# **RF Pulse Design** *Multi-dimensional Excitation I*

M229 Advanced Topics in MRI Kyung Sung, Ph.D. 2018.04.10

# **Class Business**

- Office hours
  - Instructors: Fri 10-12pm
     TAs: Xinran Zhong and Zhaohuan Zhang (time: TBD)
  - Emails beforehand would be helpful
- Homework 1 will be out on 4/12 (due on 4/26)
- Papers and Slides

# **Today's Topics**

- Recap of adiabatic pulses
- Small tip approximation
- Excitation k-space interpretation
- Design of 2D excitation pulses
  - Spiral pulse design

# **Recap of Adiabatic Pulses**

- A special class of RF pulses that can achieve uniform flip angle
- Flip angle is independent of the applied B1 field

$$\theta \neq \int_0^T B_1(\tau) d\tau$$

- Slice profile of an adiabatic pulse is obtained using Bloch simulations
- Does not follow the small tip approximation

#### **Adiabatic Pulses**

$$\theta \neq \int_0^\tau \gamma B_1(s) ds$$

- Amplitude and frequency modulation
- Long duration (8-12 ms)
- High B1 amplitude (>12 µT)
- Generally NOT multipurpose (inversion pulses cannot be used for refocusing, etc.)

$$heta = \int_0^ au \gamma B_1(s) ds$$

- Amplitude modulation
- Short duration (0.3-1 ms)
- Low B1 amplitude
- Generally multi-purpose (inversion pulses can be used for refocusing, etc.)

## **Adiabatic Pulses**

Frequency modulated pulses:

 $B_{1}(t) = A(t)e^{-i\omega_{1}(t)t}$ envelop
sweep

$$\frac{d\vec{M}}{dt} = \vec{M} \times \gamma \vec{B}_{eff}$$

where 
$$\vec{B}_{eff} = \left( egin{array}{c} A(t) \\ 0 \\ B_0 - rac{\omega}{\gamma} + rac{\omega_1(t)}{\gamma} \end{array} 
ight)$$



### Adiabatic Excitation: BIR Pulses

- BIR: B1 Insensitive Rotation
- Most popular: BIR-4 Pulse







### Adiabatic Excitation: BIR4 Pulse



### **Bloch Simulator**

#### - http://mrsrl.stanford.edu/~brian/blochsim/

[mx,my,mz] = bloch(bl,gr,tp,t1,t2,df,dp,mode,mx,my,mz)

Bloch simulation of rotations due to B1, gradient and off-resonance, including relaxation effects. At each time point, the rotation matrix and decay matrix are calculated. Simulation can simulate the steady-state if the sequence is applied repeatedly, or the magnetization starting at m0.

INPUT:

```
%%% User inputs:
mu = 5; % Phase modulation parameter [dimensionless]
beta1 = 672; % Frequency modulation parameter [rad/s]
pulseWidth = 10.24; % RF pulse duration [ms]
A0 = 0.12; % Peak B1 amplitude [Gauss].
****
                       % number of samples in the RF pulse
nSamples = 512;
dt = pulseWidth/nSamples/1000; % time step, [seconds]
tim sech = linspace(-pulseWidth/2,pulseWidth/2,nSamples)./1000';
% time scale to calculate the RF waveforms in seconds.
% Amplitude modulation function B1(t):
B1 = A0.* sech(beta1.*tim sech);
% Carrier frequency modulation function w(t):
w = -mu.*beta1.*tanh(beta1.*tim sech)./(2*pi);
% The 2*PI scaling factor at the end converts the unit from rad/s to Hz
% Phase modulation function phi(t):
phi = mu .* log(sech(beta1.*tim sech));
% Put together complex RF pulse waveform:
rf pulse = B1 .* exp(li.*phi);
% Generate a time scale for the Bloch simulation:
tim bloch = [0:(nSamples-1)]*dt;
```

```
%% The Bloch simulator requires a gradient input. For our simulation,
% gradient will be zero, as we are simulating a non-selective RF pulse.
T1_value = 10000; % [ms]
T2_value = 10000; % [ms]
f_max = 4000; % off-resonance frequency range [Hz]
freq_range = linspace(-f_max,f_max,1000); % off-resonance frequency range [Hz]
grad_pulse = zeros(1,length(rf_pulse));
mod = 0;
[mx1,my1,mz1] = bloch(rf_pulse, grad_pulse, dt, ...
T1_value/1000, T2_value/1000, freq_range, 0, mod, 0, 0, 0);
% Plotting the longitudinal magnetization to see the inversion profile as a
% function of resonance frequency
figure(2);
```

```
plot(freq_range, mz1,'k','LineWidth',LineWidthVal);
title('Inversion Profile'); xlabel('Frequency (Hz)'); ylabel('M_z'); grid on;
v = axis; axis([v(1) v(2) -1.05 1.05]);
```

