Spin currents in non-collinear antiferromagnets

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MAX-PLANCK-INSTITUT FÜR CHEMISCHE PHYSIK FESTER STOFFE Rising interest in antiferromagnets for spintronics applications

Some advantages over ferromagnets:

- ✓ Fast dynamics
- ✓ No stray fields, insensitive to magnetic fields
- ✓ Wide range of antiferromagnetic materials including metals, semiconductors, insulators, multiferroics, superconductors...

but

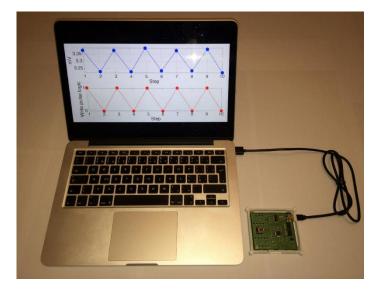
The antiferromagnetic order is difficult to manipulate and detect

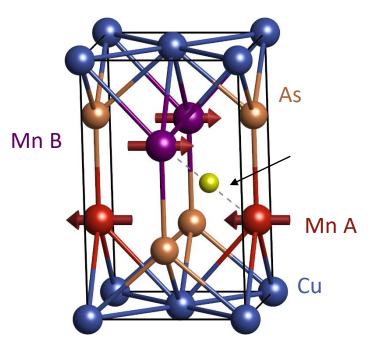
Many spintronics effects, which exist in ferromagnets are prohibited by symmetry in most antiferromagnets

Electrical control and detection of the antiferromagnetic order has been demonstrated

A key to efficient manipulation of the antiferromagnetic order is using a **staggered magnetic field**

Can be generated by spin-orbit torque

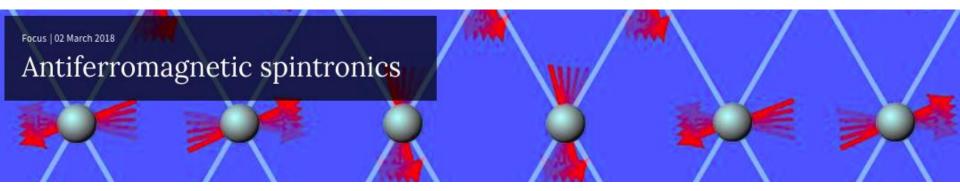




CuMnAs

Zelezny et al., PRL 113, 157201 (2014) Wadley et al., Science 351, 587–590 (2016) Bodnar et al., arXiv:1706.02482 Meinert et al., arXiv:1706.06983

Nature physics focused issue on antiferromagnetic spintronics

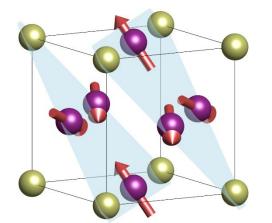


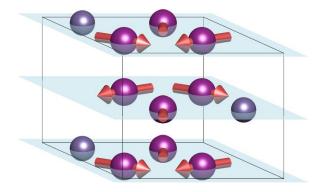
Overview - T. Jungwirth et al., Nature Physics 14, 200–203 (2018)
Transport - J. Železný et al., Nature Physics 14, 220-228 (2018)
Opto-spintronics - P. Němec et al., Nature Physics 14, 229-241 (2018)
Spin textures and dynamics - O. Gomonay et al., Nature Physics 14, 213-216 (2018)
Topology - L. Šmejkal et al., Nature Physics 14, 242-251 (2018)
Synthetic antiferromagnets - R. A. Duine et al., Nature Physics 14, 217-219 (2018)

