PHY2053 Lecture 17

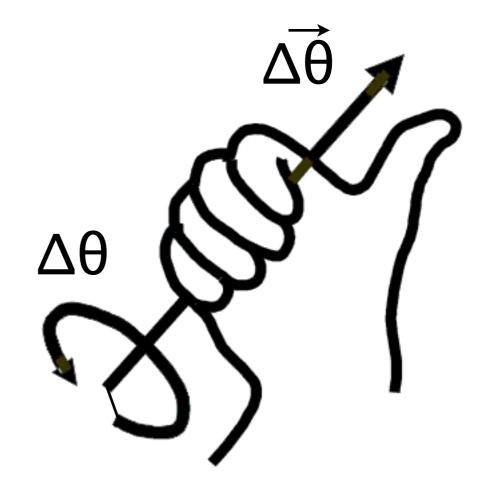
Ch. 8.4, 8.6 (skip 8.5) Rotational Equilibrium, Rotational Formulation of Newton's Laws

Linear vs Rotational Variables Transformation Table

Linear Variable	→ Rotational Variable
distance	angle
velocity	angular velocity
acceleration	angular acceleration
mass [inertia]	moment of inertia
momentum p = mv	angular momentum, L = Iω
force, F = ma	torque, $T = I\alpha$
$K_{\text{trans}} = \frac{1}{2} \text{mv}^2$	$K_{\text{rotation}} = \frac{1}{2} \omega^2 $

"Right Hand Rule" for rotational systems

- change of angle $(\Delta\theta)$ is a vector along the axis of rotation, and its direction is determined by the right hand rule:
- orient your right hand so that the fingers point in the direction in which the angle is changing; your thumb points in the direction of the rotational vector



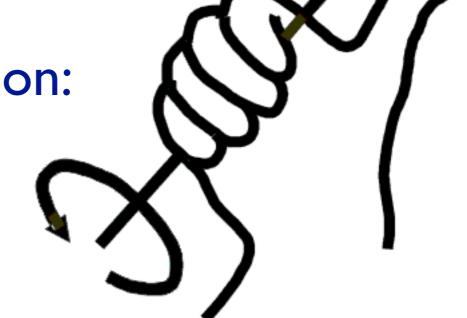
Reminder: from angle to angular acceleration

instantaneous angular velocity:

$$\vec{\omega} = \lim_{\Delta t \to 0} \frac{\Delta \vec{\theta}}{\Delta t}$$

• instantaneous angular acceleration:

$$\vec{\alpha} = \lim_{\Delta t \to 0} \frac{\Delta \vec{\omega}}{\Delta t}$$



Torque

 The equivalent of force, for rotational systems (somewhat hand-waving derivation of equivalence)