# **Small Tip Approximation**

## Bloch Equation (at on-resonance)

$$\frac{d\vec{M}_{rot}}{dt} = \vec{M}_{rot} \times \gamma \vec{B}_{eff}$$
where  $\vec{B}_{eff} = \begin{pmatrix} B_1(t) & 0 \\ 0 & 0 \\ B_0 & \frac{\omega}{\gamma} + G_z z \end{pmatrix}$ 

When we simplify the cross product,

$$\frac{d\vec{M}}{dt} = \begin{pmatrix} 0 & \omega(z) & 0\\ -\omega(z) & 0 & \omega_1(t)\\ 0 & -\omega_1(t) & 0 \end{pmatrix} \vec{M}$$
$$\omega(z) = \gamma G_z z \quad \omega_1(t) = \gamma B_1(t)$$

## **Small Tip Approximation**

$$\frac{d\vec{M}}{dt} = \begin{pmatrix} 0 & \omega(z) & 0\\ -\omega(z) & 0 & \omega_1(t)\\ 0 & -\omega_1(t) & 0 \end{pmatrix} \vec{M}$$

 $M_z \approx M_0$  small tip-angle approximation

 $\sin \theta \approx \theta$  $\cos \theta \approx 1$  $M_z \approx M_0 \rightarrow \text{constant}$ 

$$\frac{dM_z}{dt} = 0$$

$$\frac{dM_{xy}}{dt} = -i\gamma G_z z M_{xy} + i\gamma B_1(t) M_0$$

$$M_{xy} = M_x + iM_y$$

First order linear differential equation. Easily solved.

$$\frac{dM_{xy}}{dt} = -i\gamma G_z z M_{xy} + i\gamma B_1(t) M_0$$

Solving a first order linear differential equation:

$$M_{xy}(t,z) = i\gamma M_0 \int_0^t B_1(s) e^{-i\gamma G_z z \cdot (t-s)} ds$$

 $M_{r}(\tau, z) = iM_{0}e^{-i\omega(z)\tau/2} \cdot \mathcal{FT}_{1D}\{\omega_{1}(t + \frac{\tau}{2})\}|_{f = -(\gamma/2\pi)G_{z}z}$ 

(See the note for complete derivation)





- For small tip angles, "the slice or frequency profile is well approximated by the Fourier transform of B1(t)"
- The approximation works surprisingly well even for flip angles up to 90°



small-angle approximation still works reasonably well for flip angles that aren't necessarily "small"

### **Multi-Dimensional Excitation RF Pulses**

- 1D vs. N-D RF pulses
- Small tip angle approximation revisited
- Excitation k-space interpretation
- 1D examples in excitation k-space
- Excitation k-space integrals
- 2D excitation pulse design steps
- 2D spiral pulse design example
- EPI pulse design, spectral-spatial pulses (next lecture)

### What is Multi-Dimensional Excitation?

Multi-dimensional excitation occurs when using multi-dimensional RF pulses in MRI/NMR, i.e. 2D or 3D RF pulses



- 1D pulses are selective along 1 dimension, typically the slice select dimension
- 2D pulses are selective along 2 dimensions
  - So, a 2D pulse would select a long cylinder instead of a slice
  - The shape of the cross section depends on the 2D RF pulse

## Excitation k-space Interpretation

## **Small Tip Approximation**

**Small Tip Approximation** 

$$M_{xy}(t,\vec{r}) = i\gamma M_0 \int_0^t B_1(s) e^{-i\gamma \int_s^t \vec{G}(\tau) d\tau \cdot \vec{r}} ds$$