Non-collinear antiferromagnets

Non-collinear order is common in antiferromagnets

Triangular antiferromagnets Mn3X





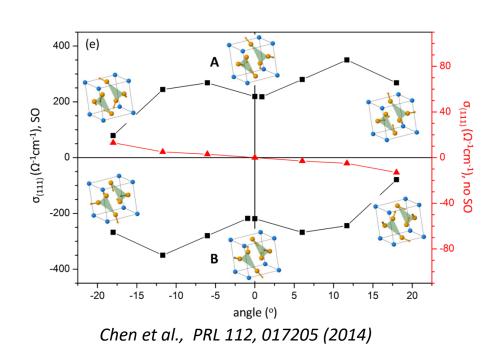
Mn₃Rh, Mn₃Pt, Mn₃Ir

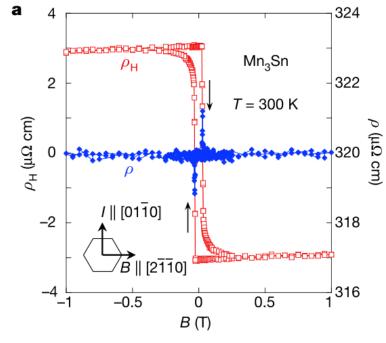
Mn₃Ga, Mn₃Ge, Mn₃Sn

• Simplest example of a non-collinear magnetic system

Anomalous Hall effect

A large anomalous Hall effect exist in the non-collinear antiferromagnets





Naktsuji et al. Nature 527, 212 (2015)

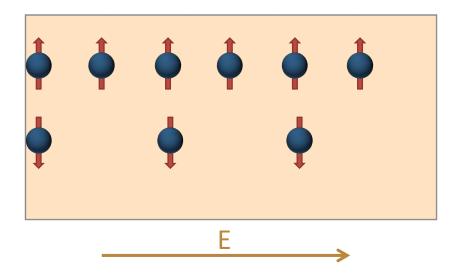
Small magnetic moment present, but is not the origin of the AHE

The AHE exists because the magnetic order breaks symmetry

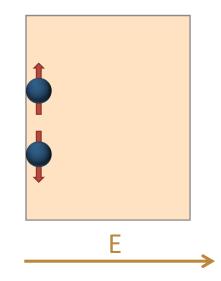
Kubler et al., EPL 108, 67001 (2014) Nayak et al. Sci. Adv. 2016;2:e1501870 Feng et al., PRB 92, 144426 (2015)

Spin currents

Spin-polarized current in ferromagnets



Spin-Hall effect



Odd under time-reversal Non-relativistic origin

Even under time-reversal Relativistic origin

Spin currents are responsible for giant and tunneling magnetoresistance, spin-transfer torque, spin-orbit torque,...

Definition of spin current

Spin current operator is normally defined as: $\hat{j}_{ij}^s = \frac{1}{2} \{ \hat{s}_i, \hat{v}_j \}$

Issues with this definition:

1. Spin is not conserved so the spin current does not satisfy the continuity equation

$$\frac{\partial S_z}{\partial t} + \nabla \cdot \mathbf{J}_s = \mathcal{T}_z$$

Shi et al., PRL 96, 076604 (2006)

2. Spin current can occur in equilibrium

Rashba, Phys. Rev. B 68, 241315(R) 2003

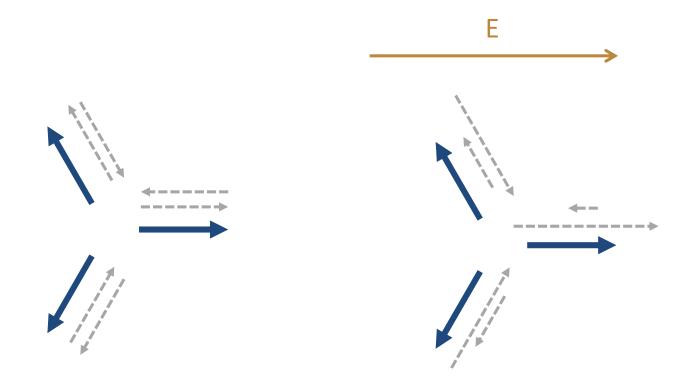
- 3. Spin currents also can exist in insulators
- 4. Spin current cannot be directly measured

But it's a well defined (and useful) physical quantity!

For accurate estimates and comparison with experiment it's best to calculate directly spin accumulation, torque...

