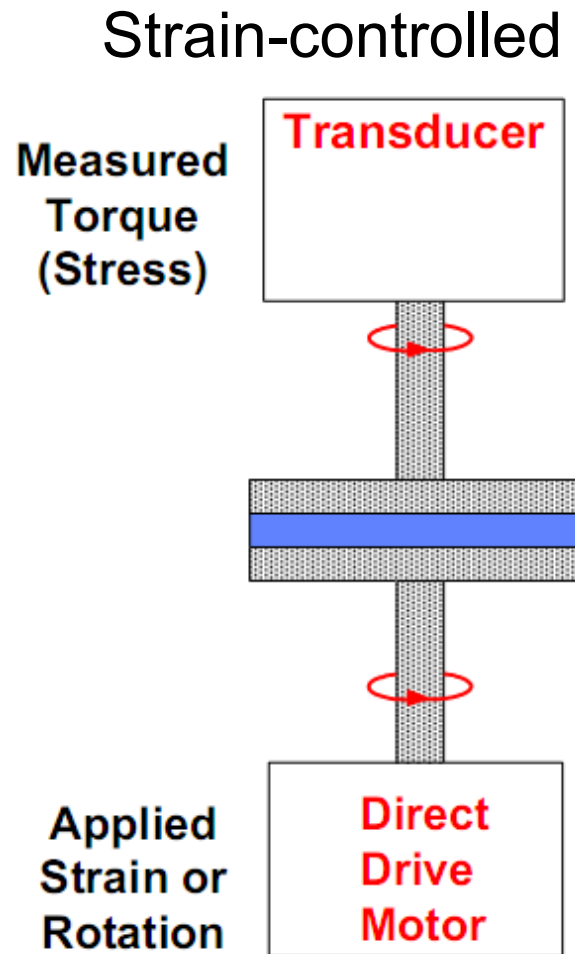
$$\eta = \frac{G''}{\omega}$$

$$\tan(\delta) = \frac{G''}{G'}$$

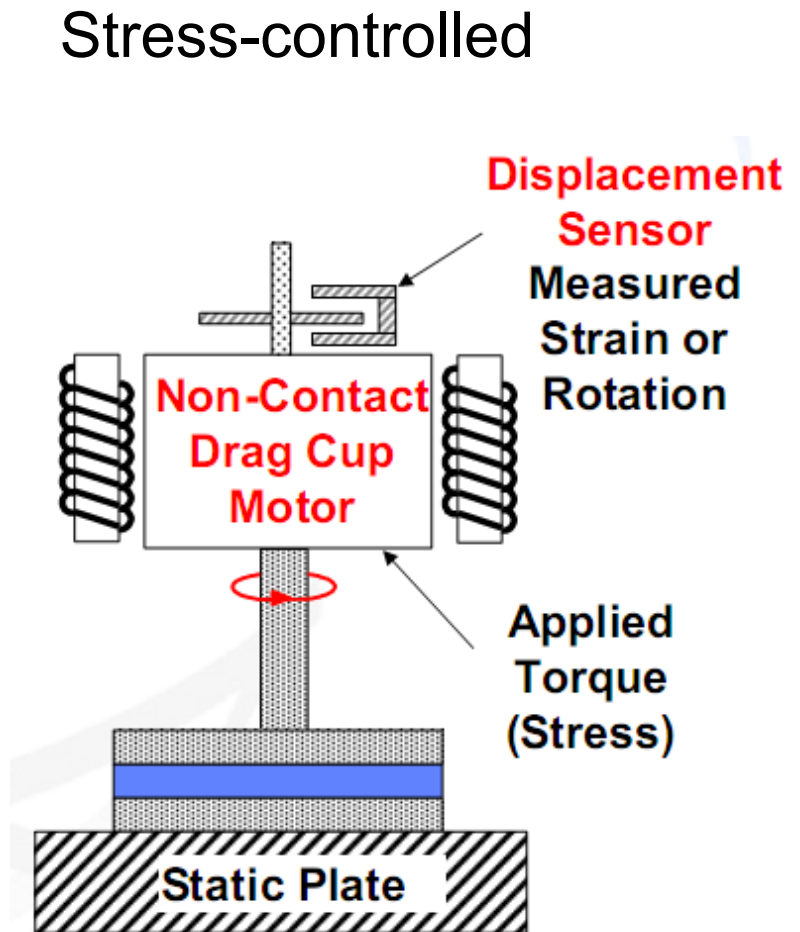
Lissajou plots



Strain-control vs stress-control



ARES



Bohlin, AR-G2, Anton Paar

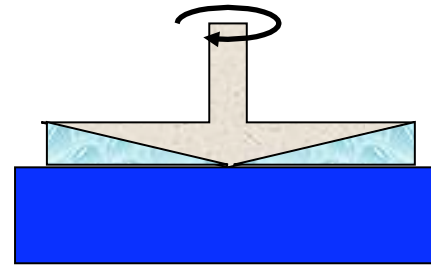
Strain-control vs stress-control

- *Strain-controlled state typically considered better defined*
- Stress-controlled rheometers have better torque sensitivity
- Strain-controlled rheometers can probe higher frequencies
- BUT... nowadays, feedback loops are fast enough that most rheometers can operate OK in both modes

Rheometer geometries

☐ Cone-plate

- uniform strain / strain-rate
- fixed gap height

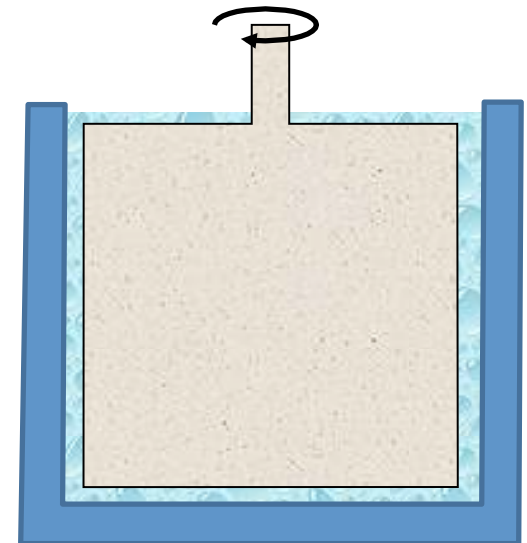
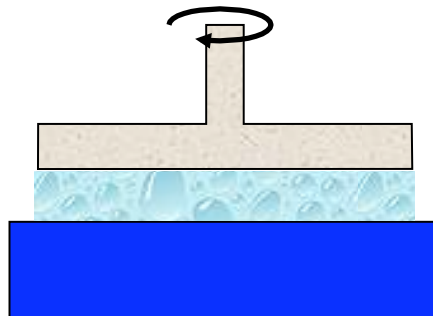


☐ Plate-plate

- non-uniform strain
- adjustable gap height
- good for testing boundary effects like slip

