# Motion Analysis of a Foldable Barrel Vault Based on Regular and Irregular Yoshimura Origami 

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#### Abstract

This paper investigates the geometry of a foldable barrel vault with Yoshimura Origami patterns during the motion. On the base of the geometry analysis of the origami unit, the radius, span, rise, and longitudinal length of the foldable barrel vault with regular Yoshimura Origami pattern in all configurations throughout the motion are determined. The results show that the radius of curvature and the span increase during deployment. But the rise increases first, followed by a decrease with increasing fold angle. Furthermore, the influence of the apex angle of the origami unit and the numbers of triangular


[^0]plates in the span direction on the geometric parameters is also investigated. Finally, the method to obtain the rise and span of the barrel vault with irregular origami pattern is also given.

## INTRODUCTION

Deployable structures have the ability to transform themselves from a small, closed or stowed configuration to a much larger, open or deployed configuration [1, 2]. Large aerospace structures, e.g. antennas and masts, are prime examples of deployable structures. Due to their size, they often need to be packaged for transportation and expanded at the time of operation [3]. On the other hand, temporary accommodation, or specially tents, is also one of the applications of deployable structures [4]. The deployment of most of such structures relies on the concept of mechanisms, i.e. the structures are assemblies of mechanisms whose mobility is retained for the purpose of deployment.

Origami, the Japanese and Chinese traditional paper craft, has been proved as a valuable tool to develop various engineering applications in numerous fields [5-7]. Several designs of origami structures have been proposed for deployable structures from around 1970s, such as the Miura-ori, which is a well-known rigid origami structure utilized in the packaging of deployable solar panels for use in space or in the folding of maps [8]. Cai et al. [9] have investigated the geometry and motion of a deployable shelter structures based on Miura-ori. Moreover, the geometry of a foldable barrel vault with modified Miura-ori patterns, which displays a curvature change during the motion, was also studied by Cai et al. [10].

The most famous origami pattern is the Yoshimura or diamond pattern [11-13]. The best way to understand the background of this problem is to consider the case of post-
buckling behavior of a thin cylindrical shell under axial loading. A family of foldable, portable structures, based on the Yoshimura origami pattern, has been presented by Foster and Krishnakumar [14]. These structures have considerable shape flexibility (multiple degrees of freedom), but once erected they possess significant stiffness. These foldable plate structures consist of a series of triangular plates, connected at their edges by continuous joints, allowing each plate to rotate relative to its neighboring plate. The plates can fold into a flat stack and unfold into a predetermined three-dimensional configuration, with a corrugated surface. Tonon $[15,16]$ has further expanded this concept to study the geometry of single and double curvature foldable plate structures, such as domes, conics, paraboloids and hyperboloids. Although they are compactly foldable into a compact stack of plates, some fold patterns cannot be developed in a plane, as opposed to the configurations previously given by Foster.

Based on the Yoshimura origami pattern, the concept for a mobile shelter was proposed by De Temmerman et al. [17]. The folding pattern and the associated design parameters were also discussed, along with their influence on the geometry behavior of the system. A historical review of origami-like deployable shelters developed by the US military has given by Thrall and Quaglia [18]. Moreover, they also proposed a novel erection strategy for origami-inspired deployable structures [19]. However, most of these researches mainly focus on the geometry of the deployable configuration. The geometry of the system during the motion is rarely studied. In this paper, the motion of the foldable barrel vault based on the regular and irregular Yoshimura origami pattern will be studied. The folding process of the barrel vault with regular Yoshimura origami pattern is shown in Fig.1. Then the radius, span, rise, and longitudinal length of the foldable barrel vault in
all configurations throughout the motion, which are very important to the design of the system, are obtained.

## GEOMETRY ANALYSIS OF A YOSHIMURA ORIGAMI UNIT

## Geometry description

Figure 2 shows a Yoshimura origami unit, the basic origami pattern with six creases, where the continuous lines indicate mountain folds while the dotted lines indicate valley folds. Figure 2 (a) demonstrates the fully unfolded configuration which is also the planar state of the basic unit and Fig. 2(b) demonstrates the partly foldable configuration.