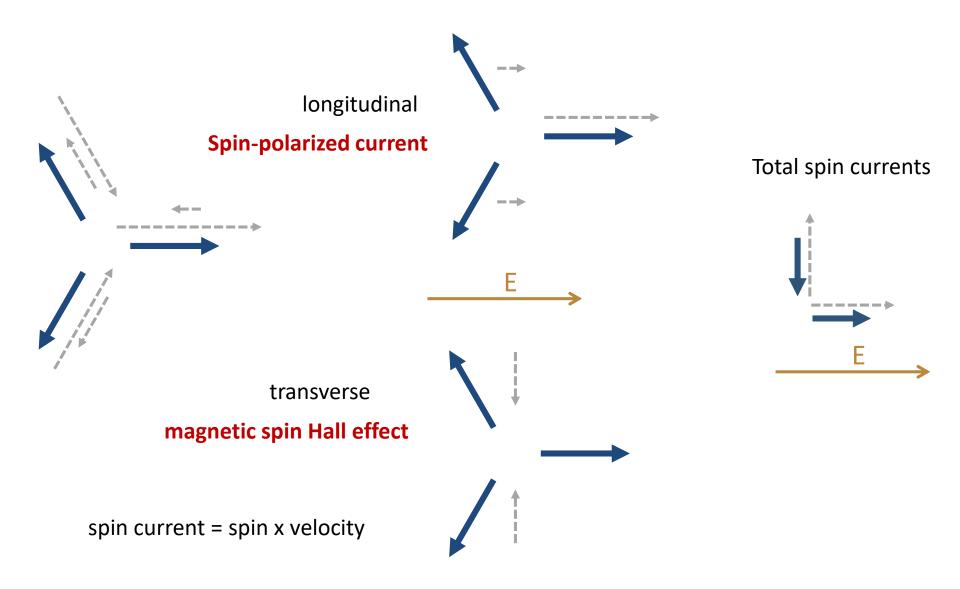
In most antiferromagnets spin-polarized current is not allowed by symmetry

In the Mn₃X non-collinear antiferromagnets spin-polarized current is present



Zelezny et al., PRL 119, 187204 (2017)

Origin of the odd spin currents



Zelezny et al., PRL 119, 187204 (2017)

$$\delta A = \chi^{I} \mathbf{E} + \chi^{II} \mathbf{E}$$
$$\chi^{I} = -\frac{e\hbar}{\pi} \sum_{\mathbf{k},n,m} \frac{\Gamma^{2} \operatorname{Re} \left(\langle n\mathbf{k} | \, \hat{A} \, | m\mathbf{k} \rangle \, \langle m\mathbf{k} | \, \hat{\mathbf{v}} \cdot \hat{\mathbf{E}} \, | n\mathbf{k} \rangle \right)}{[(E_{F} - \varepsilon_{n\mathbf{k}})^{2} + \Gamma^{2}][E_{F} - \varepsilon_{m\mathbf{k}})^{2} + \Gamma^{2}]},$$
$$\chi^{II} = -2\hbar e \sum_{\mathbf{k},n\neq m}^{n \text{ occ.}} \frac{\operatorname{Im} \left(\langle n\mathbf{k} | \, \hat{A} \, | m\mathbf{k} \rangle \, \langle m\mathbf{k} | \, \hat{\mathbf{v}} \cdot \hat{\mathbf{E}} \, | n\mathbf{k} \rangle \right)}{(\varepsilon_{n\mathbf{k}} - \varepsilon_{m\mathbf{k}})^{2}},$$

These two parts have precisely opposite transformation under time-reversal

Under time-reversal:
$$\langle n\mathbf{k} | \hat{A} | m\mathbf{k} \rangle \rightarrow \langle n\mathbf{k} | T^{-1} \hat{A} T | m\mathbf{k} \rangle^*$$

We use a constant band broadening to simulate disorder

		"Real" part	"Imaginary" part
Conductivity:	$\hat{A} = -e\hat{\mathbf{v}}$	even - ordinary conductivity	odd - Anomalous Hall effect
Spin-orbit torque:	$\hat{A} = \hat{\mathbf{T}}$	Odd - field-like torque	<mark>even</mark> – antidamping- like torque
Spin current:	$\hat{A} = \frac{1}{2} \{ \hat{\mathbf{s}}, \hat{\mathbf{v}} \}$	odd – spin-polarized current	even - spin Hall effect

