

MODELING OF INTEGRATED PLASMA PROCESSING: PLASMA PHYSICS, PLASMA CHEMISTRY AND SURFACE KINETICS

**Mark J. Kushner
University of Illinois
Department of Electrical and Computer Engineering
Urbana, IL 61801
mjk@uiuc.edu <http://uigelz.ece.uiuc.edu>**

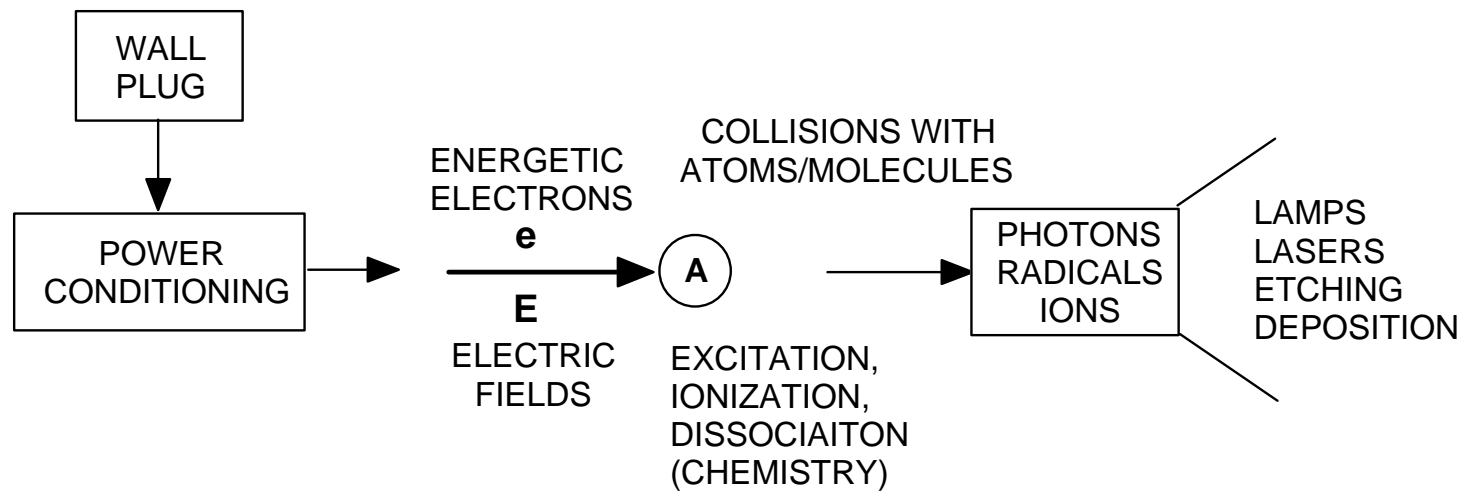
May 2003

AGENDA

- **Integration in Plasma Processing**
- **Modeling Requirements:**
 - **Plasma Physics**
 - **Plasma Chemistry**
 - **Surface Kinetics**
- **Integrated process modeling of etching and cleaning of porous silica; and metal deposition for interconnect wiring.**
- **Concluding Remarks**

COLLISIONAL LOW TEMPERATURE PLASMAS

- Partially ionized plasmas are gases containing neutral atoms and molecules, electrons, positive ions and negative ions. These systems are the plasmas of every day technology.

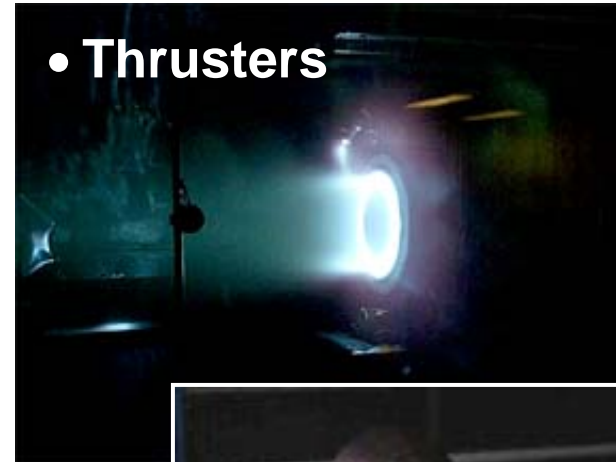


- Electrons transfer power from the "wall plug" to internal modes of atoms / molecules to "make a product", very much like combustion.
- The electrons are "hot" (several eV or 10-30,000 K) while the gas and ions are cool, creating "non-equilibrium" plasmas.

COLLISIONAL LOW TEMPERATURE PLASMAS



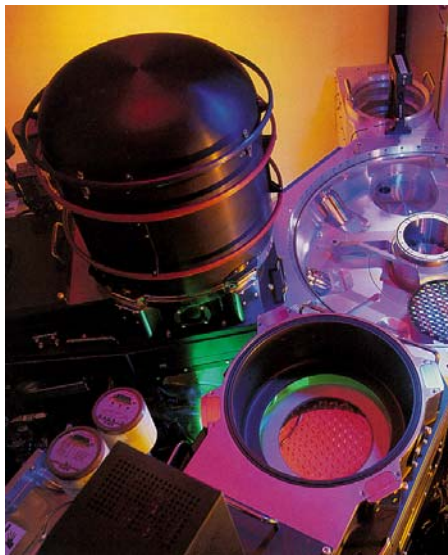
• **Lighting**



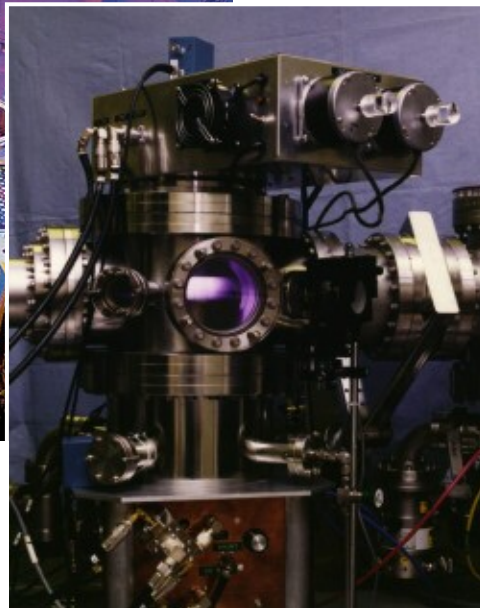
• **Thrusters**



• **Spray Coatings**



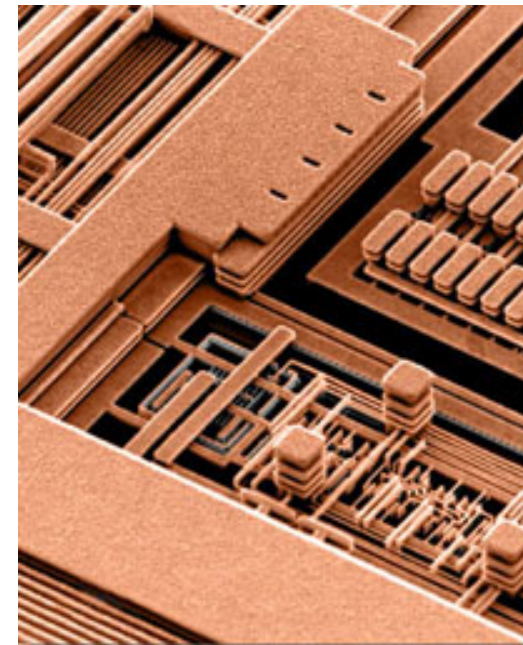
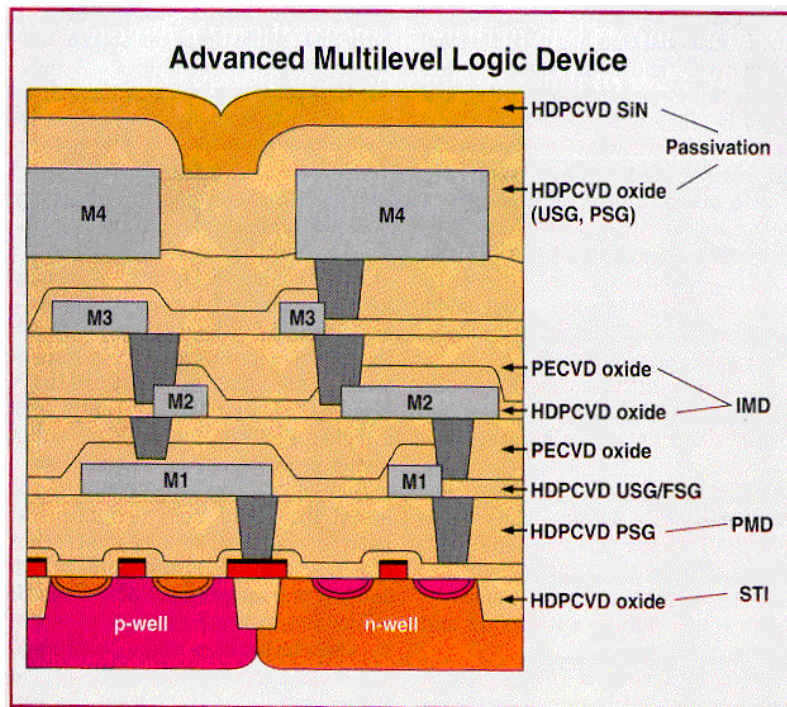
• **Materials Processing**



• **Displays**

PLASMAS IN MICROELECTRONICS FABRICATION

- The striking improvement in the functionality of microelectronics devices results from shrinking of individual components and increasing complexity of the circuitry

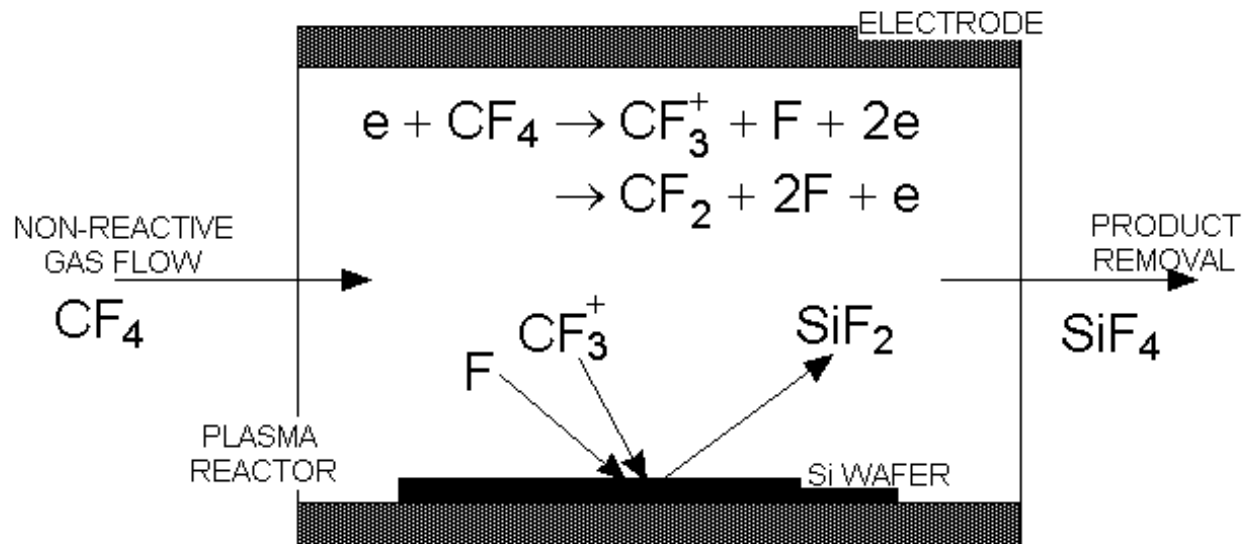


Ref: IBM Microelectronics

- Plasmas are absolutely essential to the fabrication of microelectronics.

PLASMAS IN MICROELECTRONICS FABRICATION

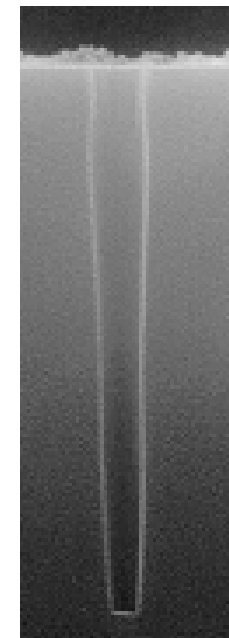
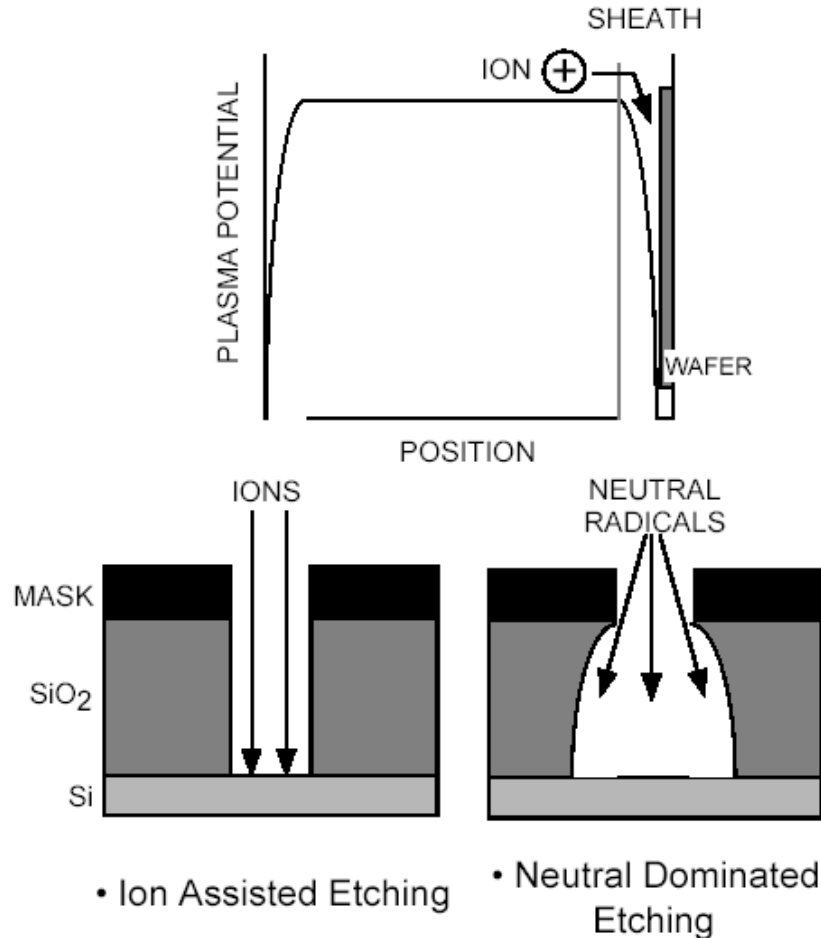
- Plasmas play a dual role in microelectronics fabrication.
- First, electron impact on otherwise unreactive gases produces neutral radicals and ions.



- These species then drift or diffuse to surfaces where they add, remove or modify materials.

PLASMAS IN MICROELECTRONICS FABRICATION

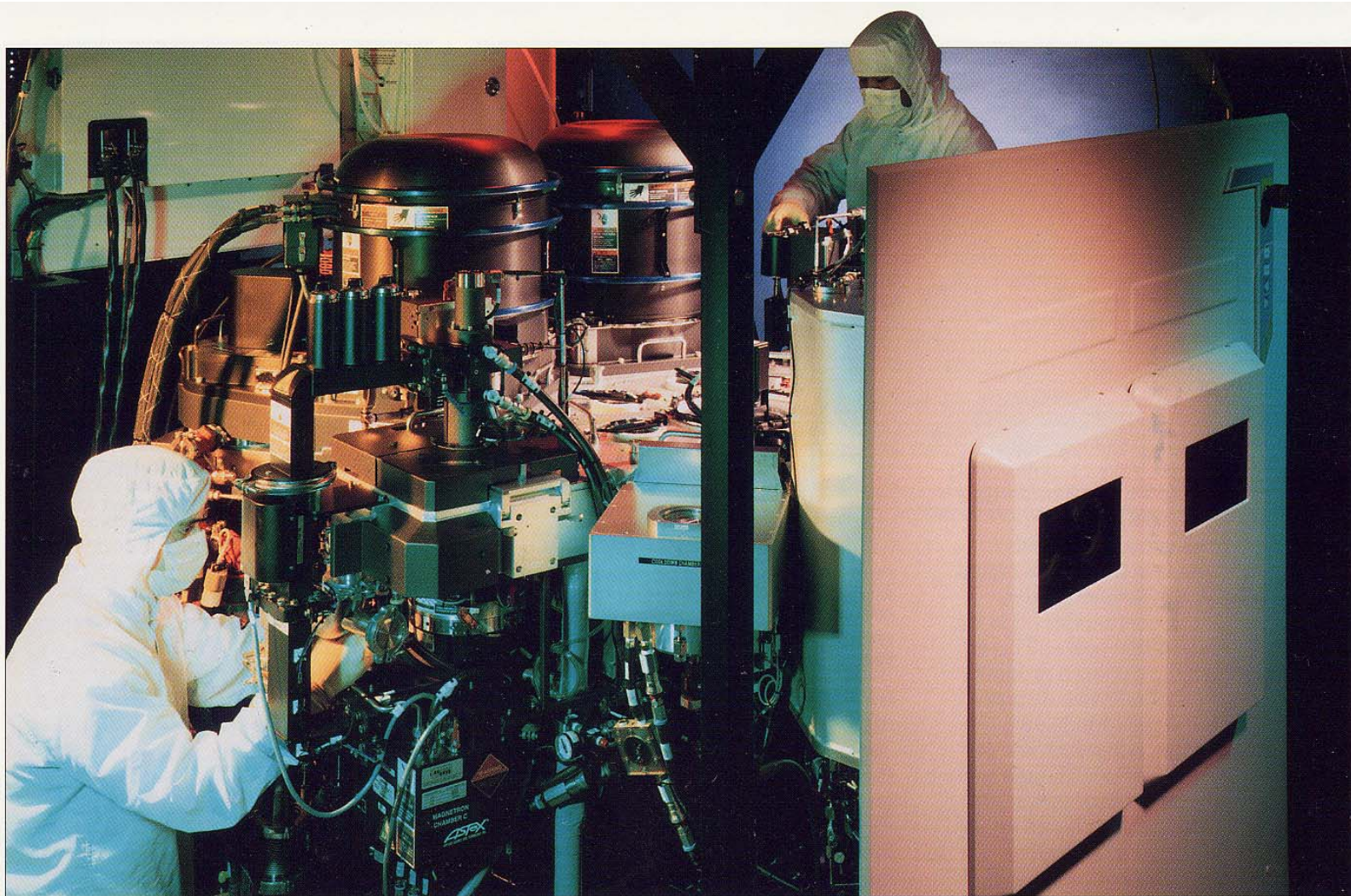
- Second, ions deliver directed activation energy to surfaces fabricating fine having extreme and reproducable tolerances.



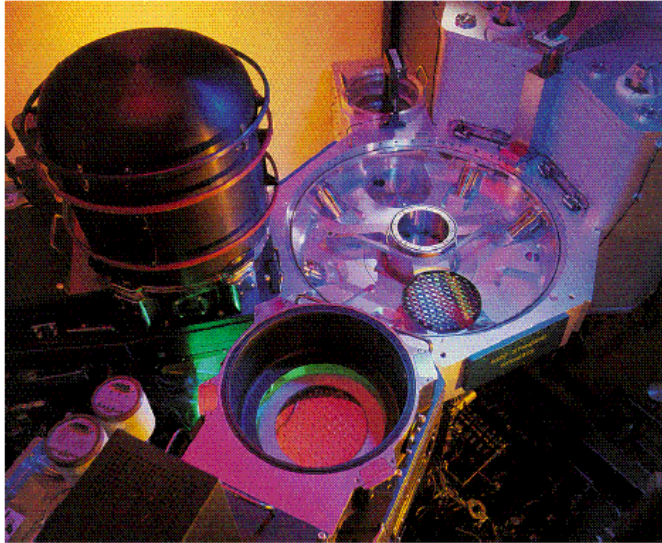
- 0.25 μm Feature (C. Cui, AMAT)

University of Illinois
Optical and Discharge Physics

APPLIED MATERIALS DECOUPLED PLASMA SOURCES (DPS)

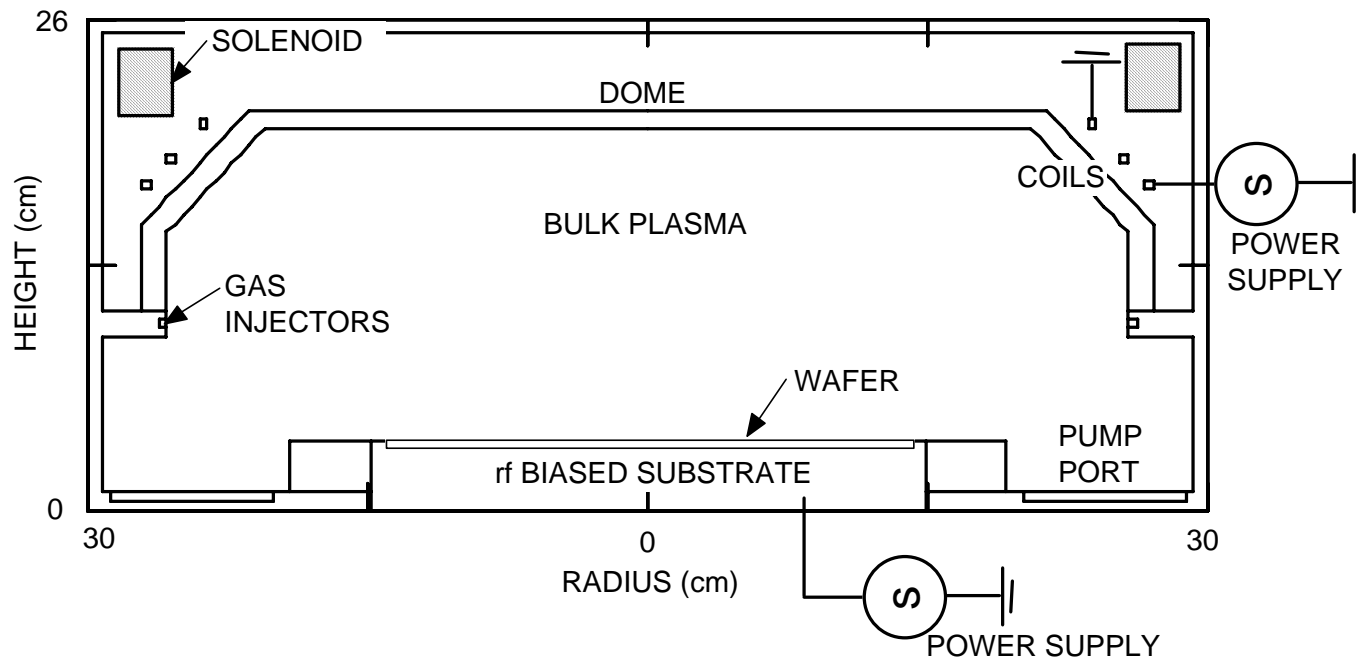


University of Illinois
Optical and Discharge Physics



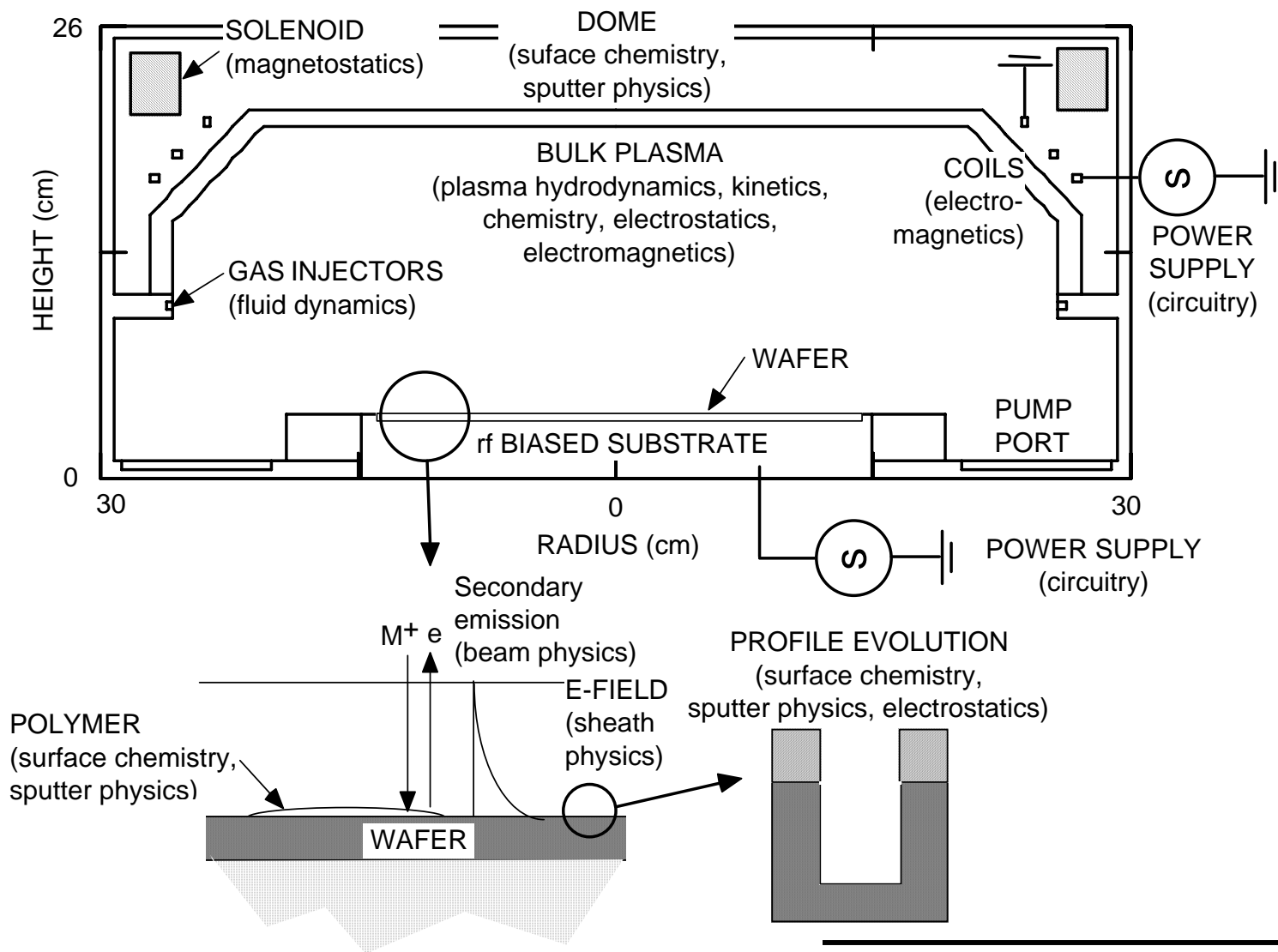
rf BIASED INDUCTIVELY COUPLED PLASMAS

- Inductively Coupled Plasmas (ICPs) with rf biasing are used here.
- $< 10\text{s mTorr}$, 10s MHz , $100\text{s W} - \text{kW}$, electron densities of $10^{11}\text{-}10^{12} \text{ cm}^{-3}$.



