

Miami 2010

Probing Dark Matter with Neutrinos

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Erkoca, Gelmini, Reno and Sarcevic, Phys. Rev D81 (2010).
Erkoca, Reno and Sarcevic, Phys. Rev. D 82 (2010).

Dark Matter Searches

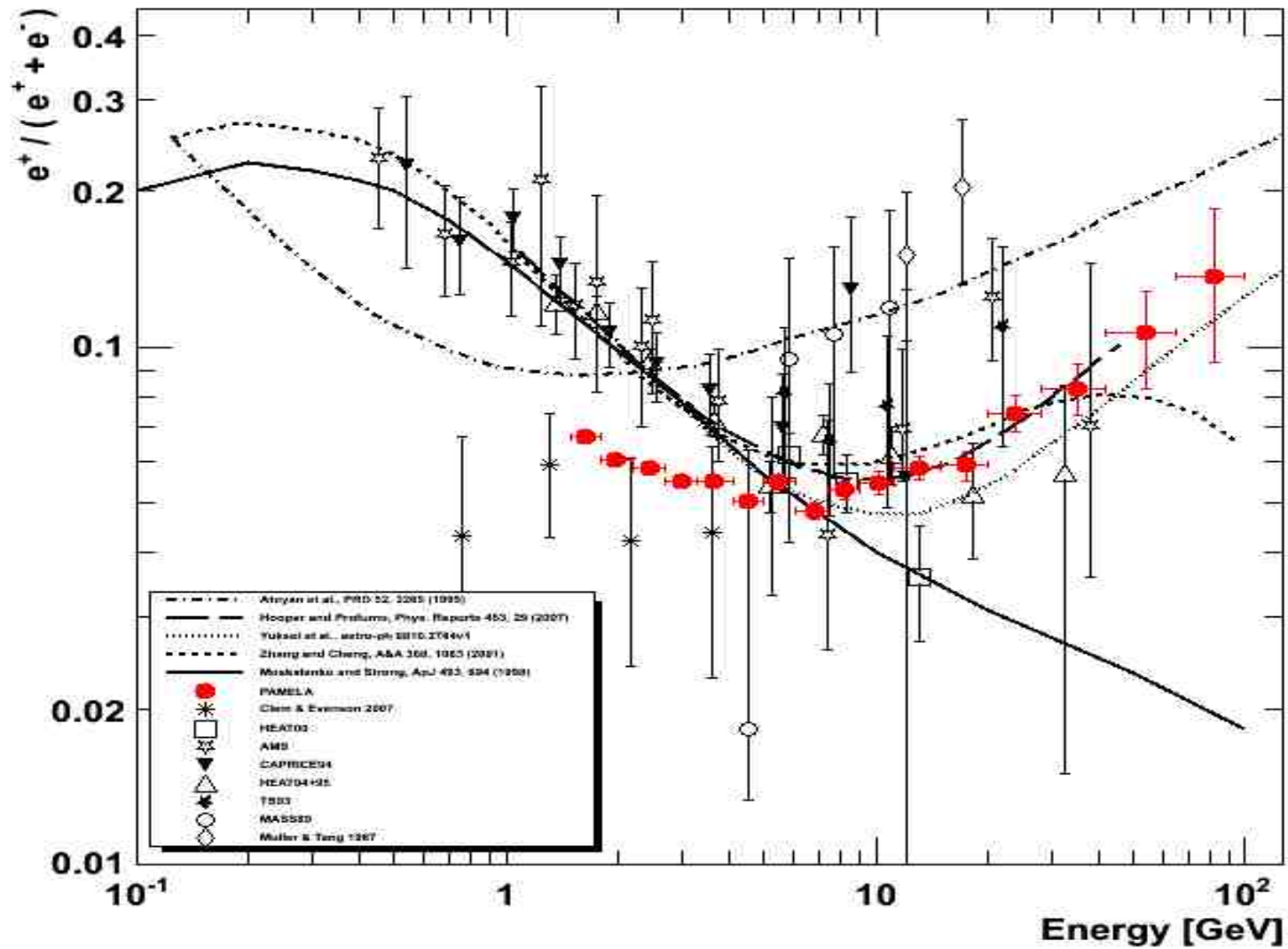
- Direct searches:

look for DM interactions with target nuclei (XENON, CDMS, DAMA, CoGeNT)

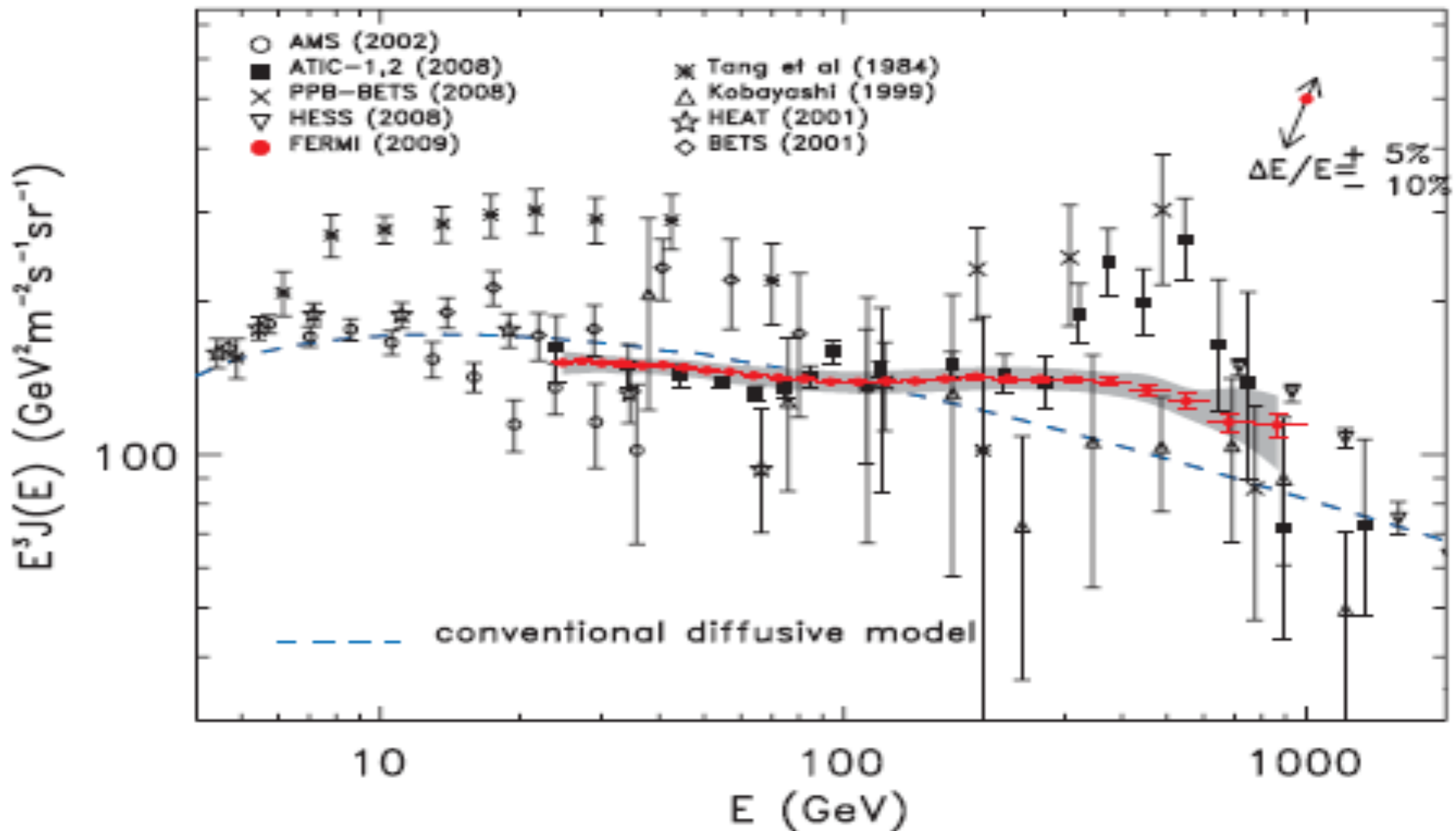
- Indirect searches:

DM annihilation producing electrons, positrons, gamma-rays (PAMELA, ATIC, FERMI/LAT, HESS ...) and neutrinos (IceCube, KM3Net...)