Let us define: 
$$ec{k}(s,t) = -rac{\gamma}{2\pi}\int_s^tec{G}( au)d au$$

$$M_{xy}(t,\vec{r}) = i\gamma M_0 \int_0^t B_1(s) e^{i2\pi \vec{k}(s,t)\cdot\vec{r}} ds$$

## **One-Dimensional Example**

$$\vec{k}(s,t) = -\frac{\gamma}{2\pi} \int_s^t \vec{G}(\tau) d\tau$$



# **One-Dimensional Example**



- This gives magnetization at t = t<sub>0</sub>, the end of the pulse
- Looks like you scan across k-space, then return to origin

# Evolution of Magnetization During Pulse

- RF pulse goes in at DC  $(k_z = 0)$
- Gradients move previously applied weighting around
- Think of the RF as "writing" an analog waveform in k-space
- Same idea applies to reception



# **Other 1D Examples**





# **Multiple Excitations**

- Most acquisition methods require several repetitions to make an image
  - e.g., 128 phase encodes
- Data is combined to reconstruct an image
- Same idea works for excitation!

## Simple 1D Example



Same profile as slice selective pulse, but zero echo time

# Small Tip Approximation Solution as k-space Integral

$$M_{xy}(t,\vec{r}) = i\gamma M_0 \int_0^t B_1(s) \underline{e^{i2\pi \vec{k}(s,t) \cdot \vec{r}}} ds$$

$$e^{i2\pi \vec{k}(s,t)\cdot \vec{r}} = \int_{\vec{k}} {}^3 \delta(\vec{k}(s,t) - \vec{k}) e^{i2\pi \vec{k}\cdot \vec{r}} d\vec{k}$$

Substituting and changing the order of integration:

$$M_{xy}(t,\vec{r}) = iM_0 \int_{\vec{k}} \left[ \int_{-\infty}^t \gamma B_1(s)^3 \delta(\vec{k}(s,t) - \vec{k}) ds \right] e^{i2\pi \vec{k} \cdot \vec{r}} d\vec{k}$$
$$p(\vec{k})$$

 $M_{xy}(t,\vec{r}) = iM_0 \int_{\vec{k}} \left[ \int_{-\infty}^t \gamma B_1(s)^3 \delta(\vec{k}(s,t) - \vec{k}) ds \right] e^{i2\pi \vec{k} \cdot \vec{r}} d\vec{k}$ 

$$M_{xy}(t,\vec{r}) = iM_0 \int_{\vec{k}} p(\vec{k}) e^{i2\pi\vec{k}\cdot\vec{r}} d\vec{k}$$

where 
$$p(\vec{k}) = \int_{-\infty}^{t} \gamma B_1(s)^3 \delta(\vec{k}(s,t) - \vec{k}) ds$$

- The magnetization is the inverse transform of p(k)
- We want this to be <u>a unit delta</u>, multiplied by a weighting function

## Small Tip Approximation Solution as k-space Integral

Multiply and divide by |k'(s,t)|:

$$p(\vec{k}) = \int_{-\infty}^{t} \frac{\gamma B_1(s)}{|k'(s,t)|} \frac{{}^3\delta(\vec{k}(s,t) - \vec{k})|k'(s,t)|ds}{\textit{Unit Delta}}$$

If we assume W(k) is single-valued

$$p(\vec{k}) = W(\vec{k}) \int_{-\infty}^{t} {}^{3}\delta(\vec{k}(s,t) - \vec{k})|k'(s,t)|ds$$
$$p(\vec{k}) = W(\vec{k})S(\vec{k})$$

## Small Tip Approximation Solution as k-space Integral

$$M_{xy}(t,\vec{r}) = iM_0 \int_{\vec{k}} p(\vec{k}) e^{i2\pi \vec{k} \cdot \vec{r}} d\vec{k}$$
$$p(\vec{k}) = W(\vec{k}) S(\vec{k})$$

$$egin{aligned} W(ec{k}) &= rac{\gamma B_1(s)}{|k'(s,t)|} & \textit{k-space weighting} \ S(ec{k}) &= \int_{-\infty}^t {}^3 \delta(ec{k}(s,t) - ec{k}) |k'(s,t)| ds \ & \textit{k-space sampling} \end{aligned}$$

## Small Tip Approximation Solution as k-space Integral

$$M_{xy}(t,\vec{r}) = iM_0 \int_{\vec{k}} W(\vec{k}) S(\vec{k}) e^{i2\pi \vec{k} \cdot \vec{r}} d\vec{k}$$