University of Illinois
Optical and Discharge Physics

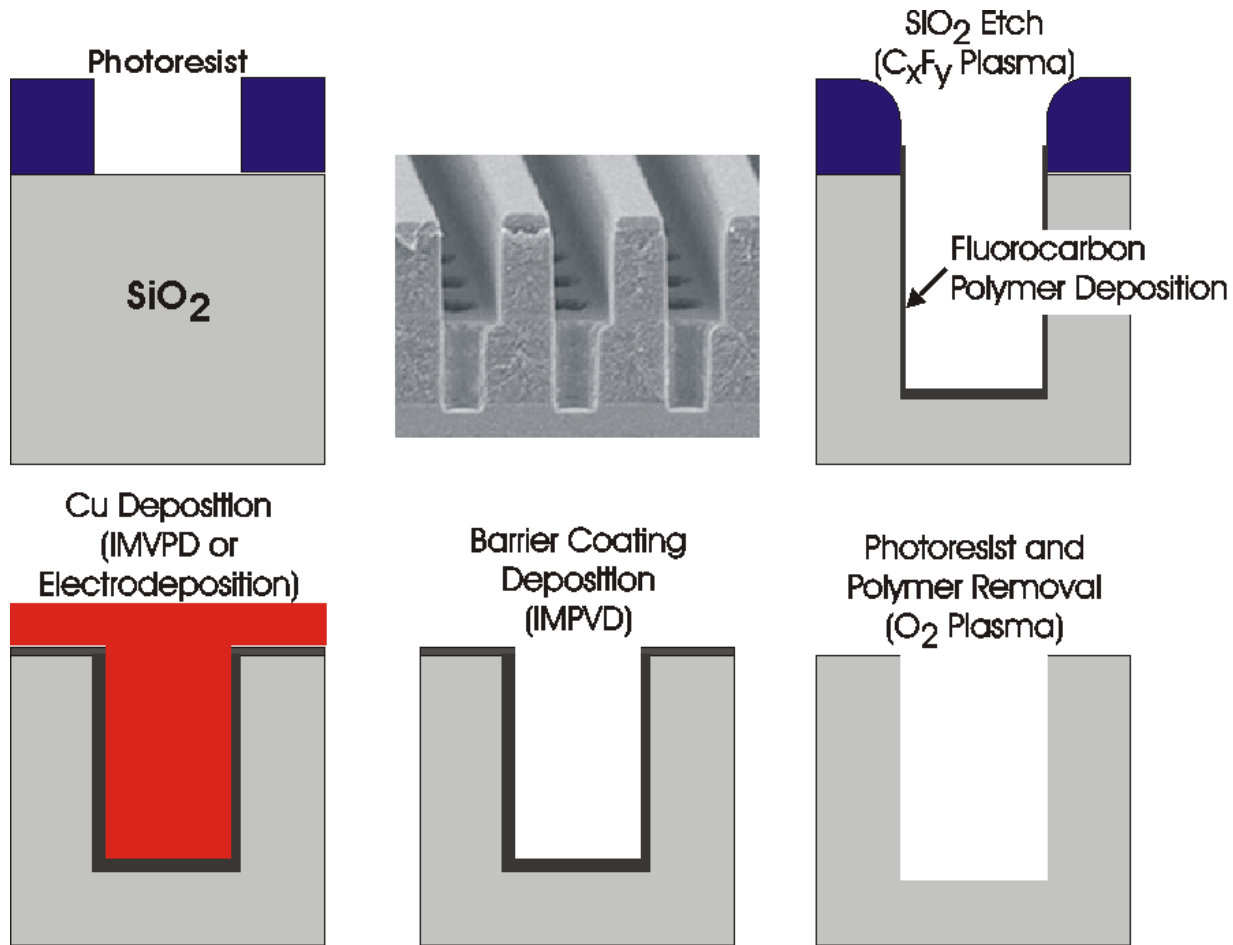
PHYSICAL PROCESSES IN REACTOR



University of Illinois
Optical and Discharge Physics

GOAL FOR PROCESS MODELING: INTEGRATION

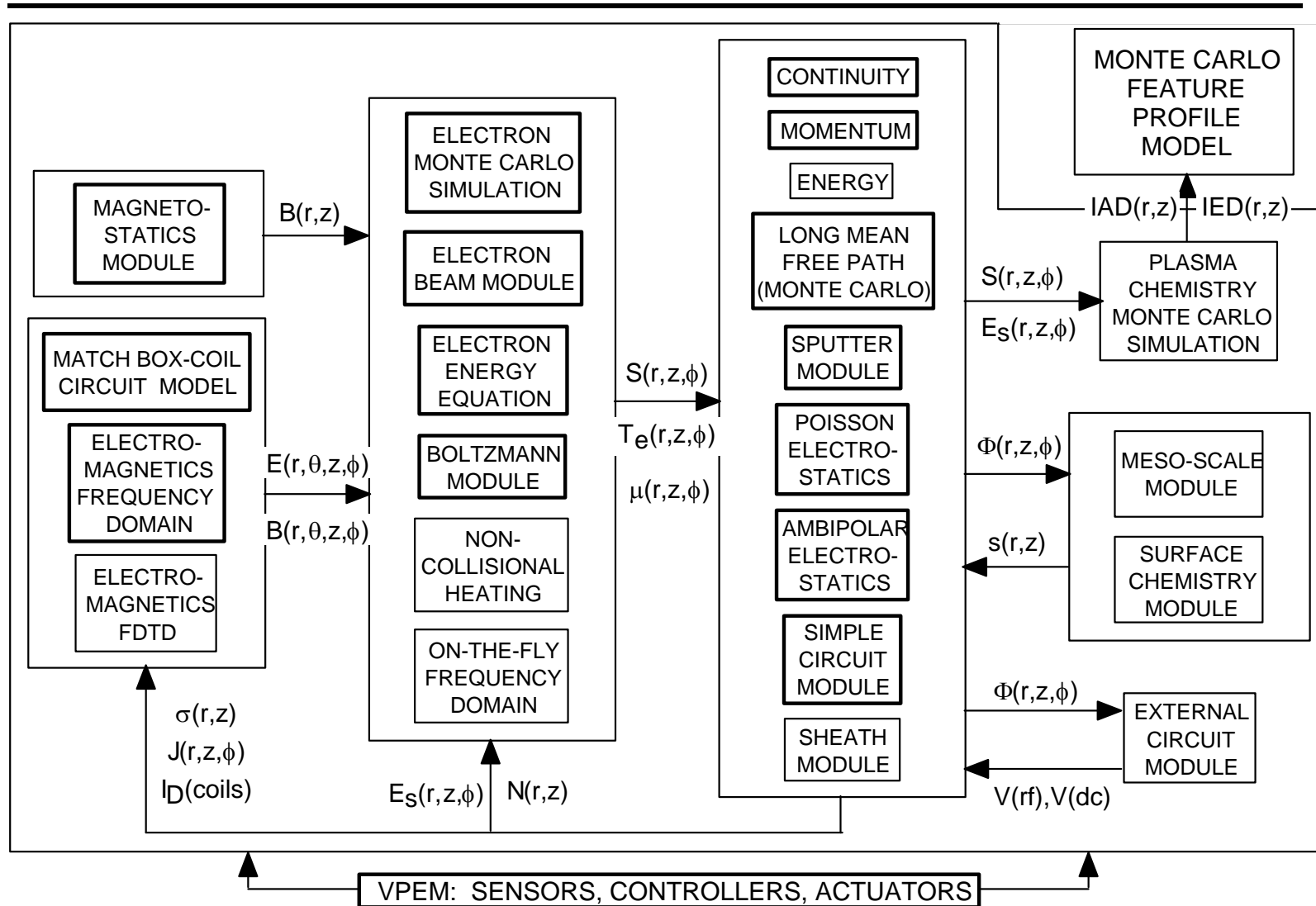
- Plasma processing involves an integrated sequence of steps, each of which depends on the quality of the previous steps.



GOAL FOR PROCESS MODELING: INTEGRATION

- To address these complexities, modeling platforms must integrate:
 - *Plasma Physics*
 - *Plasma Chemistry*
 - *Surface Kinetics*

HYBRID PLASMA EQUIPMENT MODEL



University of Illinois
Optical and Discharge Physics

ELECTROMAGNETICS MODEL

- The wave equation is solved in the frequency domain using sparse matrix techniques (2D,3D):

$$-\nabla \cdot \left(\frac{1}{\mu} \nabla \cdot \bar{\mathbf{E}} \right) + \nabla \cdot \left(\frac{1}{\mu} \nabla \bar{\mathbf{E}} \right) = \frac{\partial^2 (\epsilon \bar{\mathbf{E}})}{\partial t^2} + \frac{\partial (\bar{\boldsymbol{\sigma}} \cdot \bar{\mathbf{E}} + \bar{\mathbf{J}})}{\partial t}$$

$$\vec{\mathbf{E}}(\vec{r}, t) = \vec{\mathbf{E}}'(\vec{r}) \exp(-i(\omega t + \varphi(\vec{r})))$$

- Conductivities are tensor quantities (2D,3D):

$$\bar{\boldsymbol{\sigma}} = \sigma_o \frac{m v_m}{q \alpha} \frac{1}{\left(\alpha^2 + |\vec{\mathbf{B}}|^2 \right)} \begin{pmatrix} \alpha^2 + B_r^2 & \alpha B_z + B_r B_\theta & -\alpha B_\theta + B_r B_z \\ -\alpha B_z + B_r B_\theta & \alpha^2 + B_\theta^2 & \alpha B_r + B_\theta B_z \\ -\alpha B_\theta + B_r B_z & -\alpha B_r + B_\theta B_z & \alpha^2 + B_z^2 \end{pmatrix}$$

$$\vec{\mathbf{j}} = \bar{\boldsymbol{\sigma}} \cdot \vec{\mathbf{E}} \quad \alpha = \frac{(i\omega + v_m)}{q/m}, \quad \sigma_o = \frac{q^2 n_e}{m v_m}$$

ELECTROMAGNETICS MODEL (cont.)

- The electrostatic term in the wave equation is addressed using a perturbation to the electron density (2D).

$$\nabla \cdot \bar{E} = \frac{\rho}{\epsilon} = \frac{q\Delta n_e}{\epsilon}, \quad \Delta n_e = -\nabla \cdot \left(\frac{\bar{\sigma} \cdot \bar{E}}{q} \right) / \left(\frac{1}{\tau} + i\omega \right)$$

- Conduction currents can be kinetically derived from the Electron Monte Carlo Simulation to account for non-collisional effects (2D).

$$J_e(\vec{r}, t) = J_o(\vec{r}) \exp(i(\omega t + \phi_v(\vec{r}))) = -qn_e(\vec{r})\vec{v}_e(\vec{r}) \exp(i(\omega t + \phi_v(\vec{r})))$$

ELECTRON ENERGY TRANSPORT

- **Continuum (2D,3D):**

$$\partial \left(\frac{3}{2} n_e k T_e \right) / \partial t = S(T_e) - L(T_e) - \nabla \cdot \left(\frac{5}{2} \Phi k T_e - \bar{\kappa}(T_e) \cdot \nabla T_e \right) + S_{EB}$$

where $S(T_e)$	=	Power deposition from electric fields
$L(T_e)$	=	Electron power loss due to collisions
Φ	=	Electron flux
$\kappa(T_e)$	=	Electron thermal conductivity tensor
S_{EB}	=	Power source source from beam electrons

- Power deposition has contributions from wave and electrostatic heating.
- **Kinetic (2D,3D):** A Monte Carlo Simulation is used to derive $f(\varepsilon, \vec{r}, t)$ including electron-electron collisions using electromagnetic fields from the EMM and electrostatic fields from the FKM.

PLASMA CHEMISTRY, TRANSPORT AND ELECTROSTATICS

- Continuity, momentum and energy equations are solved for each species (with jump conditions at boundaries) (2D,3D).

$$\frac{\partial N_i}{\partial t} = -\nabla \cdot (N_i \vec{v}_i) + S_i$$

$$\frac{\partial (N_i \vec{v}_i)}{\partial t} = \frac{1}{m_i} \nabla (k N_i T_i) - \nabla \cdot (N_i \vec{v}_i \vec{v}_i) + \frac{q_i N_i}{m_i} (\vec{E} + \vec{v}_i \times \vec{B}) - \nabla \cdot \bar{\mu}_i - \sum_j \frac{m_j}{m_i + m_j} N_i N_j (\vec{v}_i - \vec{v}_j) \nu_{ij}$$

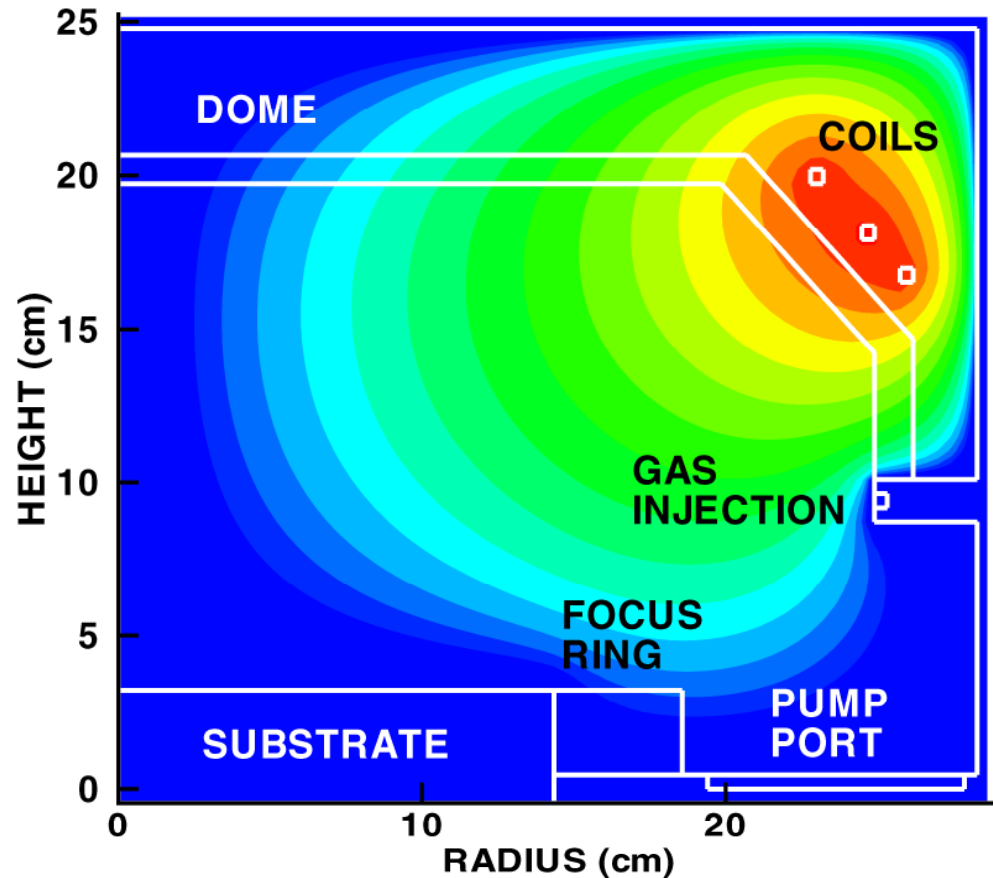
$$\begin{aligned} \frac{\partial (N_i \varepsilon_i)}{\partial t} + \nabla \cdot \mathbf{Q}_i + P_i \nabla \cdot \mathbf{U}_i + \nabla \cdot (N_i \mathbf{U}_i \varepsilon_i) &= \frac{N_i q_i^2 \nu_i}{m_i (\nu_i^2 + \omega^2)} E^2 \\ &+ \frac{N_i q_i^2}{m_i \nu_i} E_s^2 + \sum_j 3 \frac{m_{ij}}{m_i + m_j} N_i N_j R_{ij} k_B (T_j - T_i) \pm \sum_j 3 N_i N_j R_{ij} k_B T_j \end{aligned}$$

- Implicit solution of Poisson's equation (2D,3D):

$$\nabla \cdot \varepsilon \nabla \Phi(t + \Delta t) = - \left(\rho_s + \sum_i q_i N_i - \Delta t \cdot \sum_i (q_i \nabla \cdot \vec{\phi}_i) \right)$$

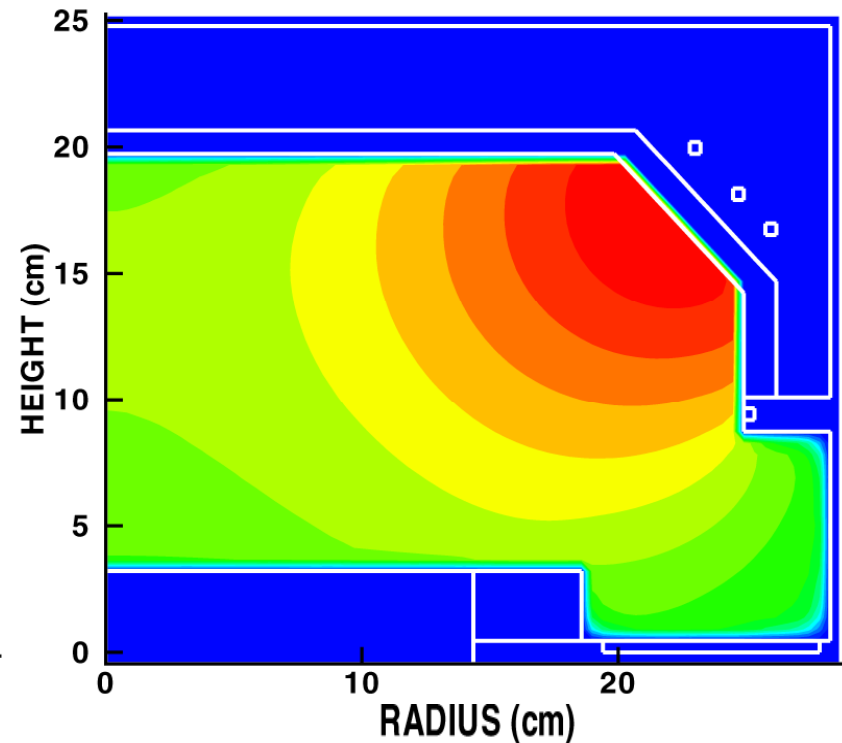
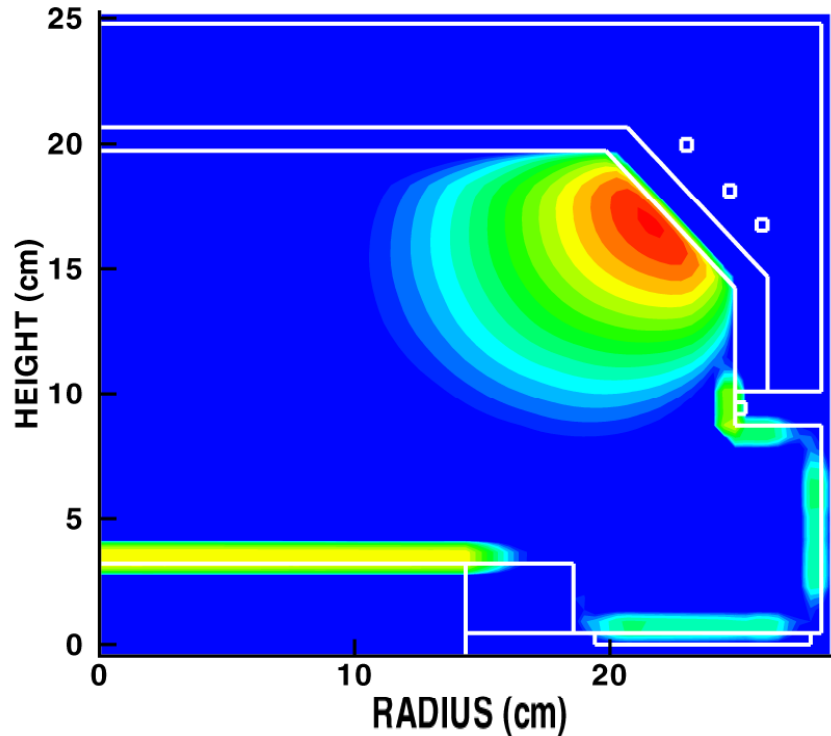
WALK THROUGH: Ar/Cl₂ PLASMA FOR p-Si ETCHING

- The inductively coupled electromagnetic fields have a skin depth of 3-4 cm.
- Absorption of the fields produces power deposition in the plasma.
- Electric Field (max = 6.3 V/cm)
 - Ar/Cl₂ = 80/20
 - 20 mTorr
 - 1000 W ICP 2 MHz
 - 250 V bias, 2 MHz (260 W)



Ar/Cl₂ ICP: POWER AND ELECTRON TEMPERATURE

- ICP Power heats electrons, capacitively coupled power dominantly accelerates ions.

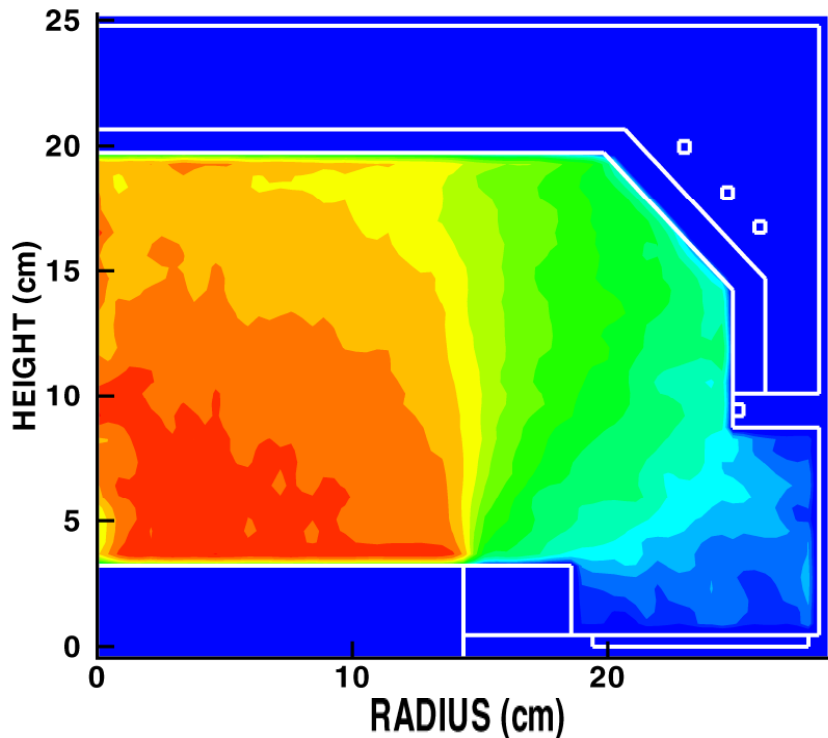


- Power Deposition (max = 0.9 W/cm³)
- Electron Temperature (max = 5 eV)

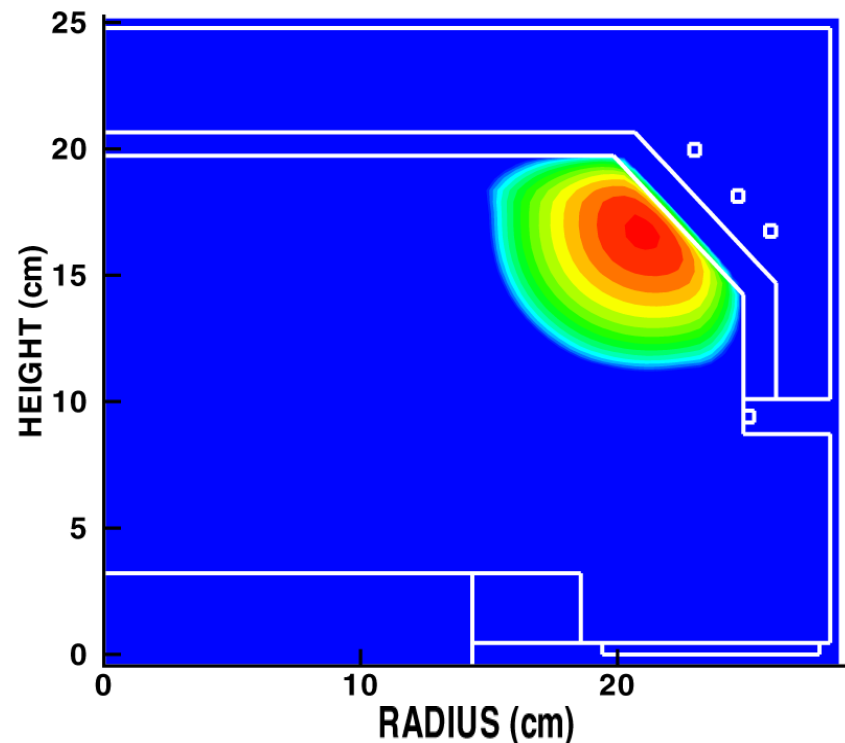
- Ar/Cl₂ = 80/20, 20 mTorr, 1000 W ICP 2 MHz, 250 V bias, 2 MHz (260 W)

Ar/Cl₂ ICP: IONIZATION

- Ionization is produced by bulk electrons and sheath accelerated secondary electrons.



- Beam Ionization
(max = $1.3 \times 10^{14} \text{ cm}^{-3}\text{s}^{-1}$)

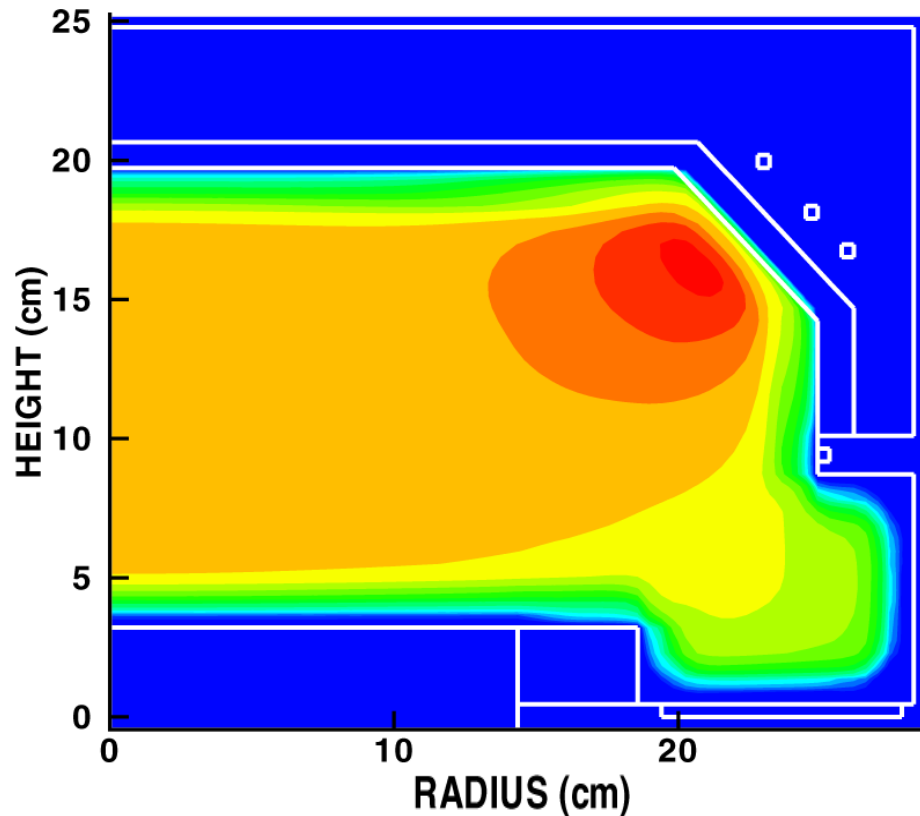


- Bulk Ionization
(max = $5.4 \times 10^{15} \text{ cm}^{-3}\text{s}^{-1}$)

- Ar/Cl₂ = 80/20, 20 mTorr, 1000 W ICP 2 MHz,
250 V bias, 2 MHz (260 W)

Ar/Cl₂ ICP: POSITIVE ION DENSITY

- Diffusion from the remote plasma source produces uniform ion densities at the substrate.