PAMELA Positron Fraction



FERMI Cosmic Ray Electron Spectrum



Neutrino Flux from DM Annihilation in the Galactic Center

Erkoca, Gelmini, Reno and Sarcevic,
Phys. Rev. D81, 096007 (2010)

- Model independent DM signals with neutrino-induced upward and contained muons and cascades (showers)
- Predictions for IceCube and Km3Net

Neutrino Flux from Dark Matter

Neutrino flux from DM annihilation/decay:

$$\left(\frac{d\phi_\nu}{dE_\nu} \right) = R \times \sum_F B_F \left(\frac{dN_\nu}{dE_\nu} \right)_F$$

here R for DM annihilation is:

$$R = B \frac{\langle \sigma v \rangle}{8\pi m_\chi^2} \int d\Omega \int_{l.o.s} \rho(l)^2 dl$$

and for DM decay:

$$R = \frac{1}{4\pi m_\chi \tau} \int d\Omega \int_{l.o.s} \rho(l) dl$$

Define $\langle J_n \rangle_\Omega$ as:

$$\langle J_n \rangle_\Omega = \int \frac{d\Omega}{\Delta\Omega} \int_{l.o.s.} \frac{dl(\theta)}{R_o} \left(\frac{\rho(l)}{\rho_o} \right)^n$$

$l(\theta)$ distance from us in the direction of the cone-half angle θ from the GC

$\rho(l)$ density distribution of dark matter halos

R_o distance of the solar system from the GC

ρ_o local dark matter density near the solar system

$$\langle \sigma v \rangle = 3 \times 10^{-26} \text{cm}^3 \text{s}^{-1}$$

$$R_o = 8.5 \text{kpc} \quad \rho_o^2 = 0.3 \text{GeV cm}^{-3}$$

Neutrino Flux (dN_ν/dE_ν) at the Production

Neutrinos can be produced directly or through decays of leptons, quarks and gauge bosons:

$$\chi\chi \rightarrow \nu_i\bar{\nu}_i$$

$$\rightarrow \tau^-\tau^+ \rightarrow (\nu_\tau l^-\bar{\nu}_l)(\bar{\nu}_\tau l^+\nu_l)$$

$$\rightarrow W^+W^- \rightarrow (l^+\nu_l)(l^-\bar{\nu}_l)$$

$$\rightarrow b\bar{b} \rightarrow (cl^-\bar{\nu}_l)(\bar{c}l^+\nu_l)$$

$$\rightarrow t\bar{t} \rightarrow bW^+\bar{b}W^- \rightarrow (cl^-\bar{\nu}_l)(l^+\nu_l)(\bar{c}l^+\nu_l)(l^-\bar{\nu}_l)$$

- Detection: neutrinos interacting below detector or in the detector producing muons
- Signals: upward and contained muons and cascade/showers
- Upward muons lose energy before reaching the detector

- Energy loss of the muons over a distance dz :

$$\frac{dE}{dz} = -(\alpha + \beta E)\rho$$

- α : ionization energy loss $\alpha = 10^{-3}\text{GeVcm}^2/\text{g}$.
- β : bremsstrahlung, pair production and photonuclear interactions $\beta=10^{-6}\text{cm}^2/\text{g}$.
- Relation between the initial and the final muon energy:

$$E_{\mu}^i(z) = e^{\beta\rho z} E_{\mu}^f + (e^{\beta\rho z} - 1) \frac{\alpha}{\beta}$$

Muon range: $R_{\mu} \equiv z = \frac{1}{\beta\rho} \log \left(\frac{\alpha + \beta E_{\mu}^i}{\alpha + \beta E_{\mu}^f} \right)$

Contained and Upward Muon Flux

Contained muon flux is given by

$$\frac{d\phi_{\mu}}{dE_{\mu}} = \int_{E_{\mu}}^{E_{max}} dE_{\nu} \left(\frac{dN}{dE_{\nu}} \right) N_A \rho \frac{d\sigma_{\nu}(E_{\nu})}{dE_{\mu}}$$

Upward muon flux is given by

$$\frac{d\phi_{\mu}}{dE_{\mu}} = \int_0^{R_{\mu}(E_{\mu}^i, E_{\mu})} e^{\beta \rho z} dz \int_{E_{\mu}^i}^{E_{max}} dE_{\nu} \left(\frac{dN}{dE_{\nu}} \right) N_A \rho \\ \times P_{surv}(E_{\mu}^i, E_{\mu}) \frac{d\sigma_{\nu}(E_{\nu})}{dE_{\mu}}$$

Hadronic Shower Flux

$$\frac{d\phi_{sh}}{dE_{sh}} = \int_{E_{sh}}^{E_{max}} dE_{\nu} \left(\frac{d\phi_{\nu}}{dE_{\nu}} \right) N_A \rho \frac{d\sigma_{\nu}(E_{\nu}, E_{\nu} - E_{sh})}{dE_{sh}}$$

Neutrino Energy Distribution

- $\chi\chi \rightarrow \nu\bar{\nu}$ channel :

$$\frac{dN_\nu}{dE_\nu} = \delta(E_\nu - m_\chi)$$

- $\chi\chi \rightarrow \tau^+\tau^-, b\bar{b}, c\bar{c}$ channels :

$$\frac{dN_\nu}{dE_\nu} = \frac{2B_f}{E_{in}}(1 - 3x^2 + 2x^3), \quad \text{where } x = \frac{E_\nu}{E_{in}} \leq 1$$

$$(E_{in}, B_f) = \begin{cases} (m_\chi, 0.18) & \tau \text{ decay} \\ (0.73m_\chi, 0.103) & b \text{ decay} \\ (0.58m_\chi, 0.13) & c \text{ decay.} \end{cases}$$

- $\chi\chi \rightarrow W^+W^-, ZZ$ channels :

$$\frac{dN_\nu}{dE_\nu} = n_f \frac{B_f}{m_\chi \beta} \quad \text{if} \quad \frac{m_\chi}{2}(1 - \beta) < E_\nu < \frac{m_\chi}{2}(1 + \beta)$$

where β is the velocity of the decaying particle (W or Z)

