## **Topological Quantum Matter,** 7.5 ECTS

#### physics.gu.se/~tfkhj/topomatter.html

Part 1: 10 double hour lectures 12/9 - 3/10, 27/10 - 31/10 (tentatively) Part 2: 5 guest lectures 10/11 - 12/12

Examination: homework problems on part 1, project based on one of the guest lectures

Literature: Lecture notes downloadable from the course homepage. Additional text/references will be made available during the course.

Masters students: To get credit for the course, please contact Bengt-Erik Mellander, <u>f5xrk@chalmers.se</u>, after completion of the examination.

## **Concepts in Topological Quantum Matter, 4.5 ECTS**

#### physics.gu.se/~tfkhj/topomatter/

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Before the era of topology in physics...

Standard paradigm of emergent order: spontaneous symmetry breaking

Some examples:



**Crystals** break the translational and rotational symmetry of free space





**Liquid crystals** break rotational but not translational symmetry



**Crystals** 



Liquid crystals



**Magnets** break time-reversal symmetry and the rotational symmetry of spin space







Liquid crystals





Superconductors break a gauge symmetry



Crystals



Liquid crystals



Magnets



Superconductors

... and many more examples





... and many more examples



Lev Landau

In 1980, the "Landau paradigm" was challenged by the discovery of the (integer) *quantum Hall effect* (von Klitzing et al.)

Electrons confined to a plane and in a strong magnetic field show, at low enough temperature, plateaus in the "Hall conductance":

$$\sigma_{xy} = n \frac{e^2}{h}$$

to a precision of at least 9 decimal places!

There is no symmetry breaking. What type of order causes this precise quantization?

## **TOPOLOGICAL ORDER!**



## Topological quantum matter

#### "Symmetry-protected topological matter"

Unique groundstate protected by a topological invariant (Chern number, Z<sub>2</sub>-index,...), ordinary electron excitations, short-range quantum entanglement:

### integer quantum Hall effect,

topological insulators, Chern insulators, topological superconductors,...

#### "Topologically ordered matter" (proper)

Groundstate degeneracies on highergenus manifolds, fractionalized excitations, long-range quantum entanglement:

fractional quantum Hall effect, quantum spin liquids,...



## Topological quantum matter

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Chern insulators, topological superconductors,...



## The topological insulators

provide the best "port of entry" to the study of topological quantum matter. This is how we shall go about it in the course! "An **electrical insulator** is a material whose internal <u>electric charges</u> do not flow freely... examples include glass, paper, Teflon, plastics, ceramic materials,... "



Ordinary insulators have nothing to do with topology, but *topological insulators* do

## 2D topological insulators...

taking off from the quantum Hall effect

### quantum Hall effect

B

"skipping current"

quantization

chiral edge states

no channel for backscattering

perfect conductance along the edge

### quantum Hall effect



Is this kind of physics possible without a magnetic field?



Well..., at least one doesn't need a net magnetic field... PRL, 1988

Duncan Haldane



In fact, one can do away with the magnetic field altogether! PRLs, 2005

Charlie Kane

Gene Mele

## To see how this is possible, consider a Gedanken experiment...

Bernevig & Zhang, PRL (2006)

To see how this is possible, consider a Gedanken experiment...



spin-orbit interaction  $(\boldsymbol{E} \times \boldsymbol{k}) \cdot \boldsymbol{\sigma} = E \sigma^z (k_y x - k_x y)$ 



## To see how this is possible, consider a Gedanken experiment...



spin-orbit interaction  $(\boldsymbol{E} \times \boldsymbol{k}) \cdot \boldsymbol{\sigma} = E\sigma^z (k_y x - k_x y)$ 



compare with an integer quantum Hall system

Lorentz force

$$\boldsymbol{A} \cdot \boldsymbol{k} \sim eB(k_y x - k_x y)$$

## **2D** topological insulator

Two copies of a quantum Hall system, bulk insulator with helical edge states



## How does Nature do it? Also by spin-orbit interactions!







## Experimental realizations...



First proposed by Kane and Mele for graphene (2005) C. L. Kane and E. J. Mele, PRL **95**, 226801 (2005)



Bernevig et al. proposal for HgTe quantum wells (2006) B. A. Bernevig, T. A. Hughes, and S. C. Zhang, Science **314**, 1757 (2006) Experimental observation by König *et al.* (2007)