Fig. 2 Basic origami unit with six creases
Angle $\beta$ between the mountain fold line and the valley fold line is defined as the apex angle of the basic origami unit. Define angle $\theta$ as the fold angle between the rigid plate with shade lines of the unit in Fig. 2(b) and the vertical plane, which is $0^{\circ}$ in the fully foldable state and $90^{\circ}$ in the fully unfolded state. It can be seen that the apex angle of the unit is a constant during the deployment of the origami, but its projections on the horizontal and vertical planes are changing.

Projection of the apex angle
The projection of the apex angle on the horizontal and vertical planes during the motion is shown in Fig. 3. The projection of the apex angle on the vertical plane is $\beta^{\prime}$, and on the horizontal plane is $\beta^{\prime \prime}$.

In the right triangle $\triangle \mathrm{ABC}$ of Fig.3, there is

$$
\begin{equation*}
H=L \tan \beta \tag{1}
\end{equation*}
$$

In the right triangle $\triangle \mathrm{OAB}$, there is

$$
\begin{equation*}
H_{1}=H \cos \theta \tag{2}
\end{equation*}
$$

In the right triangle $\triangle \mathrm{OAD}$, there is

$$
\begin{equation*}
H_{1}=L \tan \beta^{\prime} \tag{3}
\end{equation*}
$$

Substituting Eqs. (1) and (2) into Eq. (3) leads to the projection of the apex angle on the vertical plane

$$
\begin{equation*}
\tan \beta^{\prime}=\tan \beta \cos \theta \tag{4}
\end{equation*}
$$

Similarly, in the right triangle $\triangle \mathrm{OAB}$, there is

$$
\begin{equation*}
H_{2}=H \sin \theta \tag{5}
\end{equation*}
$$

In the right triangle $\triangle \mathrm{OBC}$, there is

$$
\begin{equation*}
H_{2}=L \tan \beta^{\prime \prime} \tag{6}
\end{equation*}
$$

Substituting Eqs. (1) and (5) into Eq. (6) leads to the projection of the apex angle on the horizontal plane is:

$$
\begin{equation*}
\tan \beta^{\prime \prime}=\tan \beta \sin \theta \tag{7}
\end{equation*}
$$

## MOTION ANALYSIS OF THE FOLDABLE BARREL VAULT

Foldable barrel vault
The origami pattern for the foldable barrel vault, in which the basic origami units are arranged separately in the longitudinal and span direction, is given in Fig.4. The basic units of this pattern are marked by red rectangles in the figure. The structure can also be seen as composed of a series of triangular plates connected by continuous joints, allowing each plate to rotate around the intersection lines. Assume that the number of triangular plates in the span direction is $p$, the column number in the longitudinal direction is $m$. Therefore, in Fig.4, $p$ is 7 and $m$ is 4.

Apex angles of the basic units in Fig. 4 separately are $\beta_{1}, \beta_{2}, \beta_{3}, \beta_{4}, \beta_{5}$ and $\beta_{6}$, respectively. These apex angles can be equal or unequal. This section will discuss the
situation that all origami unit with equal apex angles while the unequal apex angles situation will be considered in the next section.

Analysis of the movement process
Figure 5 shows the top view of the barrel vault in the fully unfolded state and the section view in the partly foldable state. Fig. 5(b) reveals the vertical projections of the apex angle of the basic origami units, where the connecting line of vertices of each triangle plate is a circular arc. The central angle of this arc is

$$
\begin{equation*}
\alpha=2(p-1) \beta^{\prime} \tag{8}
\end{equation*}
$$

The length of its radius is

$$
\begin{equation*}
R=\frac{L}{\sin \left(2 \beta^{\prime}\right)} \tag{9}
\end{equation*}
$$

where $L$ is length of the basic triangle plate.
In Fig. 5(b), the span of the barrel vault during the motion can be obtained as

$$
\begin{equation*}
S=2 R \sin \frac{\alpha}{2} \tag{10}
\end{equation*}
$$

Substituting Eqs. (8) and (9) into Eq. (10) leads to

$$
\begin{equation*}
S=2 \frac{L}{\sin \left(2 \beta^{\prime}\right)} \sin \left[(p-1) \beta^{\prime}\right] \tag{11}
\end{equation*}
$$