Odd components can only exist in magnetic materials

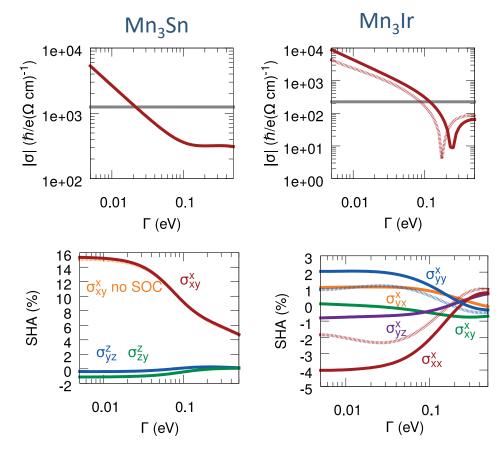
Calculations

The odd spin currents **depend strongly on disorder** unlike the intrinsic spin Hall effect

Spin current angle:
$$lpha=rac{e}{\hbar}rac{\sigma_{jk}^{i}}{\sigma_{kk}}$$

Both large longitudinal and transverse spin currents present

	α (%)	$lpha_{\perp}$ (%)
Mn ₃ Ga	13.7	13.5
Mn ₃ Ge	5.1	5.0
Mn ₃ Sn	12.4	12.4
Mn ₃ Rh	1.1	0.7
Mn ₃ Ir	3.5	1.9
Mn ₃ Pt	4.2	1.9



Spin currents are smaller than typically in ferromagnets, but still relatively large

n BCC Fe
$$oldsymbol{lpha}_{||} \sim 18\%$$

 $oldsymbol{lpha}_{\perp} \sim 1\%$

Zelezny et al., PRL 119, 187204 (2017)

Odd – magnetic spin Hall effect

$$\sigma^{x} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\sigma^{y} = \begin{pmatrix} 0 & \sigma_{yxy} & 0 \\ \sigma_{yyx} & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \sigma_{yxy} \\ 0 & 0 & 0 \\ \sigma_{yyx} & 0 & 0 \end{pmatrix}$$

Even - spin Hall effect

$$\sigma^{x} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -\sigma_{xzy} \\ 0 & \sigma_{xzy} & 0 \end{pmatrix}$$

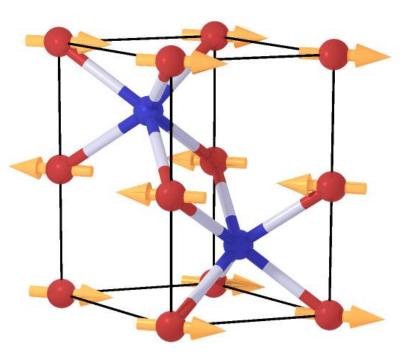
$$\sigma^{y} = \begin{pmatrix} 0 & 0 & -\sigma_{zxy} \\ 0 & 0 & -\sigma_{zxy} \\ 0 & 0 & 0 \\ \sigma_{zyx} & 0 & 0 \end{pmatrix}$$

$$\sigma^{z} = \begin{pmatrix} 0 & \sigma_{zxy} & 0 \\ \sigma_{zyx} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Symmetry and origin distinct from the conventional spin Hall effect

Odd: $\sigma(\mathbf{M}) = -\sigma(-\mathbf{M})$ Even: $\sigma(\mathbf{M}) = \sigma(-\mathbf{M})$

No experimental evidence so far

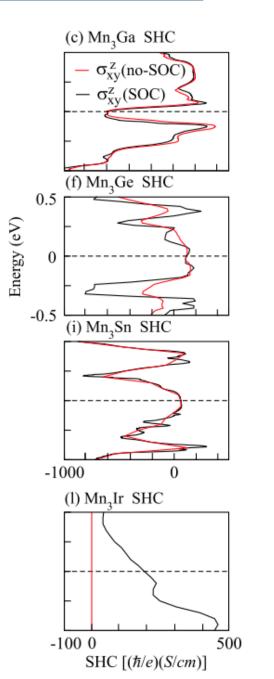


The two sublattices are not connected **by T + inversion** or **T + translation**

MnTe

Spin current angle for spin-polarized current and the magnetic SHE ${\sim}1\%$

SHE calculations

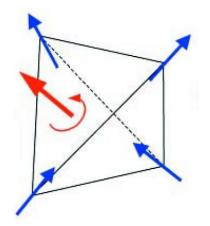


Spin Hall effect exist in these materials even without spin-orbit coupling!