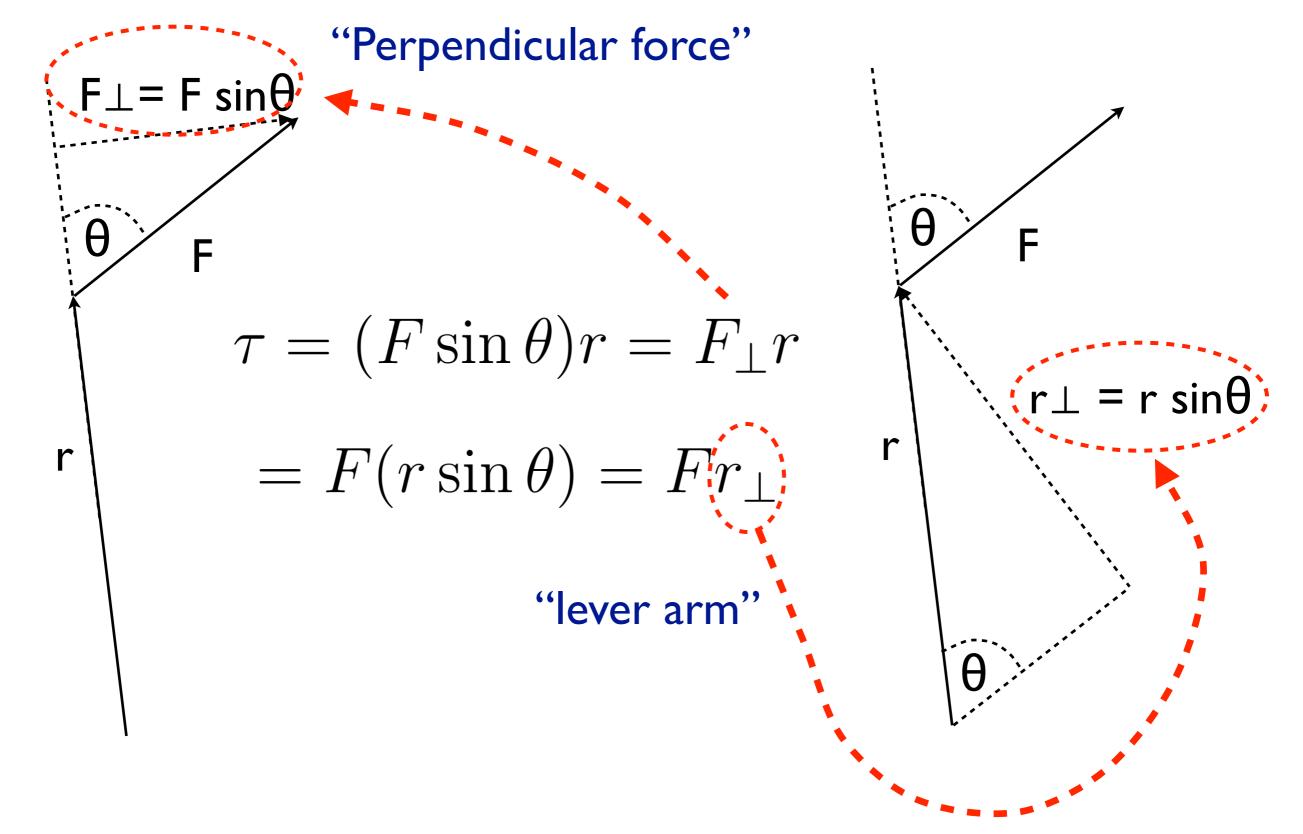
$$F_{i} = m_{i}a_{i} | r_{i}, \sum_{i} \rightarrow \sum_{i} r_{i}F_{i} = \sum_{i} m_{i}r_{i}a_{i}$$

$$[a_{i} = \alpha r_{i}] \sum_{i} r_{i}F_{i} = \sum_{i} m_{i}r_{i}^{2}\alpha$$

torque = momentum of inertia × angular acceleration

 not quite so trivial, need to account for the orientation of force wrt "lever arm" of object i, "useful force"

PHY2053, Lecture 17 Rotational Form of Newton,
$$au_i = r_i F_i \sin(\alpha_i)$$



• for the same force, the torque is directly proportional to the lever arm (projection of r perpendicular to F)

Work done by torque

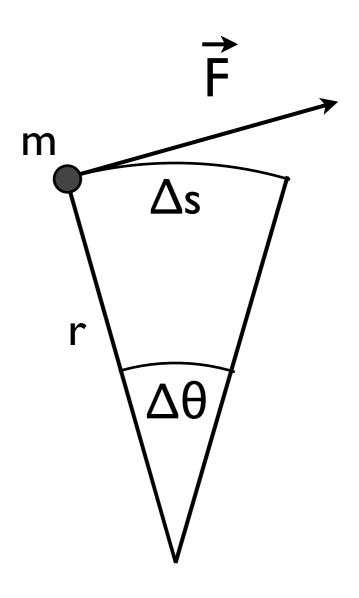
• Start from definition of work, for rotational systems:

$$W = F\Delta s; \quad [\Delta s = r\Delta\theta]$$

$$W = Fr\Delta\theta = \tau\Delta\theta$$

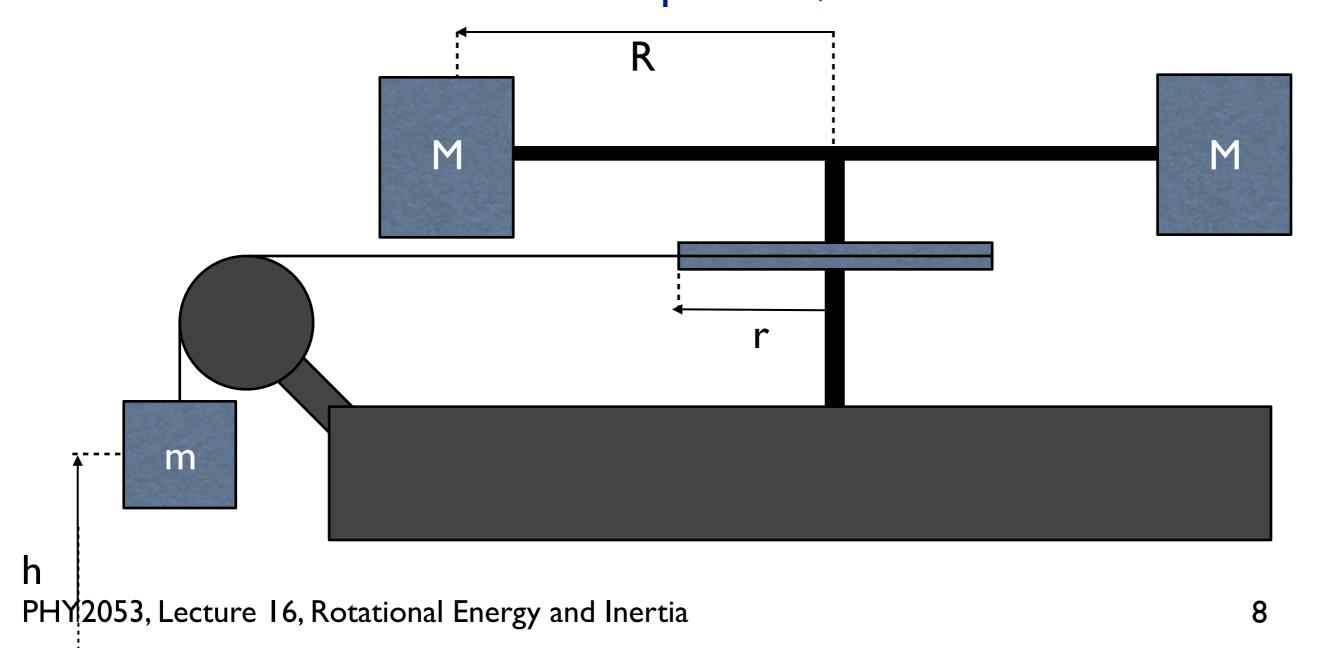
 Hooke's law, for rotational systems (torsion spring):

$$\tau = -\kappa \Delta \theta$$



Unusual Atwood's Machine

In the machine below, the small weight (mass m) is connected via an ideal pulley to the massless horizontal wheel of radius r. Ignore the moments of inertia of the connecting rods; the two masses M are both R away from the axis of rotation. What is the velocity of mass m as it falls h from its initial position, where it was at rest?