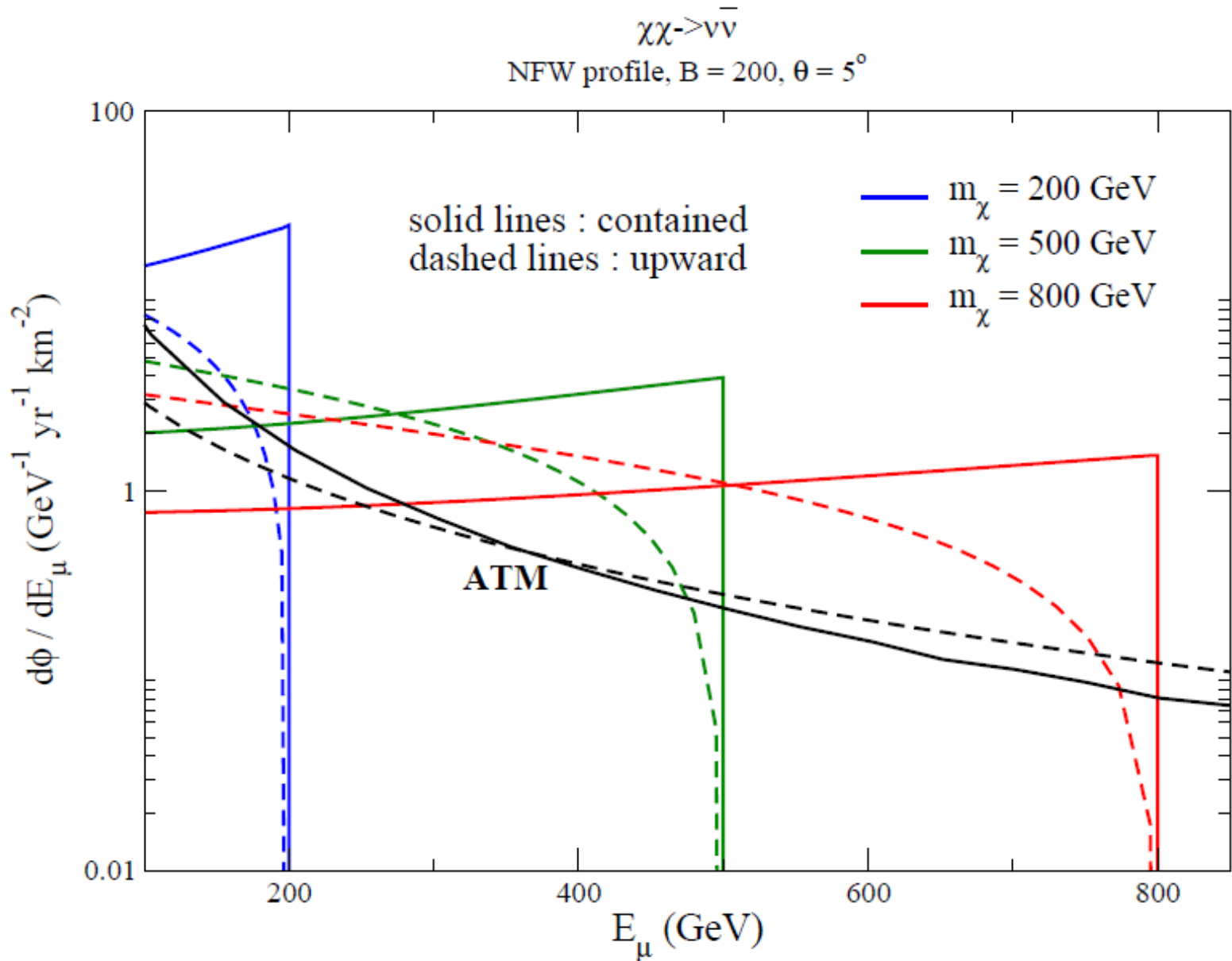
$$(n_f, B_f) = \begin{cases} (1, 0.105) & W \text{ decay,} \\ (2, 0.067) & Z \text{ decay.} \end{cases}$$

- $\chi\chi \rightarrow t\bar{t}$ channel :

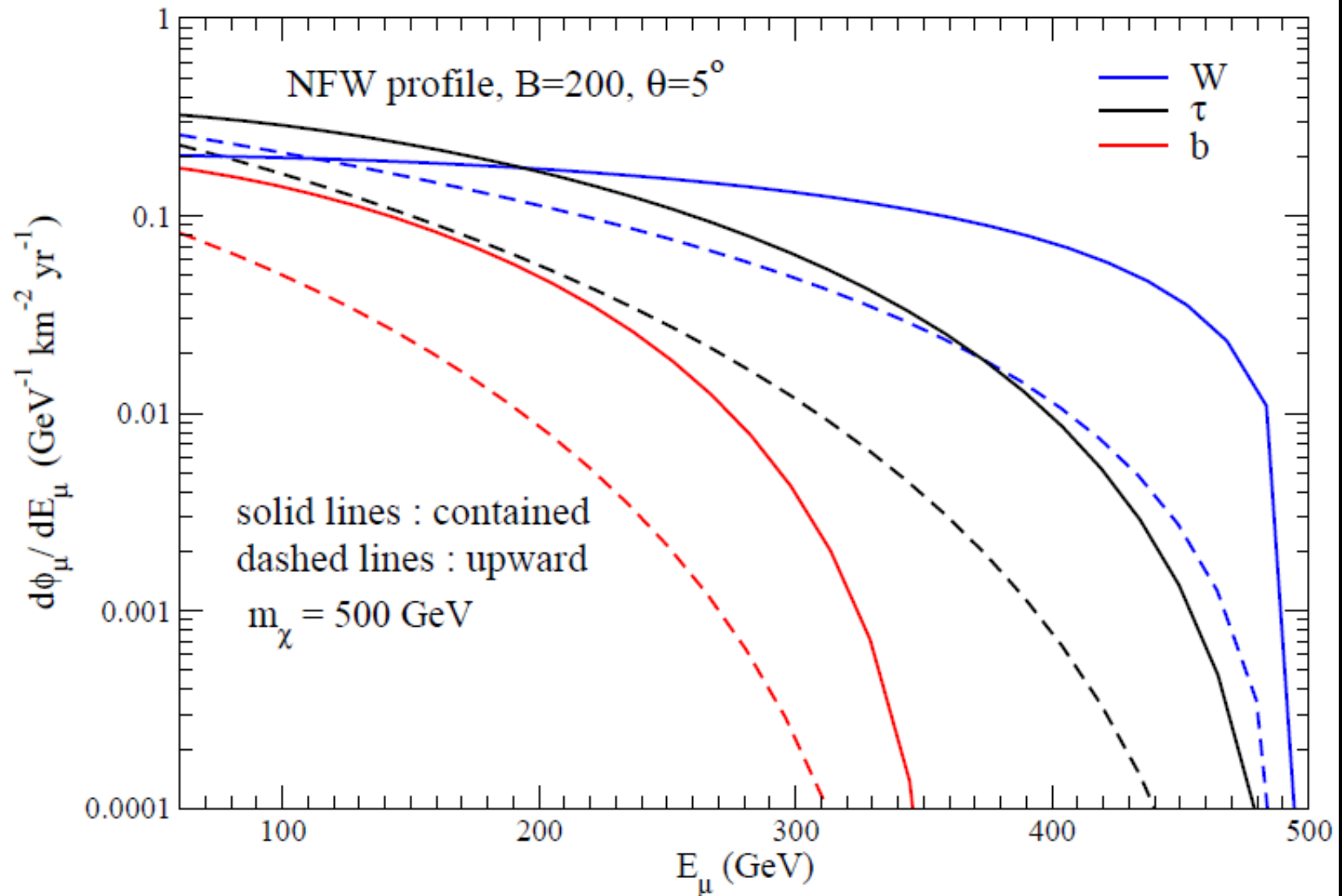
$$\left(\frac{dN_\nu}{dE_\nu}\right)_{t\bar{t}}^{\text{rest}} = \left(\frac{dN_\nu}{dE_\nu}\right)_{W+W^-} + \left(\frac{dN_\nu}{dE_\nu}\right)_{b\bar{b}}$$