M. König *et al.*, Science **318**, 766 (2007)



$$G = \frac{2e^2}{h}$$



Warm up... Gauss-Bonnet theorem

 $\frac{1}{2\pi} \int_{\mathcal{M}} K dS = 2(1-g)$ 



#### Pierre Ossian Bonnet, 1819-1892

Ancien élève de l'École polytechnique, ingénieur des Ponts et Chaussées, Bonnet préféra l'enseignement et la recherche. Répétiteur puis examinateur à l'École Polytechnique

Warm up... Gauss-Bonnet theorem



$$\frac{1}{2\pi} \int_{\mathcal{M}} K dS = 2(1-g)$$

$$g = 0 \qquad \qquad g = 1 \qquad \qquad g = 2 \qquad \qquad g = 3$$



### "Chern theorem"

## Generalized Gauss-Bonnet theorem





Shiing-Shen Chern, 1911-2004

$$\frac{1}{2\pi} \int_{\mathcal{M}} \mathcal{F} \cdot d\mathbf{S} = \mathbf{C}$$



$$\frac{1}{2\pi} \int_{\mathcal{M}} \mathcal{F} \cdot d\mathbf{S} = \mathbf{C}$$

$$\psi_{n,\boldsymbol{k}}(\boldsymbol{r}) = \boldsymbol{u_{n,\boldsymbol{k}}(\boldsymbol{r})} \mathrm{e}^{i \boldsymbol{k} \cdot \boldsymbol{r}}$$

Berry curvature Berry, RSPSA (1984)

$$\boldsymbol{\mathcal{F}}_{\boldsymbol{n}}(\boldsymbol{k}) = -i\nabla_{\boldsymbol{k}} \times \langle u_{\boldsymbol{n},\boldsymbol{k}} \mid \nabla_{\boldsymbol{k}} \mid u_{\boldsymbol{n},\boldsymbol{k}} \rangle$$



Michael Berry

$$\frac{1}{2\pi} \int_{\mathcal{M}} \mathcal{F} \cdot d\mathbf{S} = \mathbf{C}$$

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$$\frac{1}{2\pi} \sum_{n=1}^{\sharp \text{ filled bands}} \int_{\mathcal{M}=\mathrm{BZ}} \mathcal{F}_{n,z}(\mathbf{k}) \, dk_x \, dk_y = C$$

$$\psi_{n,\boldsymbol{k}}(\boldsymbol{r}) = \boldsymbol{u}_{n,\boldsymbol{k}}(\boldsymbol{r}) \mathrm{e}^{i\boldsymbol{k}\cdot\boldsymbol{r}}$$

odd under time reversal

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$$\boldsymbol{\mathcal{F}}_{n}(\boldsymbol{k}) = -i\nabla_{\boldsymbol{k}} \times \langle u_{n,\boldsymbol{k}} \mid \nabla_{\boldsymbol{k}} \mid u_{n,\boldsymbol{k}} \rangle$$

$$\frac{1}{2\pi} \sum_{n=1}^{\sharp \text{ filled bands}} \int_{\mathcal{M}=BZ} \mathcal{F}_{n,z}(\mathbf{k}) \, dk_x \, dk_y = C$$
vanishing
Chern number

broken time-reversal symmetry (like in the quantum Hall effect)

$$\psi_{n,\boldsymbol{k}}(\boldsymbol{r}) = \boldsymbol{u_{n,\boldsymbol{k}}}(\boldsymbol{r}) \mathrm{e}^{i\boldsymbol{k}\cdot\boldsymbol{r}}$$

odd under time reversal

Berry curvature Berry, RSPSA (1984)

$$\boldsymbol{\mathcal{F}}_{\boldsymbol{n}}(\boldsymbol{k}) = -i\nabla_{\boldsymbol{k}} \times \langle u_{n,\boldsymbol{k}} \mid \nabla_{\boldsymbol{k}} \mid u_{n,\boldsymbol{k}} \rangle$$



$$\frac{1}{2\pi} \sum_{n=1}^{\sharp \text{ filled bands}} \int_{\mathcal{M}=\mathrm{BZ}} \mathcal{F}_{n,z}(\mathbf{k}) \, dk_x \, dk_y = C$$

integer quantum Hall effect: "TKNN invariant"

Thouless et al., PRL (1982)

David Thouless

What if we identify timereversed points in the BZ?