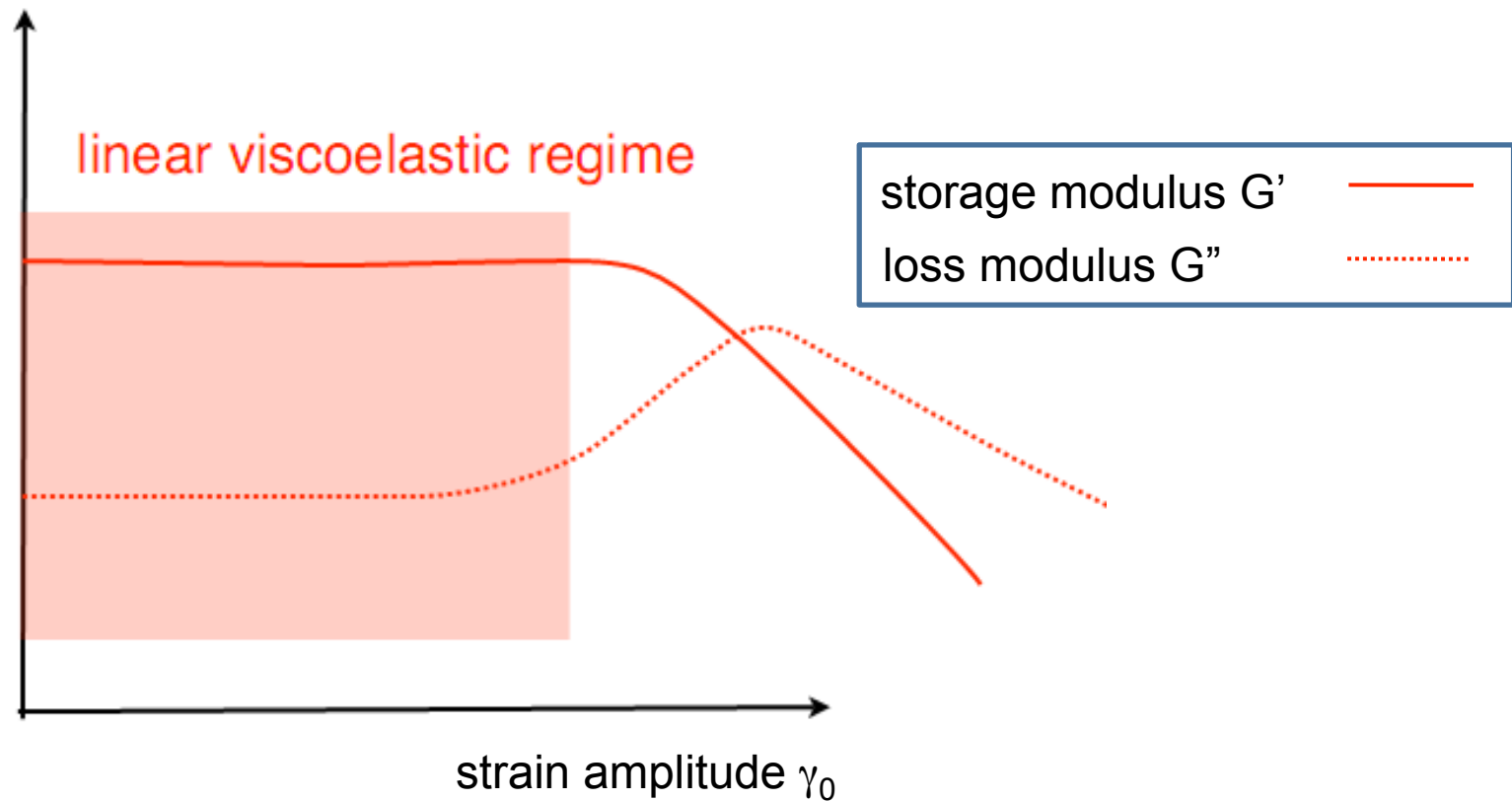
☐ Couette cell

- good sensitivity for low-viscosity fluids



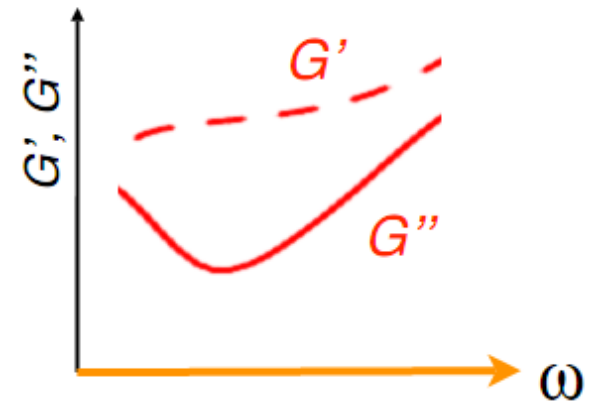
Linear viscoelasticity

Acquire data at constant frequency, increasing stress/strain



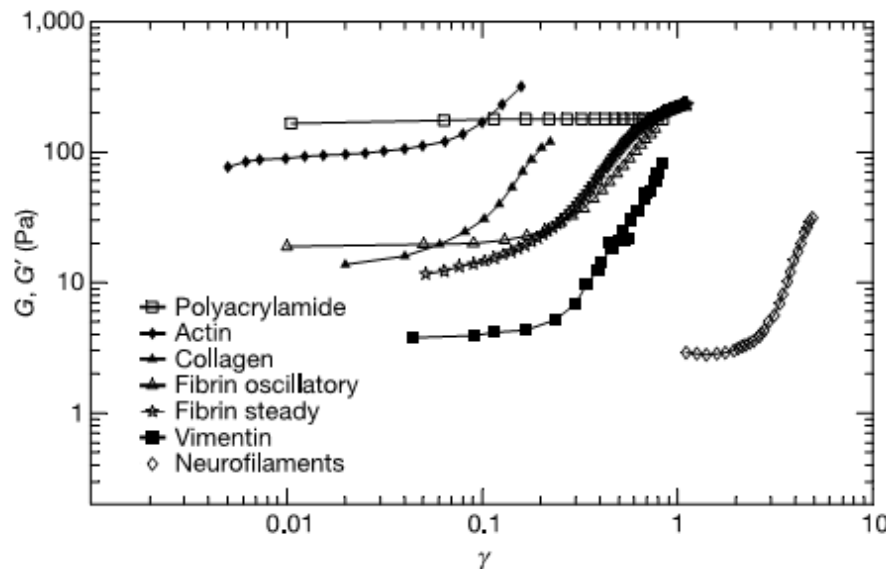
Typical protocol

- Limits of linear viscoelastic regime in desired frequency range using amplitude sweeps
=> yield stress/strain, critical stress/strain
- Test for time stability, i.e time sweep at constant amplitude and frequency
- Frequency sweep at various strain/stress amplitudes within linear regime
- Study non-linear regime

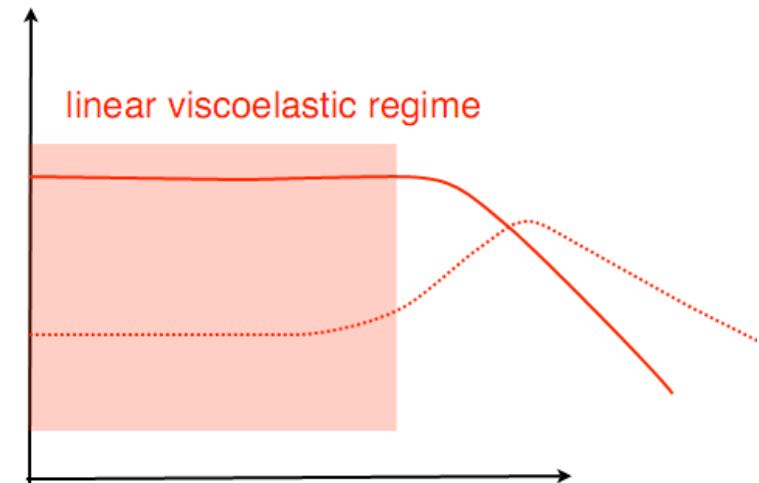


Nonlinear rheology (of biopolymers)

- “Unlike simple polymer gels, many biological materials—including blood vessels, mesentery tissue, lung parenchyma, cornea and blood clots—stiffen as they are strained, thereby preventing large deformations that could threaten tissue integrity.” (Storm et al., 2005)*