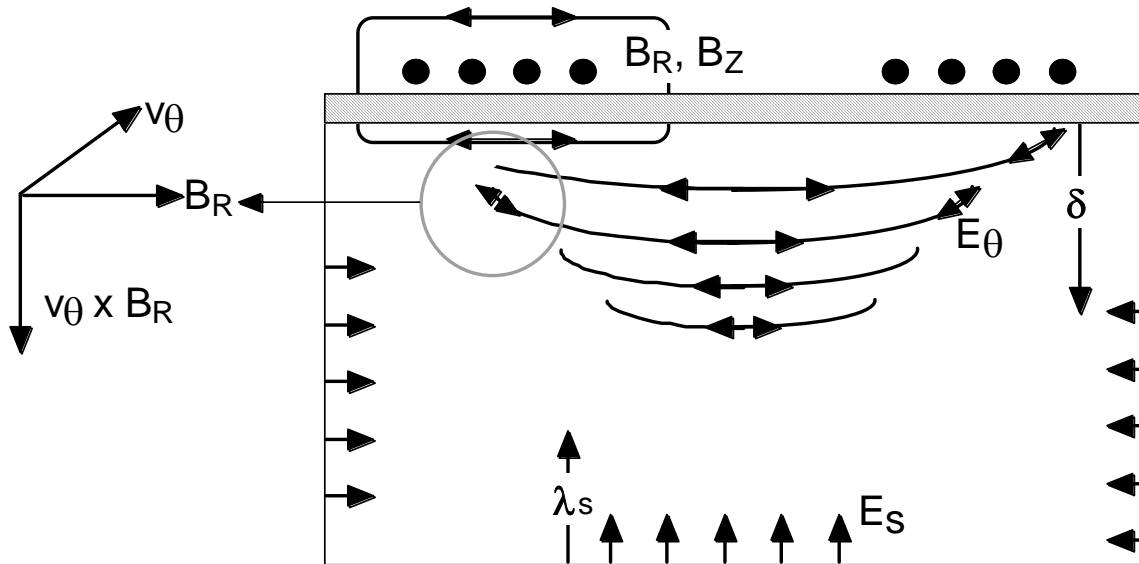


- Positive Ion Density
(max = $1.8 \times 10^{11} \text{ cm}^{-3}$)

- Ar/Cl₂ = 80/20, 20 mTorr, 1000 W ICP 2 MHz, 250 V bias, 2 MHz (260 W)

• *PLASMA PHYSICS*
(Are we getting it right?)

FORCES ON ELECTRONS IN ICPs



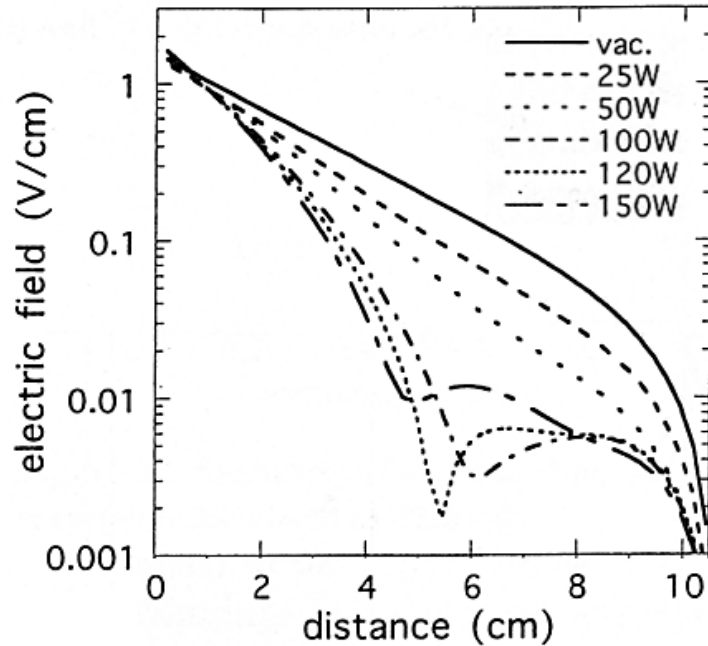
- Inductive electric field provides azimuthal acceleration; penetrates

$$\delta = \left(m_e / (e^2 \mu_o n_e) \right)^{1/2} \quad (1-3 \text{ cm})$$

- Electrostatic (capacitive); penetrates $\lambda_s \approx 10 \lambda_D$, $\lambda_D = \left(kT_e / (8\pi n_e e^2) \right)^{1/2}$
(100s μm to mm)

- Non-linear Lorentz Force $\vec{F} = v_\theta \times \vec{B}_{rf}$

ANAMOLOUS SKIN EFFECT AND POWER DEPOSITION



- Collisional heating:

$$\lambda_{mfp} < \delta_{skin}, \quad \vec{J}_e(\vec{r}, t) = \sigma(\vec{r}, t) \vec{E}(\vec{r}, t)$$

- Anomalous skin effect:

$$\lambda_{mfp} > \delta_{skin}$$

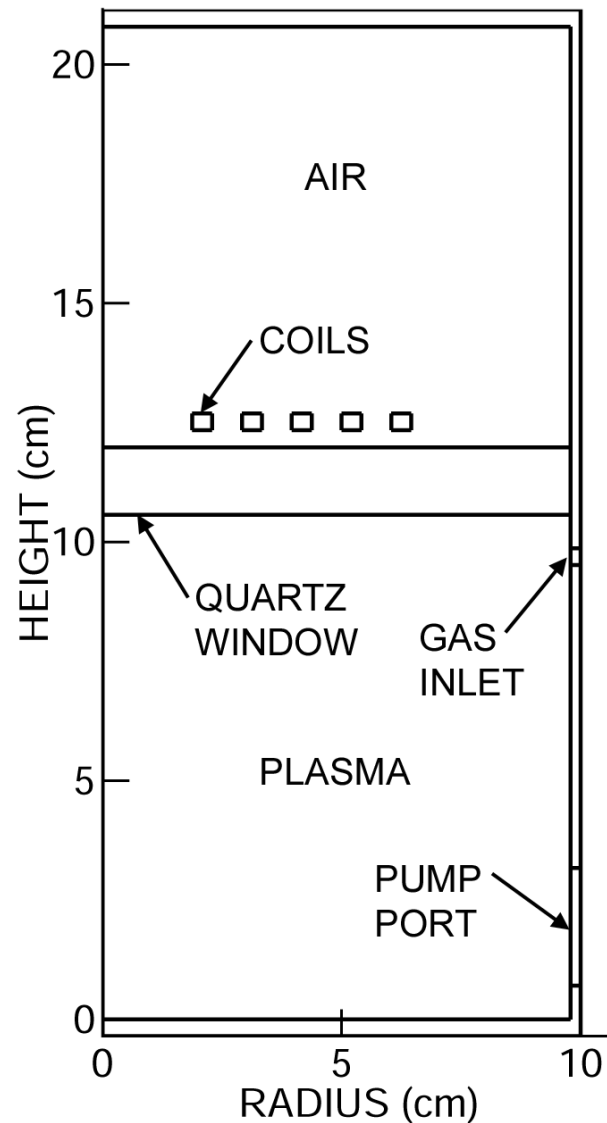
$$\vec{J}_e(\vec{r}, t) = \iint \sigma(\vec{r}, \vec{r}', t, t') \vec{E}(\vec{r}', t') d\vec{r}' dt'$$

$$\vec{F} = \vec{v} \times \vec{B}$$

- Ref: V. Godyak, “Electron Kinetics of Glow Discharges”

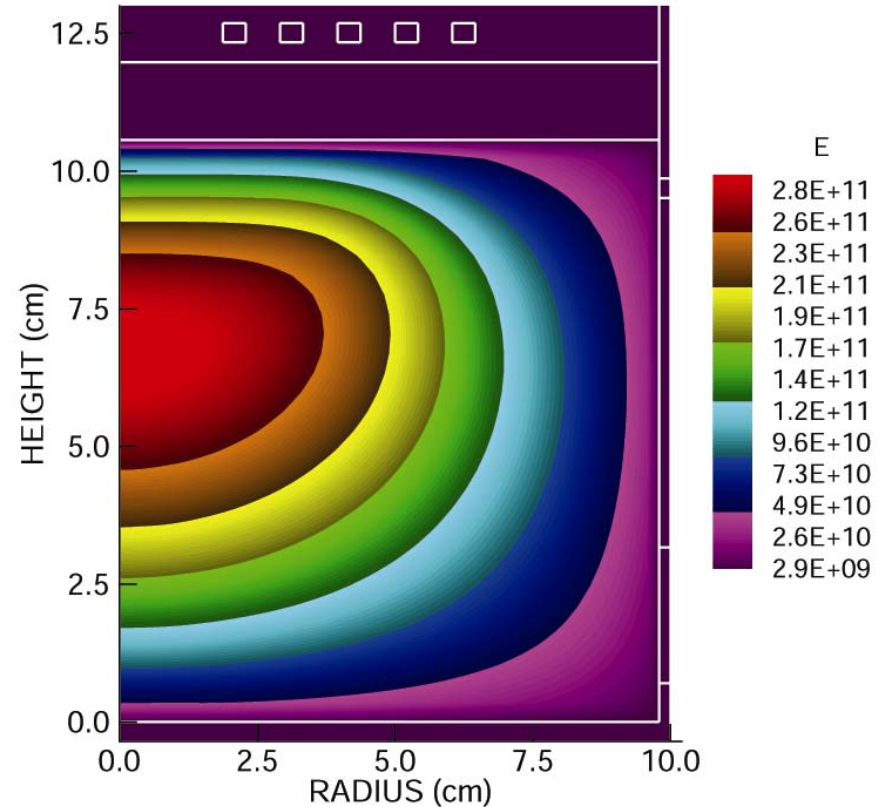
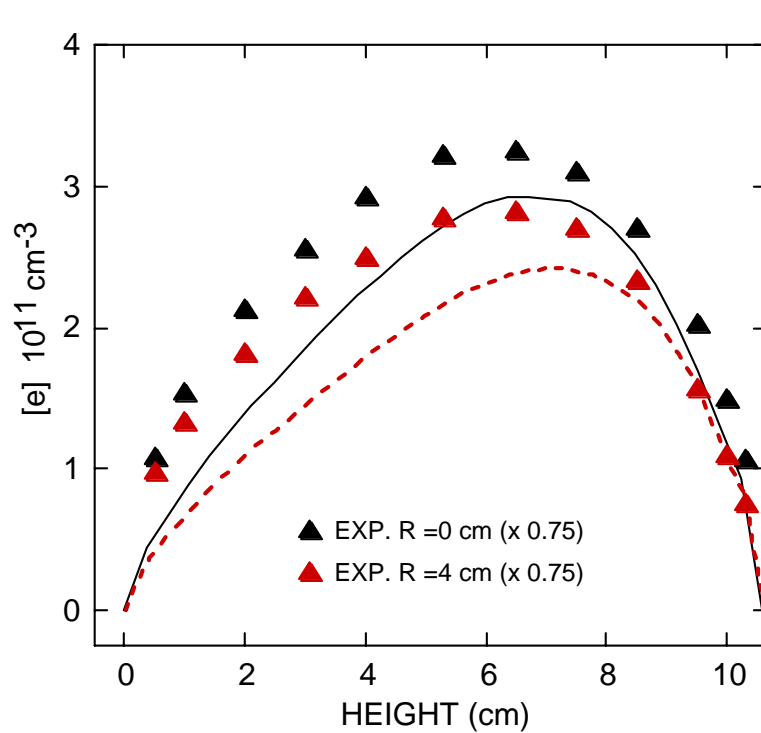
- Electrons receive (positive) and deliver (negative) power from/to the E-field.
- E-field is non-monotonic.

ICP CELL FOR VALIDATION



- Experiments by Godyak et al are used for validation.
- The experimental cell is an ICP reactor with a Faraday shield to minimize capacitive coupling.
- V. Godyak et al, J. Appl. Phys. 85, 703 (1999)

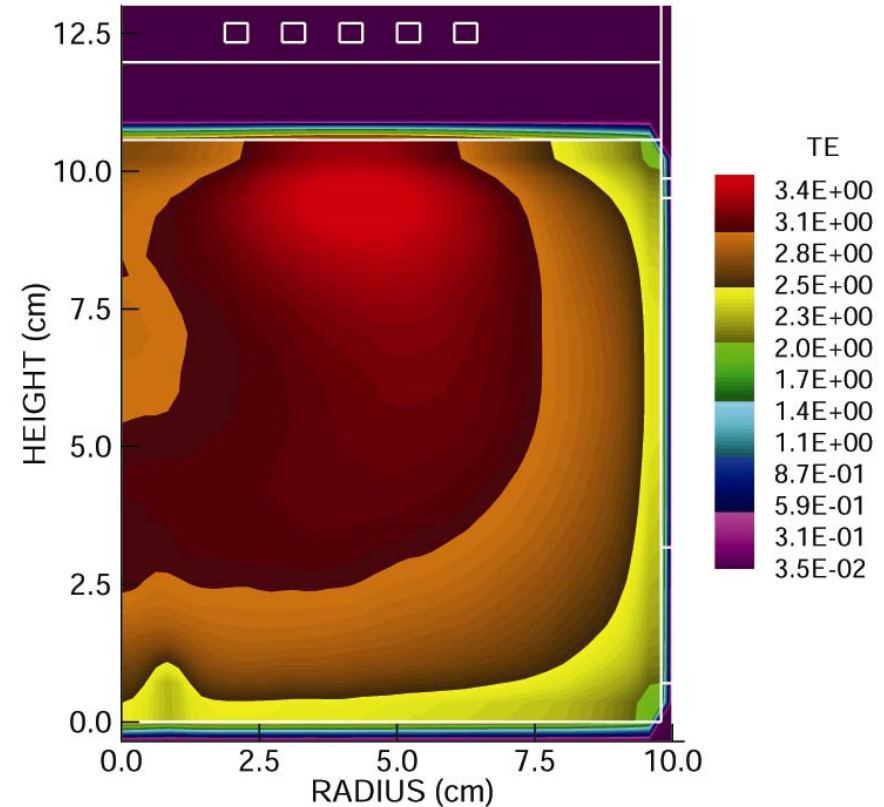
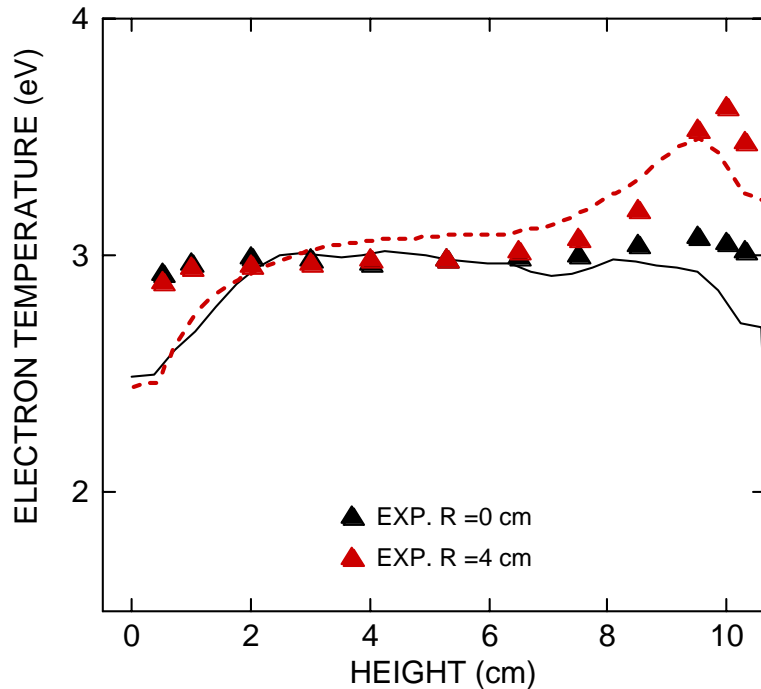
ELECTRON DENSITY: Ar, 10 mTorr, 200 W, 7 MHz



• Ref: V. Godyak, Private Comm.

- On axis peak in $[e]$ occurs inspite of off-axis power deposition.
- Model is about 30% below experiments. This likely has to do with details of the sheath model.

ELECTRON TEMPERATURE: Ar, 10 mTorr, 200 W, 7 MHz



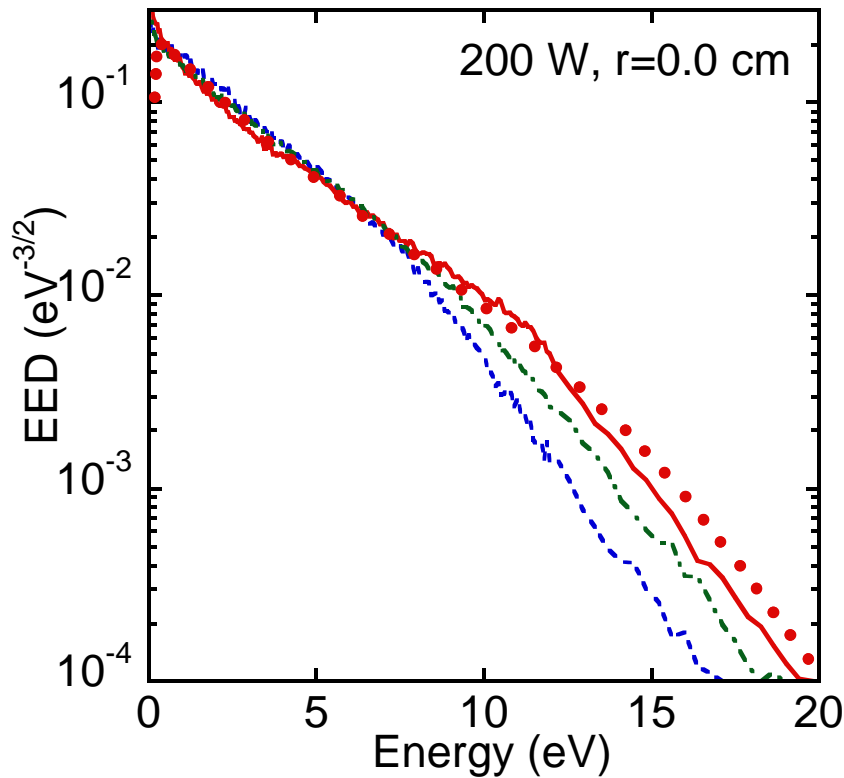
- Ref: V. Godyak, Private Comm.

- The high thermal conductivity and redistribution of energy by e-e collisions produces nearly uniform temperatures.
- T_e peaks under the coils where power deposition is largest.

University of Illinois
Optical and Discharge Physics

EEDs ALONG THE CENTERLINE OF THE REACTOR

- Godyak (1998), $z=5.0$ cm
- Model, $z=0.5$ cm
- Model, $z=5.0$ cm
- Model, $z=10.0$ cm



- The electron energy distributions show a bi-Maxwellian form, which is typical for low-pressure inductively coupled plasmas.

- Ar, 10 mTorr, 6.78 MHz, 200 W

COLLISIONLESS TRANSPORT ELECTRIC FIELDS

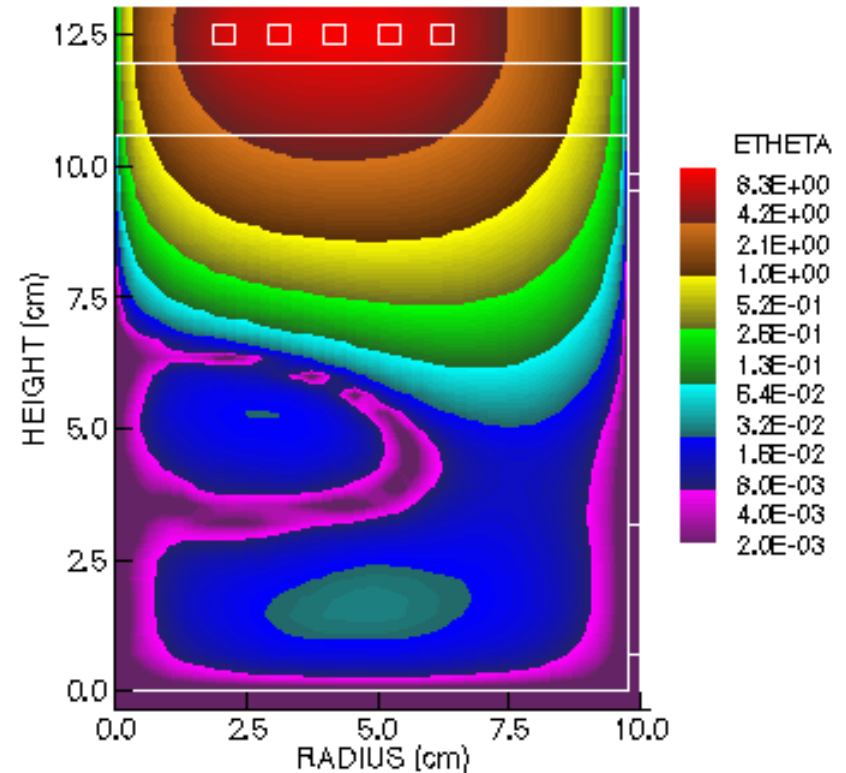
- We couple electron transport to Maxwell's equations by kinetically deriving electron current.

$$\oint \vec{j}(\vec{r}) \exp(i\omega(t-t_o)) \cdot dA = \sum_k q_k \vec{v}_k(\vec{r}) \exp(i\omega(t_k - t_o))$$

- E_θ during the rf cycle exhibits extrema and nodes resulting from this non-collisional transport.

- “Sheets” of electrons provide current sources interfering or reinforcing E_θ for the next sheet.

- Axial transport results from $\vec{v} \times \vec{B}_{rf}$ forces.

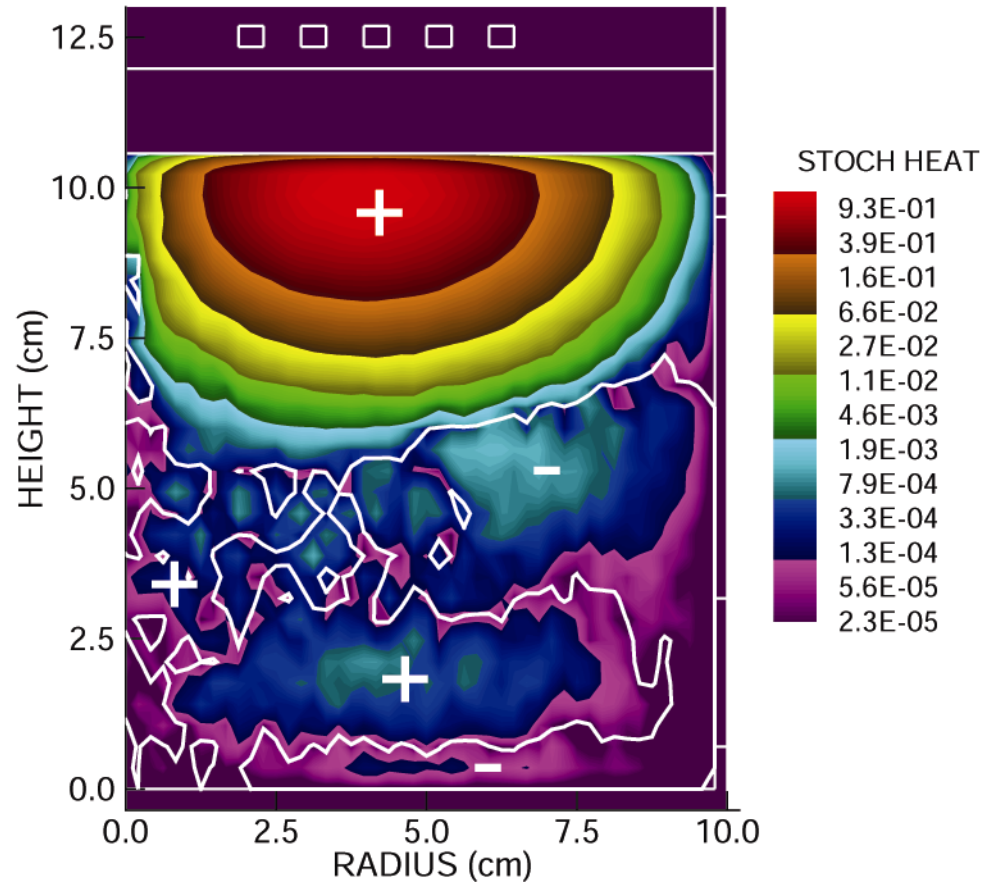
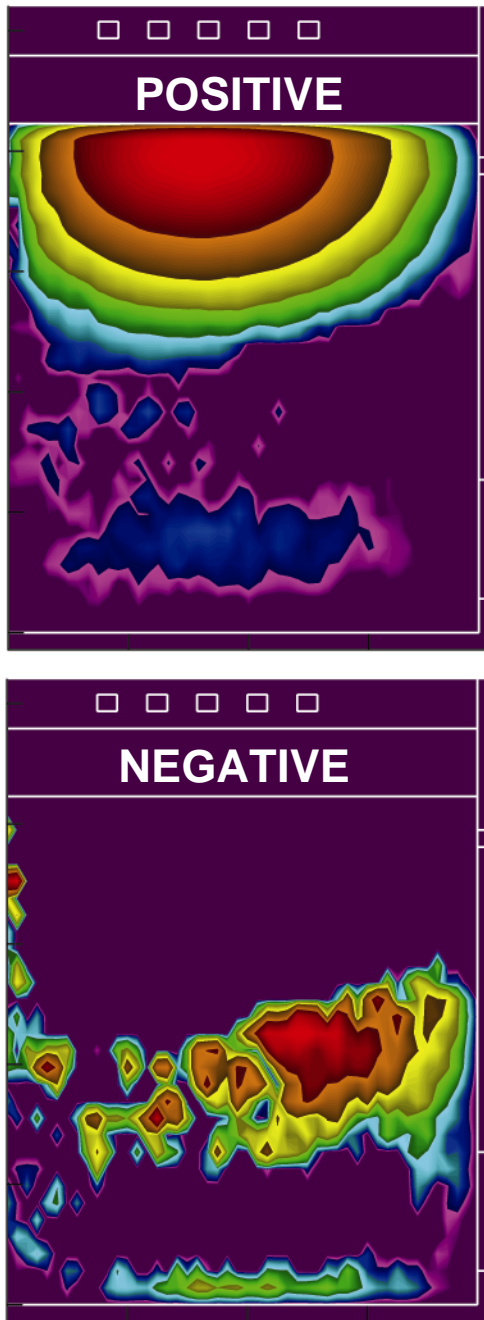


ANIMATION SLIDE

- Ar, 10 mTorr, 7 MHz, 100 W

POWER DEPOSITION: POSITIVE AND NEGATIVE

- The end result is regions of positive and negative power deposition.

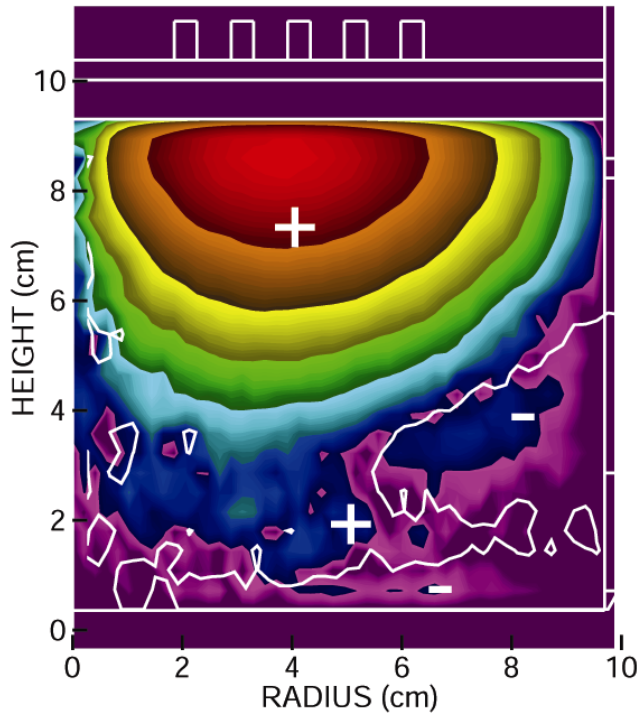


- Ar, 10 mTorr,
7 MHz, 100 W

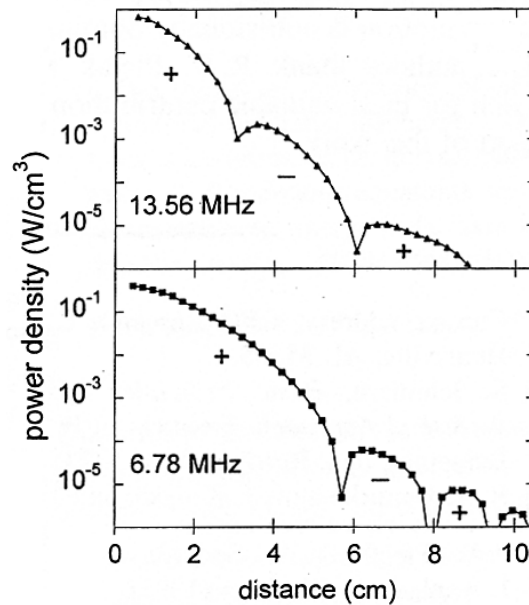
University of Illinois
Optical and Discharge Physics

POWER DEPOSITION vs FREQUENCY

- The shorter skin depth at high frequency produces more layers of negative power deposition of larger magnitude.