(No spin-orbit coupling = no relativistic effects)

Yang Zhang et al., Phys. Rev. B 95, 075128 (2017) Yang Zhang et al., New J. Phys. 20 073028 (2018) Originates from the magnetic order instead of spin-orbit coupling

Similar to AHE or orbital magnetic moment in non-coplanar magnetic systems



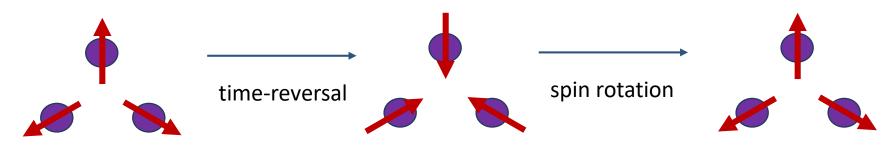
Taguchi et al., Science 291, 2573 (2001)

Can be understood in terms of symmetry and is not necessarily related to topology

Non-collinear magnetic order breaks symmetry similarly to spin-orbit coupling

Example:

In a coplanar system time-reversal + 180° spin rotation is a symmetry without SOC



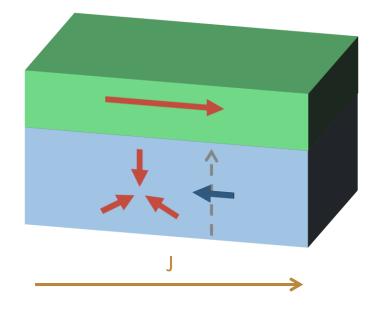
No AHE in a coplanar magnetic system

Brinkman et al., Proc. Royal Soc. London A, 294, 343-358 (1966)

Litvin et al., Physica 76, 538-554 (1974)

SOT in bilayer systems

- SHE and the magnetic SHE will generate a **spin-orbit torque** in FM/AFM bilayer systems
- The spin-orbit torque has been experimentally studied in a number of systems



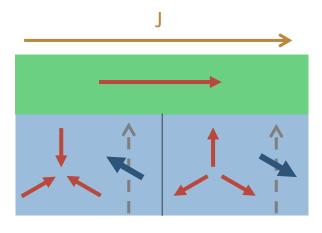
The SHE and the magnetic SHE exist even without spin-orbit coupling

The spin-orbit torque does not necessarily originate from spin-orbit coupling

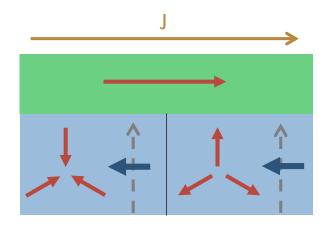
Thsitoyan et al., PRB 92, 214406 (2016) Zhang et al., Science Advances 2, (2016) Young-Wan oh et al., Nature Nanotech., (2016)

Transverse spin currents

Magnetic and conventional spin Hall effect will depend differently on presence of magnetic domains and have different symmetry



odd

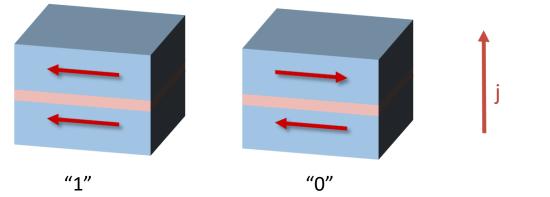


even

Sensitive to the antiferromagnetic order and presence of domains

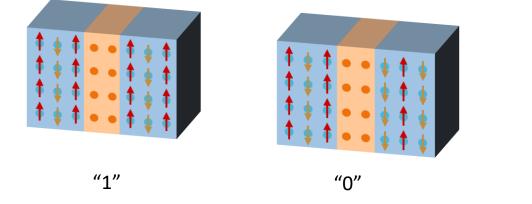
Spin current with spin-polarization along the spin current flow is possible