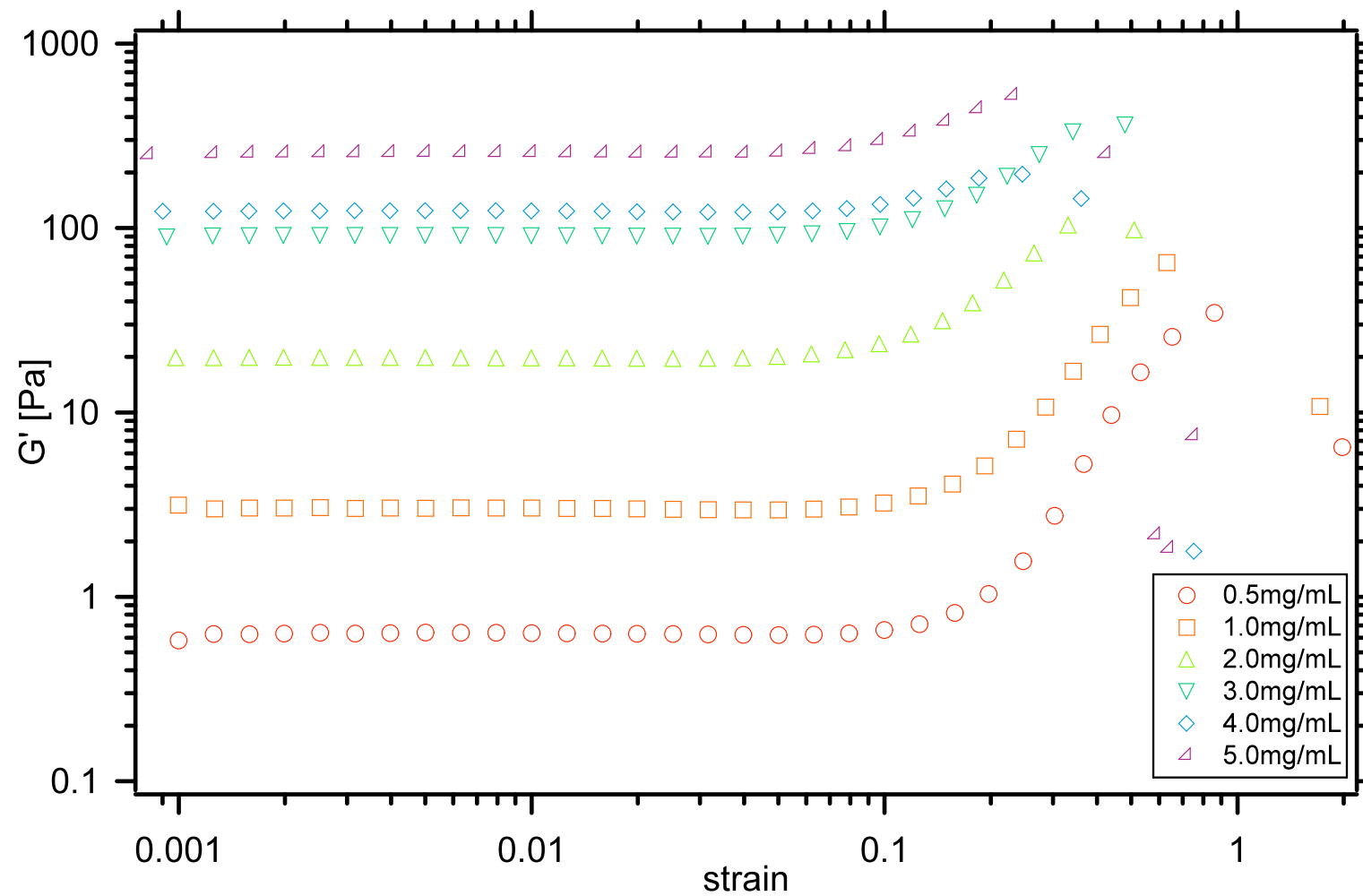


stiffening

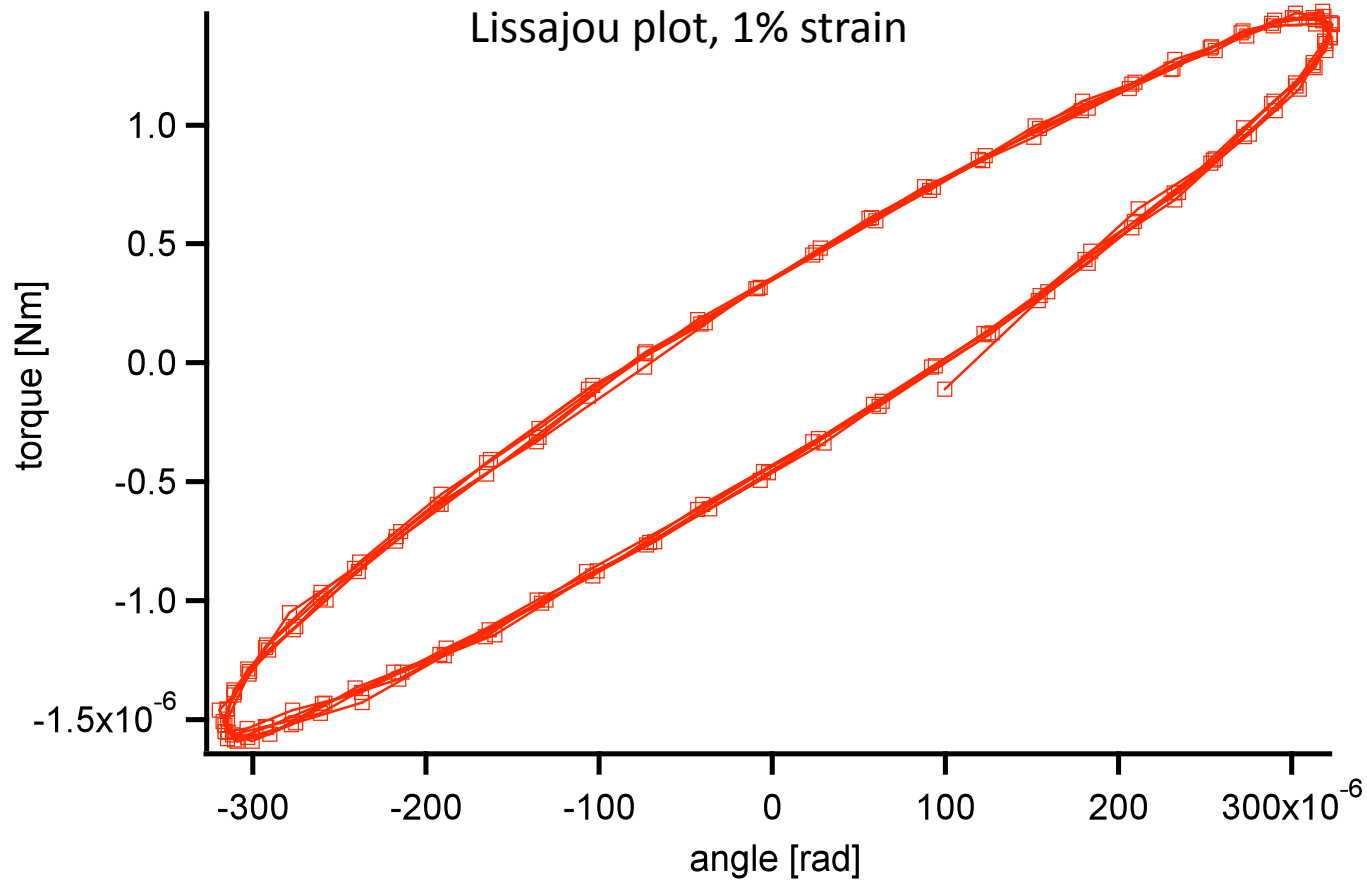


weakening

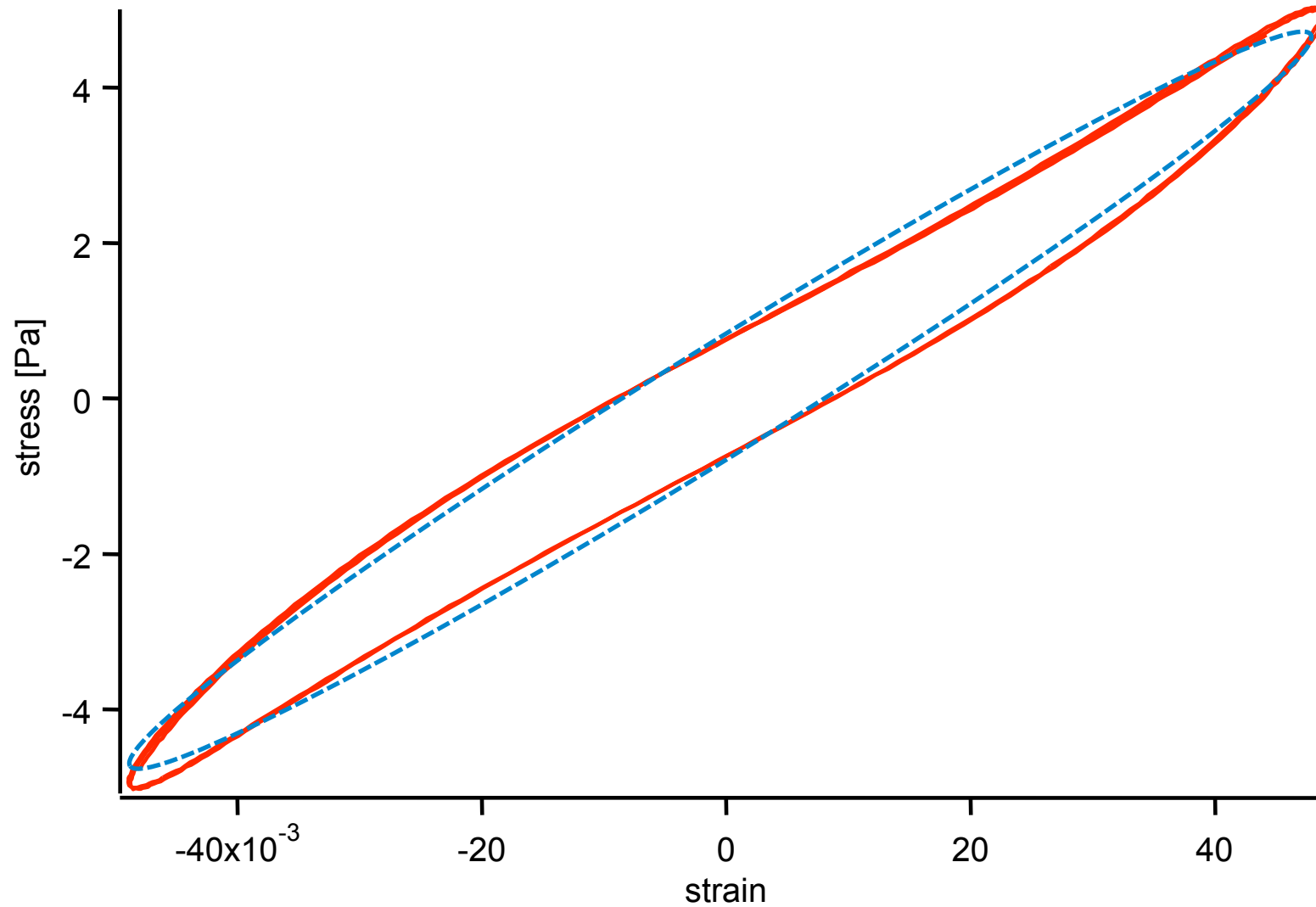