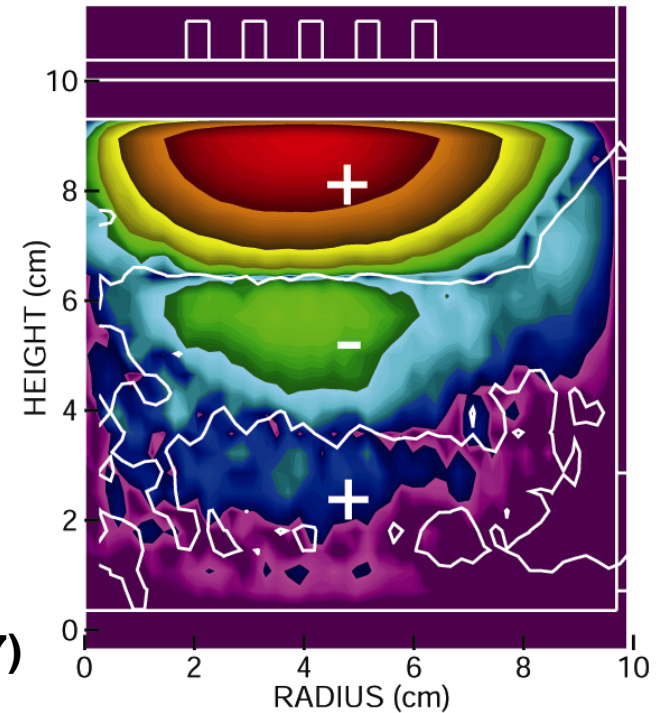


• **6.7 MHz**
 ($5 \times 10^{-5} - 1.4 \text{ W/cm}^3$)



• Ref: Godyak, PRL (1997)

• Ar, 10 mTorr, 200 W

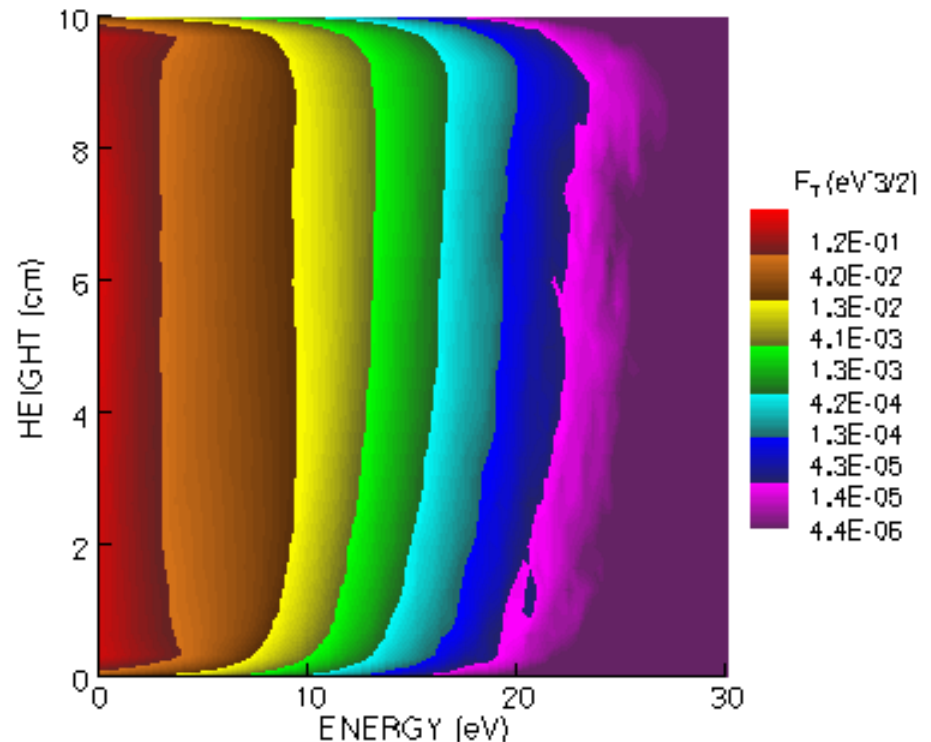


• **13.4 MHz**
 ($8 \times 10^{-5} - 2.2 \text{ W/cm}^3$)



TIME DEPENDENCE OF THE EED

- Time variation of the EED is mostly at higher energies where electrons are more collisional.
- Dynamics are dominantly in the electromagnetic skin depth where both collisional and non-linear Lorentz Forces) peak.
- The second harmonic dominates these dynamics.



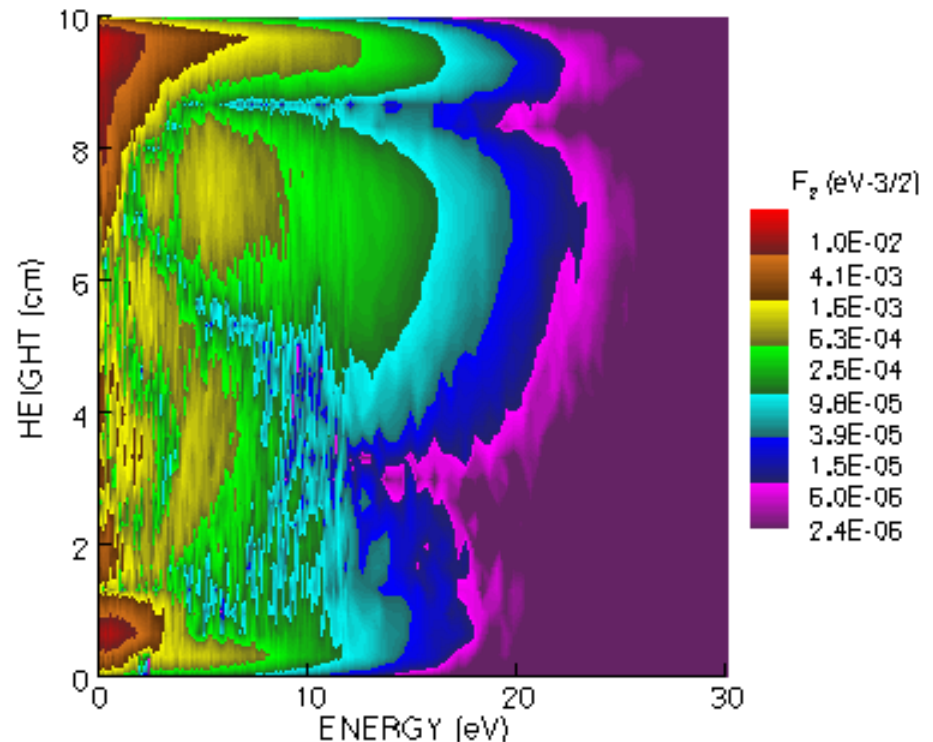
- Ar, 10 mTorr, 100 W, 7 MHz, $r = 4$ cm

ANIMATION SLIDE

University of Illinois
Optical and Discharge Physics

TIME DEPENDENCE OF THE EED: 2nd HARMONIC

- Electrons in skin depth quickly increase in energy and are “launched” into the bulk plasma.
- Undergoing collisions while traversing the reactor, they degrade in energy.
- Those surviving “climb” the opposite sheath, exchanging kinetic for potential energy.
- Several “pulses” are in transit simultaneously.



- Amplitude of 2nd Harmonic

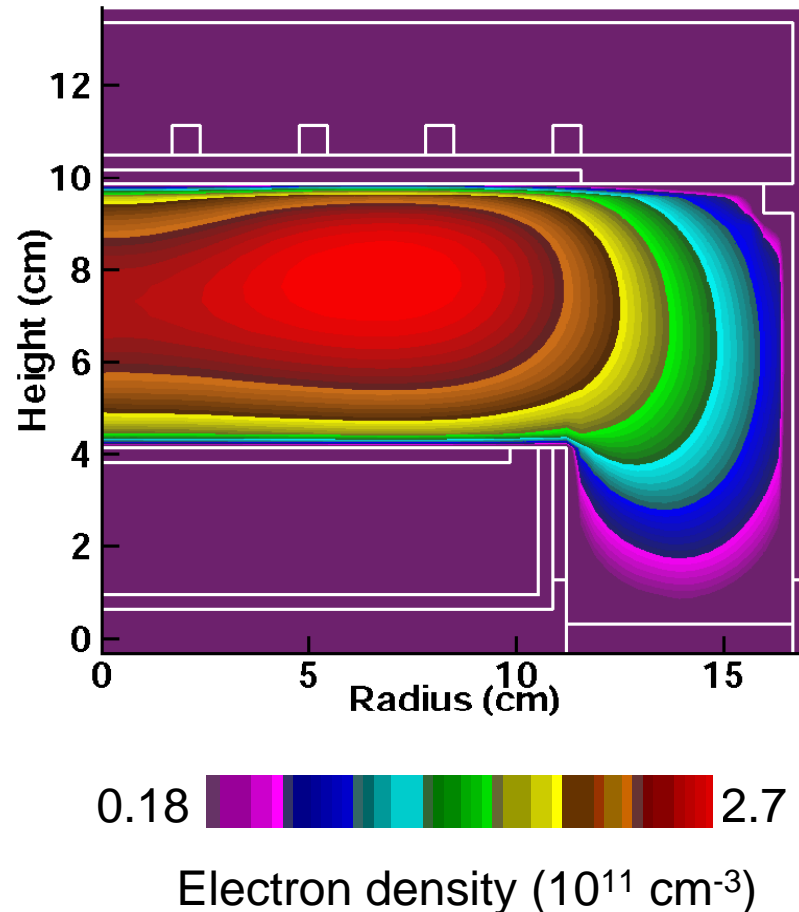
ANIMATION SLIDE

- Ar, 10 mTorr, 100 W, 7 MHz, $r = 4$ cm

University of Illinois
Optical and Discharge Physics

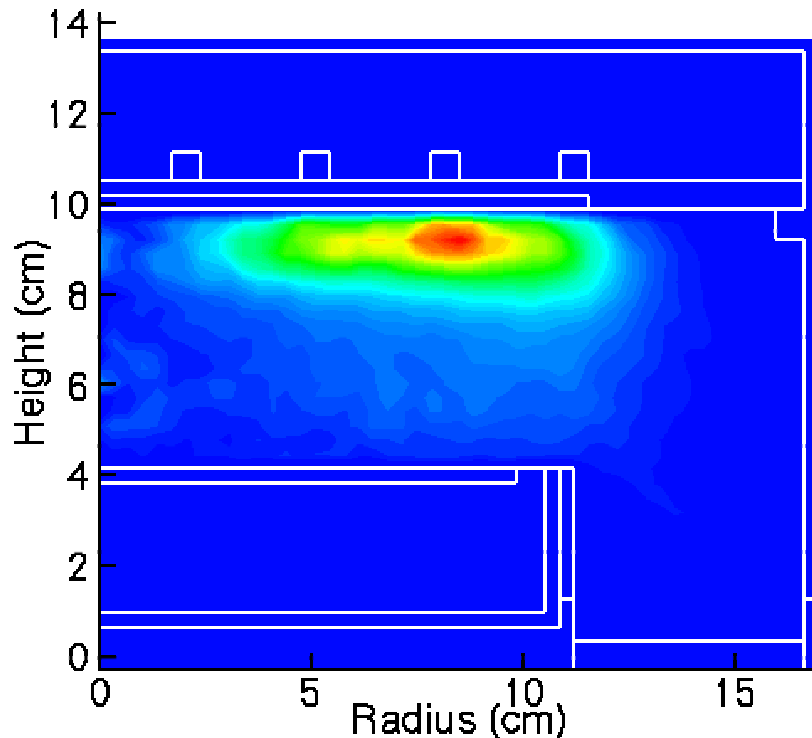
CONSEQUENCES OF ELECTRON DYNAMICS IN ICPs

- The consequences of electron dynamics were investigated for Ar/N₂ gas mixtures.
- $e^- + \text{Ar} \rightarrow \text{Ar}^+ + e^- + e^-$, $\Delta\varepsilon = 16 \text{ eV}$
High threshold reactions capture modulation in the tail of the EED.
- $e^- + \text{N}_2 \rightarrow \text{N}_2(\text{vib}) + e^-$, $\Delta\varepsilon = 0.29 \text{ eV}$
Low threshold reactions capture modulation of the bulk of the EED.
- Base case conditions:
 - Pressure: 5 mTorr
 - Frequency: 13.56 MHz
 - Ar / N₂: 90 / 10
 - Power : 650 W

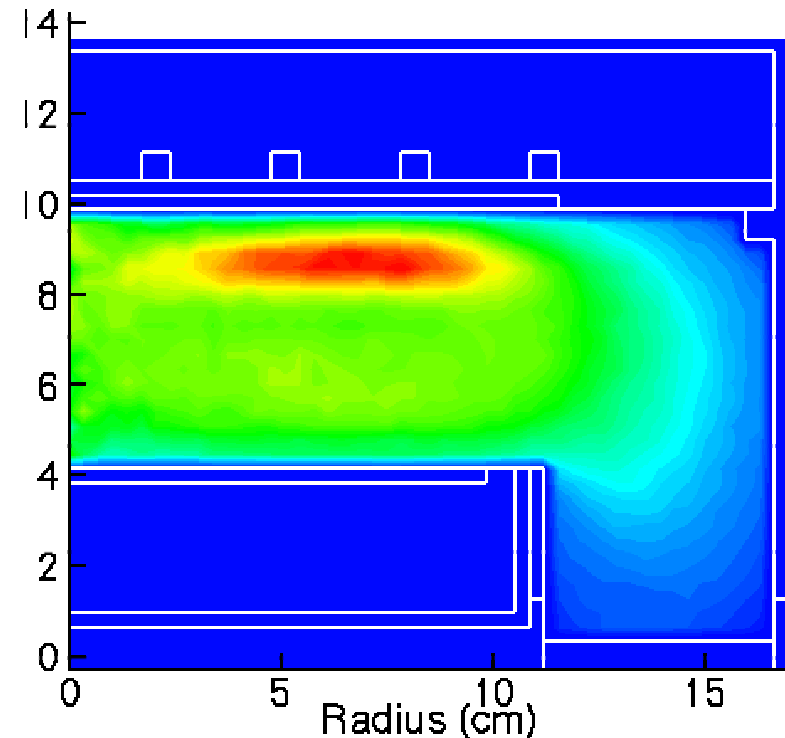


SOURCES FUNCTION vs TIME: THRESHOLD

- Ionization of Ar has more modulation than vibrational excitation of N₂ due to modulation of the tail of the EED.



• Ionization of Ar
 $6 \times 10^{14} - 3 \times 10^{16} \text{ cm}^{-3}\text{s}^{-1}$



• Excitation of N₂(v)
 $1.4 \times 10^{14} - 8 \times 10^{15} \text{ cm}^{-3}\text{s}^{-1}$

ANIMATION SLIDE

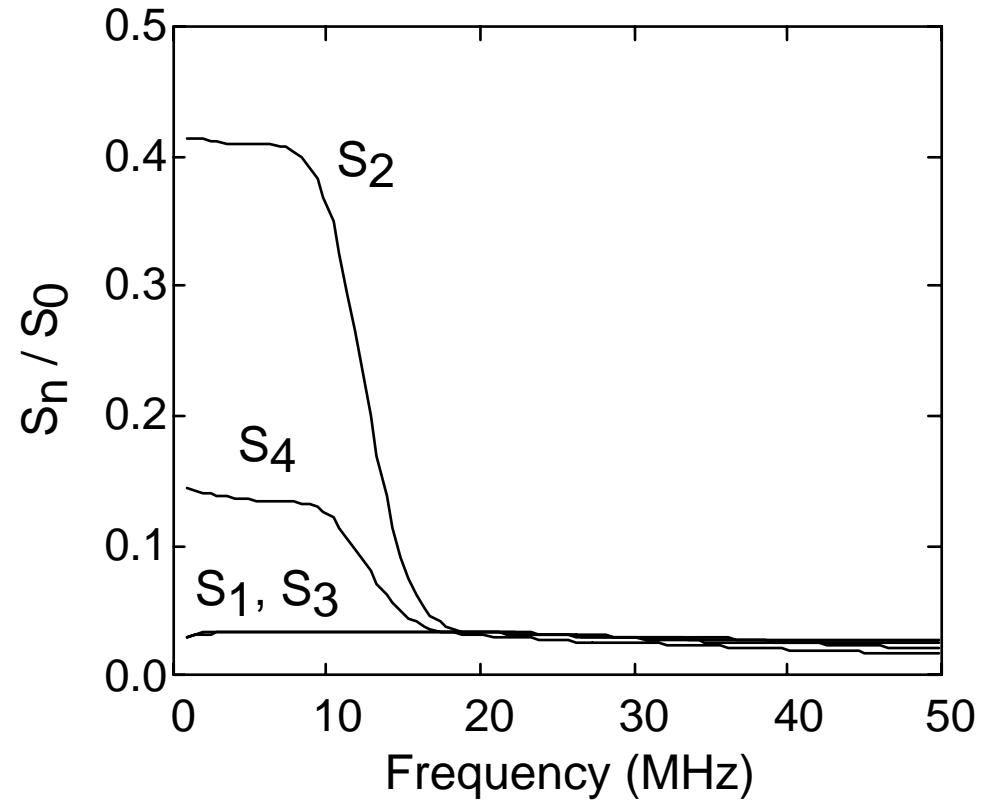
SNLA_0102_26

MIN  MAX

University of Illinois
Optical and Discharge Physics

HARMONICS OF Ar IONIZATION: FREQUENCY

- At large ω , non-linear Lorentz forces are small, and so harmonic content is also small.
- At small ω , both non-linear Lorentz forces and harmonic excitation by the electric field are large.



- Ar/N₂=90/10, 5 mTorr

- Harmonic Amplitude/Time Average

University of Illinois
Optical and Discharge Physics

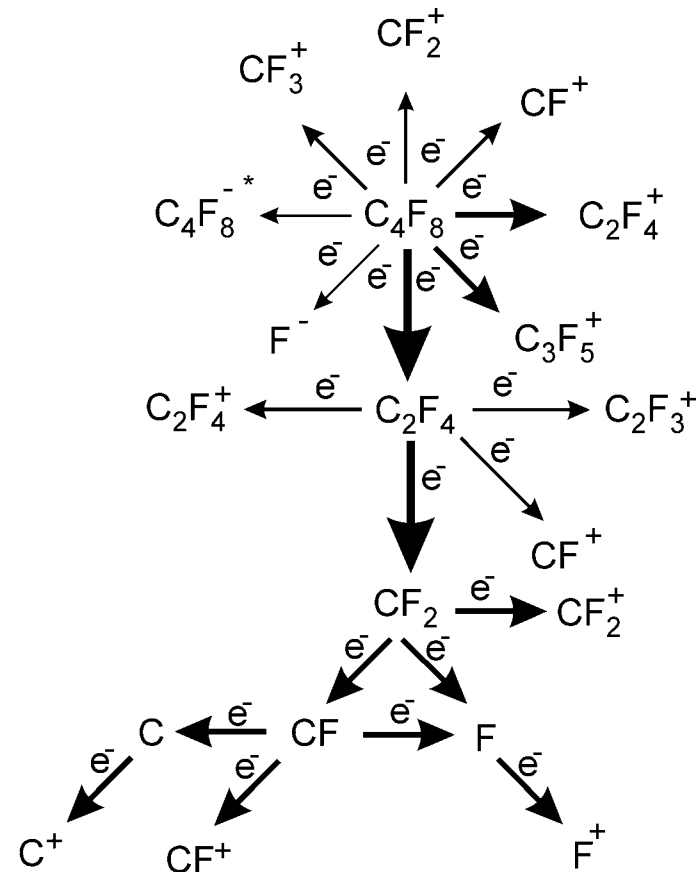
• *PLASMA CHEMISTRY*
(Are we getting this right?)

REACTION MECHANISMS FOR PLASMA ETCHING

- Recipes for plasma etching of dielectric materials (e.g., SiO_2 , Si_3N_4) often contain mixtures of many gases such as:

Ar , C_4F_8 , O_2 , N_2 , CO

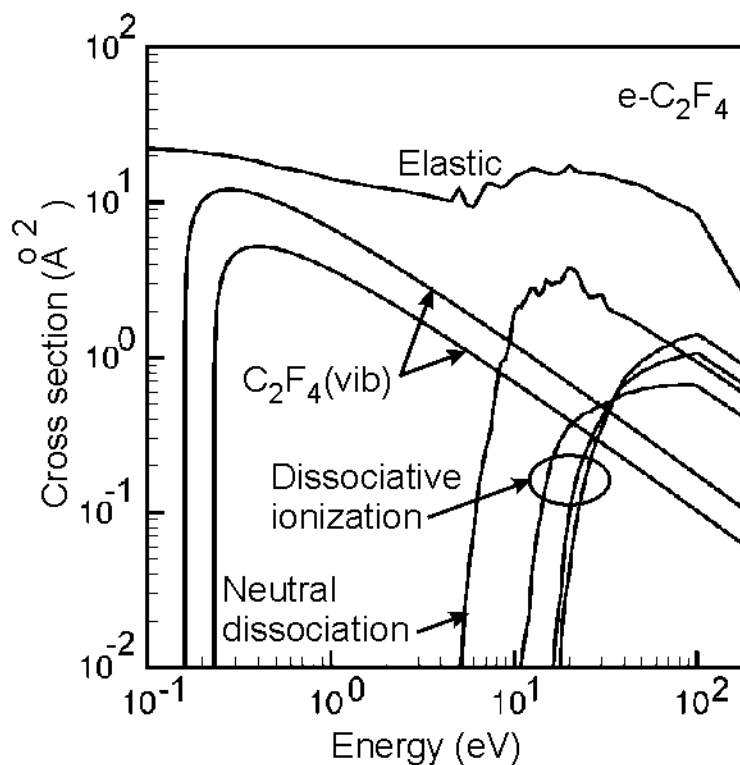
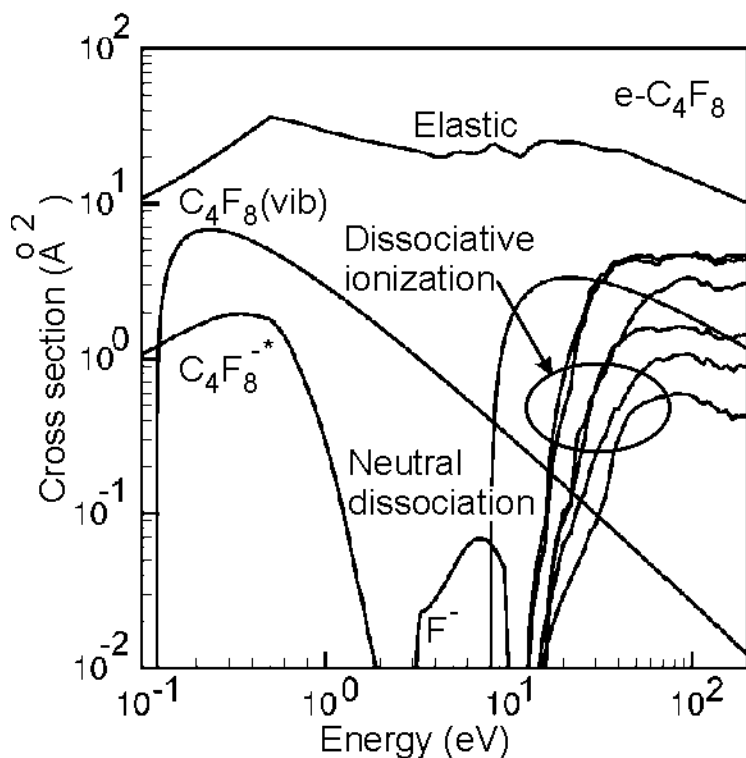
- The fluorocarbon donors are often highly dissociated, thereby requiring databases for both feedstocks and their fragments.
- For predictive modeling, reaction mechanisms must be developed for arbitrary mixtures and wide ranges of pressures.



University of Illinois
Optical and Discharge Physics

C_4F_8 , C_2F_4 CROSS SECTION SETS

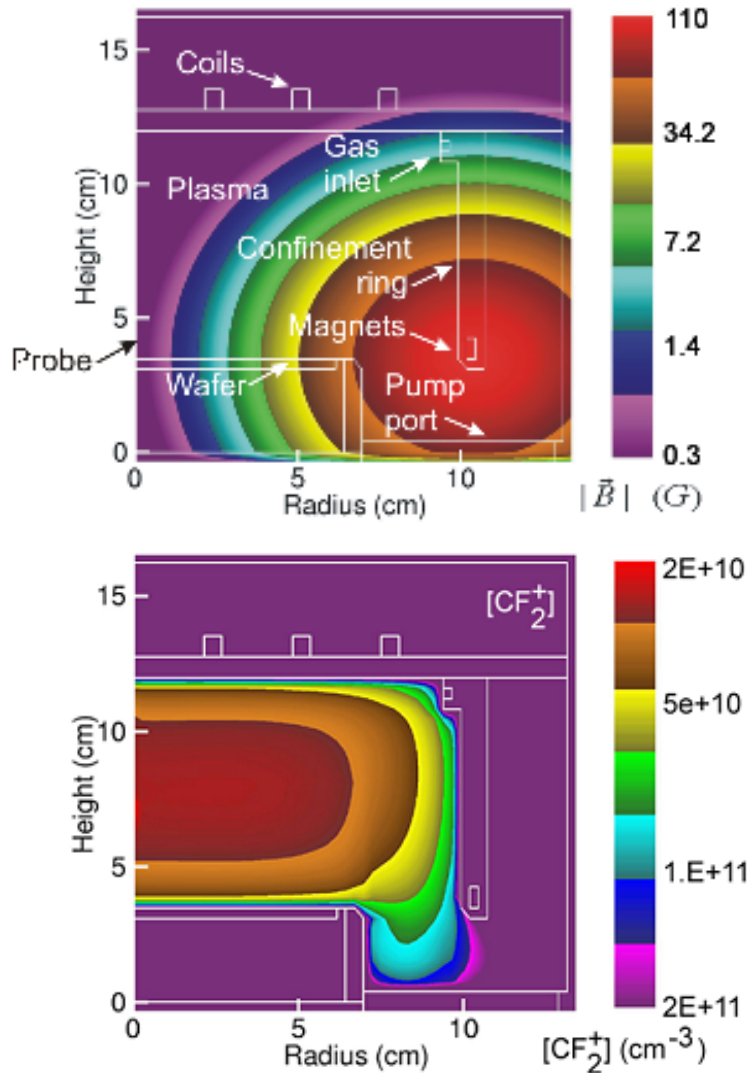
- The first step in developing a reaction mechanism is compilation of electron impact cross section sets.



