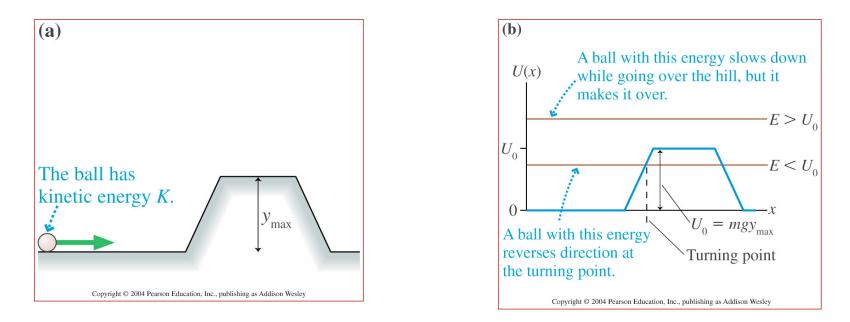
The square barrier:

Behaviour of a classical ball rolling towards a hill (potential barrier):



If the ball has energy E less than the potential energy barrier (U=mgy), then it will not get over the hill.

The other side of the hill is a **classically forbidden** region.

The square barrier:

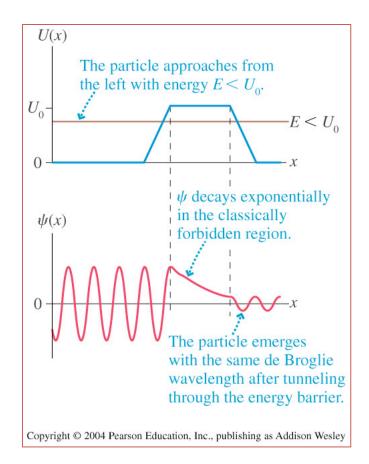
Behaviour of a quantum particle at a potential barrier

Solving the TISE for the square barrier problem yields a peculiar result:

If the quantum particle has energy E less than the potential energy barrier U, there is still a non-zero probability of finding the particle classically forbidden region !

This phenomenon is called **tunneling**.

To see how this works let us solve the TISE...



The square barrier:

Behaviour of a quantum particle at a potential barrier

To the left of the barrier (region I), U=0 Solutions are free particle plane waves:

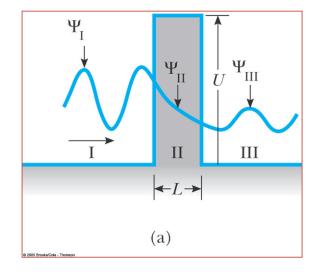
$$\phi(x) = Ae^{ikx} + Be^{-ikx}, \quad k = \frac{\sqrt{2mE}}{\hbar}$$

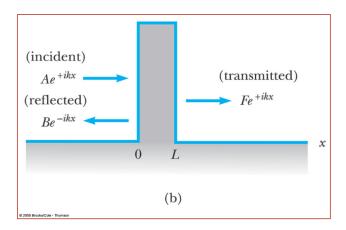
The first term is the incident wave moving to the **right**

The second term is the reflected wave moving to the **left**.

Reflection coefficient:

$$R = \frac{\left|\phi_{reflected}\right|^{2}}{\left|\phi_{incident}\right|^{2}} = \frac{\left|B\right|^{2}}{\left|A\right|^{2}}$$





The square barrier:

Behaviour of a quantum particle at a potential barrier

To the right of the barrier (region III), U=0. Solutions are free particle plane waves:

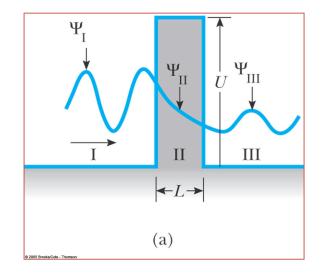
$$\phi(x) = Fe^{ikx}, \quad k = \frac{\sqrt{2mE}}{\hbar}$$