Oscillatory strain sweeps (collagen gels)



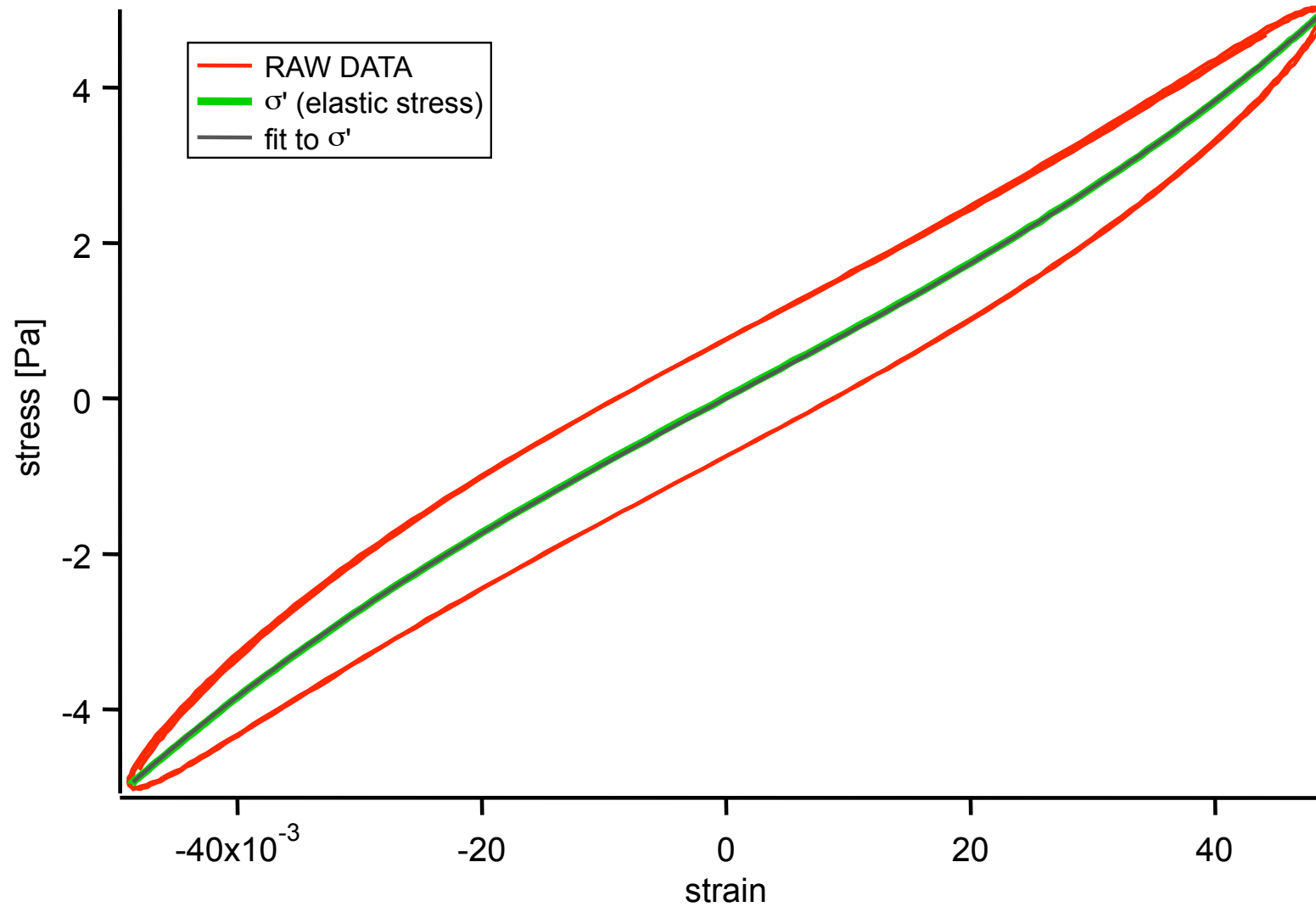
Lissajou plots from the G2 Raw data tool