Boosting this expression yields the neutrino spectrum for top quarks moving with velocity β_t

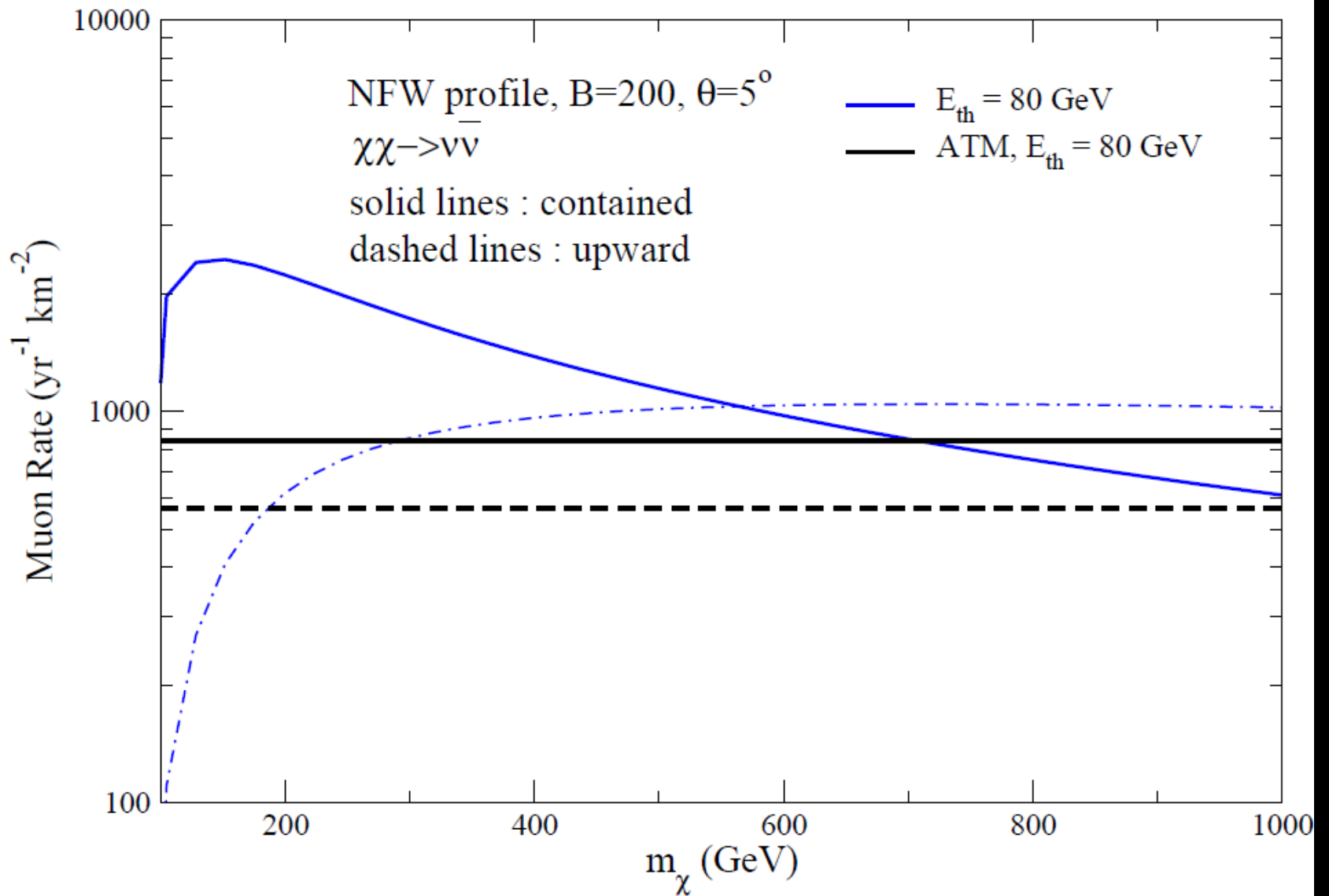