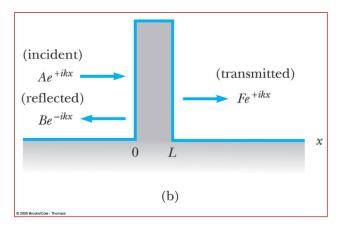
This is the transmitted wave moving to the **right**

Transmission coefficient:

$$T = \frac{\left|\phi_{transmitted}\right|^{2}}{\left|\phi_{incident}\right|^{2}} = \frac{\left|F\right|^{2}}{\left|A\right|^{2}}$$

$$T + R = 1$$





The square barrier:

Behaviour of a quantum particle at a potential barrier

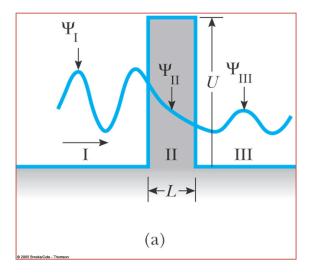
In the barrier region (region II), the TISE is

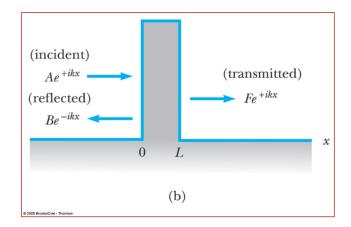
 $-\frac{\hbar^2}{2m}\frac{d^2}{dx^2}\phi(x) = (E-U)\phi(x)$

Solutions are

 $\phi(x) = Ce^{-\alpha x} + De^{\alpha x}$

$$\alpha = \frac{\sqrt{2m(U-E)}}{\hbar}$$





The square barrier:

Behaviour of a quantum particle at a potential barrier

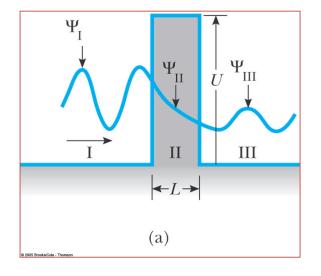
At x=0, region I wave function = region II wave function:

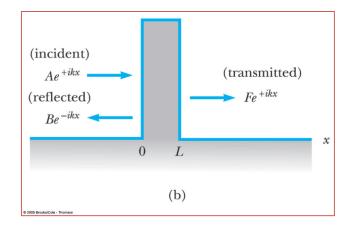
 $Ae^{ikx} + Be^{-ikx} = Ce^{-\alpha x} + De^{\alpha x}$

A + B = C + D

At x=L, region II wave function = region III wave function:

$$Ce^{-\alpha L} + De^{\alpha L} = Fe^{ikL}$$





The square barrier:

Behaviour of a quantum particle at a potential barrier

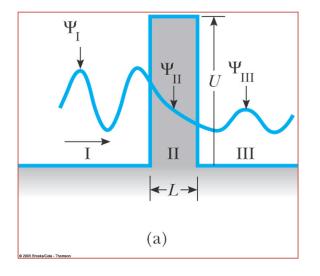
At x=0, $d\phi/dx$ in region I = $d\phi/dx$ in region II:

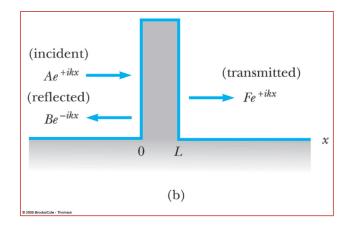
 $ikAe^{ikx} - ikBe^{-ikx} = -\alpha Ce^{-\alpha x} + \alpha De^{\alpha x}$

 $ikA - ikB = -\alpha C + \alpha D$

At x=L, $d\phi/dx$ in region II = $d\phi/dx$ in region III :

$$-\alpha C e^{-\alpha L} + \alpha D e^{\alpha L} = i k F e^{i k L}$$





The square barrier:

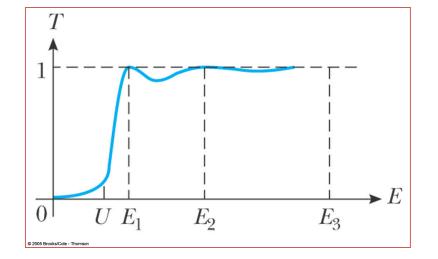
Behaviour of a quantum particle at a potential barrier