Discussion: Unusual Atwood's Machine,

via energy conservation

Utilize
$$E_i = E_f$$
 $E_i = k_i + U_i \rightarrow U_i = mgh$

$$\downarrow from rest, k_i = 0$$

$$E_f = U_f + k_c = 0 + \frac{mv^2}{2} + \frac{Tw^2}{2}, however v = \omega_i r.$$

$$= \frac{mv^2}{2} + \frac{MR^2}{2}, (\frac{v}{r})^2 \qquad \omega = \frac{v}{r} \qquad pulley$$

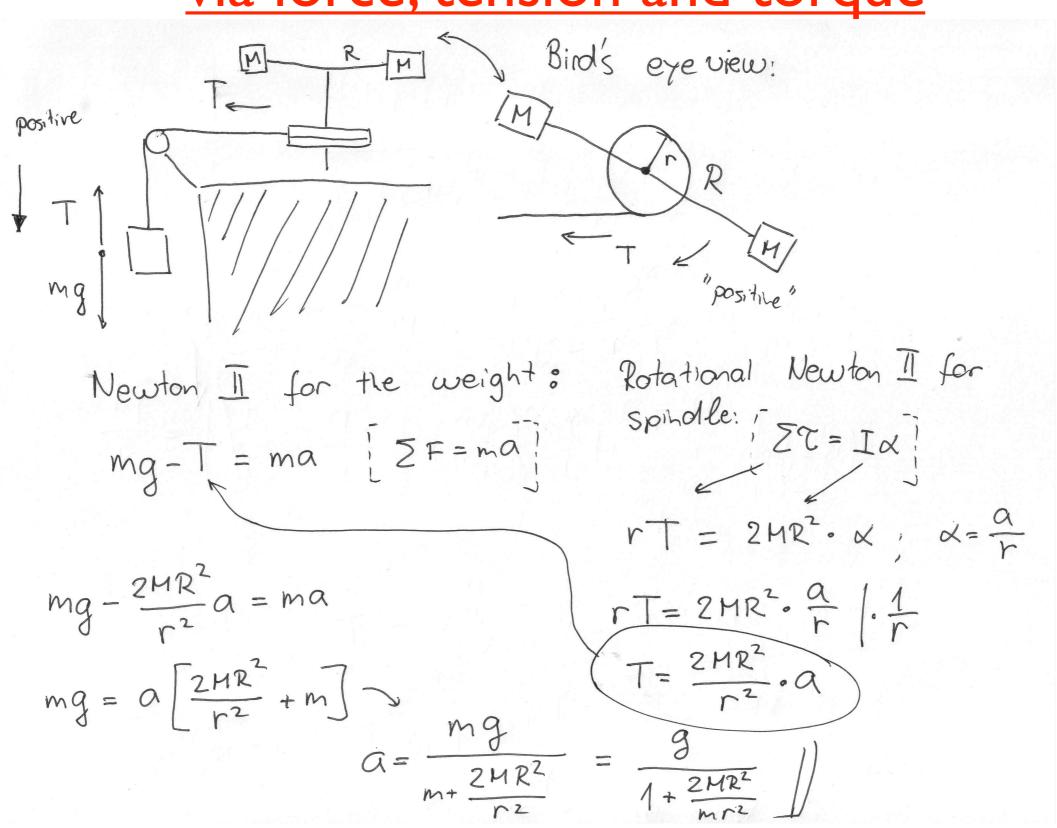
$$= \frac{mv^2}{2} + MR^2, \frac{v^2}{r^2} = \frac{mv^2}{2} + \frac{MR^2}{r^2} v^2 = v^2 \left[\frac{m}{2} + \frac{MR^2}{r^2} \right]$$

$$E_i = E_f \rightarrow U_i = k_f \rightarrow mgh = v^2 \left[\frac{m}{2} + \frac{MR^2}{r^2} \right]$$

$$mgh = mv^2 \left[\frac{1}{2} + \frac{MR^2}{mr^2} \right]$$

$$v^2 = \frac{gh}{\frac{1}{2} + \frac{MR^2}{mr^2}} \rightarrow v = \sqrt{\frac{gh}{\frac{1}{2} + \frac{MR^2}{mr^2}}}$$

Discussion: Unusual Atwood's Machine, via force, tension and torque



Discussion: Unusual Atwood's Machine, via force, tension and torque 2

We know that for constant acceleration,
$$v^2 = 2a\Delta x \rightarrow v_f^2 = 2ah$$

$$v_f^2 = 2h \cdot \frac{g}{1 + \frac{2MR^2}{mr^2}} = \frac{2gh}{1 + \frac{2MR^2}{mr^2}} \cdot \frac{\frac{1}{2}}{\frac{1}{2}} = \frac{gh}{\frac{1}{2} + \frac{MR^2}{mr^2}}$$

$$v_f = \left[\frac{gh}{\frac{1}{2} + \frac{MR^2}{mr^2}}\right] \quad \text{Same answer}$$
as when using energy conservation.