Muon Flux

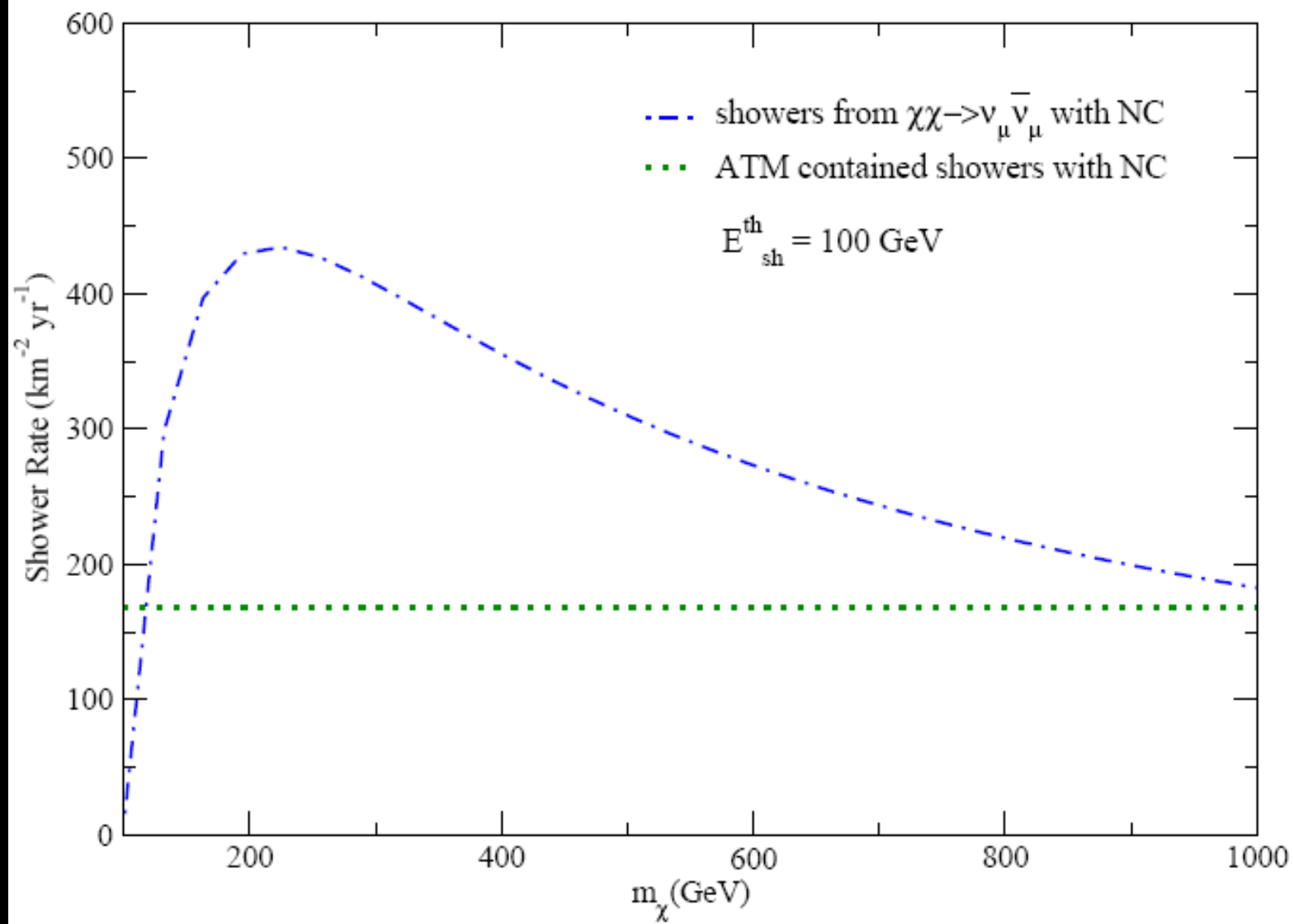


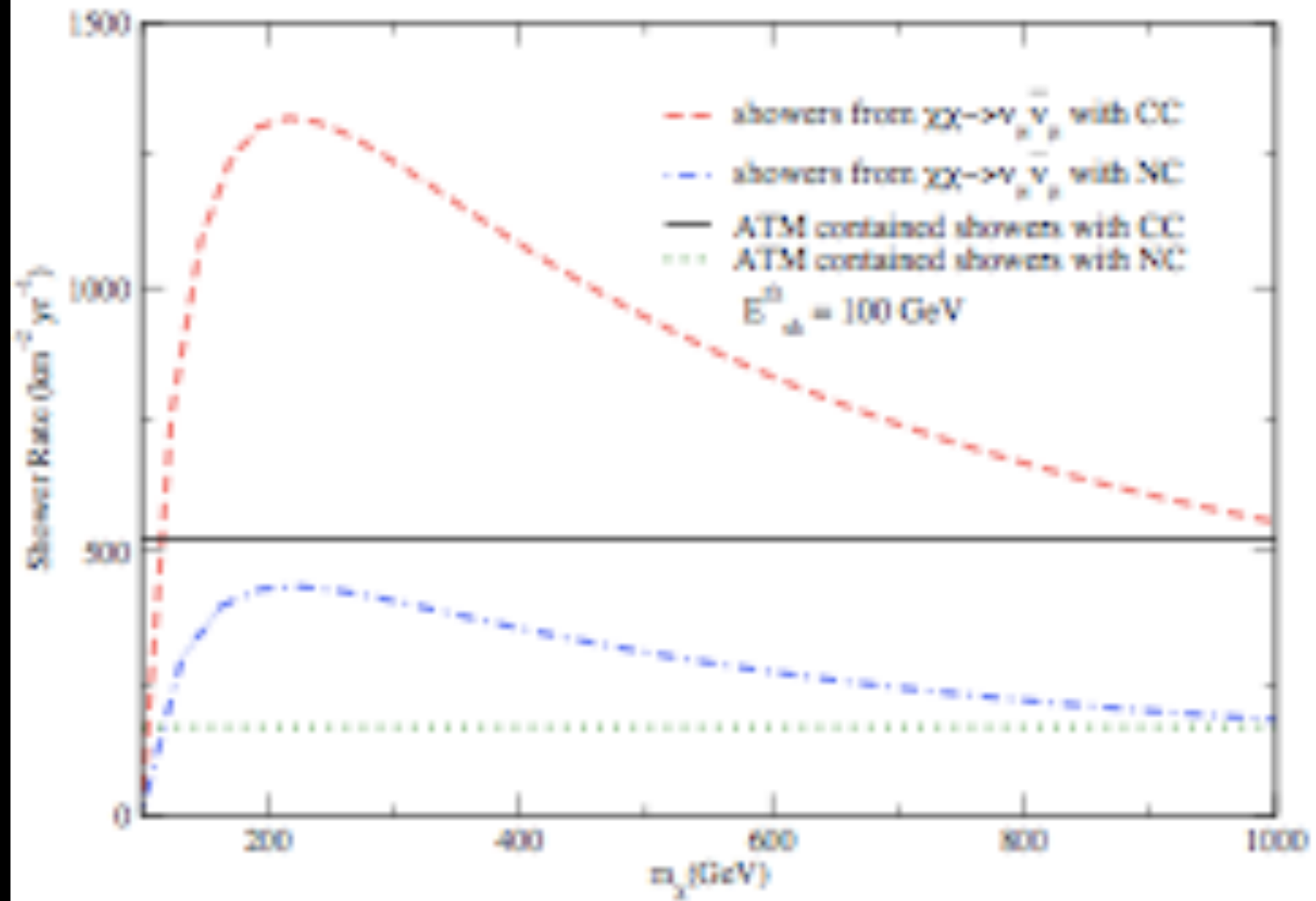
Muon Flux for Different DM Annihilation Modes