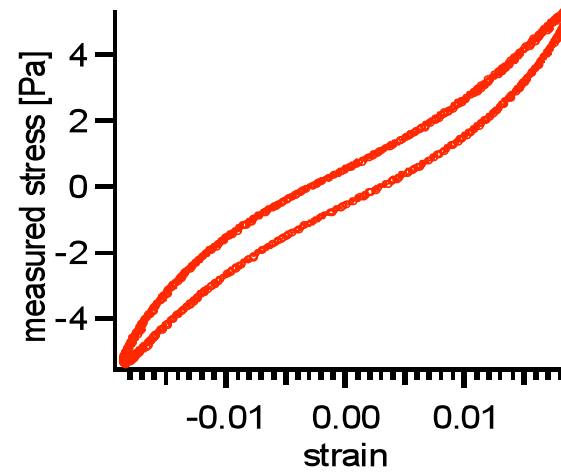
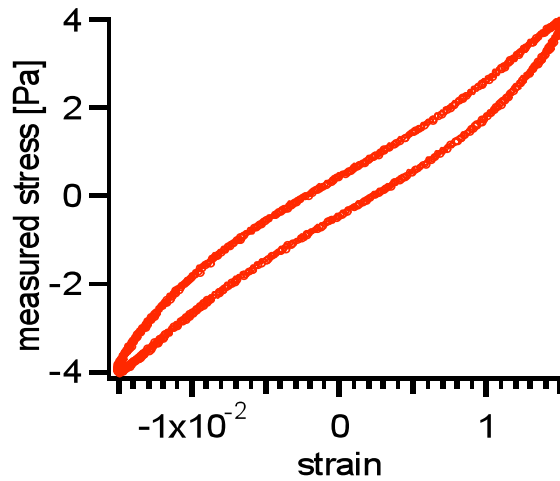
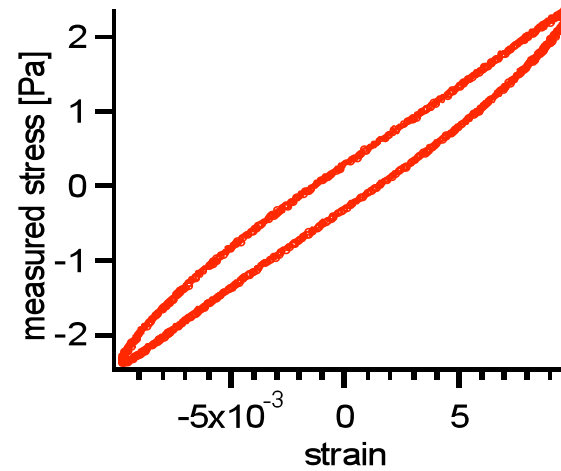
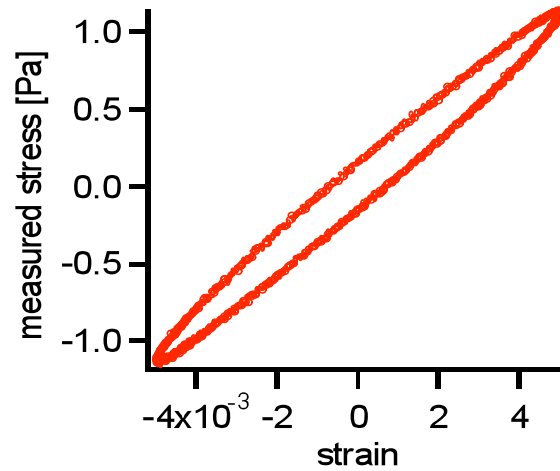
Nonlinear Lissajou plot