## Electron wave function in a crystal

$$\psi_{n,oldsymbol{k}}(oldsymbol{r}) = oldsymbol{u}_{n,oldsymbol{k}}(oldsymbol{r}) \mathrm{e}^{ioldsymbol{k}\cdotoldsymbol{r}}$$

odd under time reversal

Berry curvature Berry, RSPSA (1984)

$$\boldsymbol{\mathcal{F}}_{\boldsymbol{n}}(\boldsymbol{k}) = -i\nabla_{\boldsymbol{k}} \times \langle u_{\boldsymbol{n},\boldsymbol{k}} \mid \nabla_{\boldsymbol{k}} \mid u_{\boldsymbol{n},\boldsymbol{k}} \rangle$$

$$\frac{1}{2\pi} \sum_{n=1}^{\sharp \text{ filled bands}} \int_{\mathcal{M}=\mathrm{BZ}} \mathcal{F}_{n,z}(\mathbf{k}) \, dk_x \, dk_y = C$$

$$\frac{1}{2\pi} \sum_{n} \int \boldsymbol{\mathcal{F}}_{n}(\boldsymbol{k}) \cdot d\boldsymbol{k} = C$$

# The *parity* of the Chern number is unique!

Kane & Mele, PRL (2005)



 $C = \begin{cases} 0 \mod 2 & \text{ordinary insulator} \\ 1 \mod 2 & \text{topological insulator} \end{cases}$ 

"Z<sub>2</sub> topological invariant", counts the number of Kramers pairs at the edge of the topological insulator (bulk-boundary correspondence)

## Why all the hoopla?

#### "Who here is familiar with the concept of Topological Insulators?"

Cotte Ha Te Cotte Ha Te Cotte Ha Te Cotte Haniltonian: Te gervalves: E

TV-Series: "The Big Bang Theory", https://www.youtube.com/watch?v=HBuLMrzgbgM

BREAKTHROUGH OF THE YEAR ELECTRONS TAKE A NEW SPIN



## Why all the hoopla?

• Future electronics/spintronics:

Toward dissipationless spin transport in semiconductors

• New physics in hybrid structures: Majorana fermions, magnetic monopoles, "dyons",...



 An inroad to the general study of TOPOLOGICAL QUANTUM MATTER "Periodic tables" for symmetry-protected topological matter:

AZ class $\backslash d$	0	1	2	3
Α	$\mathbb{Z}$	0	$\mathbb{Z}$	0
AIII	0	$\mathbb{Z}$	0	$\mathbb{Z}$
AI	$\mathbb{Z}$	0	0	0
BDI	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0
D	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0
DIII	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$
AII	$2\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$
CII	0	$2\mathbb{Z}$	0	$\mathbb{Z}_2$
$\mathbf{C}$	0	0	$2\mathbb{Z}$	0
CI	0	0	0	$2\mathbb{Z}$

2D topological insulator



Refl	ection	Class	$C_q$ or $R_q$	d = 0	d = 1	d = 2	d = 3	d = 4	d = 5	d = 6	d = 7
	R	Α	$C_1$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$
j	$R^+$	AIII	$C_0$	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0
j	$R^{-}$	AIII	$C_1$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$
		AI	$R_1$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$
		BDI	$R_2$	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$	0
		D	$R_3$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$
$R^+$	$, R^{++}$	DIII	$R_4$	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0
		AII	$R_5$	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0
		CII	$R_6$	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0
		$\mathbf{C}$	$R_7$	0	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$
		$\mathbf{CI}$	$R_0$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$
	AI	$R_7$	0	0	0	$\mathbb{Z}$	0	" $\mathbb{Z}_2$ "	$\mathbb{Z}_2$	$\mathbb{Z}$	
		BDI	$R_0$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$	0	" $\mathbb{Z}_2$ "	$\mathbb{Z}_2$
		D	$R_1$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$	0	" $\mathbb{Z}_2$ "
$R^{-}$	$, R^{}$	DIII	$R_2$	" $\mathbb{Z}_2$ "	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$	0
		AII	$R_3$	0	" $\mathbb{Z}_2$ "	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$
		CII	$R_4$	$\mathbb{Z}$	0	" $\mathbb{Z}_2$ "	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0
		$\mathbf{C}$	$R_5$	0	$\mathbb{Z}$	0	" $\mathbb{Z}_2$ "	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0
	CI	$R_6$	0	0	$\mathbb{Z}$	0	" $\mathbb{Z}_2$ "	$\mathbb{Z}_2$	$\mathbb{Z}$	0	
I	2+-	BDI	$R_1$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$
F	2-+	DIII	$R_3$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0	0	$\mathbb{Z}$
F	2+-	CII	$R_5$	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$	0	0
I	?−+	CI	$R_7$	0	0	0	$\mathbb{Z}$	0	$\mathbb{Z}_2$	$\mathbb{Z}_2$	$\mathbb{Z}$
h	2-+	BDI, CII	$C_1$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$
F	2+-	DIII, CI	$C_1$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$	0	$\mathbb{Z}$

Now for some serious work on the blackboard!