And the rise of the barrel vault during the motion is

$$
\begin{align*}
H & =R-R \cos \frac{\alpha}{2} \\
& =\frac{L}{\sin \left(2 \beta^{\prime}\right)}\left\{1-\cos \left[(p-1) \beta^{\prime}\right]\right\} \tag{12}
\end{align*}
$$

The width of the basic units in the longitudinal direction is

$$
\begin{equation*}
w=L \tan \beta^{\prime \prime}=L \tan \beta \sin \theta \tag{13}
\end{equation*}
$$

Therefore, the width of the barrel vault in the longitudinal direction during the motion is

$$
\begin{equation*}
W=m w=m L \tan \beta \sin \theta \tag{14}
\end{equation*}
$$

Now the basis geometry parameter of the structure is assumed as $p=7, m=4, \beta=30^{\circ}$. During the motion, the fold angle $\theta$ is an independent variable. Therefore, the radius, span, rise and width of the structure during the motion can be investigated as follows.

The relationships between the radius and the fold angle $\theta$ are given in Fig. 6. From this figure, it can be found that with the continuous deployment of the model and the value of $\theta$ increasing from $0^{\circ}$ to $90^{\circ}$, the radius increases continuously from the initial value $R_{0}$ and trends to infinity when the structure approaches to the fully unfolded configuration. Moreover, when the fold angle $\theta$ is small, the radius changes slightly. But when $\theta$ is close to $90^{\circ}$, the radius increases exponentially. With Eqs. (4) and (9), the initial value of the radius can be obtained as

$$
\begin{equation*}
R_{0}=\frac{2 L}{\sqrt{3}} \tag{15}
\end{equation*}
$$

Figure 7 shows the relation between the span and the fold angle $\theta$. When the motion from the folded configuration to the deployment configuration, the span $S$ increases continuously and approaches to $S_{0}=(p-1) L$ until the fully unfolded state of the system.

Due to the case of $p=7$, the span in the fully deployable state is $6 L$.
As shown in Fig. 8, with the increase of the fold angle $\theta$, the rise of the barrel vault experiences two stages of variations, i.e., first increase to the maximum and then decrease to 0 . In the first stage, the rise increases to the maximum and has a small range of variation. In the second stage, the rise decreases continuously from the maximum vault to zero, which is corresponding to the fully unfolded configuration.

Figure 9 shows the longitudinal length of the barrel vault during the deployment. In the fully foldable state, the initial value of the longitudinal length is 0 , and then changes linearly. When $\theta$ is close to $90^{\circ}$, the variation rate of the longitudinal length becomes extremely slow and the value approaches to $m L \tan \beta$ in the fully deployable state.

## PARAMETRIC ANALYSIS

The influence of the apex angle $\beta$ of the origami unit and the number of triangular plates $p$ in the span direction on the radius, span, rise, longitudinal length of the barrel vault will be studied in this section.

Influence of the apex angle
The apex angle is an important factor that affects the kinematic behavior of systems. The radius, span, rise and the longitudinal length during the motion with different apex angles are shown in Figs. 10 to13 where four apex angles, $10^{\circ}, 20^{\circ}, 30^{\circ}$ and $40^{\circ}$, are considered in this study.

As shown in Fig.10, with different apex angles $\beta$, the relationships between the radius and the fold angle $\theta$ are similar. The radiuses all increase continuously with the deployment of the system. It can also be found from Fig. 10 that the radius decreases with the increase of the apex angle $\beta$ in the fully foldable state. Moreover, during the motion with the same fold angle $\theta$, the radius decreases with the increase of $\beta$ as well.

Figure 11 shows four curves of span with different apex angles. From this figure, the initial span of the barrel vault in the fully foldable state decreases continuously with the increase of the apex angle. But for the barrel vault with different apex angles, the spans in the fully deployable state are the same. It needs to point out that the value of span may be negative for the apex angle $\beta$ is $40^{\circ}$ when the fold angle is small, which means that the origami units "penetrate" through each other. However, penetration is generally not
allowable in the physical structures, which leads to a certain limit in the range of fold angle $\theta$. From the figure, the four curves present the same change rules with the deployment of the system, but with the increase of the apex angle $\beta$, the variation amplitude of span is more obvious.