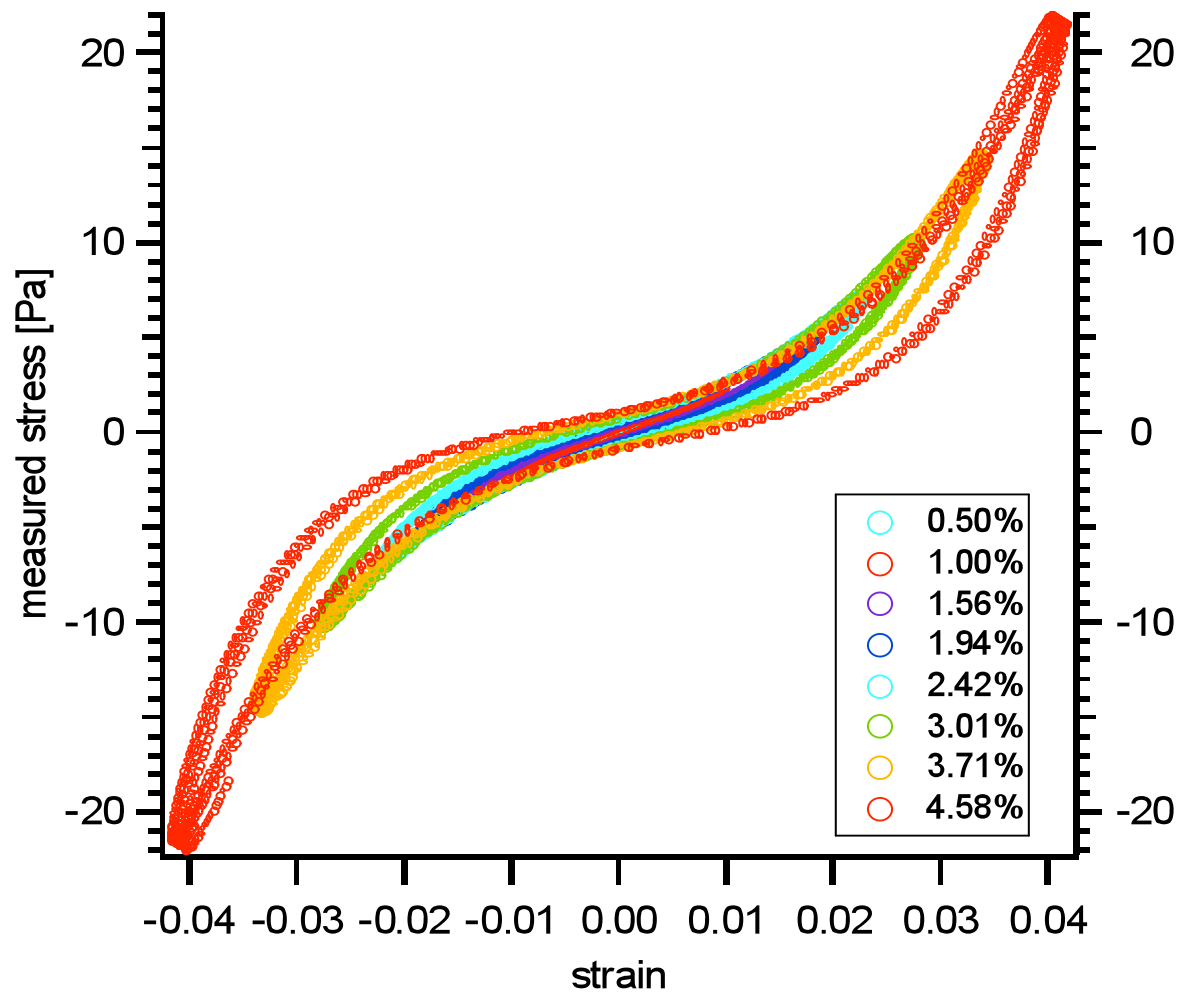
Nonlinear Lissajou plot



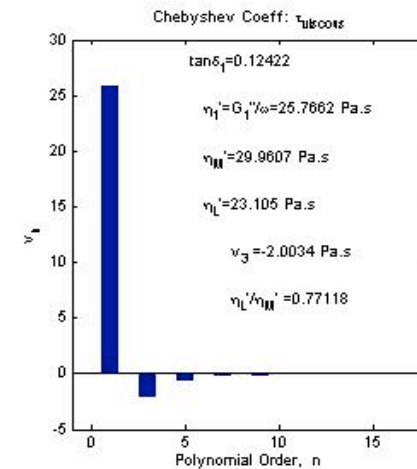
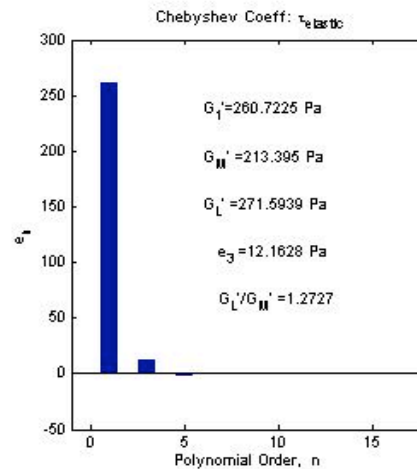
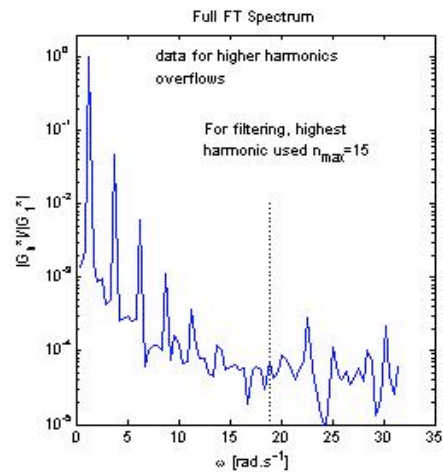
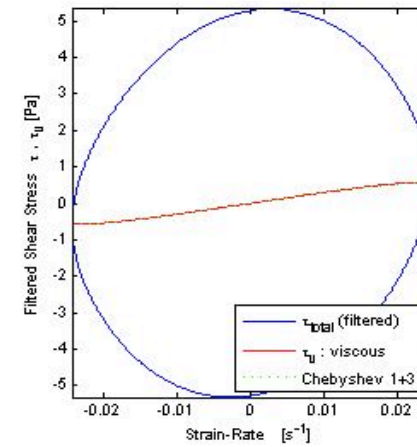
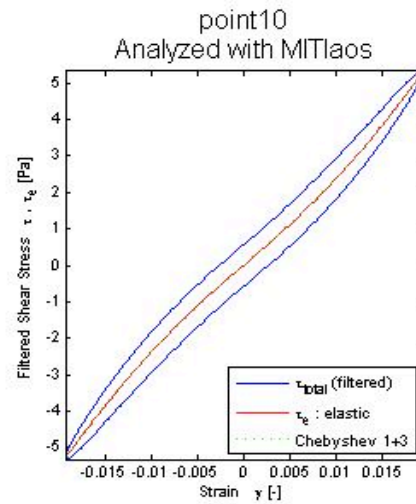
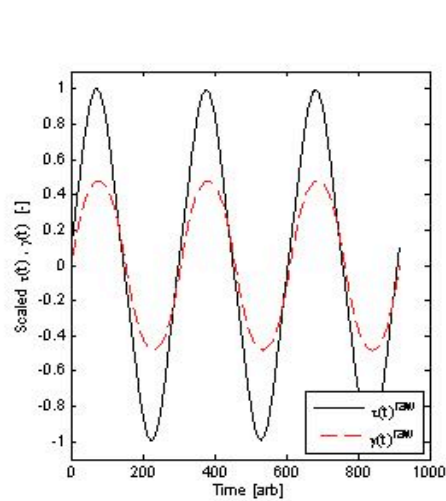
2.4mg/mL cone-plate