- Ref: V. McKoy and W. L. Morgan

University of Illinois
Optical and Discharge Physics

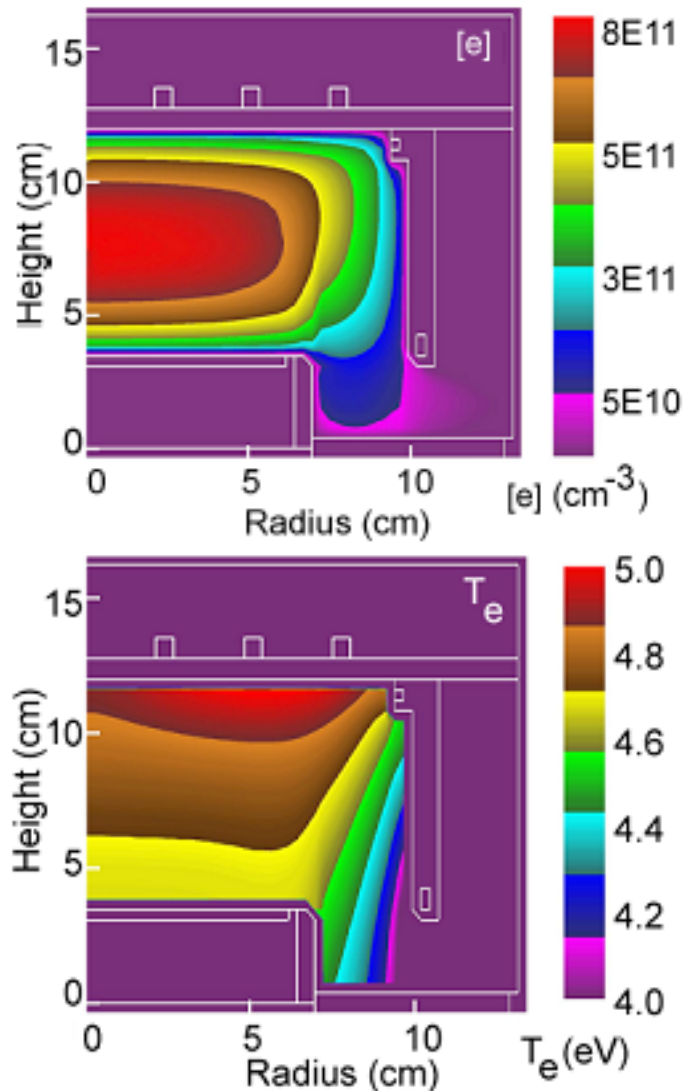
ICP CELL AND $[\text{CF}_2^+]$ FOR C_4F_8 , 10 mTorr



- An ICP reactor patterned after Oeherlein, et al. was used for validation.
- Reactor has a metal ring with magnets to confine plasma.
- CF_2^+ is one of the dominant ions in C_4F_8 plasmas due to large dissociation.
- The major path for the CF_2^+ is:
 - $\text{C}_4\text{F}_8 + e \rightarrow \text{C}_2\text{F}_4 + \text{C}_2\text{F}_4 + e$
 - $\text{C}_2\text{F}_4 + e \rightarrow \text{CF}_2 + \text{CF}_2 + e$
 - $\text{CF}_2 + e \rightarrow \text{CF}_2^+ + e + e$

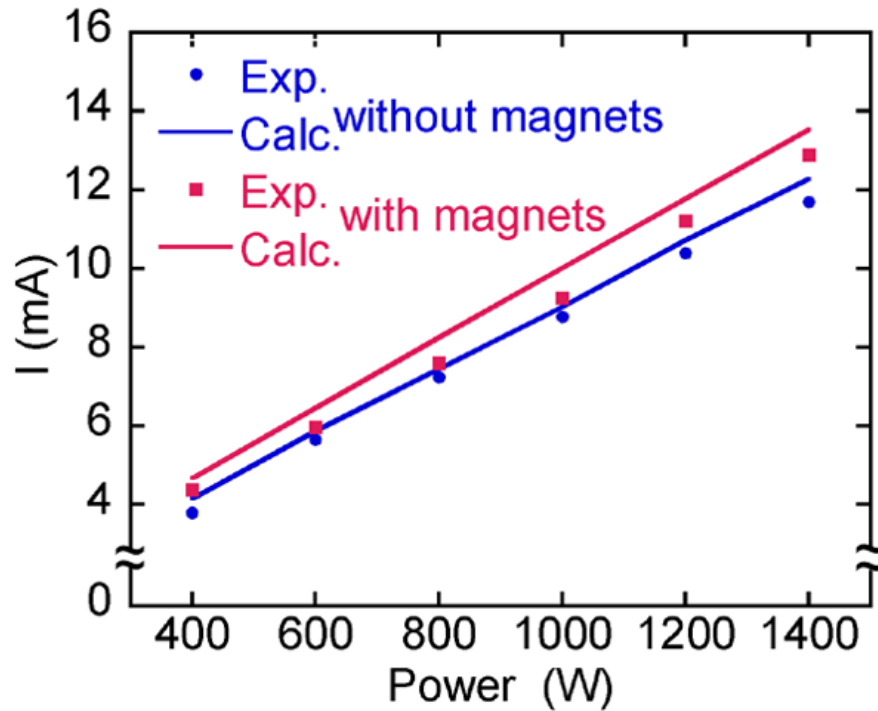
- C_4F_8 , 10 mTorr, 1.4 kW, 13.56 MHz

$[n_e]$ and T_e FOR C_4F_8 , 10 mTorr

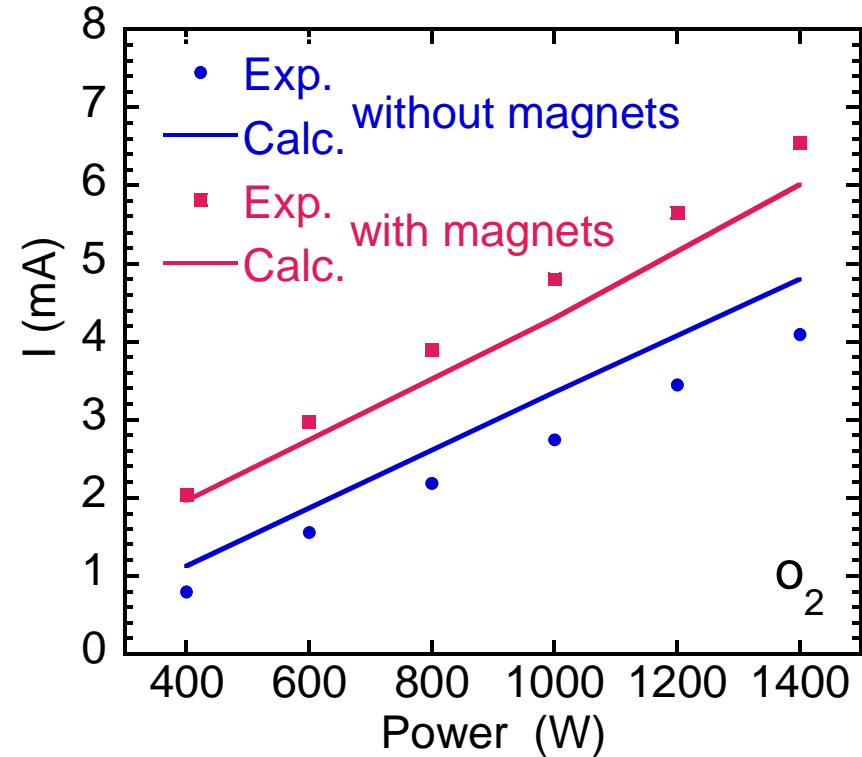


- Electron density peaks at $\approx 10^{12} cm^{-3}$.
- The peak in T_e occurs in the skin layer due to collisionless electron heating by the large electric field.
- T_e is rather uniform in the bulk plasma where electrons thermalize through e-e collisions.

I_p (PROBE CURRENT) IN ICPs SUSTAINED IN Ar, O₂



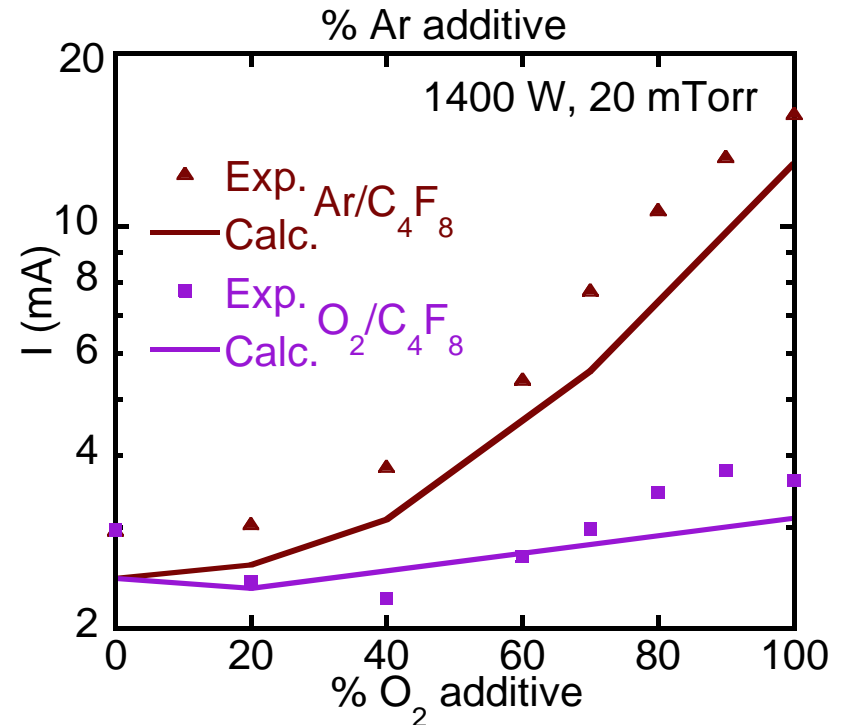
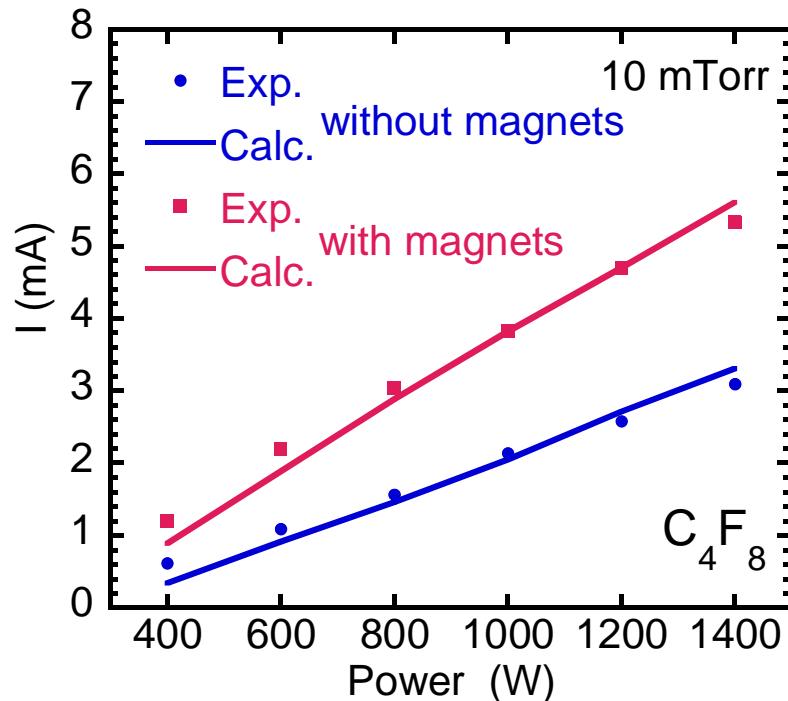
- Ar, 10 mTorr



- O₂, 10 mTorr

- Magnetic confinement is generally more effective in electronegative plasmas with a larger variety of ions.

I_p VERSUS POWER FOR ICPs IN C_4F_8 , Ar/C_4F_8 , O_2/C_4F_8

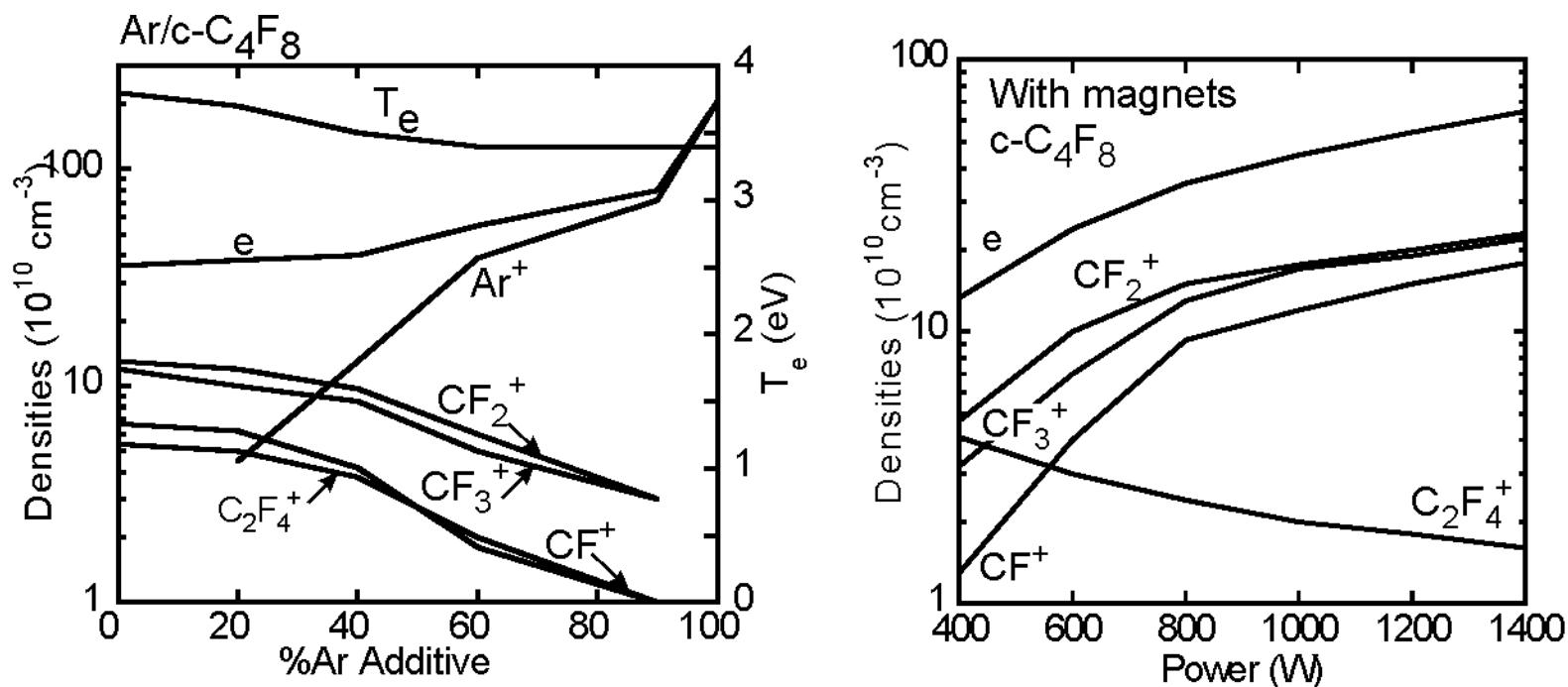


- The differences in I_p with and without magnets increases with power due to increased non-linear Lorentz force.
- I_p increases with Ar addition in Ar/C_4F_8 compared to Ar/O_2 due to higher dissociation of C_4F_8 and lower electronegativity.

• 13.56 MHz, -100 V probe bias.

University of Illinois
Optical and Discharge Physics

ION COMPOSITION IN C_4F_8 , Ar/ C_4F_8

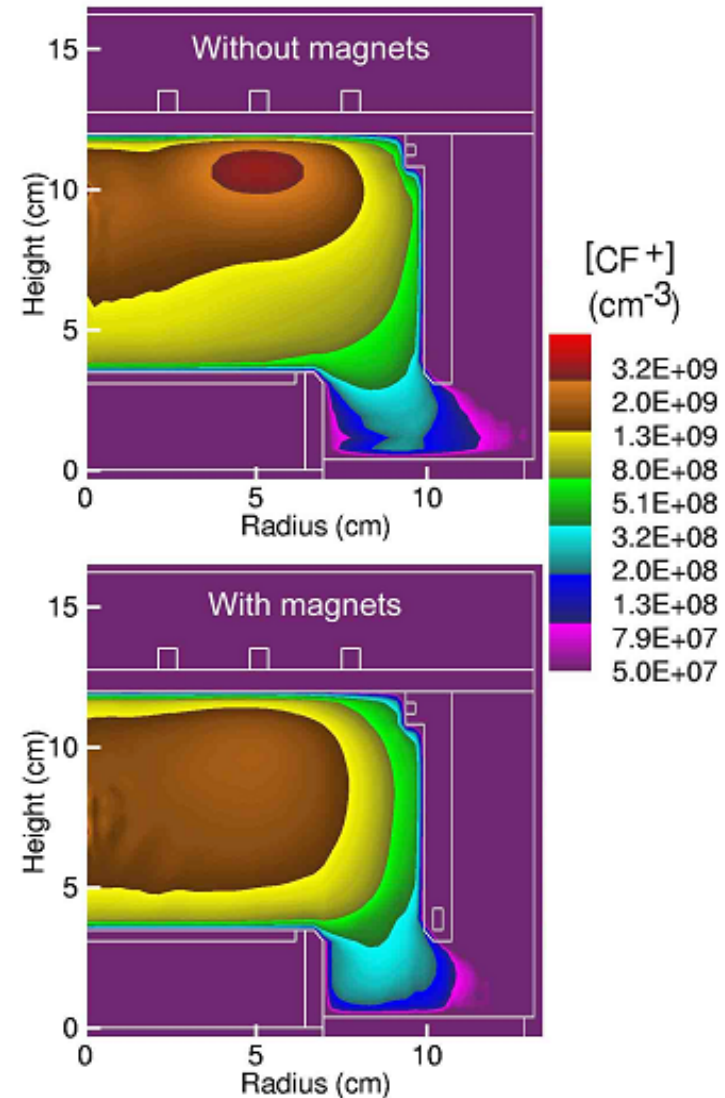


- Optimization of processing conditions on, for example, power critically depends on the composition of the radical and ion fluxes.

- 10 mTorr, 13.56 MHz

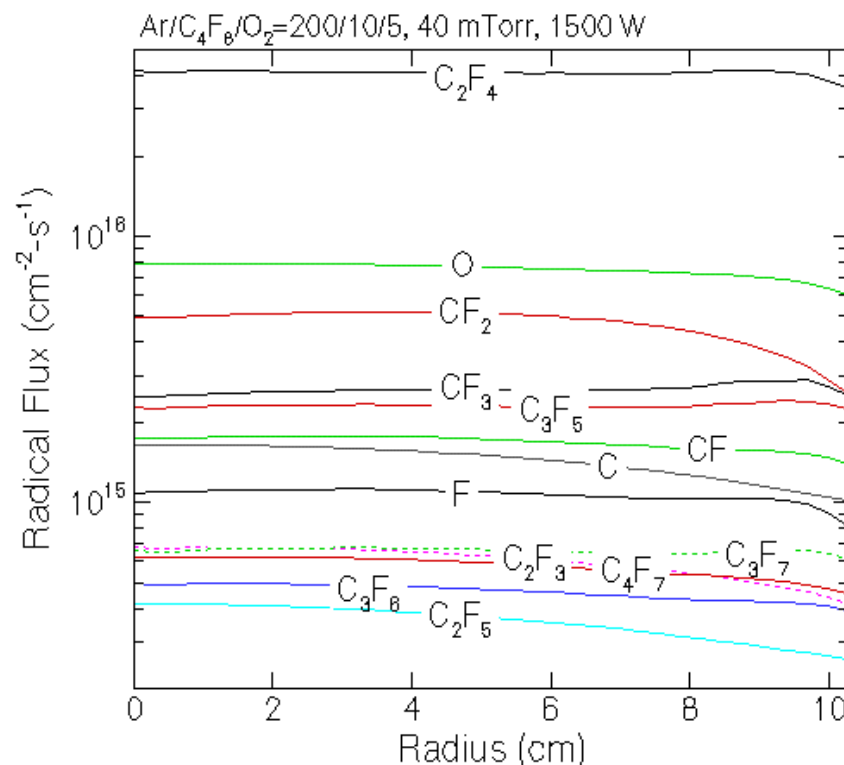
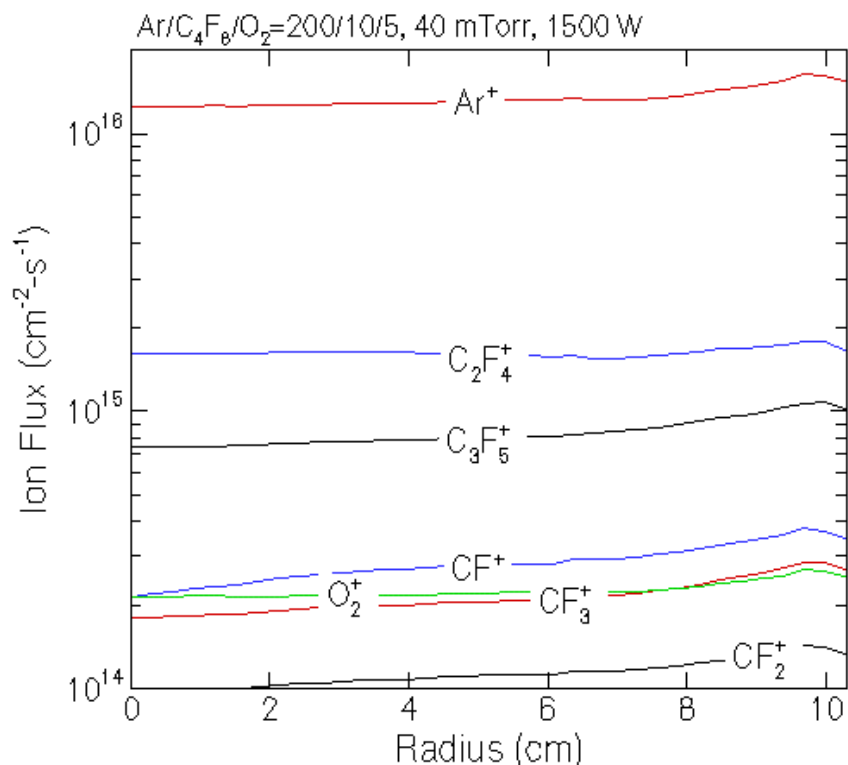
EFFECT OF MAGNETS ON [CF⁺]

- Without magnets [CF⁺] has a maximum at the edge of the classical skin depth where the electron impact ionization is the largest.
- The static magnetic fields broaden the production of [CF⁺] in the radial direction.



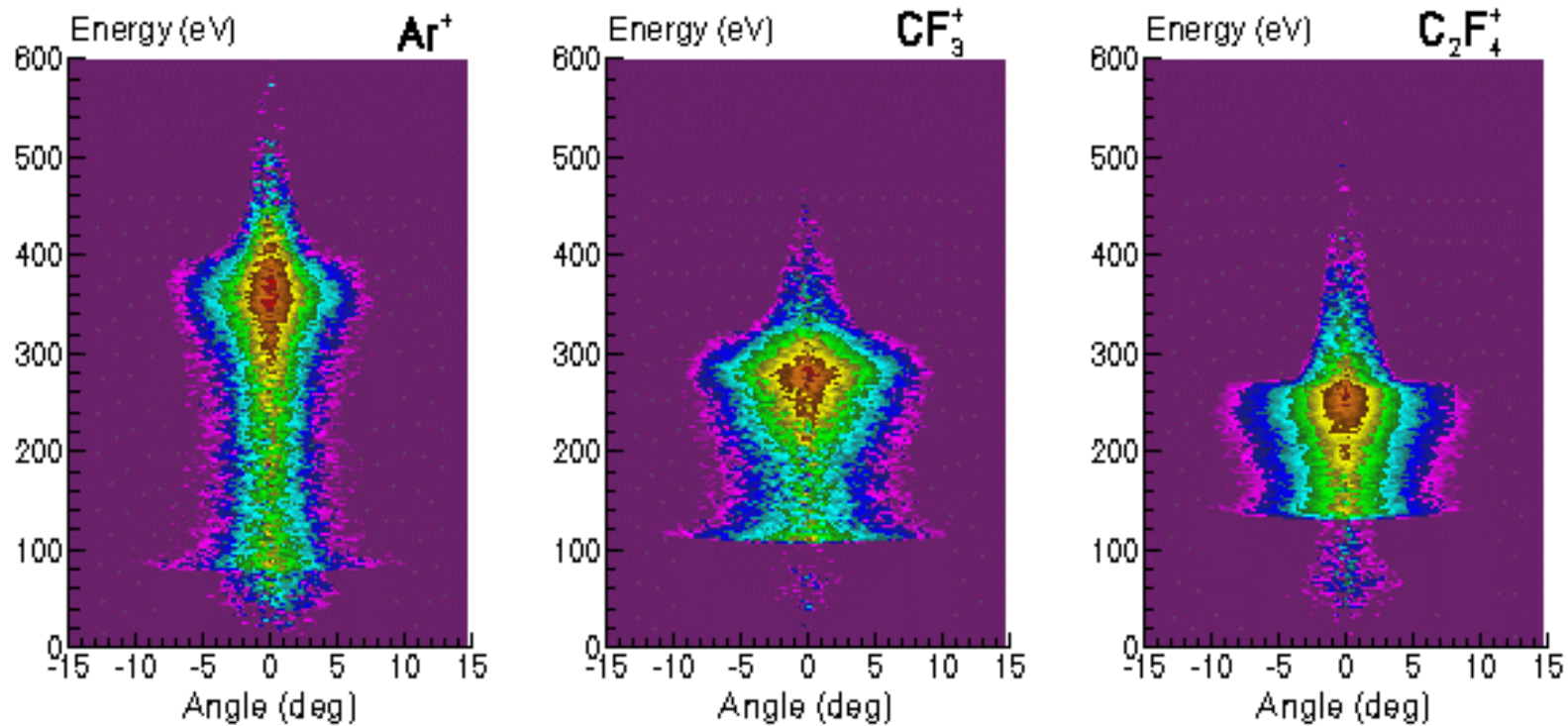
- Ar/C₄F₈=20/80, 3 mTorr, 13.56 MHz, 400 W.

TEL-DRM: Ar / C₄F₈ / O₂



- With reaction mechanisms developed for Ar / C₄F₈ / O₂ and improved ability to model MERIE systems, parameterizations were performed for TEL-DRM like conditions.
- Ar / C₄F₈ / O₂ = 200/10/5 sccm, 40 mTorr, 1500 W.

TEL-DRM: Ar / C₄F₈ / O₂

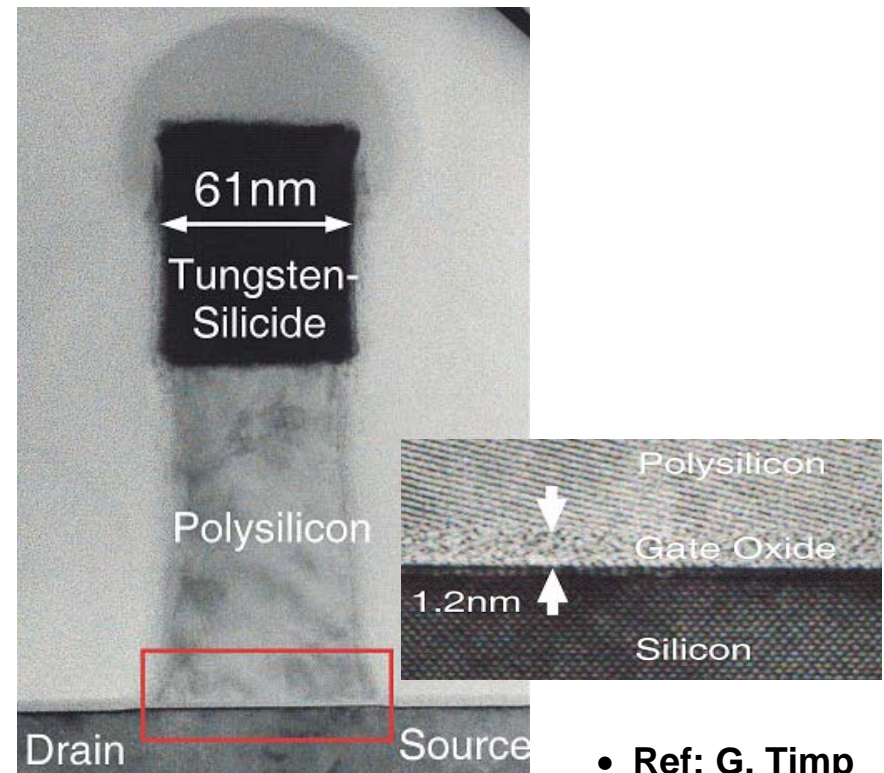


- The large variety of ion masses produces vastly different IEADs.
- Ar / C₄F₈ / O₂ = 200/10/5 sccm, 40 mTorr, 1500 W.