The effect of the apex angle $\beta$ on the rise is shown in Fig.12. When $\beta$ is $30^{\circ}$ or $40^{\circ}$, the rise of the model experiences the process of increasing first and then decreasing to 0 with the increase of fold angle $\theta$. When $\beta$ is $10^{\circ}$ or $20^{\circ}$, the rise decreases continuously with the deployment of the model. And the rises of all models are 0 in the fully deployable state. The rises are $1.46 L, 2.33 L, 2.31 L, 1.52 L$ corresponding to the apex angles are $10^{\circ}, 20^{\circ}, 30^{\circ}$ and $40^{\circ}$ in the fully foldable state. It should be noted that when $\beta$ is $40^{\circ}$, the fold angle $\theta$ can not be 0 which also means insignificance of the corresponding rise. When $\beta$ is $30^{\circ}$, in the fully foldable state, the central angle is just $360^{\circ}$ and the rise is $2 R$ or 2.31 L .

Figure 13 shows the effect of apex angles $\beta$ on the longitudinal length. As shown in the figure, the variation trends of all curves are similar. The longitudinal lengths increase continuously with the increase of the fold angle $\theta$. The longitudinal lengths of all models are 0 in the fully folded state and are $m L \tan \beta$ in the fully deployable state. Influence of the number of triangular plates

The geometry parameters of the barrel vault with different numbers of triangular plates during the motion are shown in Figs. 14-17. The numbers of triangular plates correspond to 5, 7, 9 and 11, respectively.

Figure 14 shows the relations between radius and the fold angle $\theta$ with different numbers of triangular plates $p$. From this figure, the number of triangular plates $p$ has no effect on the radius $R$ of the system. The relations between the span $S$ and the fold angle $\theta$
with different numbers of triangular plates $p$ are shown in Fig. 15. It can be seen that the four curves present the same change trends with the deployment of the system. The span increases with the increase of the fold angle. Moreover, with the increase of the number of triangular plates $p$, the change in range of the curves is more obvious. It should be noted that when the fold angle is small, the value of span may be negative for the system with $p=9$ or 11. It means that the units "penetrate" through each other, which leads to a limit deployment range of a physical structure.

The effect of the number of triangular plates $p$ on the rise is shown in Fig. 16.When numbers of triangular plates $p$ are 7,9 and 11 , the rises of the barrel vault experience the process of increasing first and then decreasing to 0 with the increase of fold angle $\theta$. When $p$ is 5 , the rise decreases continuously with the deployment of the barrel. And the rises of all models are 0 in the fully deployable state. The rises are $1.73 L, 2.31 L, 1.73 L$, $0.58 L$ corresponding to the number of triangular plates $p$ are 5, 7, 9 and 11 in the fully foldable state. It needs to point out that when $p$ is 9 or 11 , the fold angle $\theta$ cannot be 0 which also means insignificance of the corresponding rise. Figure 17 shows the relation between longitudinal length and the fold angle $\theta$ with different numbers of triangular plates $p$. It can be found that the number of triangular plates $p$ has no effect on the longitudinal length $w$ of the system.

## MOVEMENT BOUNDARY

From the previous section, the central angles of some barrel vaults are greater than $360^{\circ}$ when these systems have a small fold angle $\theta$, which means that the origami units will cross over each other. However, this will never happen for a physical model. Therefore, when the central angle is $360^{\circ}$, the barrel vault cannot be able to be further folded. Then the motion boundary of the barrel vault, i.e. the critical value of the fold
angle $\theta$ when the model cannot be folded, will be studied. From Eq. (8), the central angle is $360^{\circ}$ which can be given as

$$
\begin{equation*}
2(p-1) \beta^{\prime}=2 \pi \tag{16}
\end{equation*}
$$

which leads to

$$
\begin{equation*}
\beta^{\prime}=\frac{\pi}{p-1} \tag{17}
\end{equation*}
$$

Substituting Eq. (17) into Eq. (4) leads to

$$
\begin{equation*}
\tan \left(\frac{\pi}{p-1}\right)=\tan \beta \cos \theta \tag{18}
\end{equation*}
$$