Antiferromagnetic junctions



Spin-transfer torque and giant or tunneling magnetoresistance

FM junction



Similar effects were theoretically predicted, but are very sensitive to disorder and have never been observed experimentally

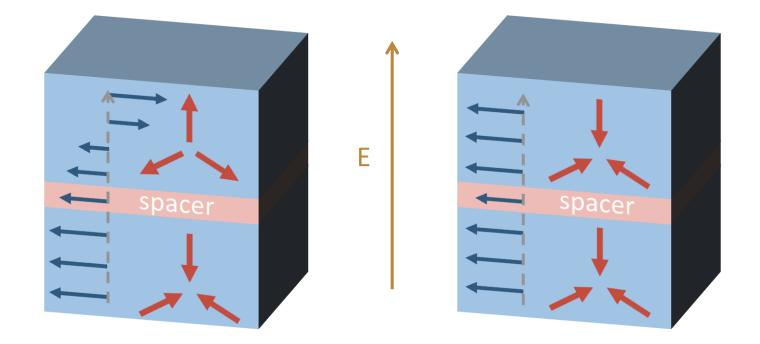
AFM junction

J. Železný, Nature Phys. 14, 220–228 (2018)

In non-collinear antiferromagnets, current is spin-polarized and thus same approach can be used as in ferromagnets



Robust giant or tunneling magnetoresistance and spin-transfer torque should be present



Significance of the AHE and the spin-polarized current is in part because they allow to detect **180° reversal** of all magnetic moments

180° reversal is equivalent to time-reversal

$$\mathbf{j} = \sigma \mathbf{E}$$
 $\mathbf{j}_{\mathbf{s}} = \sigma_{\mathbf{s}} \mathbf{E}$

$$\sigma^{\text{even}} = \sigma(\mathbf{M}_1, \mathbf{M}_2, \dots) + \sigma(-\mathbf{M}_1, -\mathbf{M}_2, \dots) \qquad \sigma_s^{\text{even}} = \sigma_s(\mathbf{M}_1, \mathbf{M}_2, \dots) + \sigma_s(-\mathbf{M}_1, -\mathbf{M}_2, \dots)$$

$$\sigma^{\text{odd}} = \sigma(\mathbf{M}_1, \mathbf{M}_2, \dots) - \sigma(-\mathbf{M}_1, -\mathbf{M}_2, \dots) \qquad \sigma_s^{\text{odd}} = \sigma_s(\mathbf{M}_1, \mathbf{M}_2, \dots) - \sigma_s(-\mathbf{M}_1, -\mathbf{M}_2, \dots)$$

$$\sigma^{odd} = AHE$$
 $\sigma_s^{odd} = spin-polarized current and magnetic SHE$

In AFM with T + inversion or T + translation symmetry both $\sigma^{\rm odd}$ and $\sigma_{\rm s}^{\rm od}$ vanish

In AFM with T + inversion or T + translation 180° reversal cannot be detected with **linear transport effects**

Possible with second order, however!