• *SURFACE CHEMISTRY*
(The most ill defined but
perhaps most important step.)

SELECTIVITY IN MICROELECTRONICS FABRICATION: PLASMAS AND POLYMERS

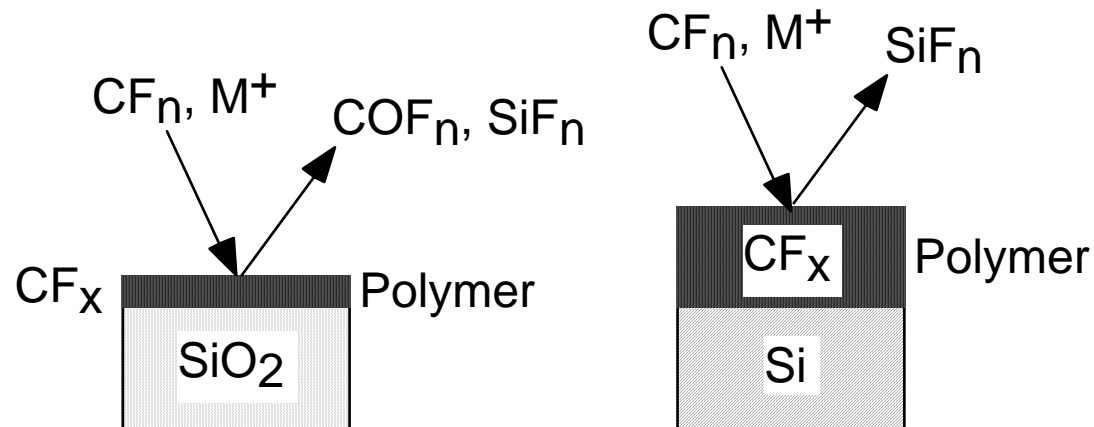
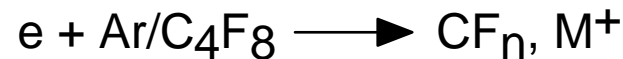
- Fabricating complex microelectronic structures made of different materials requires extreme *selectivity* in, for example, etching Si with respect to SiO_2 .
- Monolayer selectivity is required in advanced etching processes.
- These goals are met by the unique plasma-polymer interactions enabled in fluorocarbon chemistries.



University of Illinois
Optical and Discharge Physics

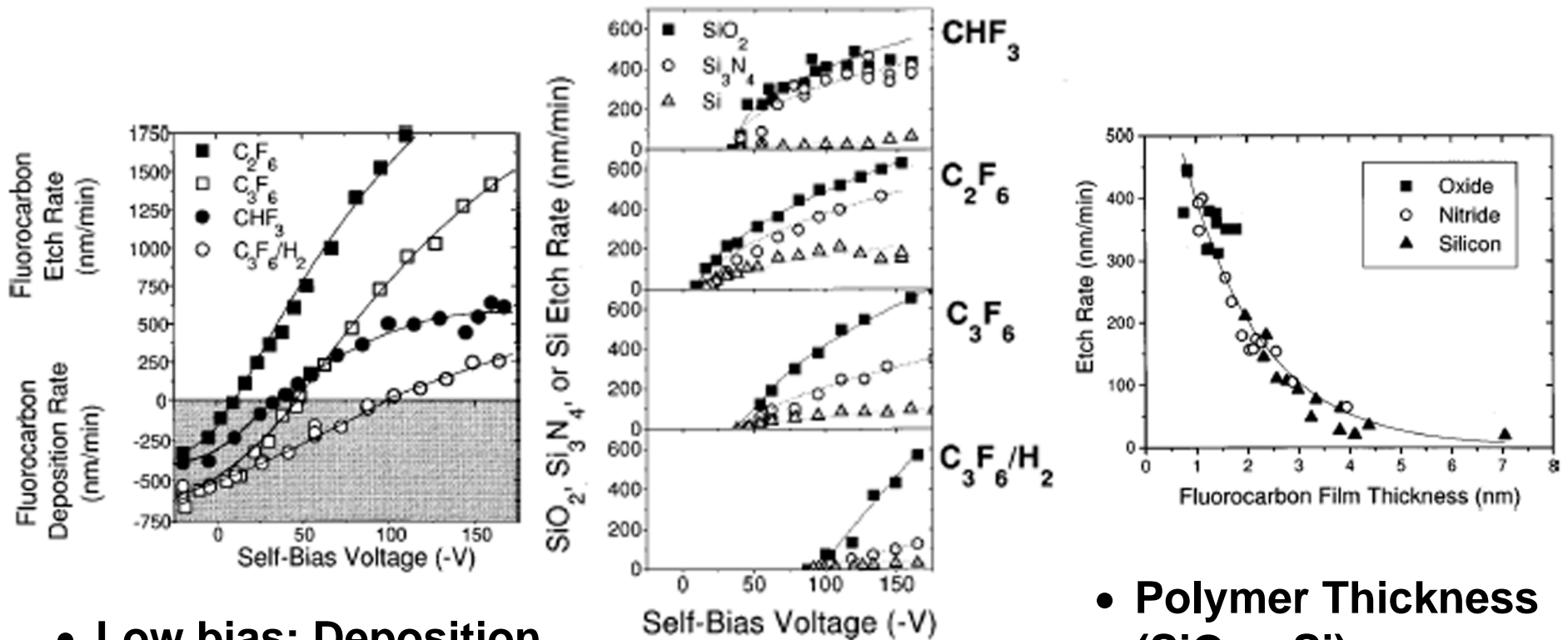
FLUOROCARBON PLASMA ETCHING: SELECTIVITY

- Selectivity in fluorocarbon etching relies on polymer deposition.
- Electron impact dissociation of feedstock fluorocarbons produce polymerizing radicals and ions, resulting in polymer deposition.



- Compound dielectrics contain oxidants which consume the polymer, producing thinner polymer layers.
- Thicker polymer on non-dielectrics restrict delivery of ion energy (lower etching rates).

FLUOROCARBON PLASMA ETCHING: SELECTIVITY



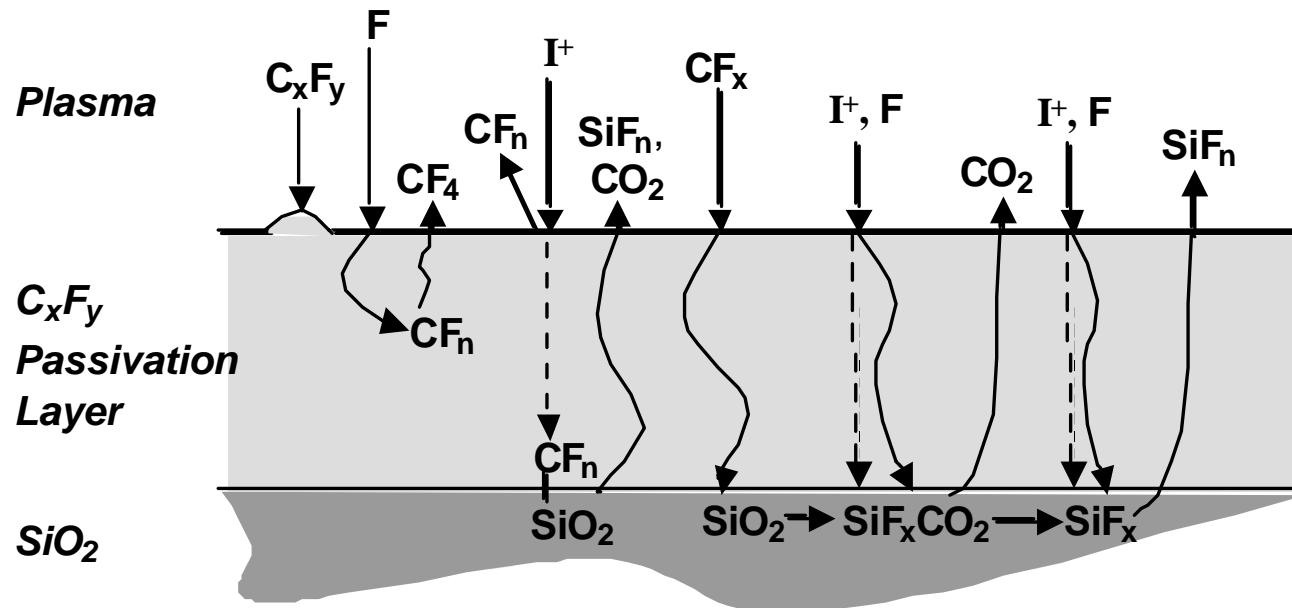
- Low bias: Deposition
- High bias: etching
- Etch Rate ($SiO_2 > Si$)

- Polymer Thickness ($SiO_2 < Si$)

• G. Oerhlein, et al., JVSTA 17, 26 (1999)

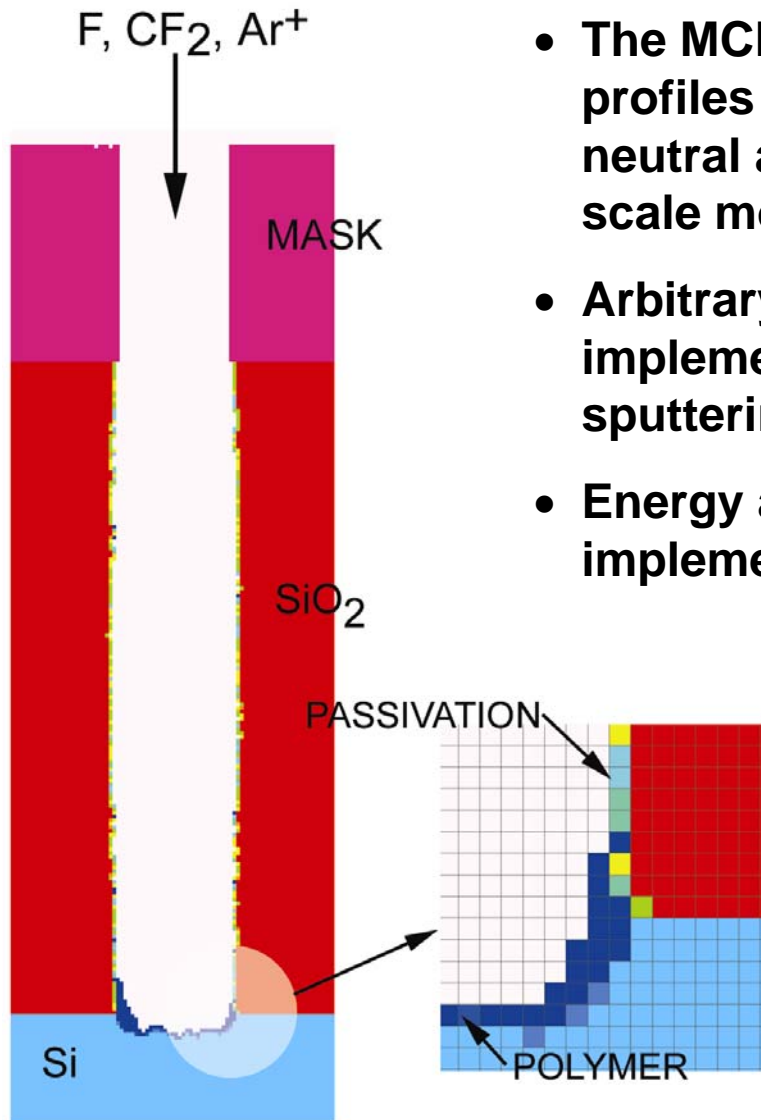
SURFACE KINETICS: FLUOROCARBON PLASMA ETCHING Si/SiO₂

- C_xF_y passivation regulates delivery of precursors and activation energy.
- Chemisorption of CF_x produces a complex at the oxide-polymer interface.
- 2-step ion activated (through polymer layer) etching of the complex consumes the polymer. Activation scales inversely with polymer thickness.
- Etch precursors and products diffuse through the polymer layer.



- In Si etching, CF_x is not consumed, resulting in thicker polymer layers.

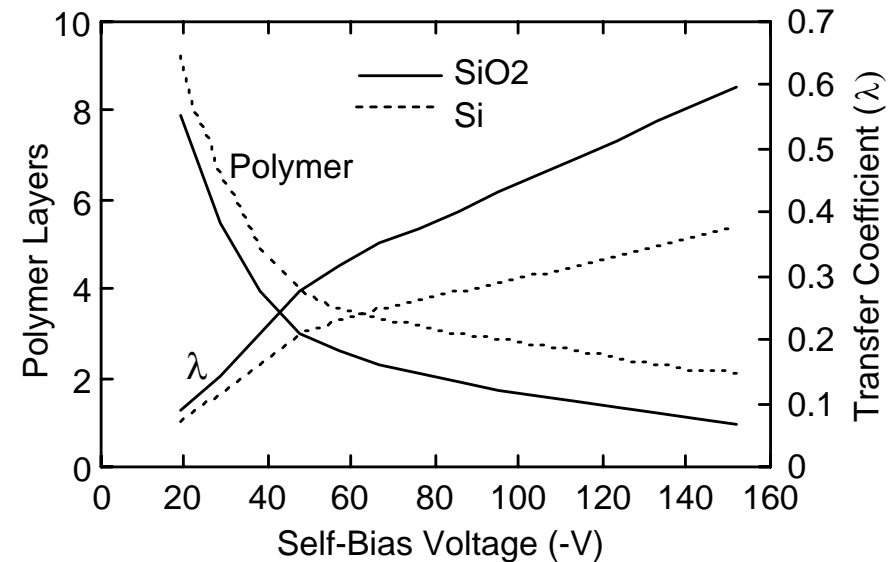
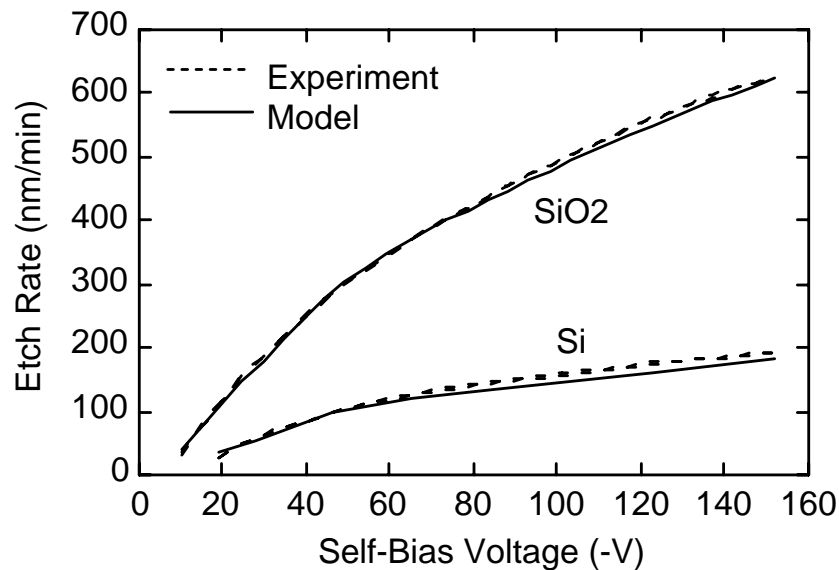
MONTE CARLO FEATURE PROFILE MODEL (MCFPM)



- The MCFPM predicts time and spatially dependent profiles using energy and angularly resolved neutral and ion fluxes obtained from equipment scale models.
- Arbitrary chemical reaction mechanisms may be implemented, including thermal and ion assisted, sputtering, deposition and surface diffusion.
- Energy and angular dependent processes are implemented using parametric forms.
- Mesh centered identify of materials allows “burial”, overlayers and transmission of energy through materials.

ETCH RATES AND POLYMER THICKNESS

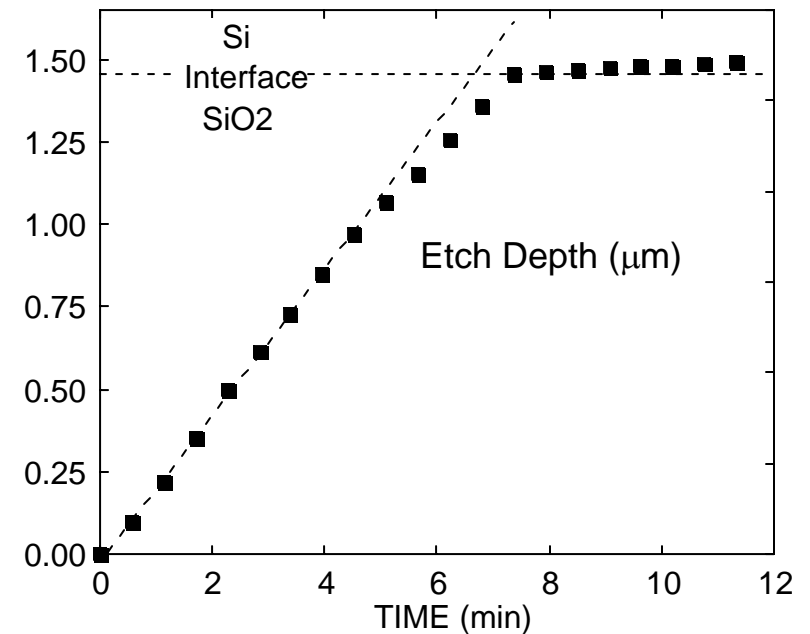
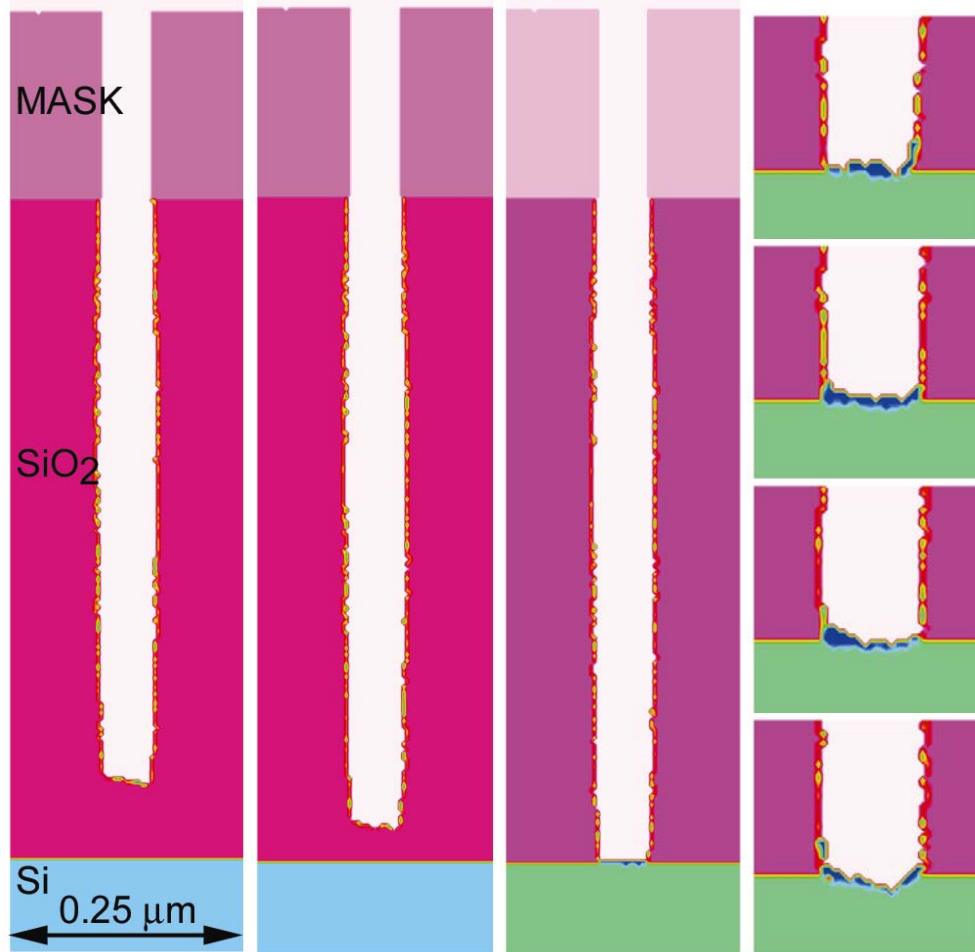
- Etch rates for Si and SiO₂ increase with increasing bias due, in part, to a decrease in polymer thickness.
- The polymer is thinner with SiO₂ due to its consumption during etching, allowing for more efficient energy transfer through the layer and more rapid etching.



- **C₂F₆, 6 mTorr, 1400 W ICP, 40 sccm**
- **Exp. Ref: T. Standaert, et al.
J. Vac. Sci. Technol. A 16, 239 (1998).**

POLYMERIZATION AIDS SELECTIVITY

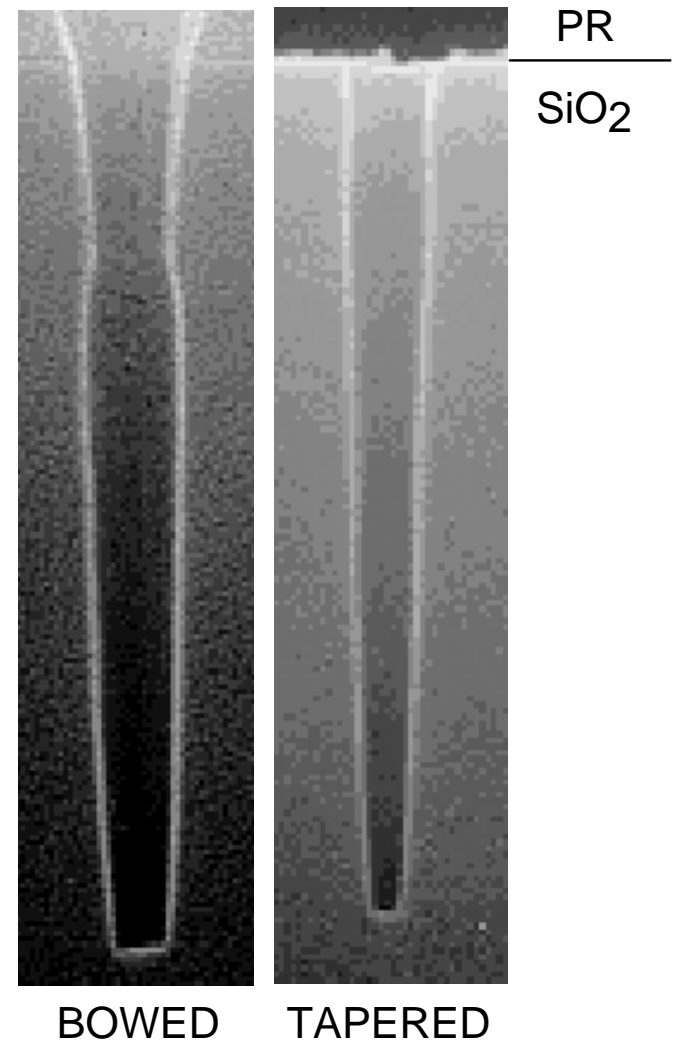
- Less consumption of polymer on Si relative to SiO₂ slows and, in some cases, terminates etching, providing high selectivity.



University of Illinois
Optical and Discharge Physics

TAPERED AND BOWED PROFILES

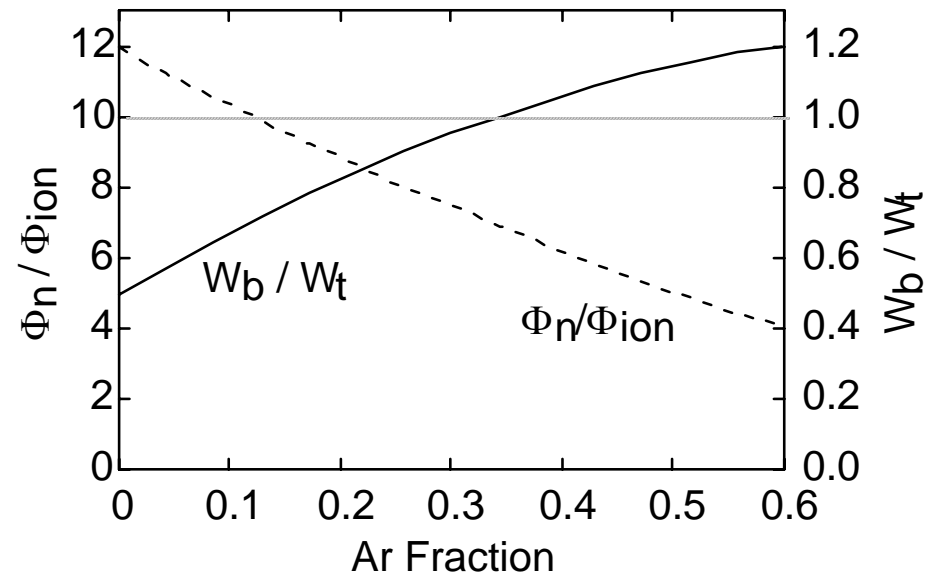
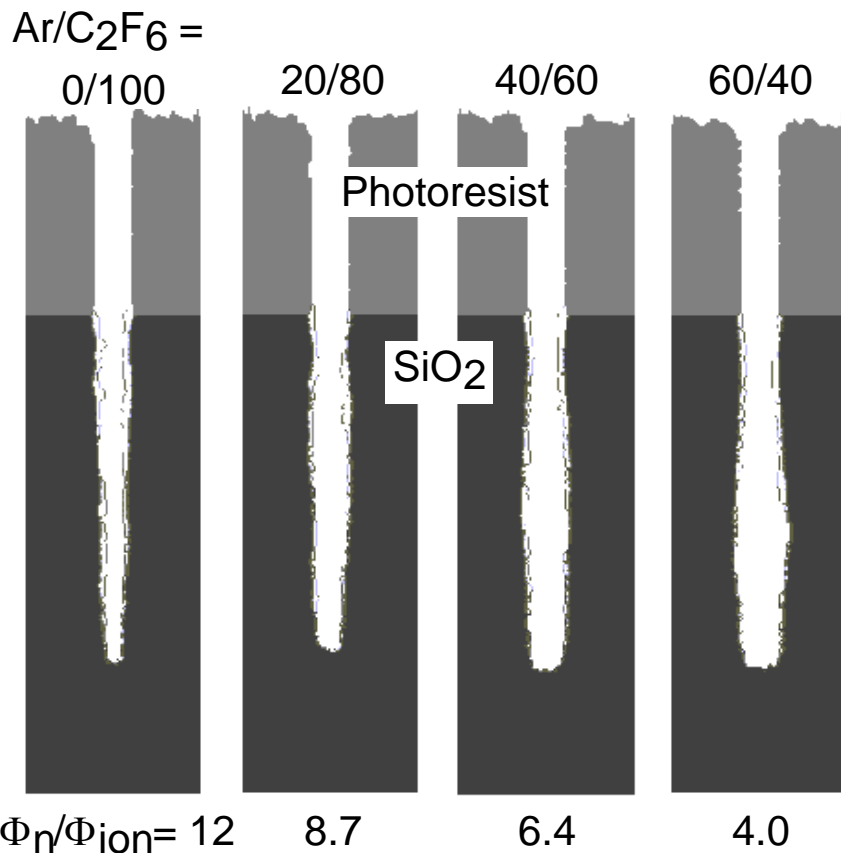
- In high aspect ratio (HAR) etching of SiO_2 the sidewall of trenches are passivated by neutrals (CF_x , $x \leq 2$) due to the broad angular distributions of neutral fluxes.
- Either tapered or bowed profiles can result from a non-optimum combination of processing parameters including:
 - Degree of passivation
 - Ion energy distribution
 - Radical/ion flux composition.