Therefore, the critical value of the fold angle $\theta$ is

$$
\begin{equation*}
\theta_{l}=\arccos \frac{\tan \left(\frac{\pi}{p-1}\right)}{\tan \beta} \tag{19}
\end{equation*}
$$

Now considering a special case, when the system is fully folded or the critical value of the fold angle $\theta$ is 0 , the central angle is $360^{\circ}$. Substituting $\theta=0$ into Eq. (19) to obtain the apex angle as

$$
\begin{equation*}
\beta=\frac{\pi}{p-1} \tag{20}
\end{equation*}
$$

It can be found from Eq. (20), the limit apex angle is only related to number of triangular plates $p$. Table 1 shows the corresponding relationship between $p$ and the apex angles $\beta$. For a barrel vault with a definitive value of $p$, when the apex angle $\beta$ of the origami unit is less than the value given in Eq.(20) or Table 1, the central angle of the barrel vault is always less than $360^{\circ}$. In other words, the folding and deployment of the system will not be restricted. However, when $\beta$ is greater than the value given in Eq. (20) or Table 1, the barrel vault cannot be fully folded. The critical value of the fold angle $\theta$
can be calculated with Eq. (19). Figure 18 shows the relationship between the critical value of the fold angle and the apex angle $\beta$. It can be seen from this figure that the critical value of the fold angle increases with the increase of the apex angle. That is to say the foldable range of the barrel vault also decreases with the increase of the apex angle. Moreover, when the apex angle of the basic unit is close to the value given in Eq. (20), the critical value of the fold angle $\theta$ changes significantly.

## MOTION ANALYSIS OF THE BARREL VAULT WITH IRREGULAR PATTERNS

The structures discussed above are all based on the regular Yoshimura origami pattern, which have consistent geometric parameters of every origami unit. In Section 2, we have mentioned that the connecting line of vertices of triangle plates is a circular arc in the partly deployable state as shown in Fig. 5(b). If the connecting line of these vertices constitutes an ellipse or parabola, the system has a larger interior space with the same rise of the system.

Figure 19 shows an origami pattern with irregular units. But the system also has symmetry of the central line. The connecting line of vertices of triangle plates is an ellipse in the partly deployable state. The folding process of this foldable plate structure is shown in Fig. 20. A graphing method is used to obtain the side projection of the barrel vault during the motion. The steps are shown in Fig. 21. The side projections from the central Yoshimura origami unit to the boundary Yoshimura origami unit are given in from Fig. 21 (a) to Fig. 21(d), where the rectangular with red color shows the basic Yoshimura origami unit. Finally, the side projection of the whole barrel vault is given in Fig. 21(e).

The rise and span of the barrel vault with irregular origami units is given in Fig. 22.
In triangle $\triangle \mathrm{ABC}$ of Fig. 22, we have

$$
\begin{equation*}
\frac{A C}{\sin \angle A B C}=\frac{A B}{\sin \angle A C B}=\frac{B C}{\sin \angle B A C} \tag{21}
\end{equation*}
$$

As shown in Fig. 21, the length of AC is $L_{1}+L_{2}$, thus the lengths of AB and BC are

$$
\left\{\begin{array}{l}
A B=\frac{\left(L_{1}+L_{2}\right) \sin \left(\beta_{2}{ }^{\prime}\right)}{\sin \left(\beta_{1}{ }^{\prime}+\beta_{2}{ }^{\prime}\right)}  \tag{22}\\
B C=\frac{\left(L_{1}+L_{2}\right) \sin \left(\beta_{1}{ }^{\prime}\right)}{\sin \left(\beta_{1}{ }^{\prime}+\beta_{2}{ }^{\prime}\right)}
\end{array}\right.
$$