Statics, Rotational Equilibrium

 Condition I: Newton's 1st law: a stationary object that has no net force acting on it remains stationary

$$\sum_{i} \vec{F}_{i} = 0$$

 Condition 2: For the object not to start rotating about an axis, the net torque on the object must be zero

$$\sum_{i} \vec{\tau}_{i} = 0$$

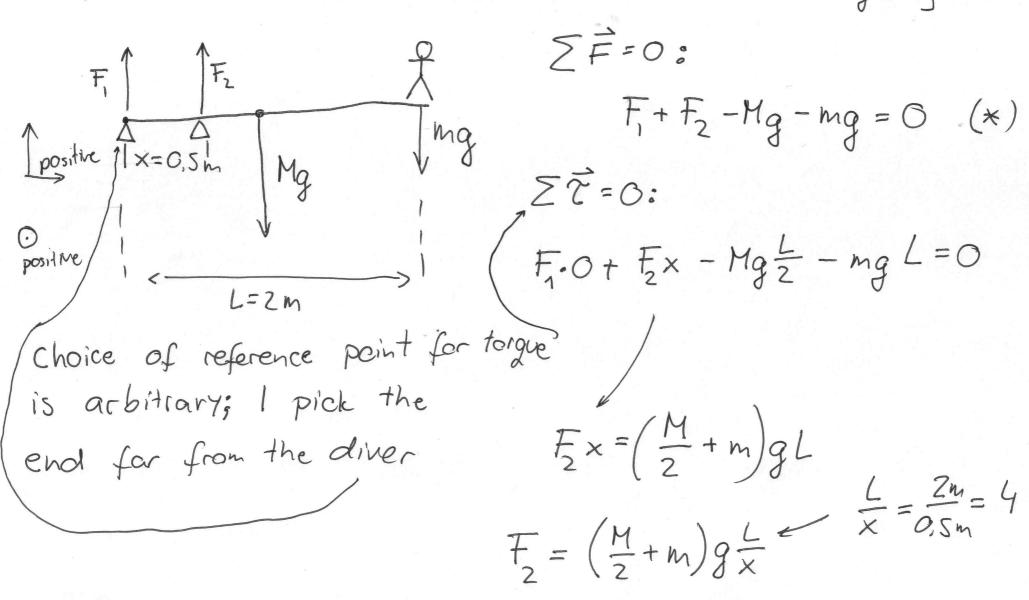
 Important: Conditions of static equilibrium, once established for one axis, apply to any arbitrary axis

Example: Springboard

• A diver weighing 80 kg is standing on the end of a 2 m long springboard which weighs 240 kg. The springboard is supported by two springs on the other end. One spring is at the very end of the board, and the other is 0.5 m toward the center of the board. Compute the forces generated by the springs if the system is in static equilibrium. Are the springs compressed or stretched?

Discussion: Springboard

First, we need a schematic with all the forces. I will make a guess about their directions, we will find out if it was correct [positive values mean correct guess]



Discussion: Springboard 2

$$F_2 = \left(\frac{240 \text{ kg}}{2} + 80 \text{ kg}\right) \cdot 9 \cdot 4 = 4 \cdot 2000 = 800 = 800 = 8 \text{ kg}$$

Fz is positive, I guessed the right direction. Spring 2 is therefore compressed.