 $j_i = \sigma_{ijk}^{(2)} E_j E_k$

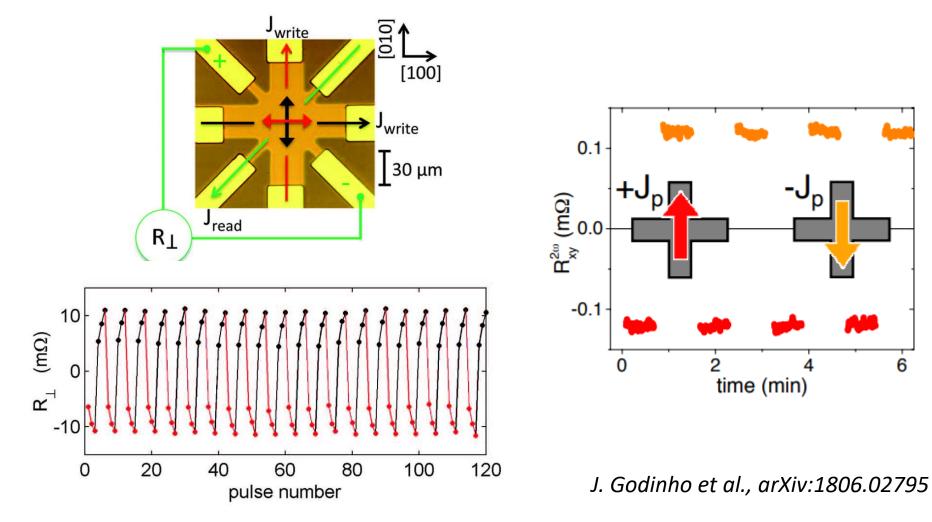
 $\sigma^{(2),}$ allowed in systems with broken inversion symmetry

The odd component of $\sigma^{(2)}$ allowed even in AFMs with PT symmetry

try Mn B Mn A Cu

180° switching in CuMnAs

Switching in CuMnAs is normally 90° measured by AMR

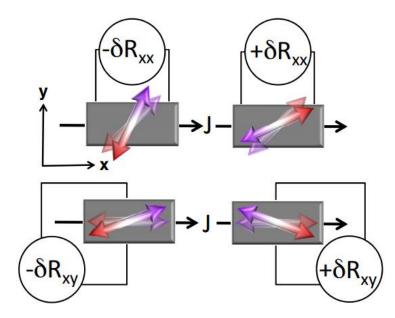


Wadley et al., Science 351, 587–590 (2016)

180° switching

6

Origin of the second-order signal



The original idea behind the experiment was that the second-order signal is due to tilting of the magnetic moments by the spin-orbit torque

But the second order signal could also simply be a second order conductivity, unrelated to a spin-polarization or a torque

Using the Boltzmann equation within the constant relaxation time approximation:

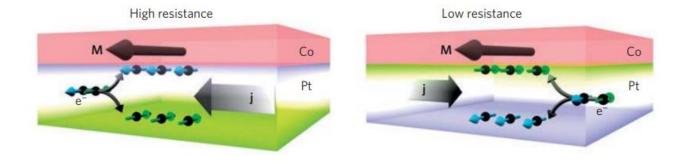
$$\sigma_{ij}^{(1)} = -e^2 \tau \int d^3k \left(v_i v_j \frac{\partial f_{\rm FD}}{\partial \varepsilon} \right)$$
$$\sigma_{ijk}^{(2)} = -e^3 \tau^2 \int d^3k \left(v_i v_j v_k \frac{\partial^2 f_{\rm FD}}{\partial \varepsilon^2} + \frac{1}{\hbar} v_i \frac{\partial v_k}{\partial k_j} \frac{\partial f_{\rm FD}}{\partial \varepsilon} \right)$$

The second-order term can be understood as originating from $\mathbf{k} \leftrightarrow -k$ band assymetry

Requires broken inversion and time-reversal + spin-orbit coupling (or non-collinear magnetic order)

Second harmonics commonly used to measure spin-orbit torques

Unidirectional magnetoresistance



Olejnik et al., Phys Rev B 91, 180402(R) (2015) Avci et al., Nature Physics 11, 570–575 (2015)

Summary

- Noncollinear antiferromagnets are very attractive for spintronics since they combine advantages of ferromagnets and antiferromagnets
 - ✓ Spin-polarized current
 - ✓ AHE, Kerr effect
 - ✓ spin Hall effect (without SOC)
 - ✓ Detecting 180° switching possible

- ✓ Fast magnetic dynamics
- Insensitive to external magnetic fields
- ✓ No stray fields

- Spin-polarized current not limited to ferromagnets, but can exist in collinear and non-collinear antiferromagnets
- Magnetic spin Hall effect: in magnetic materials a transverse current distinct from the spin Hall effect can exist
- Spin-Hall effect can exist without SOC

Zelezny et al., PRL 119, 187204 (2017)

Yang Zhang et al., New J. Phys. 20 073028 (2018)