Solving the 4 equations, we get

$$T = \frac{1}{1 + \frac{1}{4} \left[\frac{U^2}{E(U-E)} \right] \left[\frac{e^{\alpha L} + e^{-\alpha L}}{2} \right]^2}$$

For low energies and wide barriers,

$$T \approx e^{-\alpha L}$$

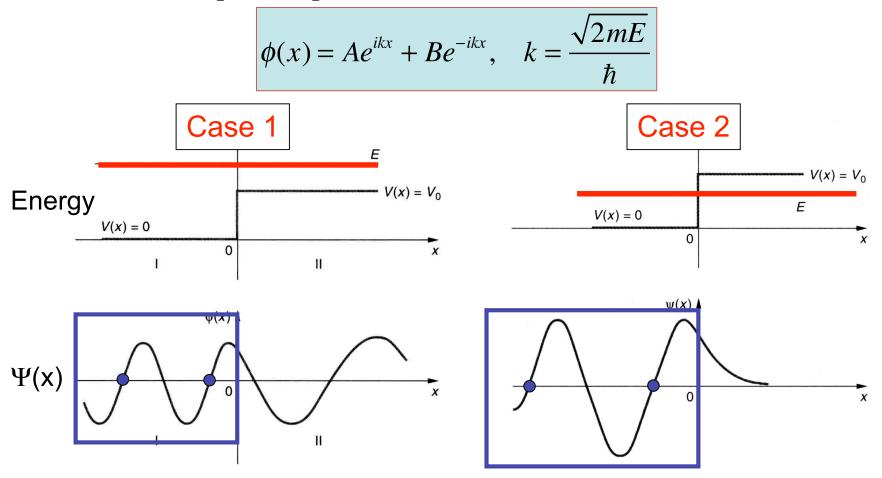


For some energies, T=1, so the wave function is fully transmitted (transmission resonances).

This occurs due to wave interference, so that the reflected wave function is completely suppressed.

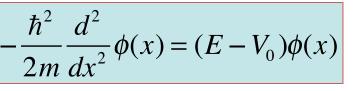
The step barrier:

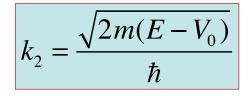
To the left of the barrier (region I), U=0. Solutions are free particle plane waves:



The step barrier:

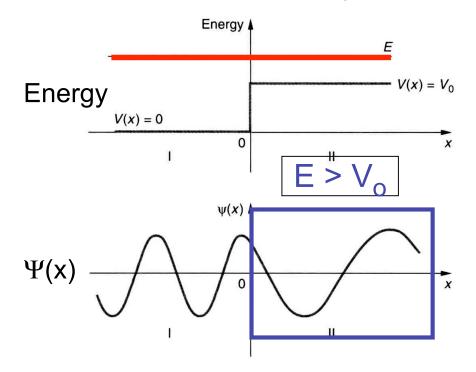
Inside Step: $U = V_o$





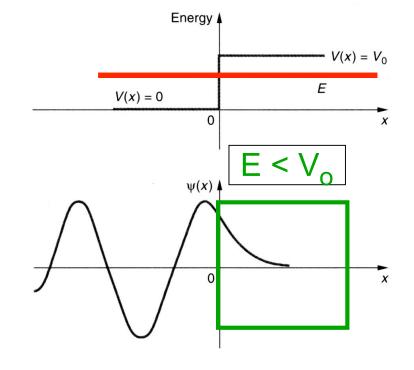
Case 2

$\Psi(x)$ is <u>decaying</u> for $E < V_o$

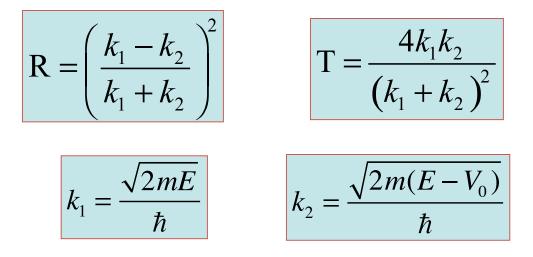


Case 1

 $\Psi(x)$ is <u>oscillatory</u> for E > V_o



The step barrier:



R(reflection) + T(transmission) = 1

Reflection occurs at a barrier ($R \neq 0$), regardless if it is step-down

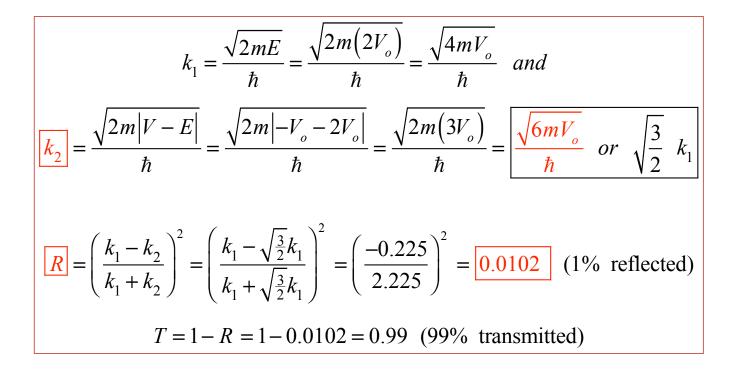
or step-up.

R depends on the wave vector difference $(k_1 - k_2)$ (or energy difference), but not on which is larger.

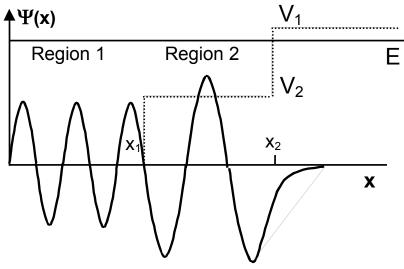
Classically, R = 0 for energy E larger than potential barrier (V_0).

The step barrier:

A free particle of mass m, wave number k_1 , and energy $E = 2V_0$ is traveling to the right. At x = 0, the potential jumps from zero to $-V_0$ and remains at this value for positive x. Find the wavenumber k_2 in the region x > 0 in terms of k_1 and V_0 . Find the reflection and transmission coefficients R and T.



Sketch the **wave function** $\psi(x)$ corresponding to a particle with energy E in the potential well shown below. Explain how and why the wavelengths and amplitudes of $\psi(x)$ are **different** in regions 1 and 2.

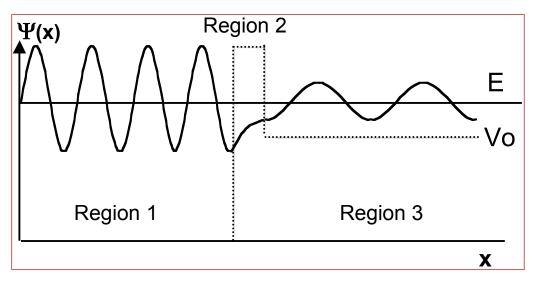


 $\psi(x)$ oscillates inside the potential well because E > V(x), and decays exponentially outside the well because E < V(x).

The **frequency** of $\psi(x)$ is **higher in Region 1** vs. Region 2 because the kinetic energy is higher $[E_k = E - V(x)]$.

The **amplitude** of $\psi(x)$ is **lower in Region 1** because its higher E_k gives a higher velocity, and the particle therefore spends less time in that region.

Sketch the wave function $\psi(x)$ corresponding to a particle with energy E in the potential shown below. Explain how and why the wavelengths and amplitudes of $\psi(x)$ are different in regions 1 and 3.



 $\psi(x)$ oscillates in regions 1 and 3 because E > V(x), and decays exponentially in region 2 because E < V(x).

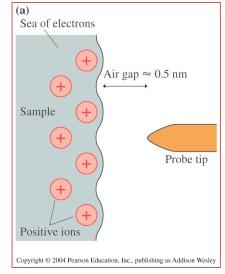
Frequency of $\psi(x)$ is **higher in Region 1** vs. 3 because kinetic energy is higher there.