In the isosceles triangle $\Delta \mathrm{O}_{1} \mathrm{AB}$, the length of $R_{1}$ is

$$
\begin{equation*}
R_{1}=O_{1} A=\frac{A B}{2 \sin \left(\beta_{1}{ }^{\prime}\right)}=\frac{\left(L_{1}+L_{2}\right) \sin \left(\beta_{2}{ }^{\prime}\right)}{2 \sin \left(\beta_{1}{ }^{\prime}+\beta_{2}{ }^{\prime}\right) \sin \left(\beta_{1}{ }^{\prime}\right)} \tag{23}
\end{equation*}
$$

In the isosceles triangle $\Delta \mathrm{O}_{2} \mathrm{BC}$, the length of $R_{2}$ is

$$
\begin{equation*}
R_{2}=O_{2} B=\frac{B C}{2 \sin \left(\beta_{2}{ }^{\prime}\right)}=\frac{\left(L_{1}+L_{2}\right) \sin \left(\beta_{1}{ }^{\prime}\right)}{2 \sin \left(\beta_{1}{ }^{\prime}+\beta_{2}{ }^{\prime}\right) \sin \left(\beta_{2}{ }^{\prime}\right)} \tag{24}
\end{equation*}
$$

In $\triangle \mathrm{BCD}$, we have

$$
\begin{equation*}
\frac{C D}{\sin \angle D B C}=\frac{B D}{\sin \angle D C B} \tag{25}
\end{equation*}
$$

As shown in Fig. 21, the length of BD is $L_{2}+L_{3}$, thus the length of CD is

$$
\begin{equation*}
C D=\frac{\left(L_{2}+L_{3}\right) \sin \left(\beta_{2}{ }^{\prime}\right)}{\sin \left(\beta_{2}{ }^{\prime}+\beta_{3}{ }^{\prime}\right)} \tag{26}
\end{equation*}
$$

In the isosceles triangle $\Delta \mathrm{O}_{3} \mathrm{CD}$, the length of $R_{3}$ is

$$
\begin{equation*}
R_{3}=O_{3} C=\frac{C D}{2 \sin \left(\beta_{3}^{\prime}\right)}=\frac{\left(L_{2}+L_{3}\right) \sin \left(\beta_{2}{ }^{\prime}\right)}{2 \sin \left(\beta_{2}{ }^{\prime}+\beta_{3}{ }^{\prime}\right) \sin \left(\beta_{3}{ }^{\prime}\right)} \tag{27}
\end{equation*}
$$

As shown in Fig.22, when the lengths of $R_{1}, R_{2}$ and $R_{3}$ are obtained, the rise $H$ and the span $S$ can be easily calculated as

$$
\begin{gather*}
H=R_{3}-\left\{\left(R_{3}-R_{2}\right) \cos \left(2 \beta_{3}{ }^{\prime}\right)+\left(R_{2}-R_{1}\right) \cos \left(2 \beta_{3}{ }^{\prime}+2 \beta_{2}{ }^{\prime}\right)+R_{1} \cos \left(2 \beta_{3}{ }^{\prime}+2 \beta_{2}{ }^{\prime}+2 \beta_{1}{ }^{\prime}\right)\right\}  \tag{28}\\
S=2\left\{\left(R_{3}-R_{2}\right) \sin \left(2 \beta_{3}{ }^{\prime}\right)+\left(R_{2}-R_{1}\right) \sin \left(2 \beta_{3}{ }^{\prime}+2 \beta_{2}{ }^{\prime}\right)+R_{1} \sin \left(2 \beta_{3}{ }^{\prime}+2 \beta_{2}{ }^{\prime}+2 \beta_{1}{ }^{\prime}\right)\right\} \tag{29}
\end{gather*}
$$

It should be noted that Eqs. (28) and (29) can only be applied to the symmetric patterns with $p=7$, whose connecting lines of these vertices of the triangle plates constitutes an ellipse. However, the graphing method given in this paper can also be used to study the shape or geometric parameters of the foldable structures with other irregular origami patterns.