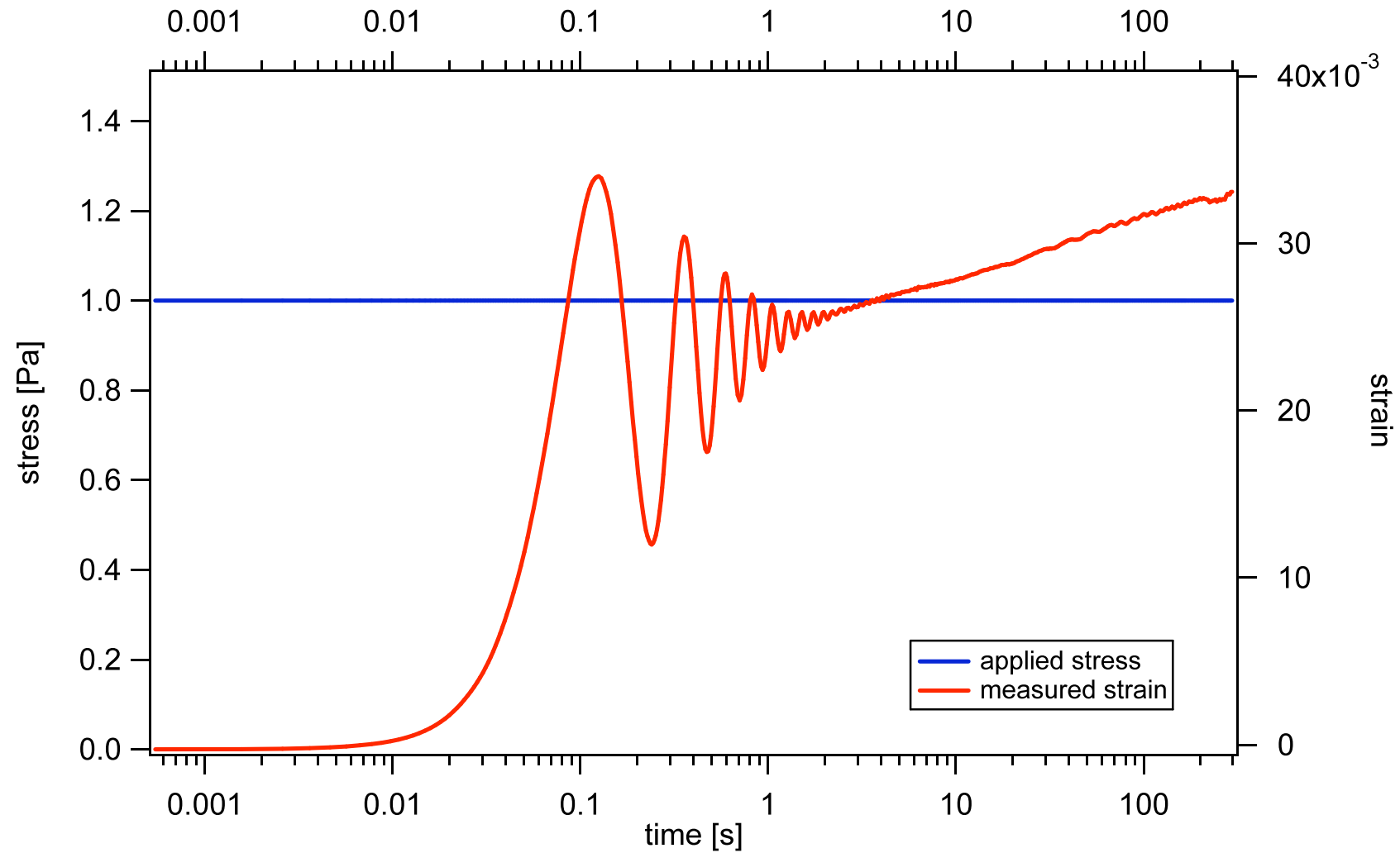
2.4mg/mL cone-plate



MIT LAOS MATLAB package

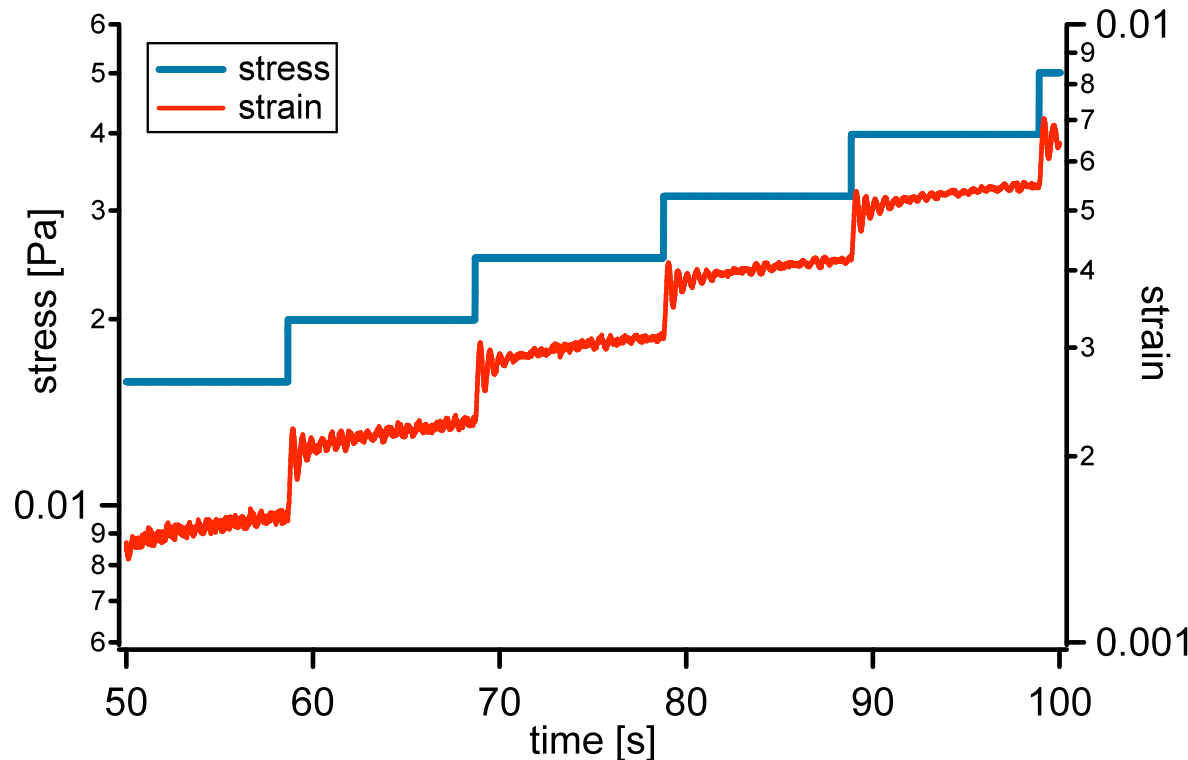


Creep-ringing



Creep-ringing

- Norman & Ryan's work here (fibrin, jamming)
- Good tutorial by Ewoldt & McKinley (MIT)



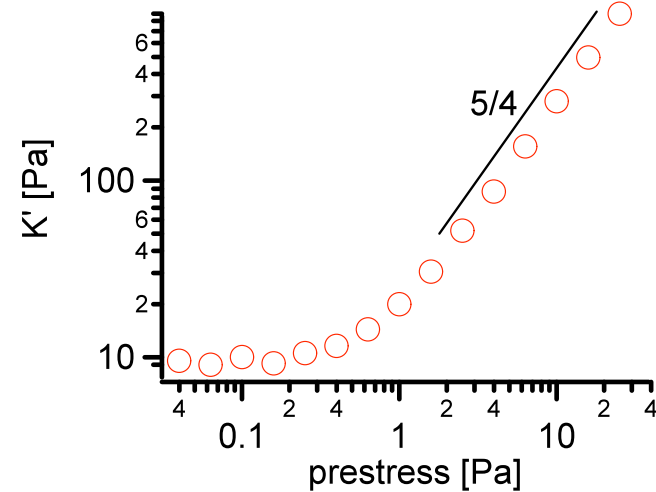
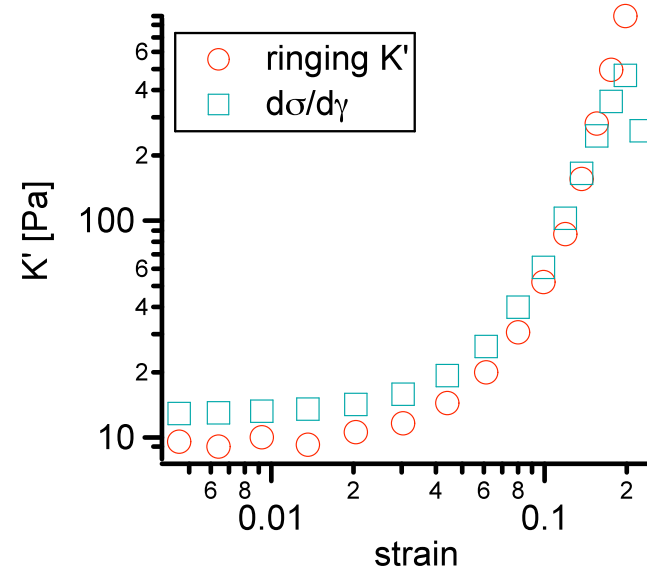
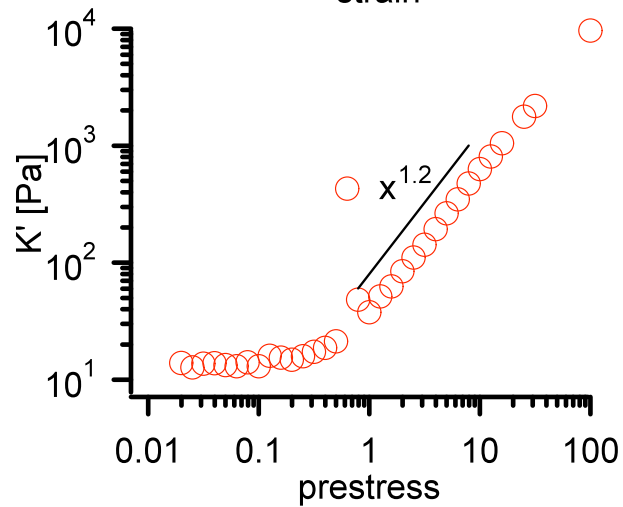
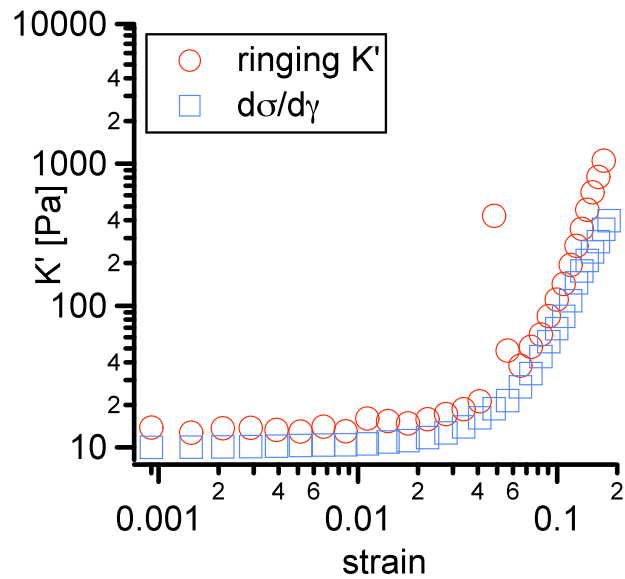
$$G' \approx \frac{I\omega^2}{b} \left(1 + \left(\Delta/2\pi \right)^2 \right)$$