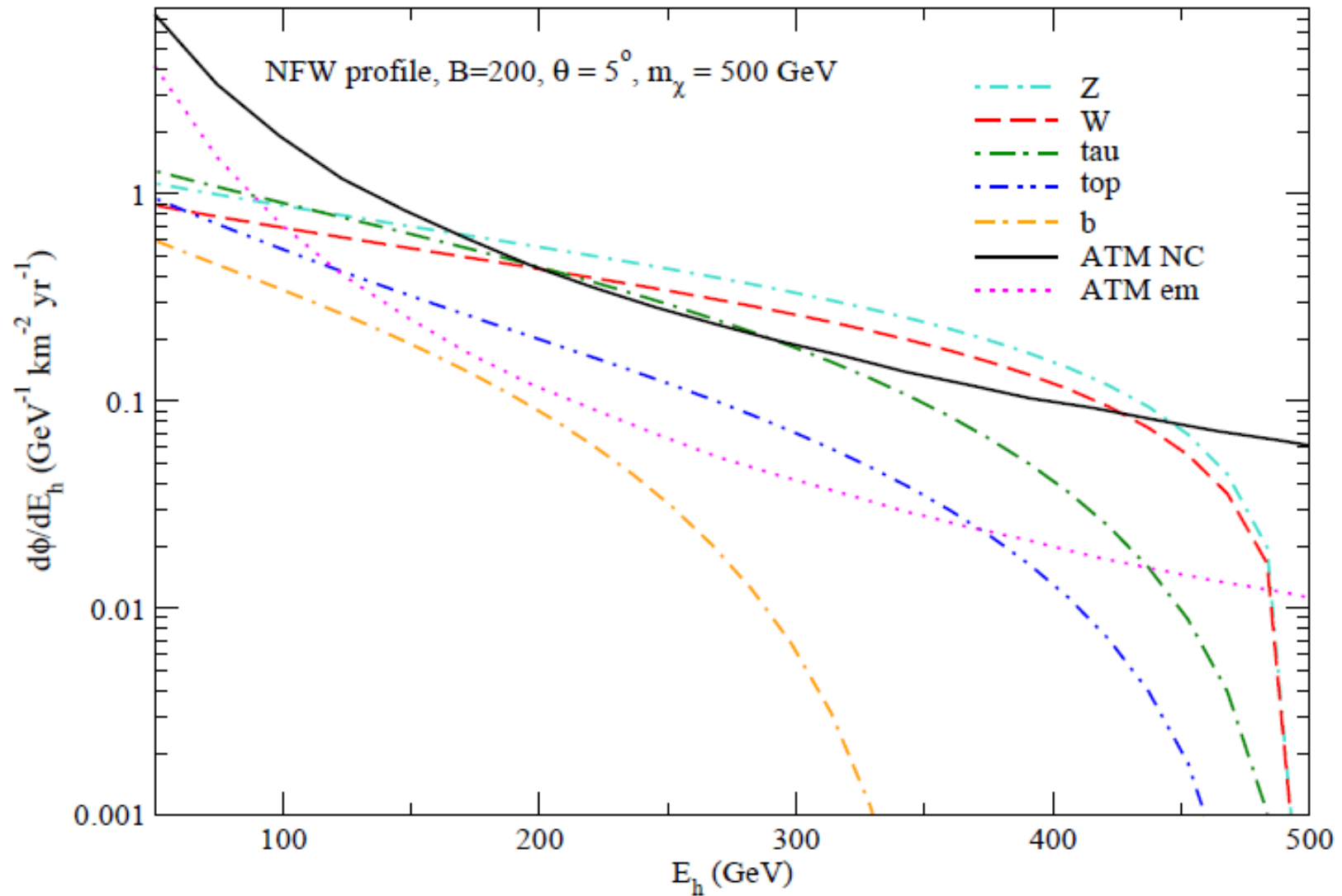
Muon Rates

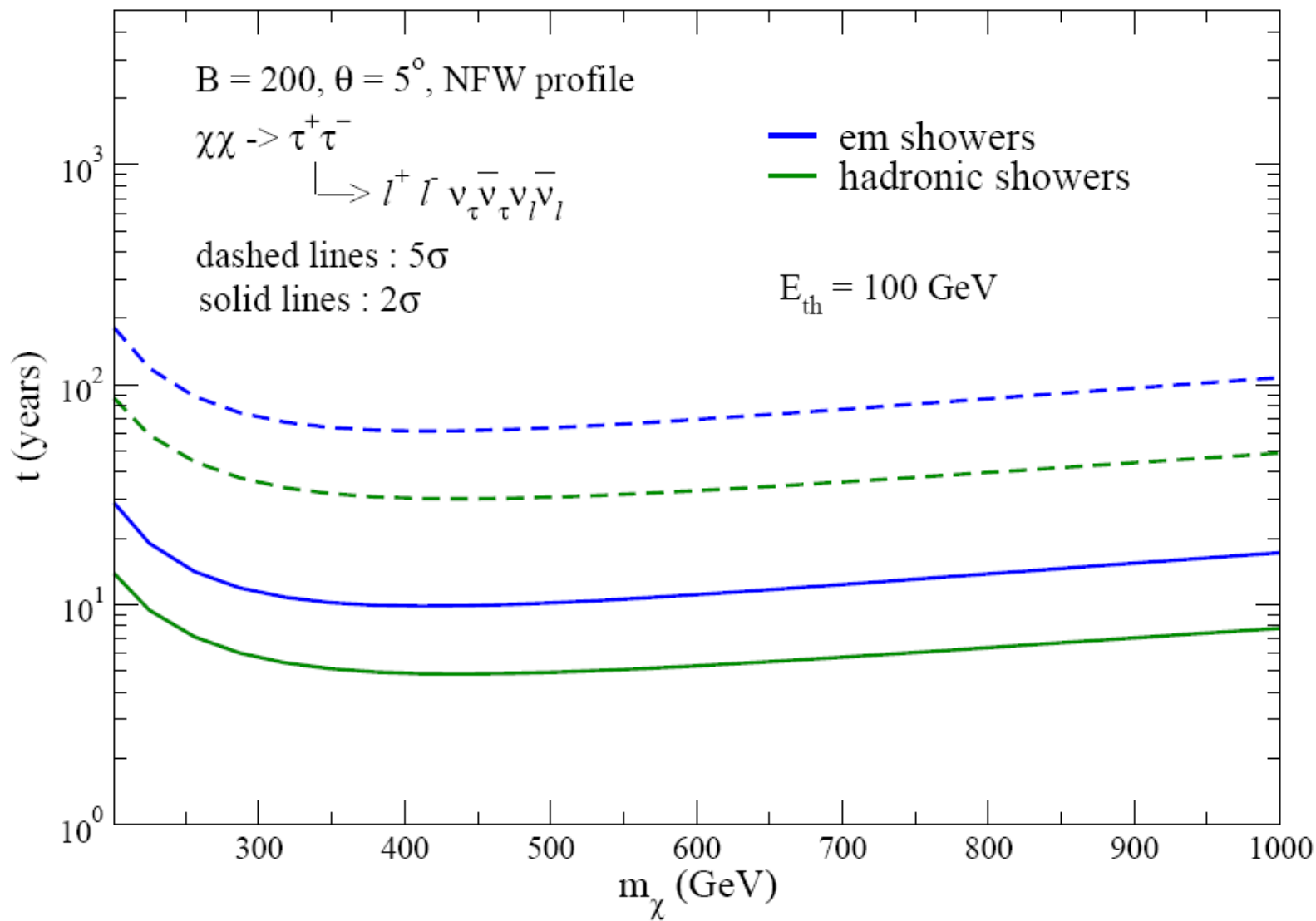






Hadronic Shower Spectra without track-like events



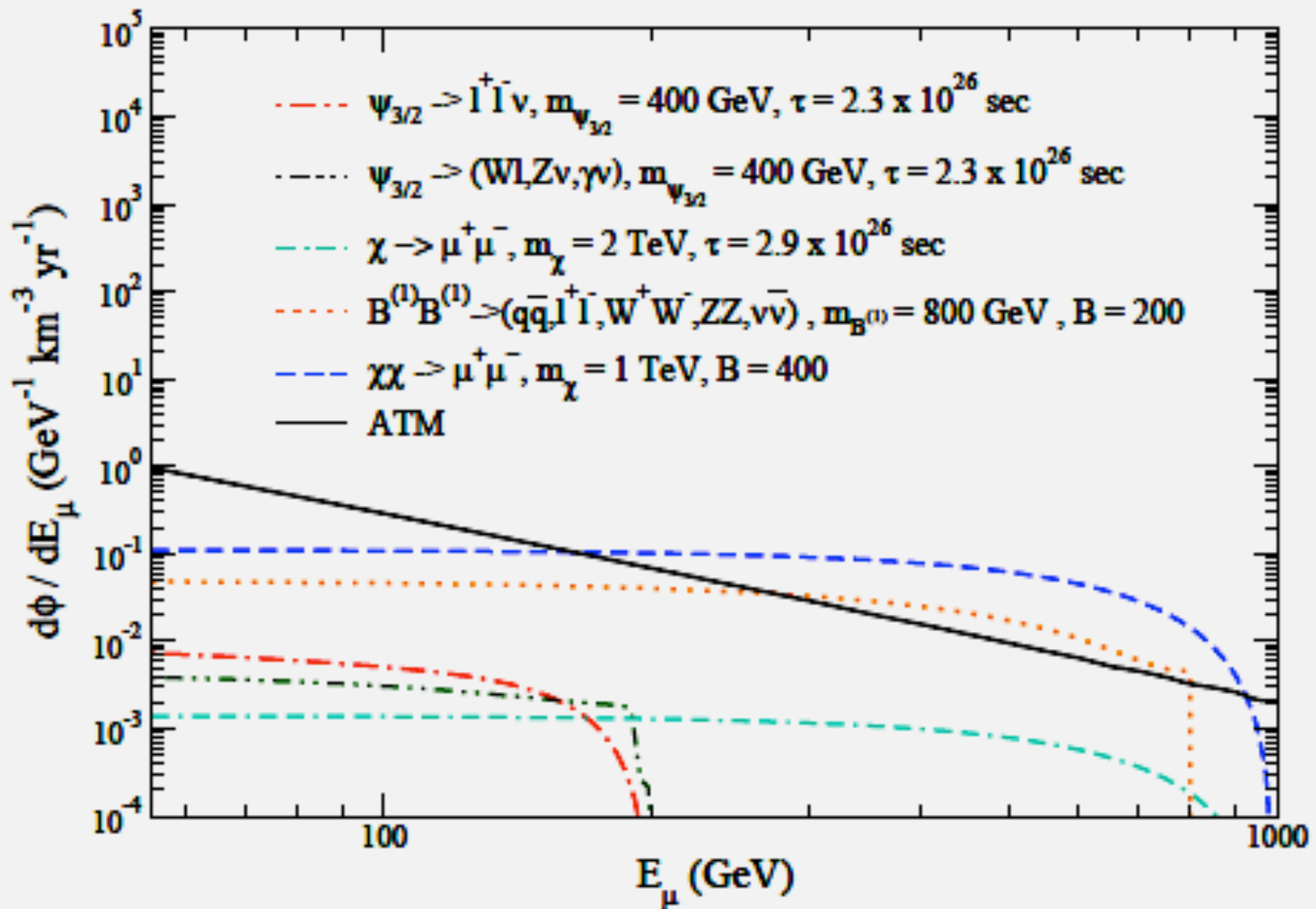


Probing the Nature of Dark Matter with Neutrinos

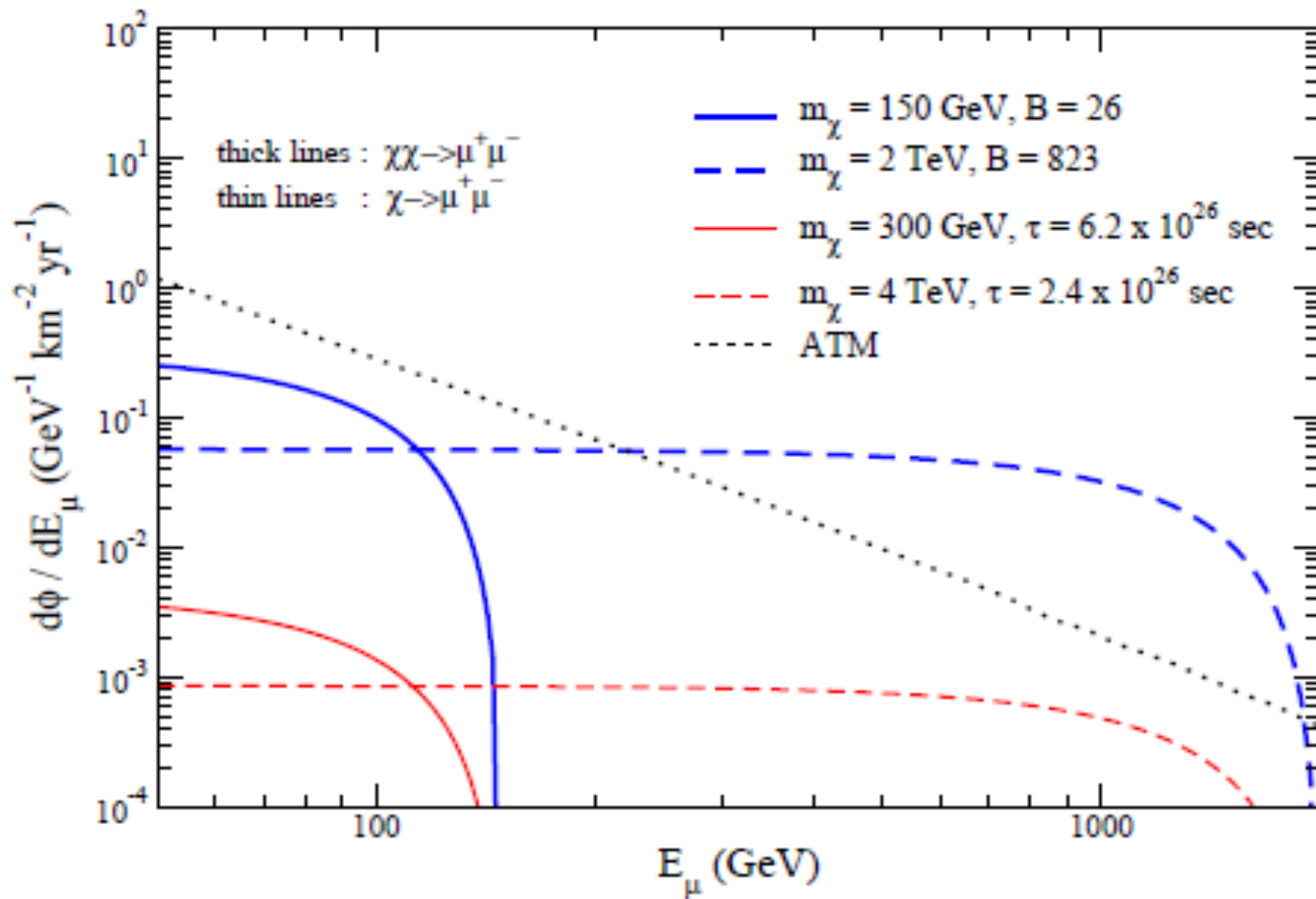
Erkoca, Reno and Sarcevic, Phys. Rev. D82,
113006 (2010)

- DM candidates: gravitino, Kaluza-Klein particle, a particle in leptophilic models.
- Dark matter signals: upward and contained muon flux and cascades (showers) from neutrino interactions
- Experimental signatures that would distinguish between different DM candidates