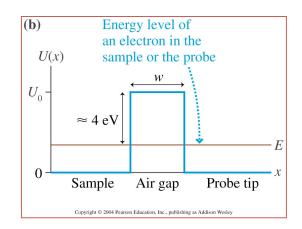
Amplitude of $\psi(x)$ in Regions 1 and 3 depends on the initial location of the wave packet. If we assume a bound particle in Region 1, then the amplitude is higher there and decays into Region 3 (case shown above).

The scanning tunneling microscope:

Scanning-tunneling microscopes allow us to see objects at the atomic level.

- A small air gap between the probe and the sample acts as a potential barrier.
- Energy of an electron is less than the energy of a free electron by an amount equal to the work function.
- Electrons can tunnel through the barrier to create a current in the probe.
- The current is highly sensitive to the thickness of the air gap.
- As the probe is scanned across the sample, the surface structure is mapped by the change in the tunneling current.

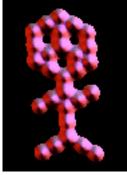




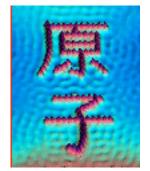
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Carbon Monoxide on Platinum



Iron on Copper

www.almaden.ibm.com/vis/stm/

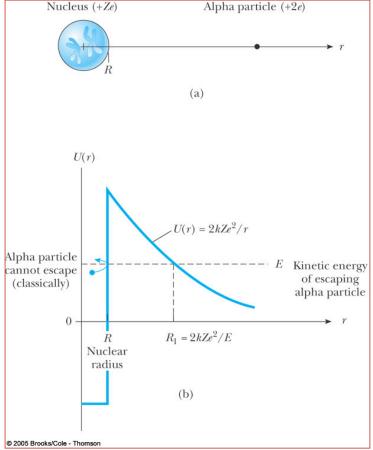
Decay of radioactive elements:

Emission of α particles (helium nucleii) in the decay of radioactive elements is an example of tunneling

- α particles are confined in the nucleus modeled as a square well
- α particles can eventually tunnel through the Coulomb potential barrier.
- Tunneling rate is very sensitive to small changes in energy, accounting for the wide range of decay times:

$$T = e^{8\sqrt{\frac{ZR}{r_0}} - 4\pi Z\sqrt{\frac{E_0}{E}}},$$

$$r_0 \approx 7.25 \, fm, E_0 = 0.0993 MeV$$



Decay of radioactive elements:

Emission of α particles (helium nucleii) in the decay of radioactive elements is an example of tunneling

• Transmission probability:

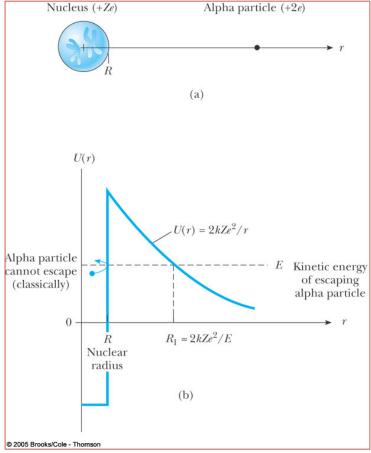
 $T = e^{8\sqrt{\frac{ZR}{r_0}} - 4\pi Z\sqrt{\frac{E_0}{E}}},$ $r_0 \approx 7.25 \, fm, E_0 = 0.0993 MeV$

• Transmission rate λ = frequency of collisions with the barrier x T

$$\lambda = fT \approx 10^{21} e^{8\sqrt{\frac{ZR}{r_0}} - 4\pi Z\sqrt{\frac{E_0}{E}}},$$

• Half life:

$$t_{1/2} = \frac{0.693}{\lambda}$$



Other applications of quantum mechanical tunneling:

- Tunneling diodes (used in digital chips in computers)
- Explanation of ammonia inversion (see text)
- Theory of black hole decay
- •