## CONCLUSIONS

The geometric parameters of a foldable barrel vault based on regular and nonregular Yoshimura origami patterns during the motion are studied in this paper. Normally, in the design of a foldable barrel vault, the designer will give the span and rise of the barrel vault when the system is under the external loads. Moreover, for a foldable structure, the volume in the fully folded configuration is also very important. Then when these parameters are known, the method given in this paper can be used to study the number of the triangle plates and the geometry of a single triangle plate. Furthermore, the deployment angle of the state when the system is under the external loads (such as the dead load, wind load or snow load etc.) can also be found. From the analysis, we can draw the following conclusions.
(1) With the deployment of the model, the radius increases gradually and at the beginning shows little change but then increases exponentially. The slope of the span increase is relatively flat at the early and final stage and shows approximate linear growth in the medium stage. With the increase of the fold angle, the rise of the model experiences two stages of variation, i.e., first increase to the maximum and then decrease to 0 . The longitudinal length increases gradually to its maximum with the deployment of the model.
(2) With the increase of the apex angle $\beta$ of the basic unit, the span decreases continuously; on the contrary, the rise changes in the opposite direction, i.e., the rise increases obviously with the increase of the apex angle. And the fold angle corresponding to the maximum of the rise also increases with the increase of $\beta$.
(3) With the increase of the numbers of triangular plates $p$, the amplitude of the span increases sharply, but the change range of the fold angle decreases. The rise increases obviously with the increase of the number of triangular plates $p$ and the fold angle corresponding to the maximum of the rise with the increase of $p$.
(4) The method to obtain the rise and span of the barrel vault with irregular origami pattern is also given.

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## Figure Captions List

Fig. 1 The folding process of a barrel vault

Fig. 2 Basic origami unit with six creases (a) Fully unfolded configuration (b)
Partially foldable configuration

Fig. 3 Projection of the apex angle

Fig. $4 \quad$ Origami pattern for a barrel vault

Fig. 5
A barrel vault (a) Top view in the fully unfolded state (b) Cross section in the partially foldable state

Fig. 6 Radius during the motion

Fig. $7 \quad$ Span during the motion

Fig. $8 \quad$ Variation of rise during motion

Fig. $9 \quad$ Variation of longitudinal length during motion

Fig. 10 Radius with different apex angles

Fig. 11 Span with different apex angles

Fig. 12 Rise with different apex angles

Fig. 13 Longitudinal length with different apex angles

Fig. $14 \quad$ Radius with different numbers of triangular plates $p$

Fig. $15 \quad$ Span with different numbers of triangular plates $p$

Fig. $16 \quad$ Rise with different numbers of triangular plates $p$

Fig. 17 Longitudinal length with different numbers of triangular plates $p$

Fig. 18 Relationship between critical value of the fold angle and apex angle $\beta$

Fig. 19 Origami pattern with irregular units

Fig. 20 The folding process of a barrel vault with irregular patterns

Fig. 21 Side projection of the barrel vault during the motion

Fig. 22 Geometric parameters of the barrel vault with irregular origami units

## Table Caption List

Table $1 \quad$ Relationship between number of triangular plates $p$ and apex angle $\beta$


Fig. 1 The folding process of a barrel vault

(a) Fully unfolded configuration (b) Partially foldable configuration Fig. 2 Basic origami unit with six creases


Fig. 3 Projection of the apex angle


Fig. 4 Origami pattern for a barrel vault

(a) Top view in the fully unfolded state (b) Cross section in the partially foldable state Fig. 5 A barrel vault


Fig. 6 Radius during the motion


Fig. 7 Span during the motion


Fig. 8 Variation of rise during motion Fig. 9 Variation of longitudinal length during motion


Fig. 10 Radius with different apex angles



Fig. 12 Rise with different apex angles
Fig. 13 Longitudinal length with different apex angles


Fig. 14 Radius with different numbers of triangular plates $p$


Fig. 15 Span with different numbers of triangular plates $p$


Fig. 16 Rise with different numbers of triangular plates $p$


Fig. 17 Longitudinal length with different numbers of triangular plates $p$


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Fig. 21 Side projection of the barrel vault during the motion


Fig. 22 Geometric parameters of the barrel vault with irregular origami units

Table 1 Relationship between number of triangular plates $p$ and apex angle $\beta$
number of triangular plates $p \quad$ apex angles $\beta\left({ }^{\circ}\right)$
$5 \quad 45$
7
30
9
11
13
22.5

18
15


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