Contained Muon Flux



Contained Muon Flux



| | | m_χ (TeV) | | | | | | | | | | |
|--|------------------------|--------------------|-----------------|-------|-------------|-------------|-------------|------------|-----------|-------|-------|----|
| | | 0.2 | 0.4 | 0.6 | 0.8 | 1 | 2 | 4 | 6 | 8 | 10 | |
| $\psi_{3/2} \rightarrow l^+ l^- \nu$ $B_\tau = 2.3$ | $N_\mu^{ct}(50^\circ)$ | 4.94 | 11.15 | 13.8 | 15.3 | 16.2 | 18.1 | 19.0 | 19.3 | 19.5 | 19.6 | |
| | $N_\mu^{up}(50^\circ)$ | 8.68 | 59.5 | 120 | 180 | 239 | 503 | 912 | 1228 | 1485 | 1704 | |
| | $N_{sh}(50^\circ)$ | 4 | 11 | 13 | 15 | 16.3 | 19 | 21 | 22 | 22 | 22 | |
| | $t_\mu^{up}(10^\circ)$ | 1.3×10^4 | 277 | 69 | 30 | 17 | 4 | 1.2 | 0.7 | 0.5 | 0.4 | |
| | $t_\mu^{up}(50^\circ)$ | 3490 | 74 | 18 | 8 | 5 | 1 | 0.32 | 0.18 | 0.12 | 0.09 | |
| | $t_{sh}(50^\circ)$ | 196 | 23 | 16 | 12 | 10 | 7 | 6.3 | 5.8 | 5.8 | 5.8 | |
| $\psi_{3/2} \rightarrow (Wl, Z\nu, \gamma\nu)$ $B_\tau = 2.3$ | $N_\mu^{ct}(50^\circ)$ | 6.1 | 8.4 | 8.9 | 9.1 | 9.15 | 9.2 | 9.2 | 9.2 | 9.2 | 9.2 | |
| | $N_\mu^{up}(50^\circ)$ | 9.9 | 50.9 | 95.6 | 139 | 181 | 364 | 638 | 844 | 1010 | 1150 | |
| | $N_{sh}(50^\circ)$ | 3.6 | 7.66 | 9.6 | 10.74 | 11.5 | 13.17 | 14.12 | 14.46 | 14.64 | 14.74 | |
| | $t_\mu^{up}(10^\circ)$ | 1×10^4 | 378 | 107 | 51 | 30 | 7.5 | 2.5 | 1.4 | 1 | 0.8 | |
| | $t_\mu^{up}(50^\circ)$ | 2693 | 101 | 29 | 14 | 8 | 2 | 0.7 | 0.4 | 0.3 | 0.2 | |
| | $t_{sh}(50^\circ)$ | 210 | 47 | 30 | 24 | 21 | 16 | 14 | 13 | 13 | 13 | |
| $\chi \rightarrow \mu^+ \mu^-$ $B_\tau = 2.9$ | $N_\mu^{ct}(50^\circ)$ | 2.13 | 6.45 | 8.43 | 9.5 | 10.2 | 11.5 | 12.2 | 12.4 | 12.5 | 12.6 | |
| | $N_\mu^{up}(50^\circ)$ | 3.14 | 29 | 62.3 | 97 | 131 | 286 | 533 | 728 | 886 | 1022 | |
| | $N_{sh}(50^\circ)$ | 1.95 | 8.22 | 12.09 | 14.55 | 16.2 | 20.2 | 22.45 | 23.27 | 23.68 | 23.94 | |
| | $t_\mu^{up}(10^\circ)$ | 1×10^5 | 1×10^3 | 252 | 104 | 57 | 12 | 3.5 | 1.9 | 1.3 | 0.97 | |
| | $t_\mu^{up}(50^\circ)$ | 2.6×10^4 | 316 | 68 | 28 | 15 | 3.2 | 0.93 | 0.5 | 0.34 | 0.26 | |
| | $t_{sh}(50^\circ)$ | 709 | 40 | 19 | 13 | 11 | 6.9 | 5.5 | 5.2 | 5 | 4.8 | |
| $B^{(1)} B^{(1)} \rightarrow \dots$ $B = 200$ | $N_\mu^{ct}(10^\circ)$ | 14.2 | 9.8 | 7.2 | 5.6 | 4.6 | 2.4 | 1.25 | 0.84 | 0.63 | 0.51 | |
| | $N_\mu^{up}(10^\circ)$ | 86.1 | 131 | 140 | 130 | 128 | 124 | 108 | 92 | 81 | 72 | |
| | $N_{sh}(10^\circ)$ | 11 | 9 | 7 | 5.7 | 4.8 | 2.6 | 1.4 | 0.9 | 0.7 | 0.6 | |
| | $t_\mu^{up}(1^\circ)$ | 1.27 | 0.63 | 0.54 | 0.65 | 0.66 | 0.7 | 0.87 | 1.14 | 1.42 | 1.72 | |
| | $t_\mu^{up}(10^\circ)$ | 1.55 | 0.68 | 0.57 | 0.71 | 0.72 | 0.76 | 1.0 | 1.36 | 1.76 | 2.2 | |
| | $t_\mu^{up}(50^\circ)$ | 5.1 | 2.2 | 1.84 | 2.29 | 2.3 | 2.44 | 3.2 | 4.5 | 5.8 | 7.2 | |
| | $t_{sh}(1^\circ)$ | 3.4 | 4.4 | 5.9 | 7.7 | 9.6 | 22 | 61 | 116 | 189 | 280 | |
| | $t_{sh}(10^\circ)$ | 1.3 | 1.9 | 2.9 | 4.3 | 5.8 | 18 | 64 | 136 | 237 | 364 | |
| | $t_{sh}(50^\circ)$ | 3.3 | 5 | 8 | 12 | 16.3 | 57 | 204 | 445 | 777 | 1202 | |
| $\chi\chi \rightarrow \mu^+ \mu^-$ $B = 400$ | $N_\mu^{ct}(10^\circ)$ | 40.19 | 29.58 | 22.01 | 17.39 | 14.3 | 7.59 | 3.90 | 2.63 | 1.98 | 1.59 | |
| | $N_\mu^{up}(10^\circ)$ | 144 | 241 | 273 | 283 | 320 | 266 | 221 | 190 | 167 | 151 | |
| | $N_{sh}(10^\circ)$ | 51.4 | 45.6 | 36.4 | 30 | 25 | 14 | 7.4 | 5 | 3.8 | 3 | |
| | $t_\mu^{ct}(1^\circ)$ | 1.11 | 1.68 | 2.55 | 3.61 | 4 | 13.64 | 44 | 92 | 156 | 238 | |
| | $t_\mu^{ct}(10^\circ)$ | 0.66 | 1.18 | 2.06 | 3.24 | 4.7 | 16.31 | 61 | 133 | 234 | 364 | |
| | $t_\mu^{ct}(50^\circ)$ | 1.93 | 3.55 | 6.38 | 10.2 | 15 | 53 | 201 | 444 | 781 | 1213 | |
| | $t_\mu^{up}(1^\circ)$ | 0.54 | 0.24 | 0.2 | 0.18 | 0.14 | 0.21 | 0.28 | 0.35 | 0.43 | 0.50 | |
| | $t_\mu^{up}(10^\circ)$ | 0.47 | 0.21 | 0.16 | 0.15 | 0.12 | 0.17 | 0.25 | 0.33 | 0.42 | 0.52 | |
| | $t_\mu^{up}(50^\circ)$ | 1.83 | 0.65 | 0.51 | 0.47 | 0.37 | 0.54 | 0.78 | 1.1 | 1.35 | 1.7 | |
| | $t_{sh}(1^\circ)$ | 0.63 | 0.72 | 0.91 | 1.12 | 1.37 | 2.58 | 5.5 | 9 | 13 | 18 | |
| | $t_{sh}(10^\circ)$ | 0.12 | 0.14 | 0.2 | 0.26 | 0.34 | 0.87 | 2.63 | 5.34 | 9 | 13.6 | |
| | | $t_{sh}(50^\circ)$ | 0.18 | 0.22 | 0.33 | 0.48 | 0.7 | 2.1 | 7.2 | 15.5 | 27 | 42 |
| | Atmospheric | N_μ^{ct} | 2.28(1°) | | | 227.5(10°) | | | 5347(50°) | | | |
| N_μ^{up} | | 28(1°) | | | 2794(10°) | | | 65668(50°) | | | | |
| N_{sh} | | 0.3(1°) | | | 28.8(10°) | | | 676(50°) | | | | |

DM Detection with Neutrino Telescopes

IceCUBE : 1 km³ neutrino detector at South Pole

- detects Cherenkov radiation from the charged particles produced in neutrino interactions
- contained and upward muon events and showers
- contained muons from GC
- showers from GC with IceCUBE+DeepCore

KM3Net : a future deep-sea neutrino telescope

- contained and upward muon events and showers
- upward muons from GC

Summary

- Neutrinos could be used to detect dark matter and to probe its physical origin
- Contained and upward muon flux is sensitive to the DM annihilation mode and to the mass of dark matter particle
- Combined measurements of cascade events and muons with IceCube+DeepCore and KM3Net look promising
- Neutrinos can probe DM candidates, such as gravitino, Kaluza-Klein DM, and a particle in leptophilic models