From (*):
$$F_1 = -F_2 + Mg + mg = (-800g + 240g + 80g) kg$$

= $[-480 kg] \cdot g \simeq -4.8 kN /l$

F₁ is negative, I guessed the direction wrong: F₁ is a downward force, and spring 1 is therefore stretched.

Angular Momentum

Momentum-impulse theorem:

$$F_i \Delta t = m_i \Delta v_i \quad | r_i, \sum_i \rightarrow \sum_i r_i F_i \Delta t = \sum_i m_i r_i \Delta v_i; \quad [v_i = \omega r_i]$$
$$\sum_i \tau_i \Delta t = \sum_i r_i F_i \Delta t = \sum_i m_i r_i^2 \Delta \omega = I \Delta \omega$$

by analogy with p = mv, define angular momentum:

$$\vec{L} = I \vec{\omega}$$

• angular momentum-impulse theorem:

$$\vec{\tau}\Delta t = \Delta \vec{L} = I\Delta \vec{\omega}$$
 $\vec{\tau} = 0 \rightarrow \vec{L} = co\vec{n}st$

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Example: "Angular Collision"

- A disk of mass m_1 and radius R is spinning with an angular frequency ω_i . A second disk, of mass $3m_1$ is dropped onto disk I so that they share the same axis of rotation. After a brief period of friction, the two disks are rotating with the same angular frequency ω_f .
- What is the ratio ω_f/ω_i ?
- What fraction of the initial kinetic energy was wasted through friction?

Discussion: Angular Collision

The angular nomentum of the system is unchanged:
$$L_{i} = L_{f} \qquad L_{i} = I_{1}\omega_{i} \qquad L_{f} = \left[I_{1}+I_{2}\right]\omega_{f}$$

$$\left[I_{1}+I_{2}\right]\omega_{f} = I_{1}\omega_{i} \quad \Rightarrow \quad \omega_{f} = \omega_{i} \quad \frac{I_{1}}{I_{1}+I_{2}}$$

$$I_{1} = m_{1}R^{2}\cdot\frac{1}{2} \quad ; \quad I_{2} = m_{2}R^{2}\cdot\frac{1}{2} \quad ; \quad m_{2} = 3m_{1}$$

$$\omega_{f} = \omega_{i}\cdot\frac{m_{1}\cdot\frac{R^{2}}{2}}{m_{1}\cdot\frac{R^{2}}{2}} = \omega_{i}\cdot\frac{m_{1}}{m_{1}+m_{2}} = \omega_{i}\cdot\frac{m_{i}}{q_{m_{i}}} = \frac{\omega_{i}}{q_{m_{i}}}$$
 The initial and final kinetic energies are:
$$K_{i} = \frac{I_{1}}{2}\omega_{i}^{2} \qquad K_{f} = (I_{1}+I_{2})\cdot\frac{1}{2}\cdot\omega_{f}^{2}$$

Discussion: Angular Collision 2

$$\frac{K_f}{K_i} = \frac{\frac{1}{2} \left[I_1 + I_2 \right] \omega_f^2}{\frac{1}{2} I_1 \omega_i^2} = \left(\frac{I_1 + I_2}{I_1} \right) \frac{\omega_f^2}{\omega_i^2}$$

we know $\omega_f = \frac{\omega_i}{4}$, $I_1 = \frac{m_1 R^2}{2}$, $I_2 = 3m_1 \frac{R^2}{2}$

$$\frac{K_f}{K_i} = \left[\frac{I_1}{I_1} + \frac{I_2}{I_1} \right] \cdot \left(\frac{\omega_f}{\omega_i} \right)^2 = \left[1 + \frac{3m_1 R^2}{2} \right] \cdot \left(\frac{1}{4} \right)^2 =$$

$$= \left[1 + 3 \right] \cdot \frac{1}{16} = \frac{4}{16} = \frac{1}{4} = 25\%$$

The second disk is dropped into friction as the second disk is dropped onto the first.

The vector nature of torque and angular momentum

Next Lecture