So, the inverse Fourier transform of the k-space weighting will give us the excitation profile!

$$W(\vec{k}) = \frac{\gamma B_1(s)}{|k'(s,t)|}$$

k-space weighting

## **Design of 2D Excitation Pulses**

## 2D Pulse Design

1. Choose a k-space trajectory

$$ec{k}(s,t) = -rac{\gamma}{2\pi}\int_{s}^{t}ec{G}( au)d au$$

2. Choose a weighting function

 $W(\vec{k})$ 

3. Design the RF pulse

$$B_1(s) = \frac{1}{\gamma} W(\vec{k}) \cdot |k'(s,t)|$$

## 1. Choose a k-space trajectory

- Select a k-space trajectory that uniformly covers k-space
  - k-space extent (−k<sub>max</sub>, k<sub>max</sub>) ⇒ spatial resolution
  - sampling density ( $\Delta k$ )  $\Rightarrow$  spatial FOV

# 1. Choose a k-space trajectory



• Spiral is common for pencil beams

• EPI is common for spectral-spatial pulses

## 2. Choose a weighting function

$$M_{xy}(t,\vec{r}) = iM_0 \int_{\vec{k}} W(\vec{k}) S(\vec{k}) e^{i2\pi\vec{k}\cdot\vec{r}} d\vec{k}$$

- An excitation profile is the inverse Fourier transform of the weighting function
- If you know what excitation profile you want, its Fourier transform will be the weighting function
- Localized excitation ⇒ low-pass k-space weighting



## 3. Design the RF pulse

$$ec{k}(s,t) = -rac{\gamma}{2\pi}\int_{s}^{t}ec{G}( au)d au$$

$$k'(s,t) = -rac{\gamma}{2\pi} \vec{G}(s)$$

$$B_1(s) = \frac{1}{2\pi} |\vec{G}(s)| W(\vec{k})$$

 $B_1$  needs to be scaled for flip angle

## 2D Spiral Pulse Design

- Two major choices:
  - Resolution  $\Delta r$

$$\Delta r = \frac{1}{2k_{max}}$$

- Smallest volume / minimum transition width
- Field-of-View (FOV)
  - Distance to center of first sidelobe

$$FOV = \frac{1}{\Delta k} = \frac{2N}{2k_{max}}$$



# 2D Spiral Pulse Design

- Spiral Gradient Design
  - Constant angular rate spiral
  - Constant slew rate spiral



# 2D Spiral Pulse Design

• Truncated Jinc Weighting

 $W(\vec{k}) = jinc(N\frac{k_r}{k_{max}}) \cdot rect(\frac{k_r}{2k_{max}})$ 



Minimum transition width, but residual ripples

# 2D Spiral Pulse Design

Windowed Jinc Weighting

$$W(\vec{k}) = jinc(N\frac{k_r}{k_{max}}) \cdot A(\frac{k_r}{2k_{max}})$$



Doubled transition width, but smoother response

## 2D Spiral Pulse Design

Calculation of the RF pulse
given W(k) and k(s,t)

$$W(\vec{k}) = \frac{\gamma B_1(s)}{|k'(s,t)|}$$
$$B_1(s) = \frac{1}{2\pi} |\vec{G}(s)| W(\vec{k})$$

- needs to be scaled for flip angle
- [k'(s,t)] is an estimate of the density compensation function d(t)



# Conclusions

- N-D RF pulses are selective in N-dimensions
- The small tip approximation can be extended to describe the excitation k-space
- The small tip approximation solution can be used to show that the excitation profile of an N-D pulse is given by the inverse Fourier transform of the excitation k-space weighting

# Conclusions

- An N-D RF pulse can be designed by:
  - Choosing a k-space trajectory
  - Choosing a k-space weighting function
  - Then calculating the  $B_1(t)$  and G(t) functions

## **Next Class**

- N-D pulses with EPI trajectory
- Spatial-spectral pulses
- Matlab demo of N-D pulses

# Thank You!

- Further reading
  - Read "Spatial-Spectral Pulses" p.153-163
- Acknowledgments
  - John Pauly's EE469b (RF Pulse Design for MRI)
  - Shams Rashid, Ph.D.

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