$$G'' \approx \frac{I\omega^2}{b} \left(\Delta/\pi \right)$$

$$\Delta = \frac{1}{n} \ln \left(\frac{A_1}{A_{n+1}} \right)$$

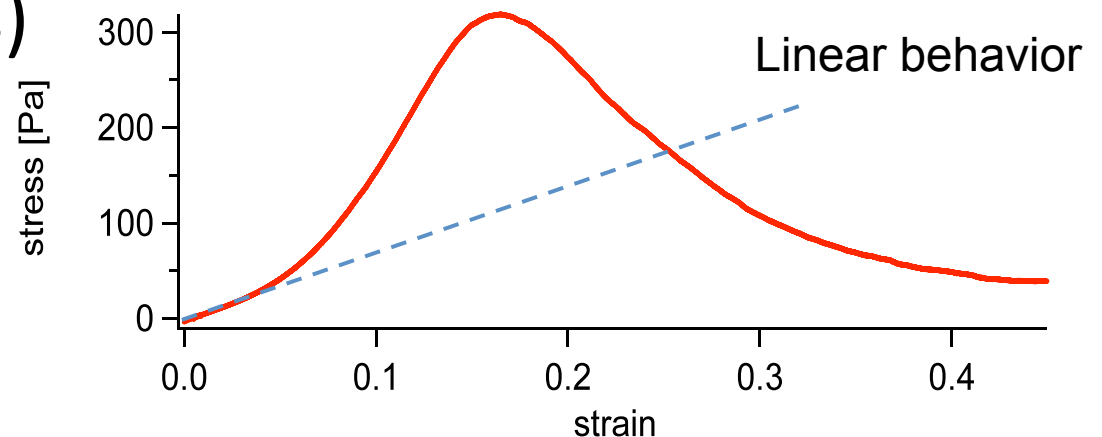
$$b = \frac{2\pi \cdot R^3}{3 \tan \theta}$$

Creep-ringing results



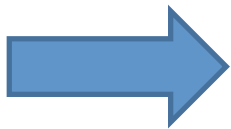
More nonlinear rheology

- Stress/strain ramps with constant rate
- Pre-stress measurements, i.e. small stress oscillations around a constant (pre-)stress
- Pre-strain measurements
- Transient responses in LAOS (talk to Stefan)
- Fourier domain analysis
- SRFS (talk to Hans)



Origin of nonlinear behavior

- Distribution of length-scales / inhomogeneities
- Rearrangement of particles / filaments
- Non-affine motion
- How do we find out?



Observation at the microscopic scale:

- Microrheology
- Microscopy

Microrheology basics

- General idea: look at the thermally-driven motion of micron-sized particles embedded in a material
- Mean-square displacement of particles as a function of time provides microscopic information on local elastic and viscous material properties as a function of frequency
- Mason and Weitz, PRL, 1995

Short and long timescales

Short time scales:
diffusive

$$\text{MSD}(\tau) \equiv \left\langle \left(\mathbf{r}_{t+\tau} - \mathbf{r}_t \right)^2 \right\rangle_t \sim 4D\tau$$

$$D \sim \frac{k_B T}{6\pi \cdot a \cdot \eta}$$

r: position vector
D: diffusion constant
 τ : lag time
kT: thermal energy
a: particle size
 η : viscosity

Long time scales:
spring-like

$$\frac{1}{2} K \cdot \text{MSD}(\tau) \sim \frac{1}{2} k_B T$$

K: effective spring-
constant, linked to
elastic properties

What about intermediate times?

Generalized Stokes-Einstein

$$D \sim \frac{k_B T}{6\pi \cdot a \cdot \eta} \quad \longrightarrow \quad \eta \sim \frac{k_B T}{6\pi \cdot a \cdot D} \sim \frac{k_B T \cdot 4\tau}{6\pi \cdot a \cdot \text{MSD}(\tau)}$$

$$\eta(\tau) \sim \frac{\tau}{\text{MSD}(\tau)}$$

Take Laplace transform of $\eta(\tau)$ numerically, to get $\eta(s)$ – with $s=i\omega$.
From earlier, we know:

$$G''(\omega) \sim \eta(\omega) \cdot \omega$$

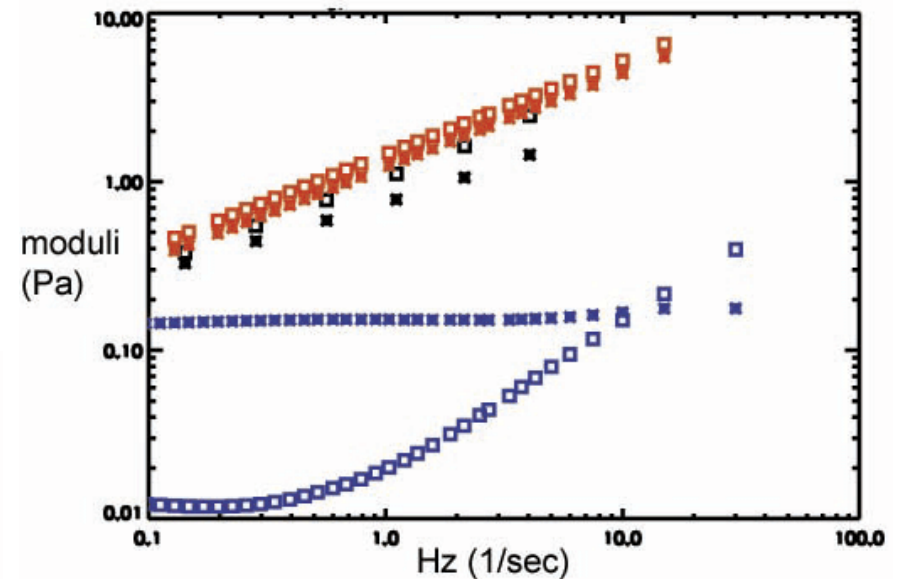
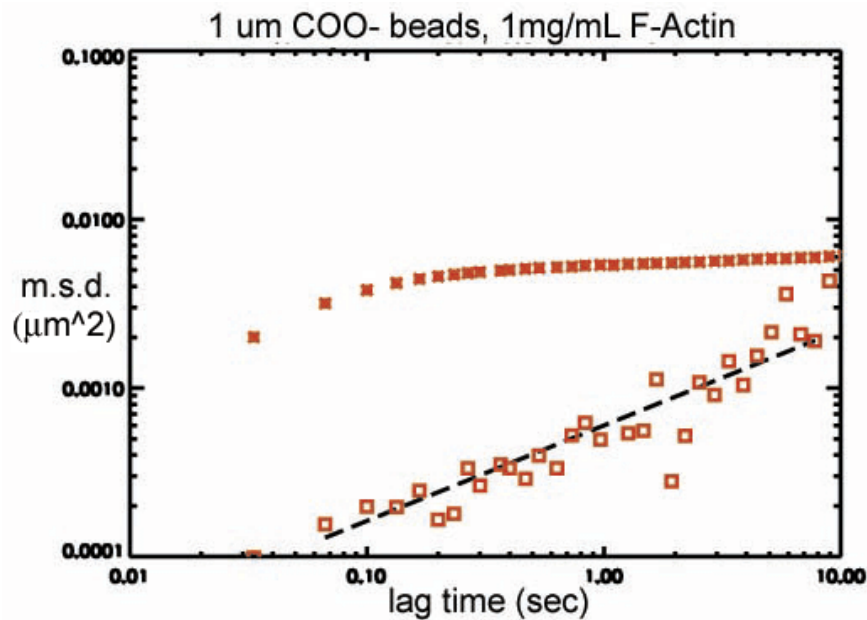
We can then get the generalized complex modulus, by analytically extending:

$$G''(s) \sim \eta(s) \cdot s$$

i.e. $G^*(i\omega) \sim \eta(i\omega) \cdot i\omega$

2-point vs 1-point microrheology

2-point microrheology calculates a mean-square displacement from the correlated pair-wise motion of particles, rather than the single-particle MSD.



Black: bulk rheology
Red: 2-point microrheology
Blue: 1-point microrheology
Open symbols: G'

Other considerations

- Non-linear regime non-trivial, but more interesting.
- Surface effects can be important.
- Imaging to figure out mechanisms.
- Richness of effects, mechanisms, time-, length- and energy- scales present in soft matter / complex fluids.
- More to explore on Weitzlab webpage.
- More at ComplexFluids meetings.