University of Illinois
Optical and Discharge Physics

PROFILE TOPOLOGY: NEUTRAL TO ION FLUX RATIO

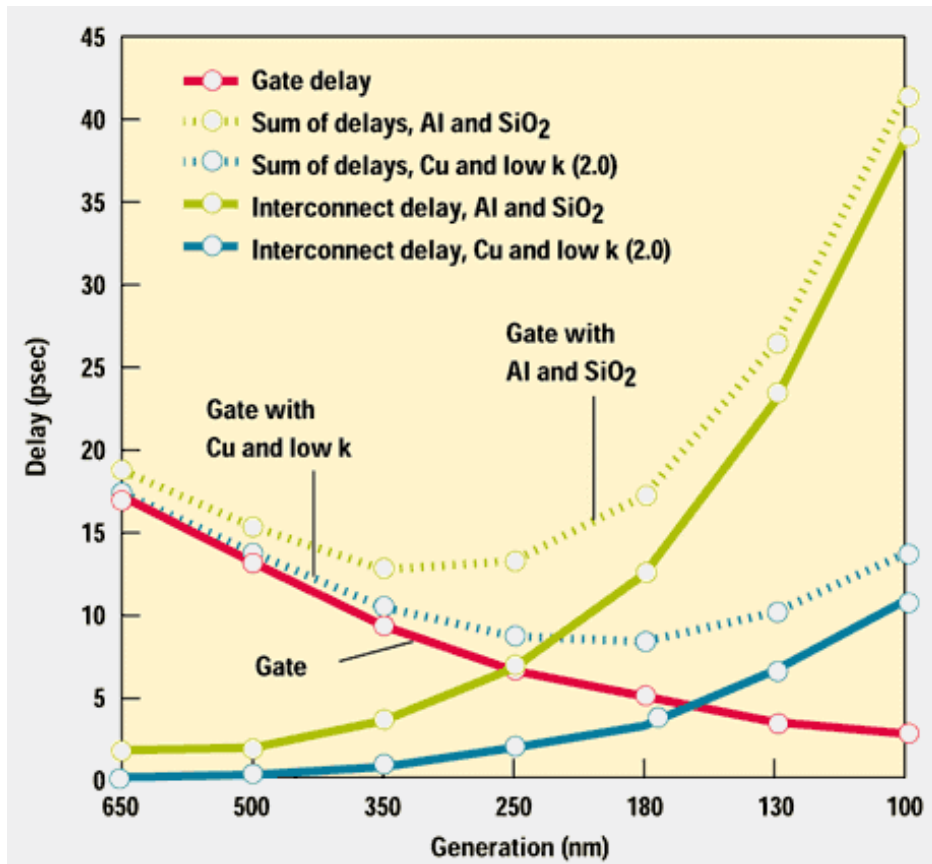
- Profiles depend on ratio of polymer forming fluxes to energy activating fluxes. Small ratios produce bowing, large ratios tapering.
- Controlling this ratio through gas mixture (e.g., Ar/C₂F₆) enables specification of profile topology.



University of Illinois
Optical and Discharge Physics

LOW-K DIELECTRICS

- As feature sizes decrease and device count increases, the diameter of interconnect wires shrinks and path length increases.



- Large RC-delay limits processor performance.
- To reduce RC-delay, low dielectric constant (low-k) materials are being investigated.

• L. Peters, *Semi. Intl.*, 9/1/1998

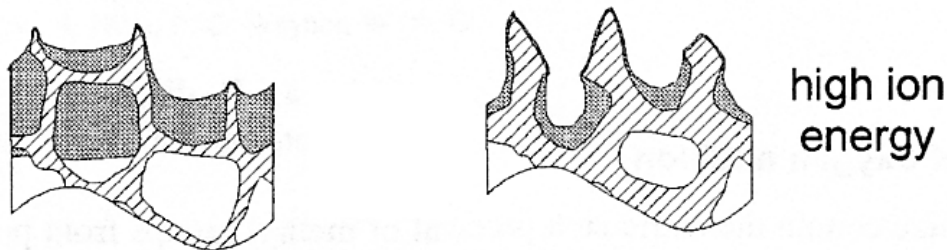
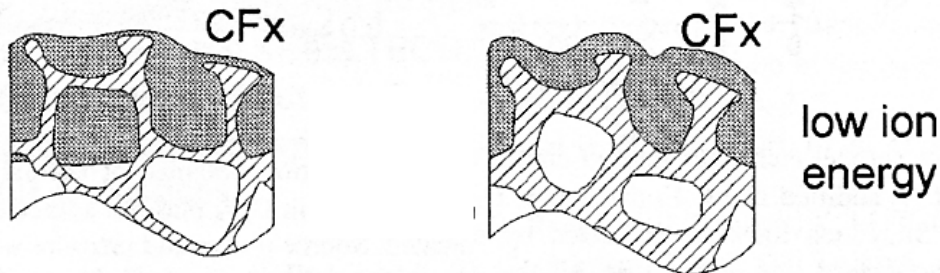
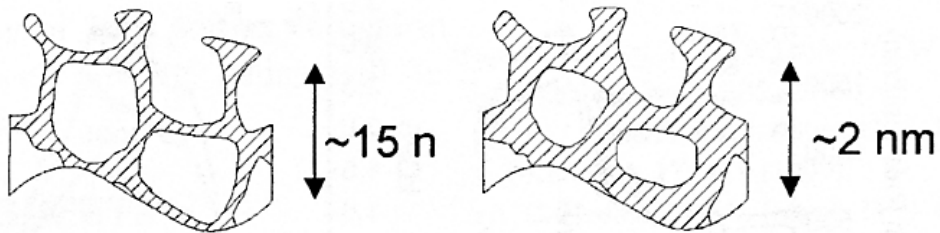
University of Illinois
Optical and Discharge Physics

POROUS SILICON DIOXIDE

- Porous SiO₂ (xerogels) have low-k properties due to their lower mass density resulting from (vacuum) pores.
 - Typical porosities: 30-70%
 - Typical pore sizes: 2-20 nm
- Porous SiO₂ (P-SiO₂) is, from a process development viewpoint, an ideal low-k dielectric.
 - Extensive knowledge base for fluorocarbon etching of conventional non-porous (NP-SiO₂).
 - No new materials (though most P-SiO₂ contains some residual organics)
 - Few new integration requirements

WHAT CHANGES WITH POROUS SiO₂?

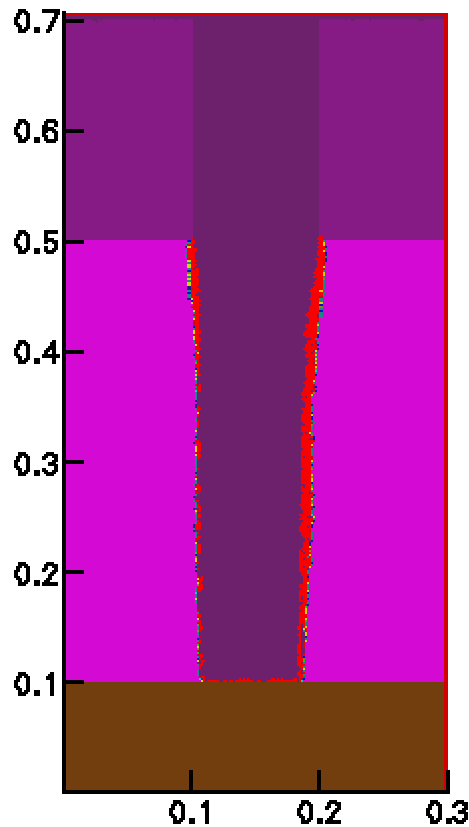
- The “opening” of pores during etching of P-SiO₂ results in the filling of the voids with polymer, creating thicker layers.
- Ions which would have otherwise hit at grazing or normal angle now intersect with more optimum angle.



- An important parameter is L/a (polymer thickness / pore radius).

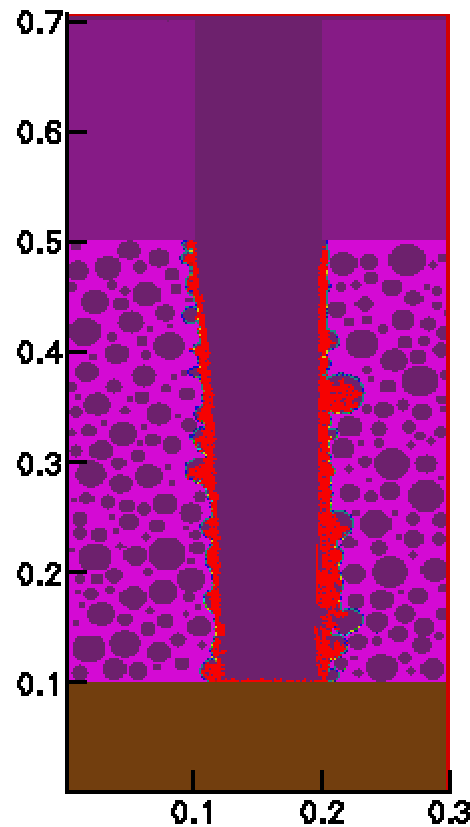
- Adapted: Standaert, JVSTA 18, 2742 (2000)

ETCH PROFILES IN SOLID AND POROUS SiO₂



• Position (μm)

• **Solid**



• Position (μm)

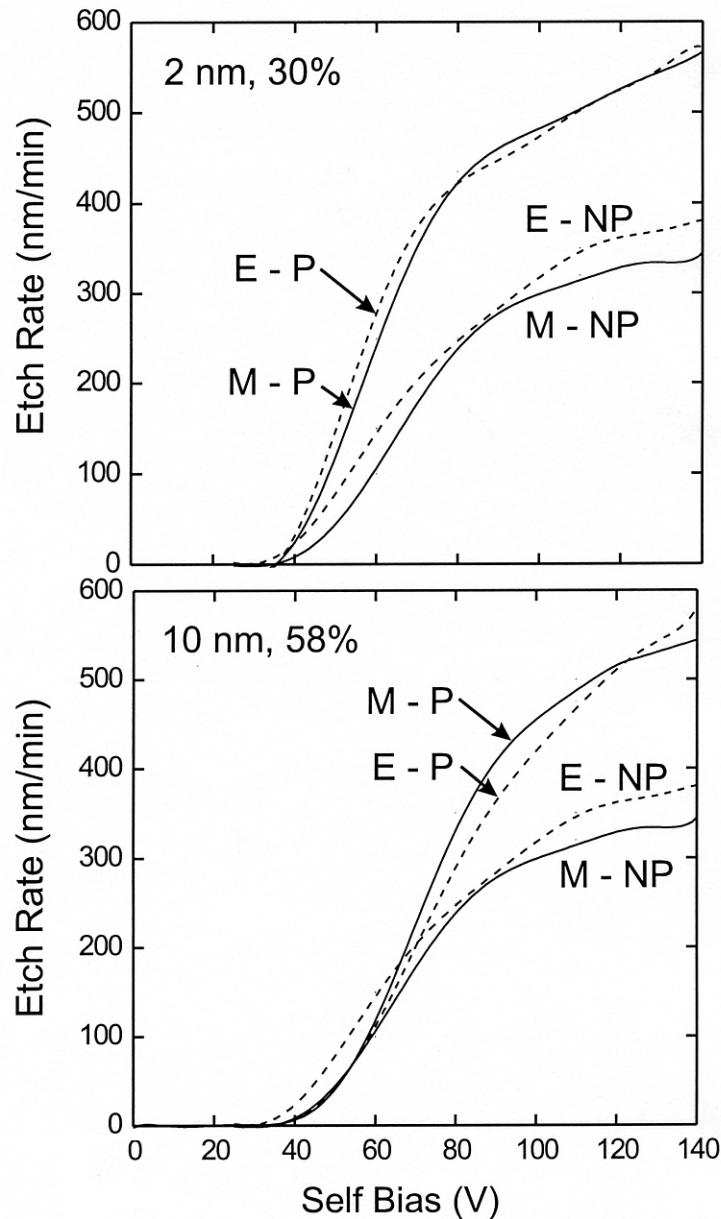
• Porosity = 45 %
Pore radius = 10 nm

- Porous SiO₂ is being investigated for low-permittivity dielectrics for interconnect wiring.
- In polymerizing environments with heavy sidewall passivation, etch profiles differ little between solid and porous silica.
- The “open” sidewall pores quickly fill with polymer.

ANIMATION SLIDE

University of Illinois
Optical and Discharge Physics

ETCHING OF POROUS SiO₂



- Etch rates of P-SiO₂ are generally higher than for non-porous (NP).
- Examples:
 - 2 nm pore, 30% porosity
 - 10 nm pore, 58% porosity
- Higher etch rates are attributed to lower mass density of P-SiO₂.
- CHF₃ 10 mTorr, 1400 W

Exp: Oehrlein et al. Vac. Sci. Technol. A **18**, 2742 (2000)

ADVMET_1002_23

University of Illinois
Optical and Discharge Physics

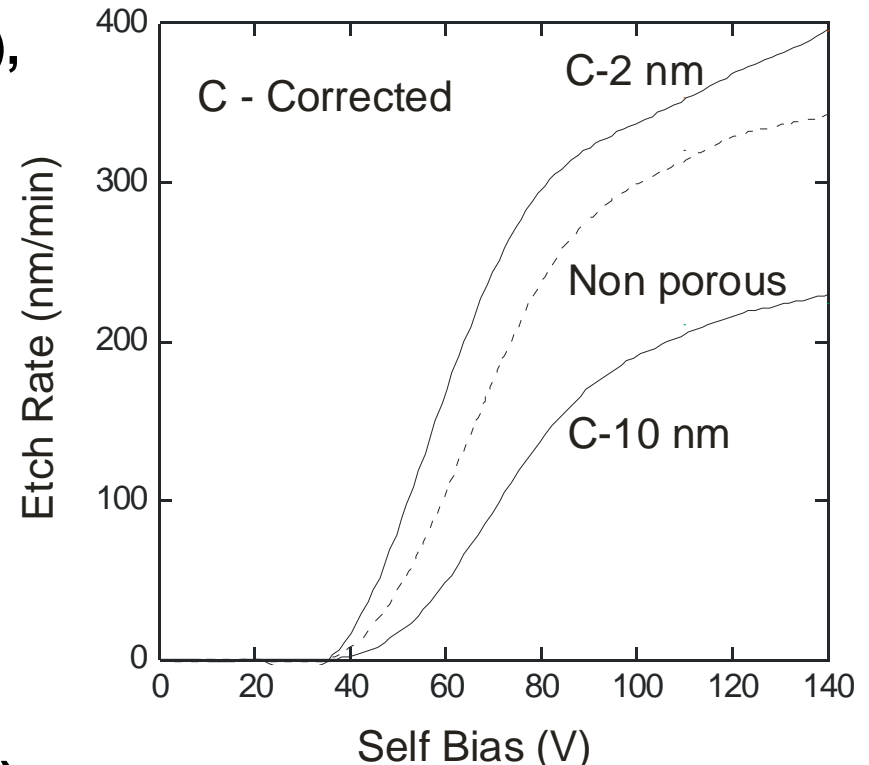
PORE-DEPENDENT ETCHING

- To isolate the effect of pores on etch rate, corrected etch rate is defined as

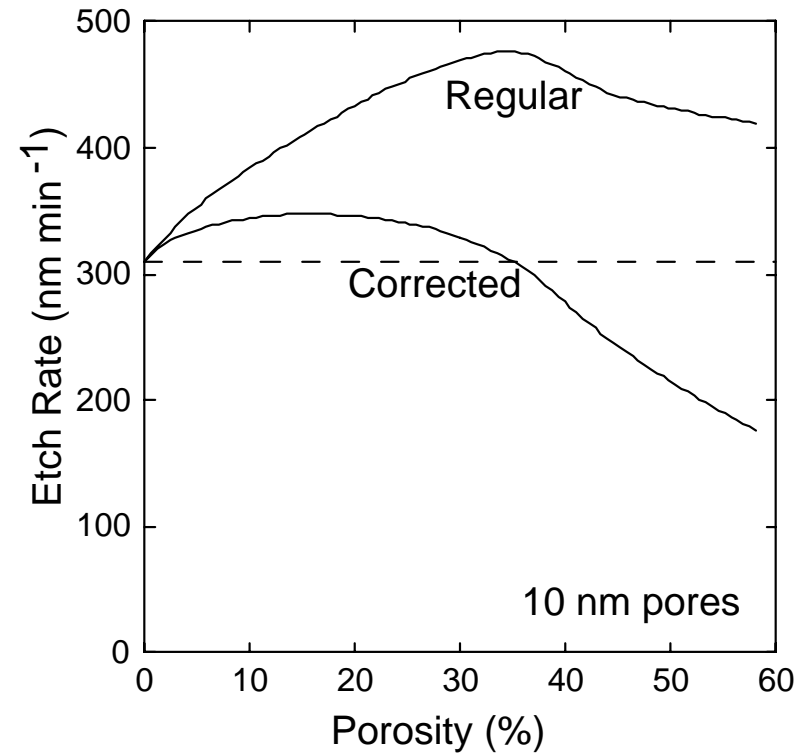
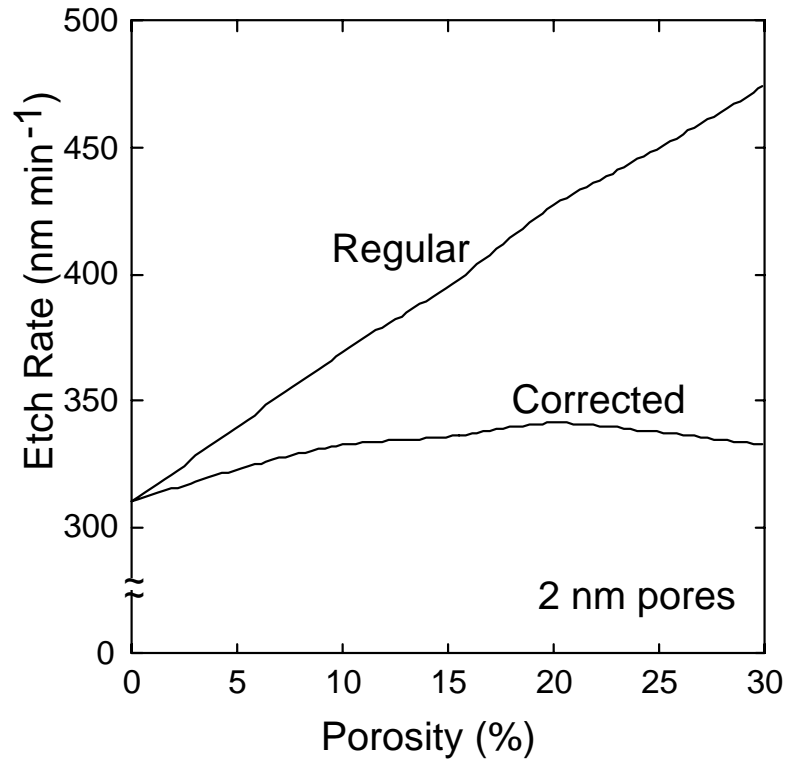
$$\text{Etch Rate (ER)}_{\text{corrected}} = \text{ER}_{\text{regular}} \times (1 - p),$$

p = porosity

- If etching depended only on mass density, corrected etch rates would equal that of NP- SiO_2 .
- 2 nm pores $L/a \geq 1$: C-ER > ER(SiO_2). Favorable yields due to non-normal incidence may increase rate.
- 10 nm pores $L/a \leq 1$: C-ER < ER(SiO_2). Filling of pores with polymer decrease rates.



EFFECT OF POROSITY ON BLANKET ETCH RATES

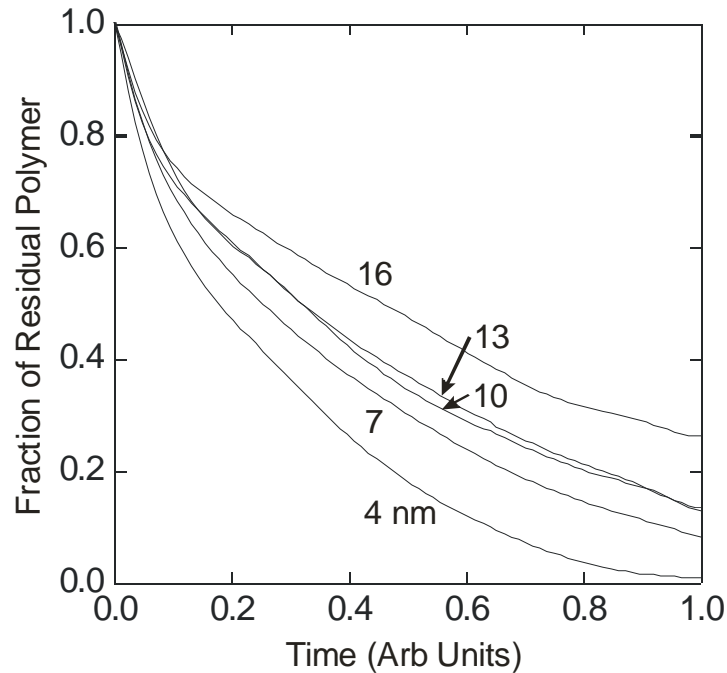


- **2 nm pores: Etch rate increases with porosity.**
- **10 nm pores: Polymer filling of pores reduces etch rate at large porosities.**

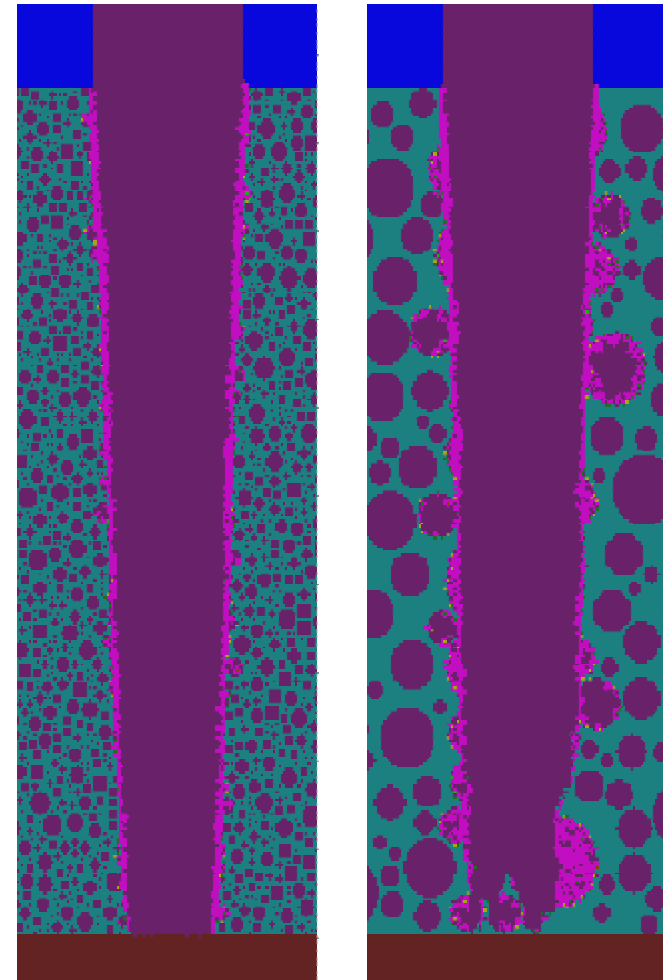
OXYGEN PLASMA CLEANING OF POLYMER

- After etching, the polymer must be removed from the feature.
- O₂ plasmas are typically used for polymer stripping, usually during photoresist mask removal.
- Unlike hydrocarbon polymers which spontaneously react with O, fluorocarbon polymers require ion activation for etching.
 - Polymer + Energetic Ion → Activated Polymer Site (P*)
 - P* + O → Volatile Products
- Removal of polymer from porous materials is difficult due to shadowing of ion fluxes caused by the pore morphology.

EFFECT OF PORE RADIUS ON CLEANING



- Larger pores are more difficult to clean due to small view angle of ion fluxes.
- Lower fluxes of less energetic ions reduce activation and lengthen cleaning time.



4 nm

16 nm

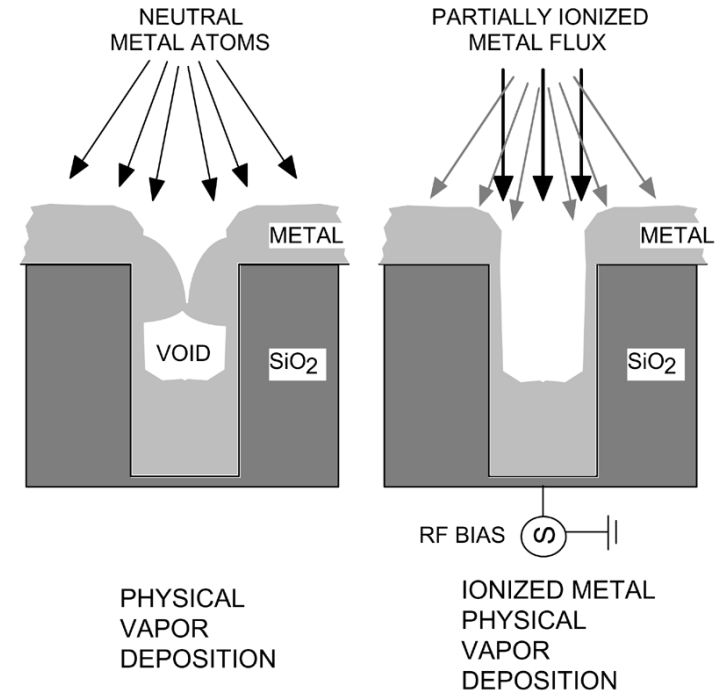
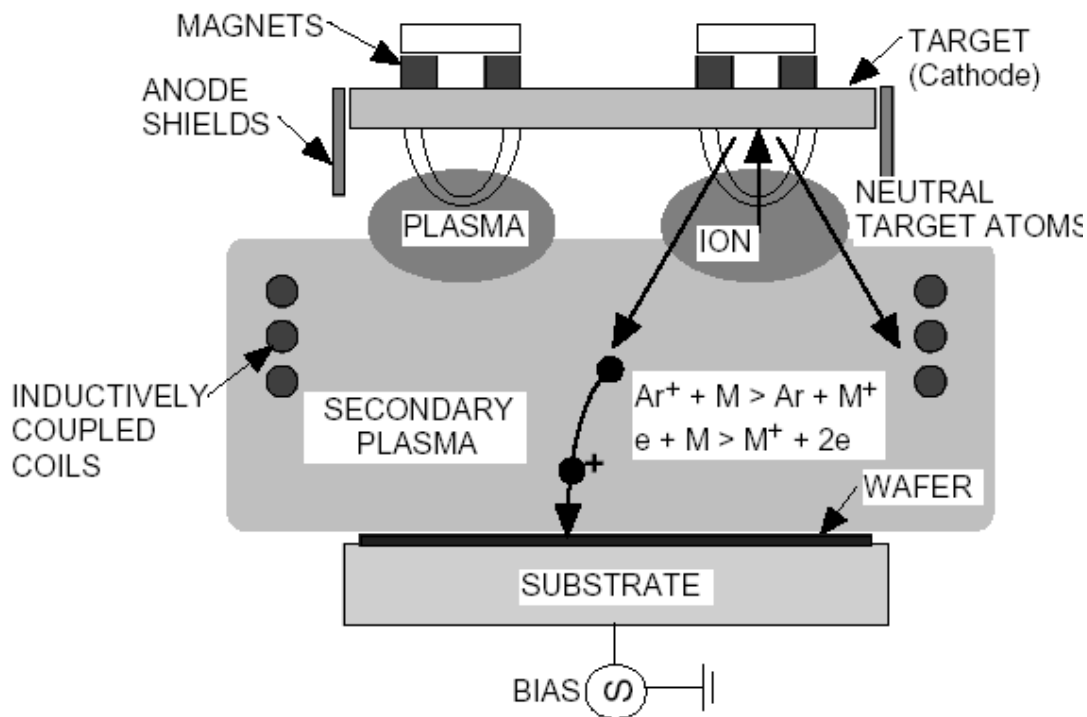
ANIMATION SLIDE

University of Illinois
Optical and Discharge Physics

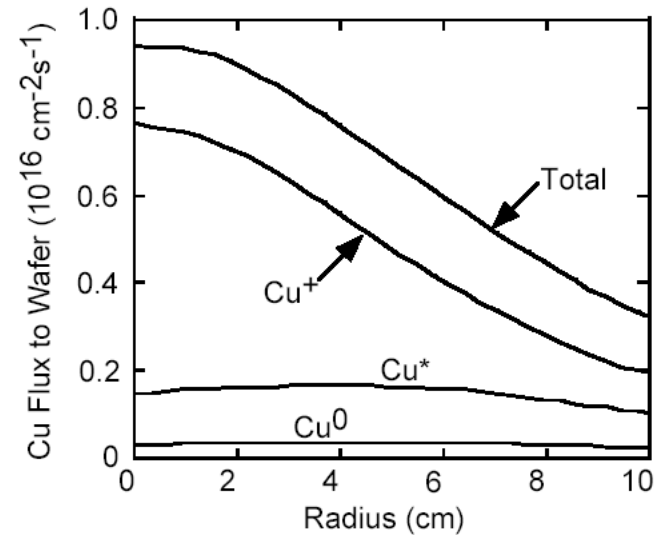
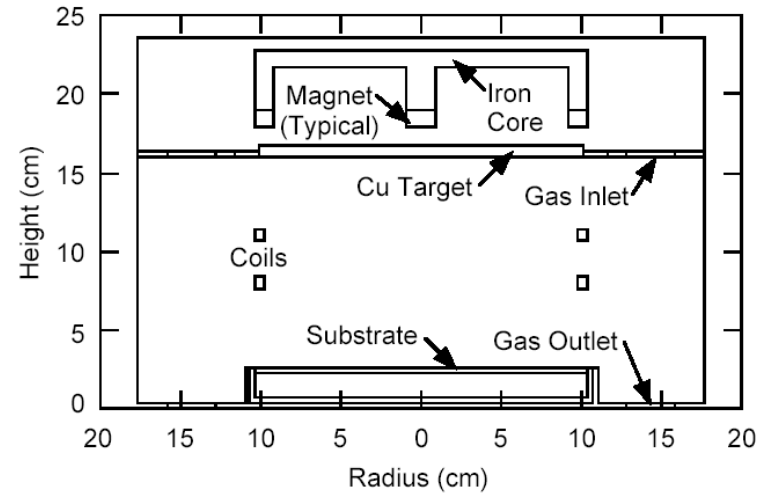
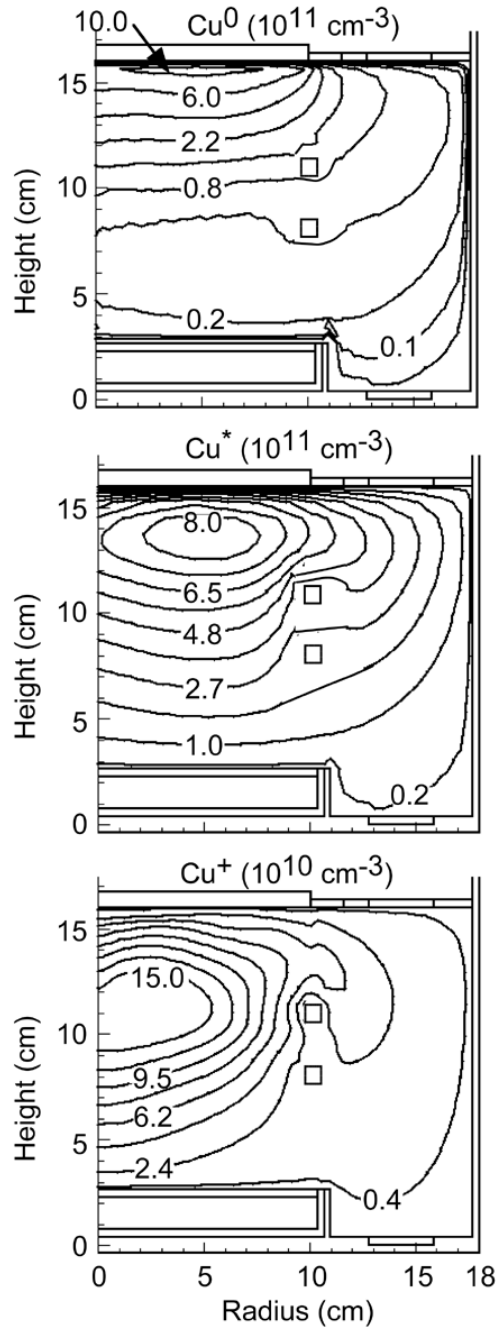
- ***TOWARDS INTEGRATED
PROCESS MODELING
(The last step...metal
deposition.)***

IONIZED METAL PHYSICAL VAPOR DEPOSITION (IMVPD)

- IMPVD is a technique to deposit seed layers and barrier coatings, and fill trenches.
- A flux of both neutral and metal atoms more uniformly produce depositions without formation of voids.



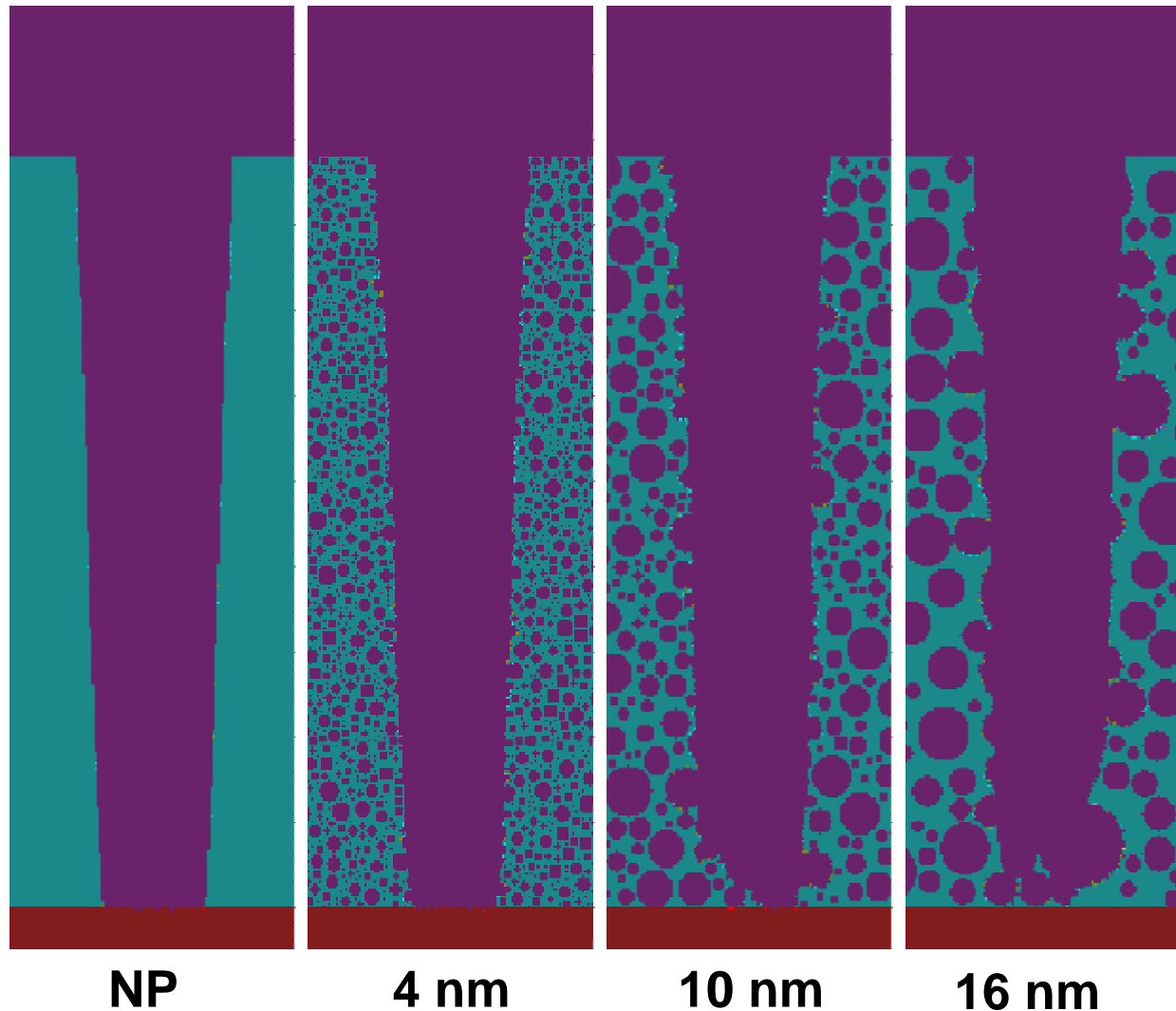
Cu IMVPD: REACTOR SCALE MODELING



- 40 mTorr Ar
- 1 kW ICP
- 0.3 kW Magnetron
- -25 V bias

University of Illinois
Optical and Discharge Physics

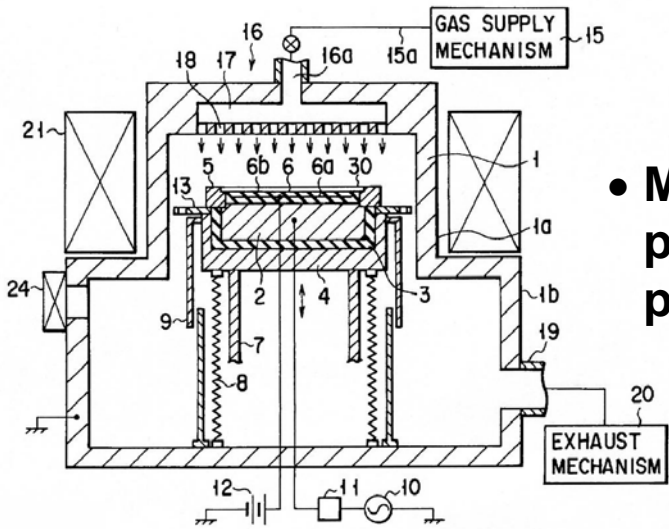
EFFECT OF PORE RADIUS ON Cu DEPOSITION



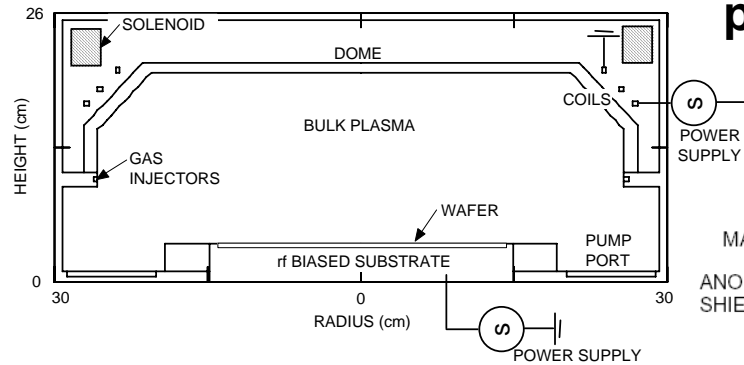
- Surrogate study for seed layer deposition and barrier coating.
- Voids are created at the pore surface or initiated due to the presence of pores.
- Presence of voids are pronounced for bigger pores.

ANIMATION SLIDE

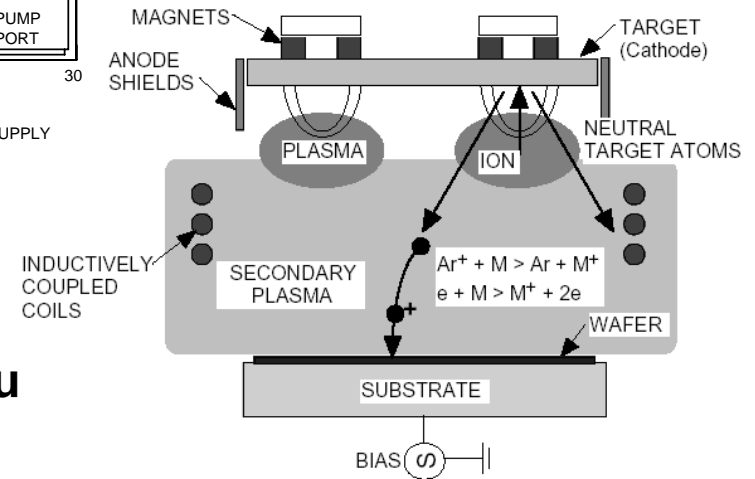
University of Illinois
Optical and Discharge Physics



- MERIE Fluorocarbon plasma etching of porous SiO₂

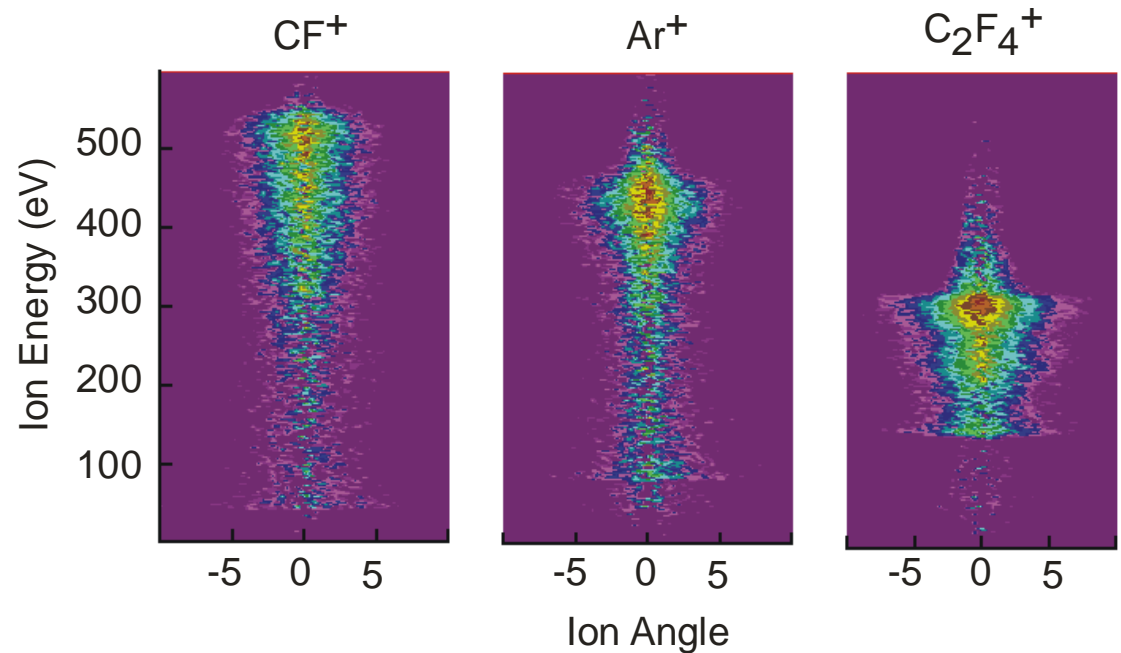
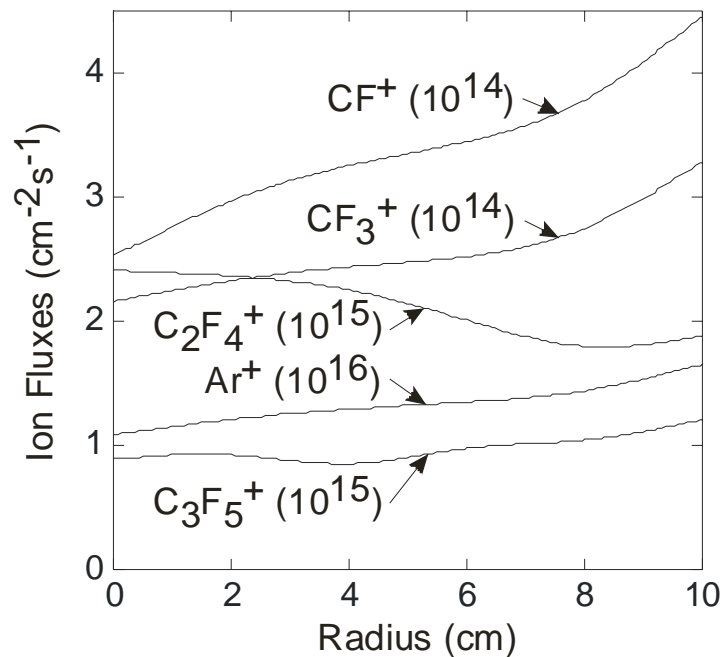


- ICP O₂ plasma cleaning of PR and polymer.



- IMPVD of Cu seed layer

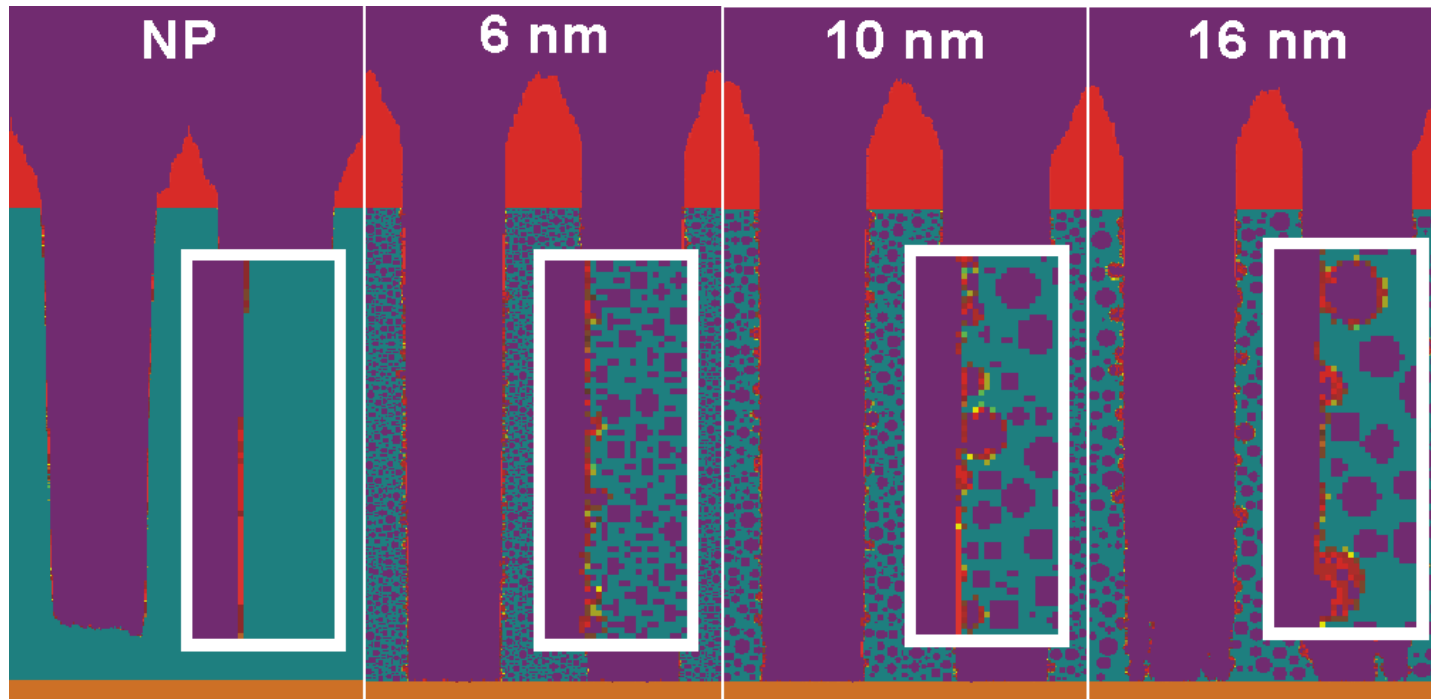
MERIE: ION FLUXES AND ENERGIES



- Ar/O₂/ C₄F₈ = 200/5/10 sccm
- 2000 W
- 40 mTorr

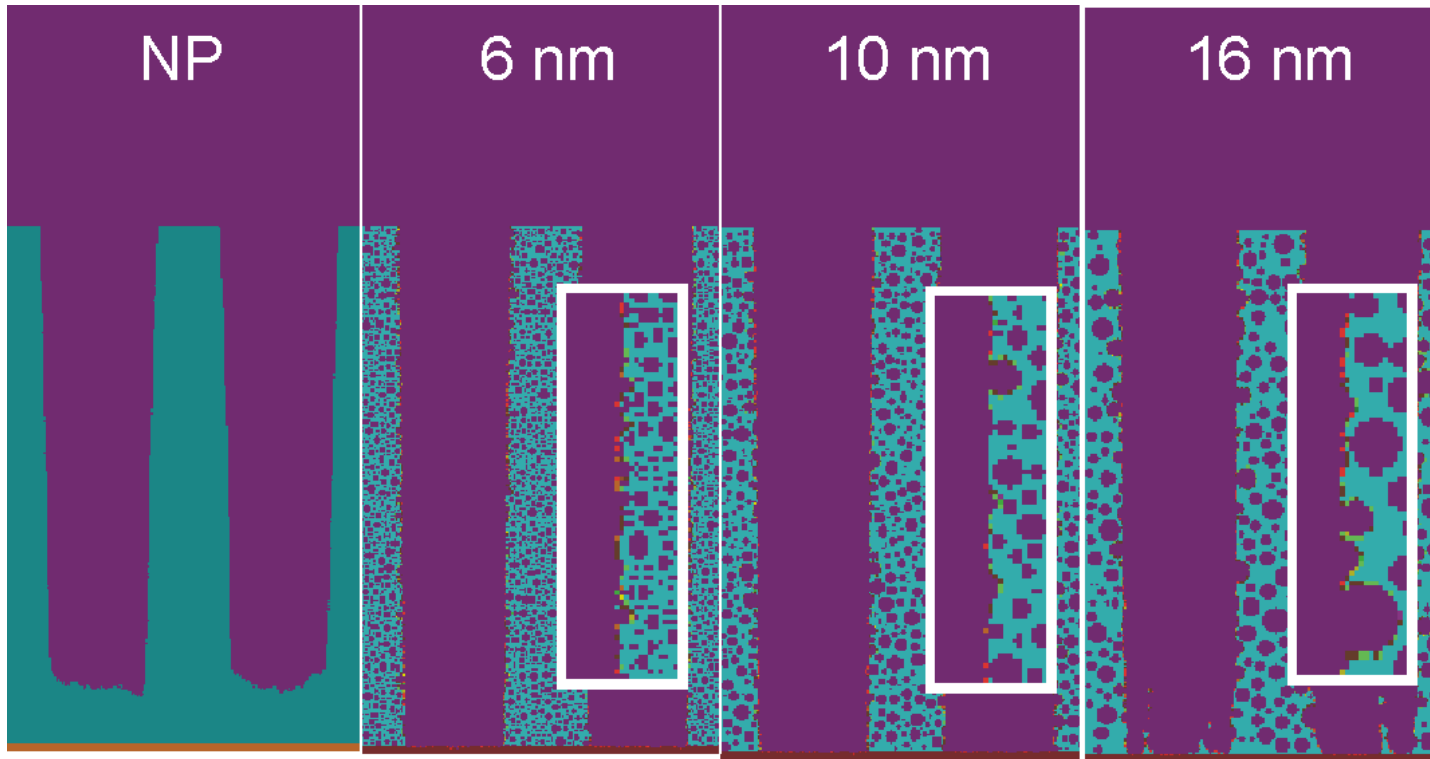
- Due to high dilution and low fractional dissociation, dominant ions are Ar⁺, C₂F₄⁺

MERIE: POROUS SiO₂ ETCH



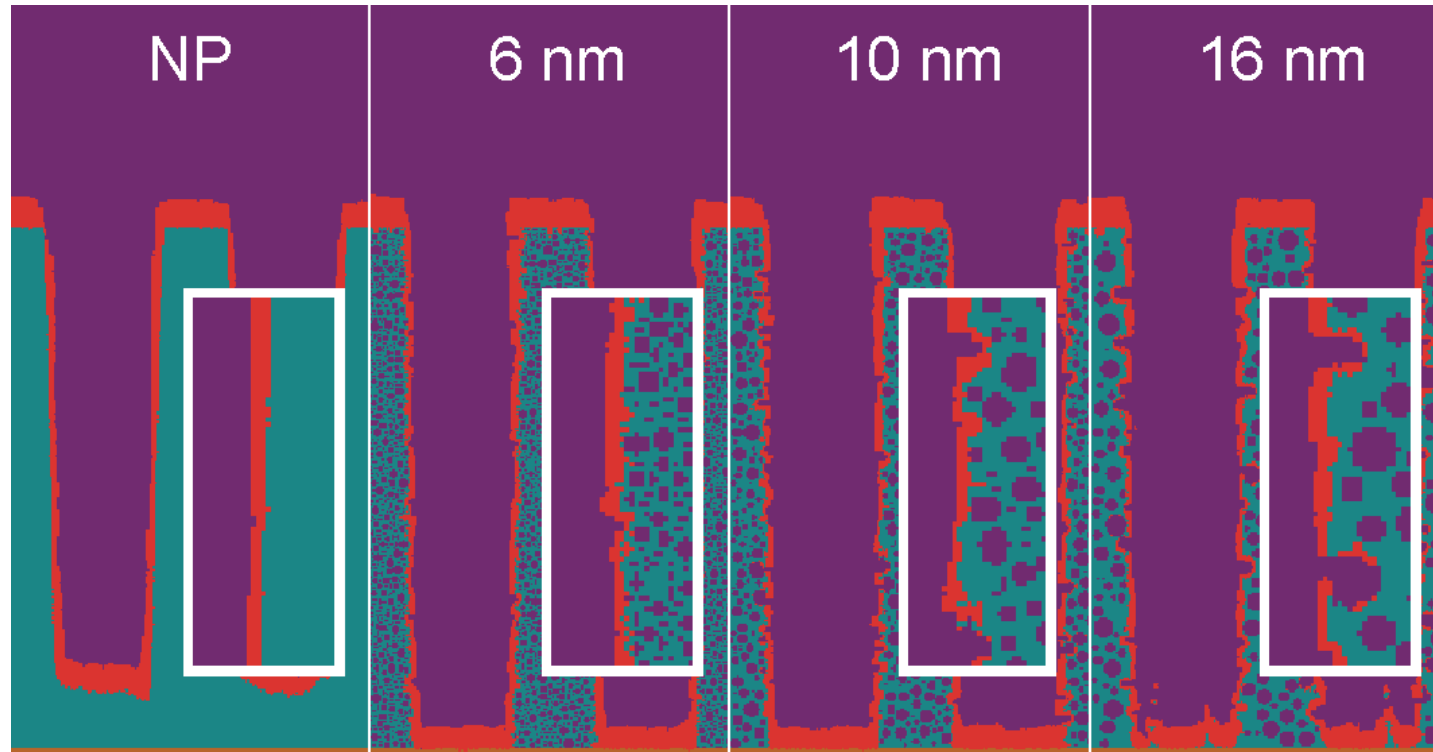
- More rapid etching with porous SiO₂ results in less mask erosion and better profile control, but more polymer filling of pores.

ICP: POROUS SiO₂ AND PHOTORESIST CLEAN



- Longer cleaning times are required with more porous materials to remove polymer which is shaded from ion flux.

IMPVD: Cu SEED LAYER DEPOSITION



- Thicker seed layers are required with large pores to cover over (or fill) gaps resulting from open structures.

CONCLUDING REMARKS

- **Integrated plasma process modeling requires addressing a wide range of physical phenomena.**
- **The large variety of gas mixtures, reactor geometries, plasma sources and materials motivates development of generalized modeling platforms with few a priori assumptions.**
- **The fundamental modeling challenges are no different than in experimental integration:**
 - **If a single module (process) is validated (optimized) in isolation, will it still be valid (optimum) when integrated with other steps?**

ACKNOWLEDGEMENTS

- **Dr. Alex V. Vasenkov**
- **Dr. Gottlieb Oherlein**
- **Mr. Arvind Sankaran**
- **Mr. Pramod Subramonium**

- **Funding Agencies:**
 - **3M Corporation**
 - **Semiconductor Research Corporation**
 - **National Science Foundation**
 - **SEMATECH**
 - **CFDRC Inc.**